

New Developments with M3D-C1

N.M. Ferraro, General Atomics

M.S. Chance, J. Chen, S.C. Jardin, PPPL

F. Delalondre, F. Zhang, RPI

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Outline

- **Resistive Wall Boundary Conditions**
- **Linear Non-Axisymmetric Response**
- **Nonlinear 3D**

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Thin-Shell Resistive Wall Boundary Conditions

$$\mathbf{B} = \nabla\psi \times \nabla\varphi - \nabla_{\perp} f' + F\nabla\varphi$$

$$F = F_0 + R^2 \nabla_{\perp}^2 f$$

- **M3D-C¹ advances ψ and F**
- **Essential (Dirichlet) condition on ψ :**

$$\frac{\partial\psi}{\partial t} = -\frac{\eta_W}{\delta} R \left(\hat{\mathbf{t}} \cdot \mathbf{B} - \hat{\mathbf{t}} \cdot \mathbf{B}^v \right) - \frac{V_L}{2\pi}$$

- **Natural (Neumann) boundary condition on F :**

$$\int dV \frac{\mu}{R^2} \frac{\partial F}{\partial t} = -\oint dS \frac{\mu}{R} \frac{\eta_W}{\delta} \left(\frac{F}{R} - \hat{\varphi} \cdot \mathbf{B}^v \right) + \int dV \nabla\varphi \cdot \nabla\mu \times \mathbf{E}$$

Vacuum Response Depends on Plasma Response Non-Locally

$$\begin{aligned}(\hat{\mathbf{n}} \cdot \mathbf{B}^v)_i &= (\hat{\mathbf{n}} \cdot \mathbf{B})_i \\(\hat{\mathbf{t}} \cdot \mathbf{B}^v)_i &= M_{ij}^t (\hat{\mathbf{n}} \cdot \mathbf{B})_j \\(\hat{\boldsymbol{\varphi}} \cdot \mathbf{B}^v)_i &= M_{ij}^\varphi (\hat{\mathbf{n}} \cdot \mathbf{B})_j\end{aligned}$$

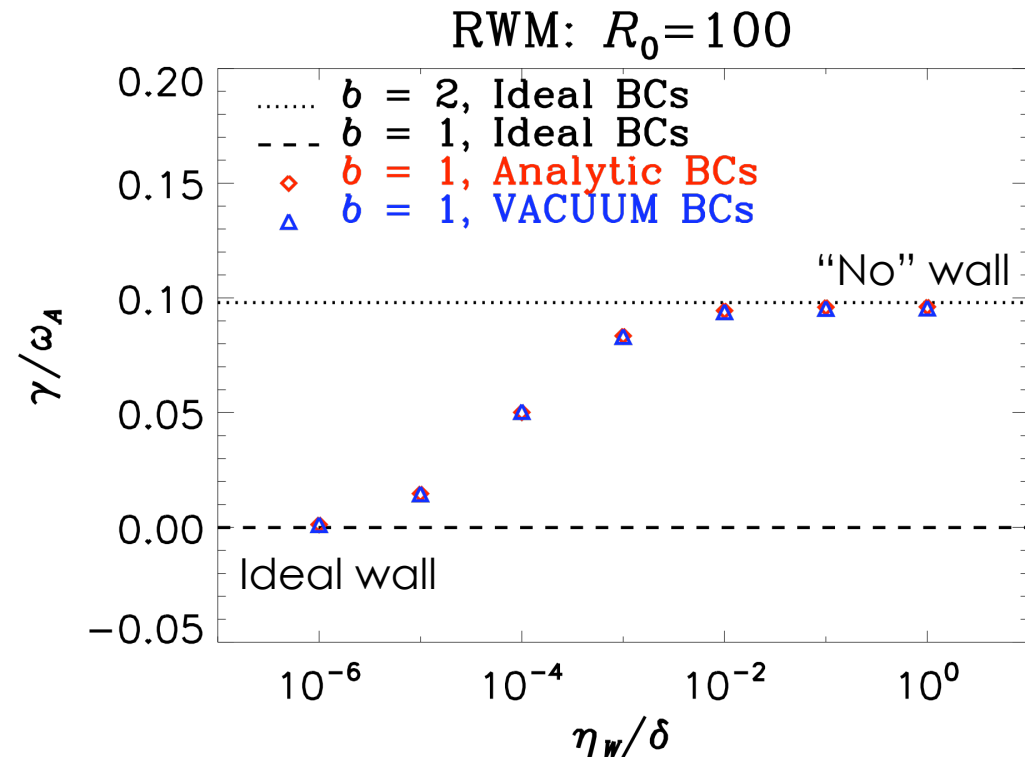
i, j range over all boundary nodes

- **VACUUM*** calculates response matrices M in arbitrary geometry
- **M is dense; all boundary nodes coupled to each other**
 - Adds communication; hurts scalability
 - Not yet supported by SCOREC libraries in parallel

