Adaptive Grid Generation for Magnetically Confined Plasmas

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Magnetic Reconnection, Final State

Magnetic Flux



Current Density



$$A = 1$$

$$M = 1/2$$

$$\eta = 10^{-4}$$

$$\mu = 10^{-4}$$

$$\epsilon = 10^{-4}$$

$$dt = 20$$

$$nx = 6$$

$$ny = 16$$

$$np = 12$$

$$nproc = 16$$

$$cpu = 3.5$$
 hr

Stream Function



Vorticity



The Need for a 3D Adaptive Field-Aligned Grid

- > An essential feature of magnetic confinement is very strong anisotropy, $\chi_{\parallel} >> \chi_{\perp}$.
- ➤ The most unstable modes are those with $k_{\parallel} \ll 1/R < 1/a \ll k_{\perp}$.
- The most effective numerical approach to these problems is a field-aligned grid packed in the neighborhood of singular surfaces and magnetic islands. NIMROD.
- Long-time evolution of helical instabilities requires that the packed grid follow the moving perturbations into 3D.
- Multidimensional oblique rectangular AMR grid is larger than necessary and does not resolve anisotropy.
- Novel algorithms must be developed to allow alignment of the grid with the dominant magnetic field and automatic grid packing normal to this field.
- Such methods must allow for regions of magnetic islands and stochasticity.

Adaptive Mesh Refinement vs. Harmonic Grid Generation

Adaptive Mesh Refinement

- 1. Coarse and fine patches of rectangular grid.
- 2. Complex data structures.
- 3. Oblique to magnetic field.
- 4. Static regrid.
- 5. Explicit time step; implicit a research problem.
- 6. Berger, Collela, Samtaney, Jardin

Harmonic Grid Generation

- 1. Harmonic mapping of rectangular grid onto curvilinear grid.
- 2. Logically rectangular
- 3. Aligned with magnetic field.
- 4. Static or dynamic regrid.
- 5. Explicit or implicit time step.
- 6. Liseikin, Winslow, Dvinsky, Brackbill

SEL Code Features

- Spectral elements: exponential convergence of spatial truncation error + adaptable grid + parallelization.
 - George Em Karniadakis and Spencer J. Sherwin, "Spectral/*hp* Element Methods for CFD," Oxford, 1999.
 - Ronald D. Henderson, "Adaptive spectral element methods for turbulence and transition," in *High-Order Methods for Computational Physics*, T.J. Barth & H. Deconinck (Eds.), Springer, 1999.
- Time step: fully implicit, 2nd-order accurate, Newton-Krylov iteration, static condensation preconditioning.
- Highly efficient massively parallel operation with MPI and PETSc.
- ➢ Flux-source form: simple, general problem setup.

Spatial Discretization

Flux-Source Form of Equations

$$\frac{\partial u^i}{\partial t} + \nabla \cdot \mathbf{F}^i = S^i$$

$$\mathbf{F}^i = \mathbf{F}^i(t, \mathbf{x}, u^j, \nabla u^j)$$

$$S^i = S^i(t, \mathbf{x}, u^j, \nabla u^j)$$

Galerkin Expansion

$$u^{i}(t, \mathbf{x}) \approx \sum_{j=0}^{n} u_{j}^{i}(t) \alpha_{j}(\mathbf{x})$$

Weak Form of Equations

$$(\alpha_i, \alpha_j)\dot{u}_j^k = \int_{\Omega} d\mathbf{x} \left(S^k \alpha_i + \mathbf{F}^k \cdot \nabla \alpha_i \right) - \int_{\partial \Omega} d\mathbf{x} \alpha_i \mathbf{F}^k \cdot \hat{\mathbf{n}}$$

