Two-Fluid NSTX Steady-States with Flow

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- We are using M3D- $C¹$ to calculate axisymmetric toroidal steady-states of a comprehensive two-fluid model.
- These steady-states are steady on all timescales self-consistently include two-fluid effects, gyroviscosity, flow, and anisotropic transport.
- In particular, we would like to understand the effects of two-fluid terms and gyroviscosity on the steady-states.
- These steady-states may be used as accurate equilibria for three-dimensional stability studies.

Physical Model

$$
\frac{\partial n}{\partial t} + \nabla \cdot (n\mathbf{u}) = \Sigma
$$
\n
$$
n\left(\frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u}\right) = \mathbf{J} \times \mathbf{B} - (\nabla p + \nabla \cdot \Pi) - \mathbf{u}\Sigma
$$
\n
$$
\frac{1}{\Gamma - 1} \left[\frac{\partial p}{\partial t} + \nabla \cdot (p\mathbf{u})\right] = -p\nabla \cdot \mathbf{u} + \frac{d_i}{\Gamma - 1} \frac{\mathbf{J}}{n} \cdot \left(\nabla p_e - \Gamma \frac{p_e}{n} \nabla n\right)
$$
\n
$$
-\nabla \cdot \mathbf{q} - \Pi : \nabla \mathbf{u} + d_i \Pi_e : \nabla \frac{\mathbf{J}}{n} + \frac{1}{2} u^2 \Sigma
$$
\n
$$
\frac{\partial \mathbf{B}}{\partial t} = -\nabla \times \mathbf{E}
$$
\n
$$
\mathbf{E} + \mathbf{u} \times \mathbf{B} = \eta \mathbf{J} + \frac{d_i}{n} (\mathbf{J} \times \mathbf{B} - \nabla p_e - \nabla \cdot \Pi_e)
$$
\n
$$
\mathbf{J} = \nabla \times \mathbf{B} \qquad \Sigma = \sigma + D\nabla^2 n
$$
\n
$$
\Pi = \Pi_o + \Pi_\wedge + \Pi_\parallel \qquad \Pi_e = \lambda n \nabla \mathbf{J}
$$
\n
$$
\mathbf{q} = -\kappa_o \nabla T - \kappa_\parallel \mathbf{B} \mathbf{B} \cdot \nabla T
$$

Digression: Interpretation of Gyroviscosity

$$
\nabla \cdot \boldsymbol{\mathsf{\Pi}}_\wedge \thickapprox -mn\mathbf{u}_\ast \cdot \nabla \mathbf{u}
$$

where

$$
\mathbf{u}_*=-\frac{1}{mn}\nabla\times\left(\frac{p}{B^2}\mathbf{B}\right)
$$

• Braginskii shows that difference between fluid drift velocity and average gyro-center drift velocity is

$$
\mathbf{u} - \langle \mathbf{v}_c \rangle = -\frac{1}{mn} \nabla \times \left(\frac{p}{B^2} \mathbf{B}\right) = \mathbf{u}_*
$$

• Therefore,

$$
\mathit{mnu}\cdot\nabla u+\nabla\cdot\Pi_\wedge\approx\mathit{mn}(u-u_*)\cdot\nabla u=\mathit{mn}\left\langle v_c\right\rangle\cdot\nabla u
$$

• "Gyroviscous cancellation" cancels fluid drifts from inertia.

- The simulation is initialized with ^a solution to the Grad-Shafranov equation.
- A loop voltage is applied by changing the flux at the boundary of the simulation domain at a constant rate ψ \cup \equiv $=V_{L}/2\pi$.
- A localized density source is included to offset diffusive flux out of the simulation domain.
- The simulation is run until ^a steady state in all hydrodynamic quantities is reached (may not be stationary).
- $\eta = \eta_0 T^{-3/2}$. The vacuum region is simply a low temperature region outside the separatrix.

Radial Flows

- Pfirsch-Schlüter theory of radial flows is well satisfied.
- Radial flows are proportional to η and V_L .

Toroidal Flow: Edge Flows

Toroidal Flow: Core Rotation

Poloidal Flow

Parallel Viscosity

- Collisional parallel viscosity damps poloidal flows.
- Does not damp toroidal flows.

Conclusions

- We have been able to obtain self-consistent steady-states of the extended-MHD equations for realistic ^plasma configurations with free boundaries.
- The flows observed in the steady-states are in good agreement with Pfirsch-Schlüter theory.
- Strong up-down antisymmetric toroidal edge flows exist in highly resistive SOLs.
- Hall term and electron pressure gradient have little effect on steady-state.
- In highly resistive, low viscosity case, gyroviscosity may lead to persistent, large-scale oscillations.

- Better modeling of edge/SOL
	- Realistic boundary shapes/conditions
	- Pedestal modeling for H-mode
	- $-$ "Rice Scaling" for spontaneous rotation $M_A \sim \beta_N$.
- Need some model for neoclassical parallel viscosity (bootstrap current).
- Coupling to realistic transport models (TGLF?)
- 3D nonlinear simulations