* M.S. Chance, *Phys. Plasmas* **4**, 2161 (1997)

Resistive Wall Mode Test

- Equilibrium is no-wall unstable
- Stable with conducting wall at $b=1$



- Growth rate should transition from ideal-wall limit to no-wall limit as η_w/δ is increased.
- In the large-aspect limit, we know response matrix analytically

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Linear Non-axisymmetric Field Response with M3D-C1

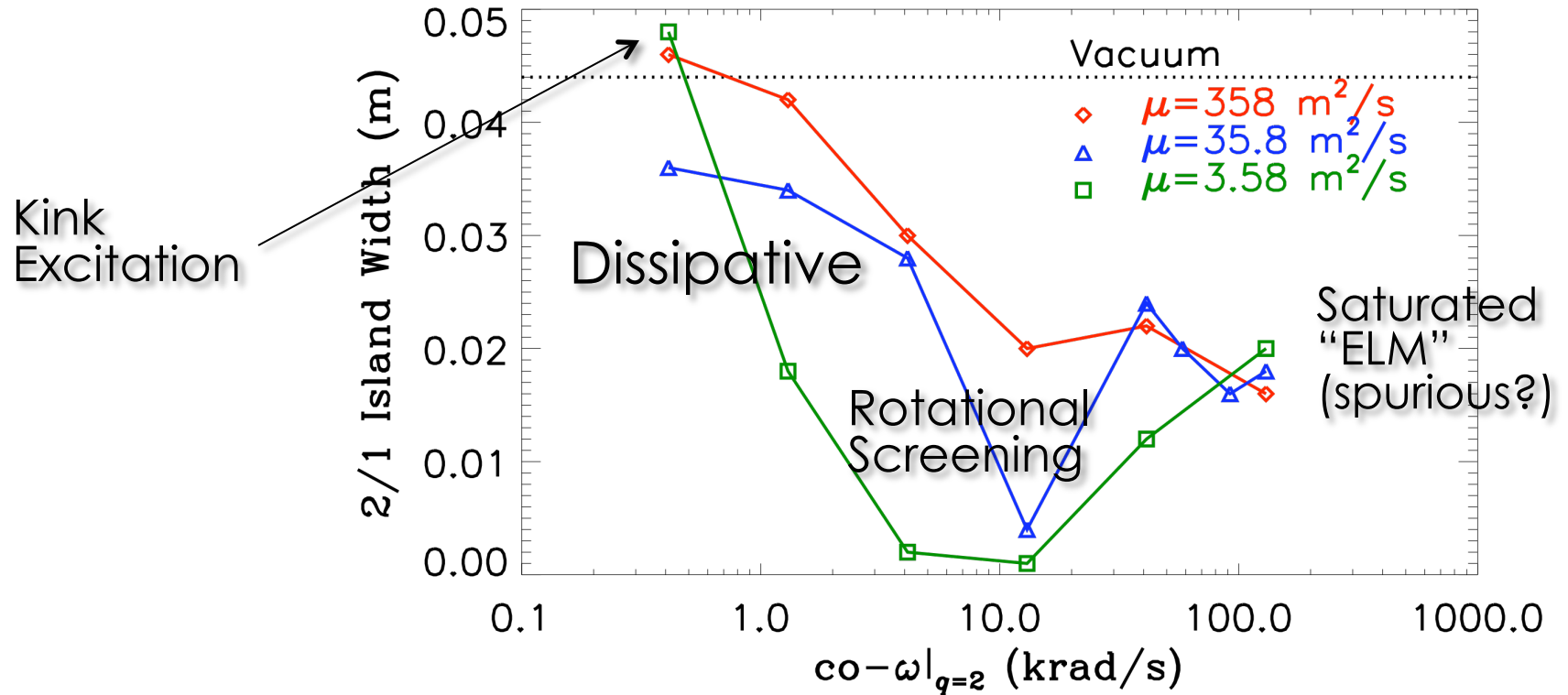
- $\mathbf{B}(t) = \mathbf{B}_0 + \mathbf{B}_1(t)$
 - \mathbf{B}_0 is the axisymmetric equilibrium field
 - $\mathbf{B}_1(0)$ is the “vacuum field” from non-axisymmetric coils (I-coils).
- **Conducting-wall boundary condition**
 - \mathbf{B} held constant in time on simulation domain boundary (approximately vacuum vessel)
- **Simulation is time-advanced until the steady-state is reached.**
- **Final \mathbf{B}_1 is applied field + plasma response.**