Fully Implicit Newton-Krylov Time Step

 $M\dot{u} = r$

$$\mathbf{M}\left(\frac{\mathbf{u}^{+}-\mathbf{u}^{-}}{h}\right) = \theta \mathbf{r}^{+} + (1-\theta)\mathbf{r}^{-}$$
$$\mathbf{R}\left(\mathbf{u}^{+}\right) \equiv \mathbf{M}\left(\mathbf{u}^{+}-\mathbf{u}^{-}\right) - h\left[\theta \mathbf{r}^{+} + (1-\theta)\mathbf{r}^{-}\right] = 0$$
$$\mathbf{J} \equiv \mathbf{M} - h\theta \left\{\frac{\partial r_{i}^{+}}{\partial u_{j}^{+}}\right\}$$

 $\mathbf{R} + \mathsf{J}\delta \mathbf{u}^{+} = \mathbf{0}, \quad \delta \mathbf{u}^{+} = -\mathsf{J}^{-1}\mathbf{R}\left(\mathbf{u}^{+}\right), \quad \mathbf{u}^{+} \to \mathbf{u}^{+} + \delta \mathbf{u}^{+}$

- Nonlinear Newton-Krylov iteration.
- Elliptic equations: $\mathbf{M} = 0$.
- Static condensation, fully parallel.
- PETSc: GMRES with Schwarz ILU, overlap of 3, fill-in of 5.

Static Condensation

$$\mathbf{L}\mathbf{u} = \mathbf{r} \tag{1}$$

Partition: (1) element edges: (2) element interiors

$$\mathbf{u} = \begin{pmatrix} \mathbf{u}_1 \\ \mathbf{u}_2 \end{pmatrix}, \quad \mathbf{r} = \begin{pmatrix} \mathbf{r}_1 \\ \mathbf{r}_2 \end{pmatrix}, \quad \mathbf{L} = \begin{pmatrix} \mathbf{L}_{11} & \mathbf{L}_{12} \\ \mathbf{L}_{21} & \mathbf{L}_{22} \end{pmatrix}$$
(2)

$$\mathbf{L}_{11}\mathbf{u}_1 + \mathbf{L}_{12}\mathbf{u}_2 = \mathbf{r}_1 \tag{3}$$

$$\mathbf{L}_{22}\mathbf{u}_2 = \mathbf{r}_2 - \mathbf{L}_{21}\mathbf{u}_1 \tag{4}$$

$$\mathbf{\bar{L}}_{11} \equiv \mathbf{L}_{11} - \mathbf{L}_{12} \mathbf{L}_{22}^{-1} \mathbf{L}_{21}
\bar{\mathbf{r}}_{1} \equiv \mathbf{r}_{1} - \mathbf{L}_{12} \mathbf{L}_{22}^{-1} \mathbf{r}_{2}$$
(5)

$$\bar{\mathbf{L}}_{11}\mathbf{u}_1 = \bar{\mathbf{r}}_1 \tag{6}$$

Equation (4) solved by local LU factorization and back substitution.
Equation (6), substantially reduced, solved by global Newton-Krylov.

Adaptive Grid Kinematics: How to Use Logical Coordinates.

$$x^{j}(\xi^{k}) = \sum_{i} x_{i}^{j} \alpha_{i}(\xi^{k}), \quad j,k = 1,2$$
$$\mathcal{J} \equiv (\hat{\mathbf{z}} \cdot \nabla \xi^{1} \times \nabla \xi^{2})^{-1} = \frac{\partial x^{1}}{\partial \xi^{1}} \frac{\partial x^{2}}{\partial \xi^{2}} - \frac{\partial x^{1}}{\partial \xi^{2}} \frac{\partial x^{2}}{\partial \xi^{1}}$$
$$\frac{\partial u^{k}}{\partial t} + \nabla \cdot \mathbf{F}^{k} = S^{k}, \quad \frac{\partial u^{k}}{\partial t} + \frac{1}{\mathcal{J}} \frac{\partial}{\partial \xi^{j}} \left(\mathcal{J} \mathbf{F}^{k} \cdot \nabla \xi^{j}\right) = S^{k}$$

$$u^{k}(t,\mathbf{x}) \approx \sum_{j=0}^{\infty} u_{j}^{k}(t)\alpha_{j}(\xi), \quad (u,v) \equiv \int_{\Omega} uv d\mathbf{x} = \int_{\Omega} uv \mathcal{J}d\xi$$

$$(\alpha_i, \alpha_j) \dot{u}_j^k = \int_{\Omega} \left(S^k \alpha_i + \mathbf{F}^k \cdot \nabla \xi^j \frac{\partial \alpha_i}{\partial \xi^j} \right) \mathcal{J} d\xi - \int_{\partial \Omega} \alpha_i \mathbf{F}^k \cdot \hat{\mathbf{n}} \mathcal{J} d\xi$$

