NIMROD Two-Fluid, Algorithm and Calculations

D. C. Barnes *NIMROD-CEMM Meeting* November 13-14, 2004





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Outline

- Two-fluid MHD
- Time-implicit method a stability theorem
- NIMROD-2F implementation and dispersion tests
- FRC application
- Conclusions





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Single/Two-Fluid "MHD" (1F/2F)

$$Mn\frac{\partial \mathbf{u}}{\partial t} = (\nabla \times \mathbf{B})/\mu_0 \times \mathbf{B} - \nabla(P_e + P_i) - Mn\mathbf{u} \cdot \nabla \mathbf{u}$$
$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times \left\{ \mathbf{u} \times \mathbf{B} - \frac{1}{e} \left[\frac{(\nabla \times \mathbf{B}) \times \mathbf{B}/\mu_0 - \nabla P_e}{n} \right] - m_e \left(\frac{\partial}{\partial t} + \mathbf{u} \cdot \nabla - \frac{\nabla \times \mathbf{B}/\mu_0}{en} \cdot \nabla \right) \left(\mathbf{u} - \frac{\nabla \times \mathbf{B}/\mu_0}{en} \right) \right] \right\}$$
$$\frac{\partial n}{\partial t} = -\nabla \cdot n\mathbf{u}$$
$$\frac{\partial P_e}{\partial t} = -\mathbf{u} \cdot \nabla P_e - \Gamma_e P_e \nabla \cdot \mathbf{u} + \frac{\nabla \times \mathbf{B}/\mu_0}{en} \cdot \nabla P_e + \Gamma_e P_e \nabla \cdot \frac{\nabla \times \mathbf{B}/\mu_0}{en}$$
$$\frac{\partial P_i}{\partial t} = -\mathbf{u} \cdot \nabla P_e - \Gamma_e P_e \nabla \cdot \mathbf{u} + \frac{\nabla \times \mathbf{B}/\mu_0}{en} \cdot \nabla P_e + \Gamma_e P_e \nabla \cdot \frac{\nabla \times \mathbf{B}/\mu_0}{en}$$

$$\frac{\partial P_i}{\partial t} = -\mathbf{u} \bullet \nabla P_i - \Gamma_i P_i \nabla \bullet \mathbf{u}$$





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Can Rewrite Hall Term

$$Mn\frac{\partial \mathbf{u}}{\partial t} = (\nabla \times \mathbf{B})/\mu_0 \times \mathbf{B} - \nabla (P_e + P_i) - Mn\mathbf{u} \cdot \nabla \mathbf{u}$$
$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times \left\{ \mathbf{u} \times \mathbf{B} - \frac{1}{e} \left[M \frac{\partial \mathbf{u}}{\partial t} + M\mathbf{u} \cdot \nabla \mathbf{u} + \frac{\nabla P_i}{n} \right] - m_e \left(\frac{\partial}{\partial t} + \mathbf{u} \cdot \nabla - \frac{\nabla \times \mathbf{B}/\mu_0}{en} \cdot \nabla \right) \left(\mathbf{u} - \frac{\nabla \times \mathbf{B}/\mu_0}{en} \right) \right\}$$
$$\frac{\partial n}{\partial t} = -\nabla \cdot n\mathbf{u}$$
$$\frac{\partial P_e}{\partial t} = -\mathbf{u} \cdot \nabla P_e - \Gamma_e P_e \nabla \cdot \mathbf{u} + \frac{\nabla \times \mathbf{B}/\mu_0}{en} \cdot \nabla P_e + \Gamma_e P_e \nabla \cdot \frac{\nabla \times \mathbf{B}/\mu_0}{en}$$
$$\frac{\partial P_i}{\partial t} = \nabla \nabla P_e - \nabla \nabla P_e \nabla \nabla \mathbf{u}$$

$$\frac{\partial P_i}{\partial t} = -\mathbf{u} \bullet \nabla P_i - \Gamma_i P_i \nabla \bullet \mathbf{u}$$





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1F→2F Plasma Features

- Waves (uniform, unbounded plasma)
 - Whistler
 - Kinetic Alfven
 - Low frequency electrostatic
- Drift waves
- ω^* stabilization of MHD modes
- Electron-Ion decoupling





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Physics of HMHD

- Full (warm 2 fluid) dispersion relation has 3 waves (cubic)
- This was given by Stringer (1963) and is also discussed by Swanson

$$w = \omega^{2}$$

$$w^{3} - A w^{2} + Bw - C = 0$$

$$A = k_{\parallel}^{2} + (1 + \hat{\beta})k^{2} + H^{2}k_{\parallel}^{2}k^{2}$$

$$B = k_{\parallel}^{2}k^{2}(1 + 2\hat{\beta} + \hat{\beta}H^{2}k^{2})$$

$$C = \hat{\beta}k^{2}k_{\parallel}^{4}$$





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Modest Parameters Give Extreme Stiffness



E.G. $\hat{\beta} = 10^{-8}$, $k_{\parallel}/k_{\perp} = 10^{-4}$, $d_i/L_x = 0.1 \Rightarrow 12$ orders of magnitude in ω