Response Calculations Require Stable Equilibria

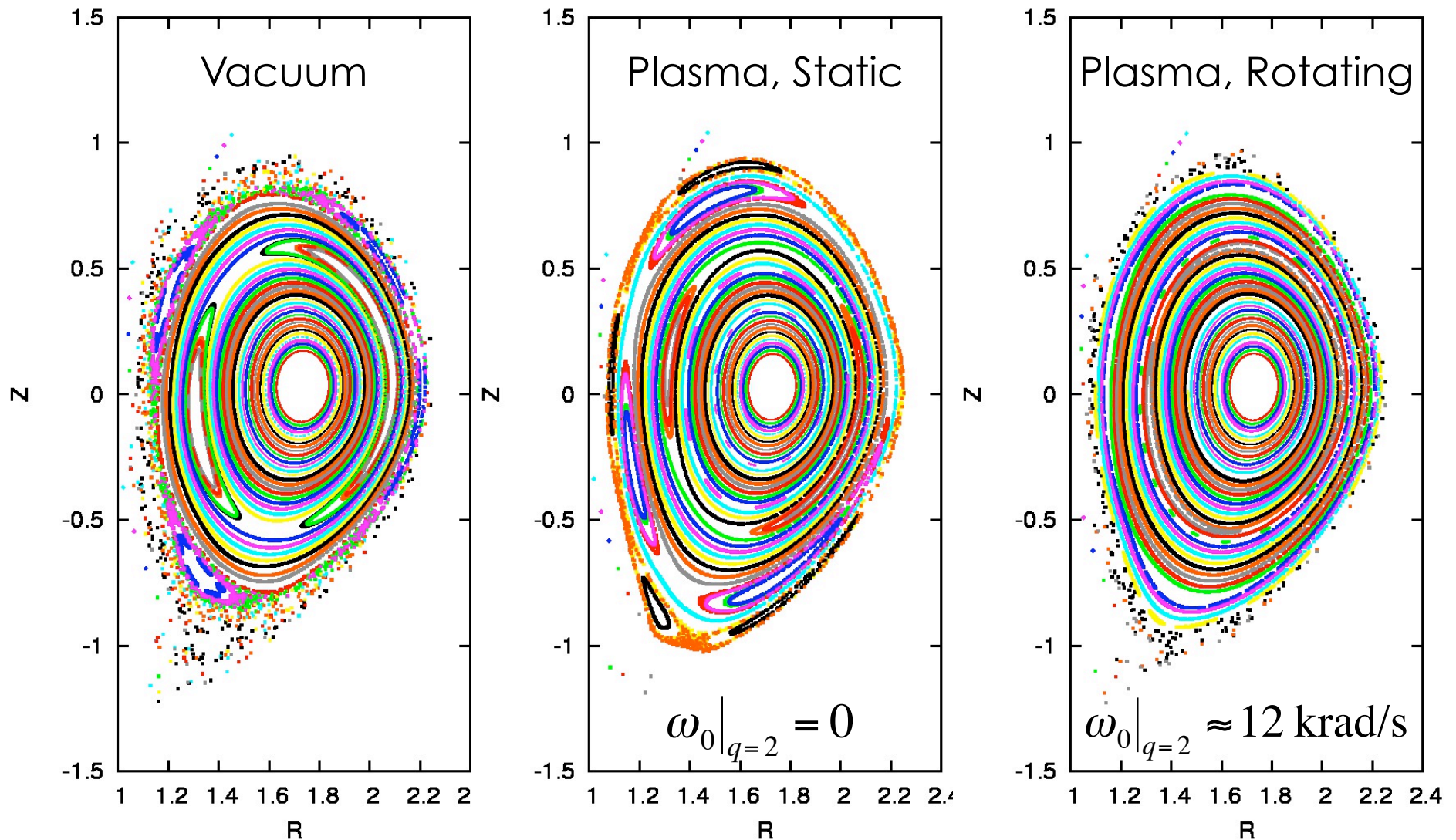
- **In practice, equilibria are almost always weakly unstable to “numerical tearing” modes**
 - Due mainly to lack of resolution around rational surfaces
 - No steady-state → This invalidates response
- **With MARS, this is usually not a problem since the response frequency is chosen *a priori***
- **With initial value code, the equilibrium must be made to be stable to these spurious modes**
 - Change equations: thermal diffusion, viscosity
 - Change equilibrium: rotation