Adaptive Grid Dynamics: How to Choose Logical coordinates.

$$\mathcal{L} \equiv \frac{1}{2} \int \left[\left(\mathbf{B} \cdot
abla \xi^j \right)^2 + \epsilon |
abla \xi^j|^2 \right] d\mathbf{x}$$

 $rac{\delta \mathcal{L}}{\delta \xi^j} = 0 \Rightarrow
abla \cdot \left(\mathbf{g} \cdot
abla \xi^j \right) = 0, \quad \mathbf{g} \equiv \mathbf{B}\mathbf{B} + \epsilon \mathbf{B}$

Beltrami equation + boundary conditions \Rightarrow logical coordinates. Alignment with magnetic field except where **B** \rightarrow 0, isotropic term dominates.

Vladimir D. Liseikin A Computational Differential Geometry Approach to Grid Generation Springer Series in Synergetics, 2003

Domains and Transformations

Used in Harmonic Grid Generation Figure by Andrei Simakov



Modified Beltrami Equation

Variational Principle

$$\mathcal{L} = \frac{1}{2} \int_{\Omega} \frac{1}{w\sqrt{g}} \mathbf{g} : \nabla \xi^i \nabla \xi^i d\mathbf{x}$$

Euler-Lagrange Equation

$$\nabla \cdot \left(\frac{1}{w\sqrt{g}}\mathbf{g} \cdot \nabla \xi^i\right) = 0$$

Expressed in Logical Coordinates (Chacon)

$$\frac{1}{\mathcal{J}}\frac{\partial}{\partial\xi^j}\left(\frac{\mathcal{J}}{w\sqrt{g}}g^{kl}\frac{\partial\xi^i}{\partial x^k}\frac{\partial\xi^j}{\partial x^l}\right) = 0, \quad \frac{\partial\xi^i}{\partial x^j} \to \frac{\partial x^i}{\partial\xi^j}$$

Metric Tensor Used for Alignment

 $\mathbf{g} = \mathbf{B}_1 \mathbf{B}_1 + \mathbf{B}_2 \mathbf{B}_2 + \epsilon \mathbf{I}, \quad \mathbf{B}_1 \equiv \hat{\mathbf{z}} \times \nabla \psi, \quad \mathbf{B}_2 = k \hat{\mathbf{z}} \times \mathbf{B}^1$

Boundary Conditions

At each boundary, one of the x^i is held fixed at the edge of the domain and $\nabla x^i \cdot \nabla x^j = 0$ for $i \neq j$.

Pure Alignment, w = 1



Magnetic flux is multiply connected; grid is simply connected. Crossings occur where $\mathbf{B} = \mathbf{z} \ \mathfrak{D} \ \forall \ \psi$ is small. Alignment Error: 0.012 max, 0.0055 RMS.

Pure Adaptation, g = I



Weight function derived from log of spatial truncation error. Grid density almost perfectly reproduces weight function, including absolute magnitudes.

Alignment + Adaptation



This is a compromise. Neither the alignment nor the adaptation is quite as good as for the pure cases. Alignment Error: 0.036 max, 0.018 RMS.

Conclusions

- Pure alignment works very well, giving a simply-connected grid which is well-aligned except where **B** is small.
- Pure adaptation works very well, concentrating the grid in regions of large spatial truncation error.
- > Alignment + adaptation works reasonably well, a good compromise.
- ➤ Imperfections are compensated by high-order spectral elements.

Next

➤ Use adapted grid for computations.

Vladimir D. Liseikin Esteemed Colleague and Good Friend