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Modest Parameters Give Extreme Stiffness



E.G. $\hat{\beta} = 10^{-2}$, $k_{\parallel}/k_{\perp} = 10^{-4}$, $d_i/L_x = 0.1 \Rightarrow 12$ orders of magnitude in ω





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Extension of 1F Method to 2F Challenging

- 1F of form $\dot{\mathbf{B}} = \dot{\mathbf{B}}(\mathbf{u}), \dot{\mathbf{u}} = \dot{\mathbf{u}}(\mathbf{B}) \Rightarrow$ $\mathbf{B}^{n+1} - \mathbf{B}^n = \Delta t \dot{\mathbf{B}}\left(\frac{\mathbf{u}^{n+1} + \mathbf{u}^n}{2}\right), \mathbf{u}^{n+1} - \mathbf{u}^n = \Delta t \dot{\mathbf{u}}\left(\frac{\mathbf{B}^{n+1} + \mathbf{B}^n}{2}\right) \Rightarrow$ $\mathbf{B}^{n+1/2} - \mathbf{B}^{n-1/2} = \Delta t \dot{\mathbf{B}}(\mathbf{u}^n), L(\mathbf{u}^{n+1} - \mathbf{u}^n) = \Delta t \dot{\mathbf{u}}(\mathbf{B}^{n+1/2})$
- 2F is not!
- Challenge is to obtain method for low dissipation
 - Real frequency physical modes should give real frequency numerical modes





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A Useful Stability Theorem

- To simulate low dissipation cases, need "spectral fidelity"
 - If physical system has only real frequency modes, numerical system will have only real frequency modes (or controlled damping would also work)
- One case of sufficient conditions

$$\dot{U} = F(U, \nabla) \Rightarrow \frac{U^{n+1} - U^n}{\Delta t} = F(U^{n+\sigma}, D)$$

$$U^{n+\sigma} \text{ is linear combination of } U^{n+1}, U^n, U^{n-1}, \dots, \text{ centered at}$$

$$\sigma \ge 1/2, D \text{ is difference operator } \ni :iD \text{ is Hermitian}$$





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Some Possible Schemes

$$U^{n+1/2} = \frac{U^{n+1} + U^{n}}{2} : (1L) \Longrightarrow Z - 1 = -i\widetilde{\omega}\Delta t \frac{Z + 1}{2}$$
$$U^{n+1/2} = \lambda \frac{U^{n+1} + U^{n}}{2} + (1 - \lambda) \frac{3U^{n+1} + U^{n-1}}{4} : (2L) \Longrightarrow$$
$$[i\widetilde{\omega}\Delta t (3 - \lambda) + 4]Z^{2} + (2i\widetilde{\omega}\Delta t - 4)Z + i\widetilde{\omega}\Delta t (1 - \lambda) = 0$$

Both have vanishing damping for low frequencies. 1L has no damping for high frequencies as well, while 2L gives strongest damping when $\lambda = \frac{3}{4} (|Z| = 1/3)$.





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Using NIMROD Machine for This Form

- Use Predictor/Corrector method
- Advance all equations (slave) except momentum (master) with trial u*~uⁿ⁺¹
- Error in momentum gives correction
- Linear operator from linearizing change in slave variables in change in u
- P/C required because not possible to incorporate exact linear change of slave variables
 - $\nabla \nabla \neq \nabla^2$ (use compact stencil for 2nd order operator)
 - Slave equations contain operator inversion





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Some Numerical Details

- Care in integration by parts because of different BC
- Correct BC is to require only $u_n = 0$

$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times \mathbf{u} \times \mathbf{B} - \frac{M}{e} \nabla \times \frac{\partial \mathbf{u}}{\partial t} + \dots$$
$$\int d\mathbf{r} \boldsymbol{\xi}^* \bullet \frac{\partial \mathbf{B}}{\partial t} \neq \int d\mathbf{r} \nabla \times \boldsymbol{\xi}^* \bullet \left(\mathbf{u} \times \mathbf{B} - \frac{M}{e} \frac{\partial \mathbf{u}}{\partial t} \right)$$





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Waves in a Box Tests







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History for KAW Calculation







Profile for KAW Calculation



Numerical Results show good Dispersion



NILINGOD Non-tdon, develop Open Discussion Project

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Convergence Acceleration of P/C

- Can consider present scheme as preconditioner
- Apply GMRES to iteration
- Some success with partial GMRES (1 lag evel)
- For larger problems, can use direct solve per block to get Schwarz preconditioner





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Extension of Linear Solve

- Instead of master/slave scheme, solve all equations simultaneously (8 unknowns/ node instead of 3)
- No P/C required
- NIMROD machinery supports this (almost)
- Almost same as scheme(s) of Chacón and Glasser





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NIMROD Initialized with FRC Equilibrium





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Conclusion

- NIMROD 2F well under way
- Waves in box test (nearly) passed
- Early version into CVS (update soon)
- All thermoelectric terms in (nonlinear?)
 - FRC application beginning
- Need better separatrix conditions
- Algorithm improvements continue in parallel with applications
 - Finite m_e
- GMRES, Block direct + Schwarz
- Alternative: fully couple all equations





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