Rotation & Dissipation Affect Stability & Screening

- **Dissipative terms inhibit screening response**
 - Magnetic islands form
- **Equilibrium rotation enhances screening**



Rotation Improves Core Screening; But Stochasticizes Edge



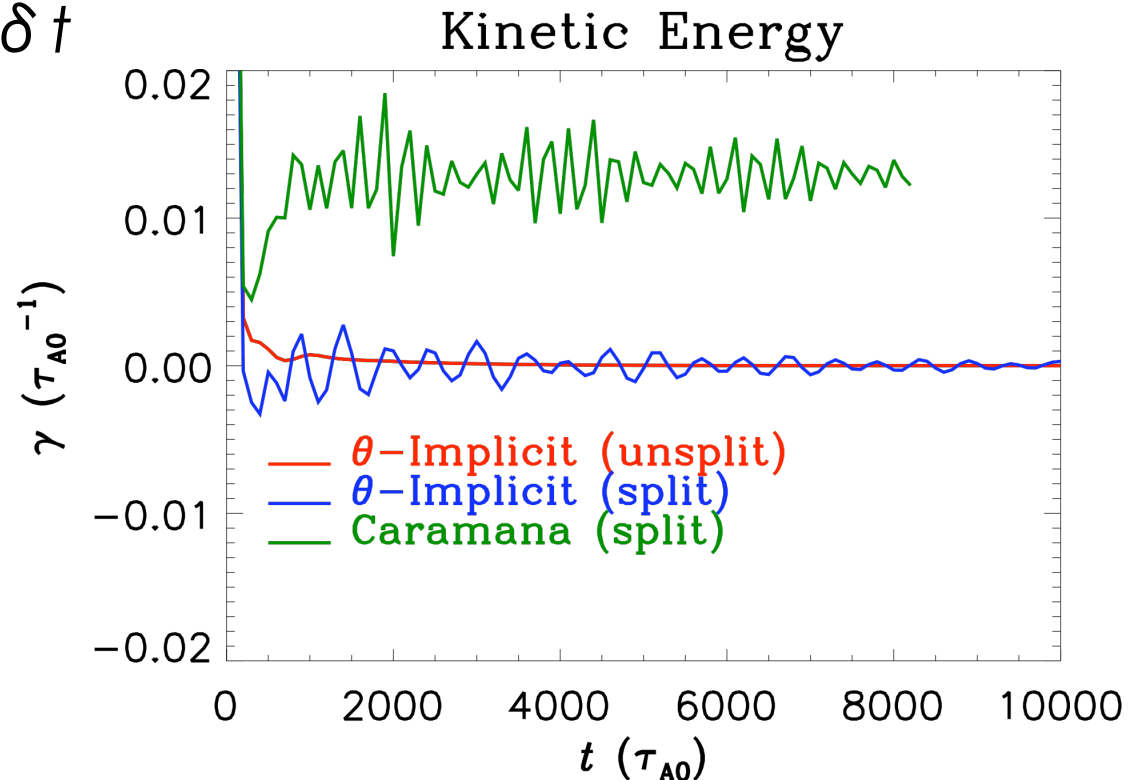
Summary of Non-Axisymmetric Results

- We are able to calculate response with **Spitzer resistivity, rotation, and two-fluid** terms
- **Initial-value calculations require dissipation to ensure stability**
 - Dissipation inhibits screening
- **Rotation enhances screening**
 - Direction of rotation is important even *in single fluid MHD*
- **Poster Tuesday morning**

Unsplit Method Superior to Split Methods for Finding Linear Perturbed Equilibrium

- **Split method has difficulty obtaining perturbed equilibrium**

- Persistent oscillations at low dissipation
- More sensitive to δt
- Caramana method more susceptible to numerical instability in under-resolved cases

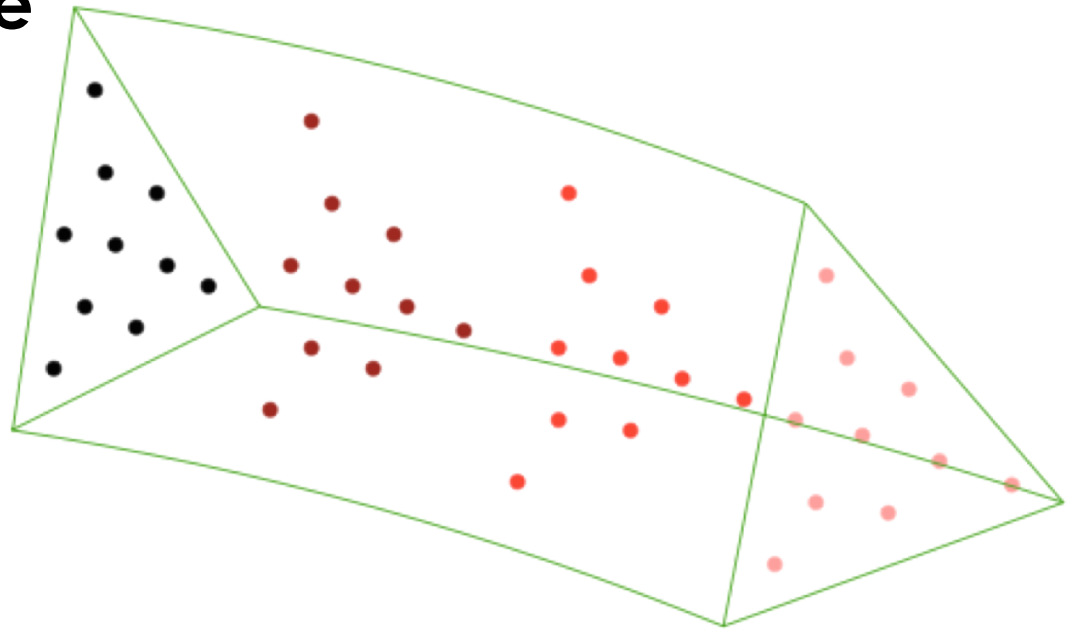


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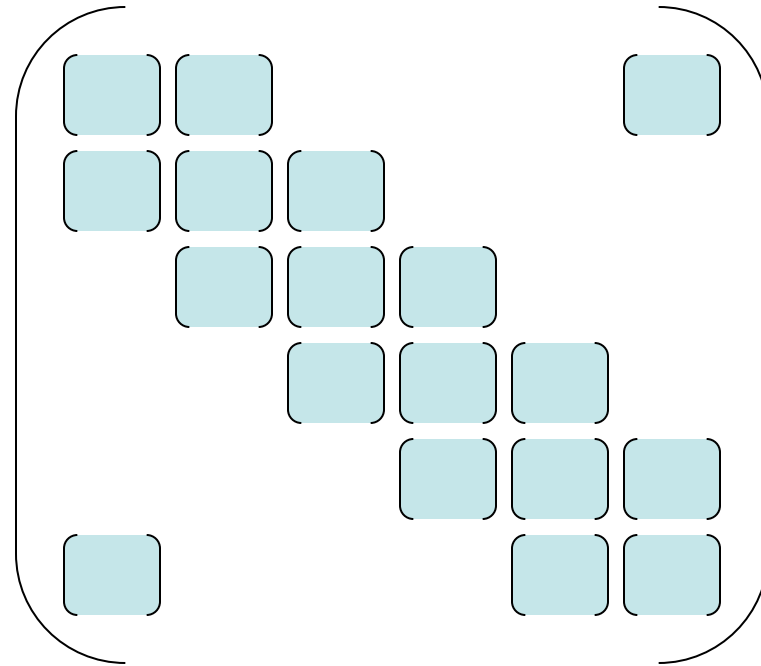
M3D-C1 Uses Wedge-Shaped Elements in 3D

- **Tensor product**
 - Poloidal: reduced quintic (C^1)
 - Toroidal: reduced cubic polynomials (C^1)
 - Integration quadrature is also tensor product
- **$6 \times 2 = 12$ DOFs/node**
- **3D mesh is series of 2D planes**
- **Allows packing in toroidal direction**



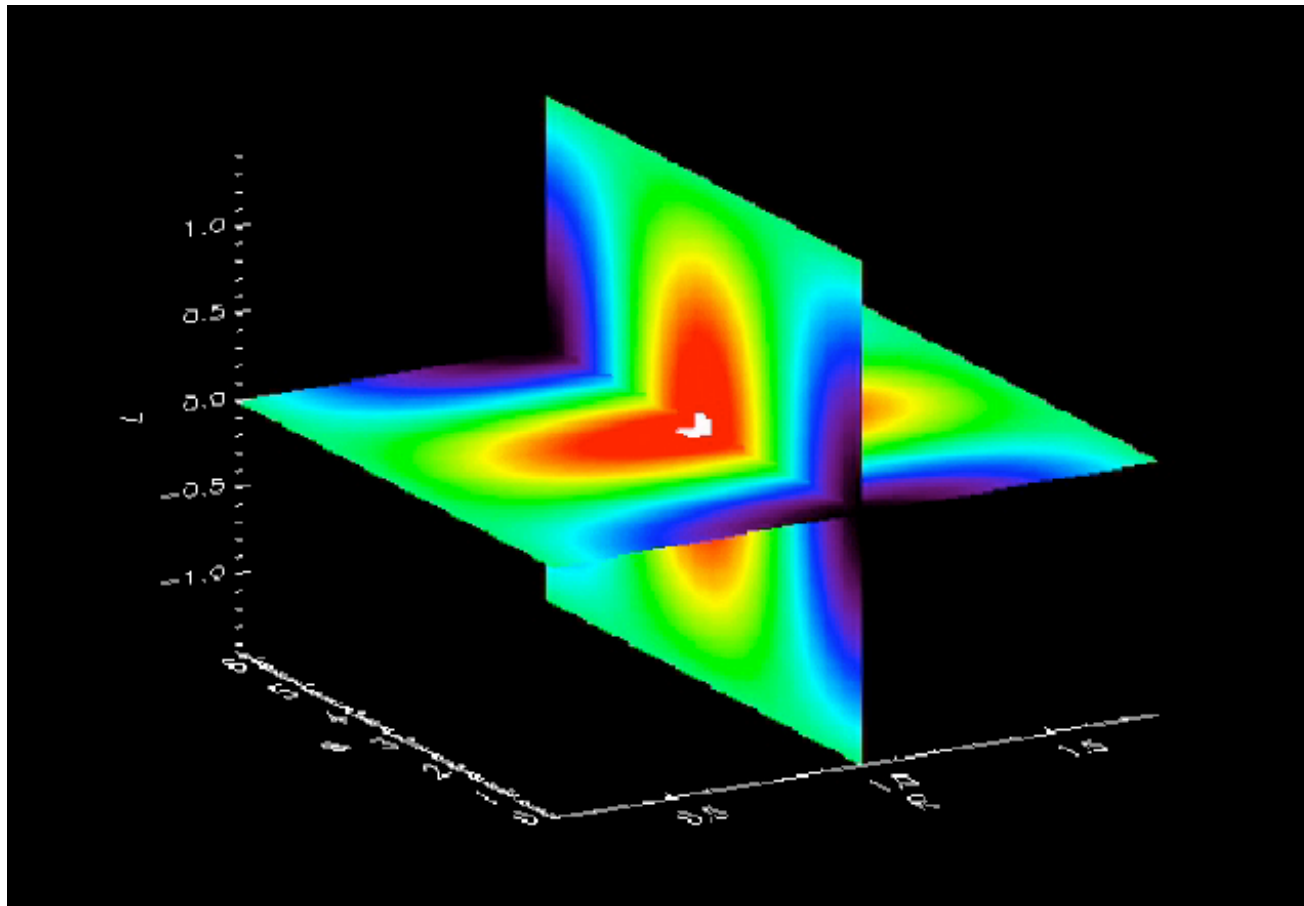
3D Matrix Has Cyclic Block-Tridiagonal Form

- Each plane corresponds to “block”
- Only nearest neighbor planes are coupled
- Presently this is solved without (further) preconditioning
- Typical problem: 10^5 DOFs/block

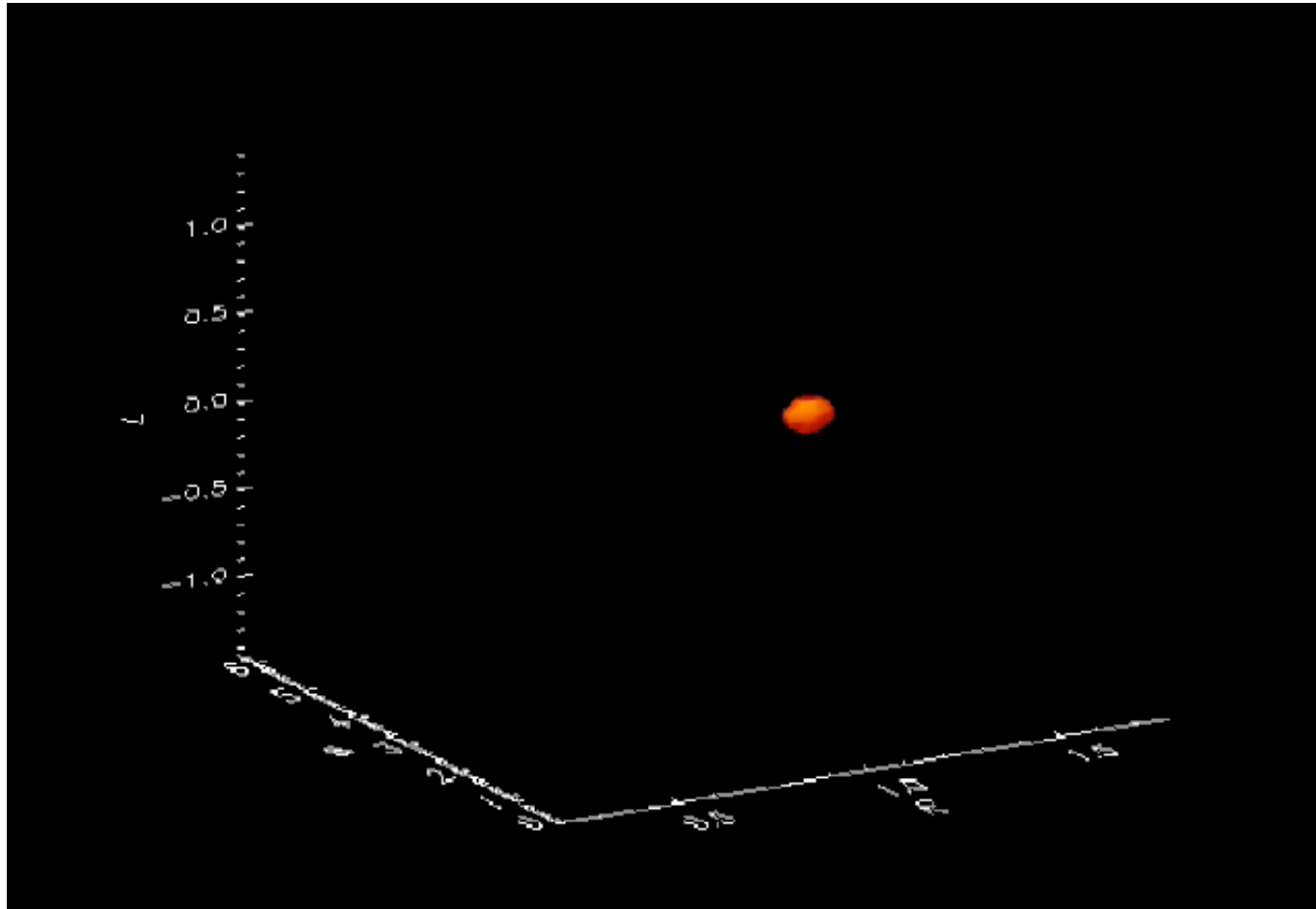


3D Results: Alfvén Wave

- With just 2 toroidal planes, ω is correct to 1 part in 10^5



3D Results: Anisotropic Diffusion in Helical Field



Summary

- **Resistive-wall boundary conditions are implemented**
 - Not yet functional in parallel
- **Linear 3D response successfully calculated with Spitzer resistivity, rotation, and two-fluid physics**
- **First fully-3D simulations have been run successfully**
 - Much future work will involve solver strategies