

M3D CDX-U Benchmark Status

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PPPL

CEMM Meeting

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Characteristics of the Current Drive Experiment Upgrade (CDX-U)

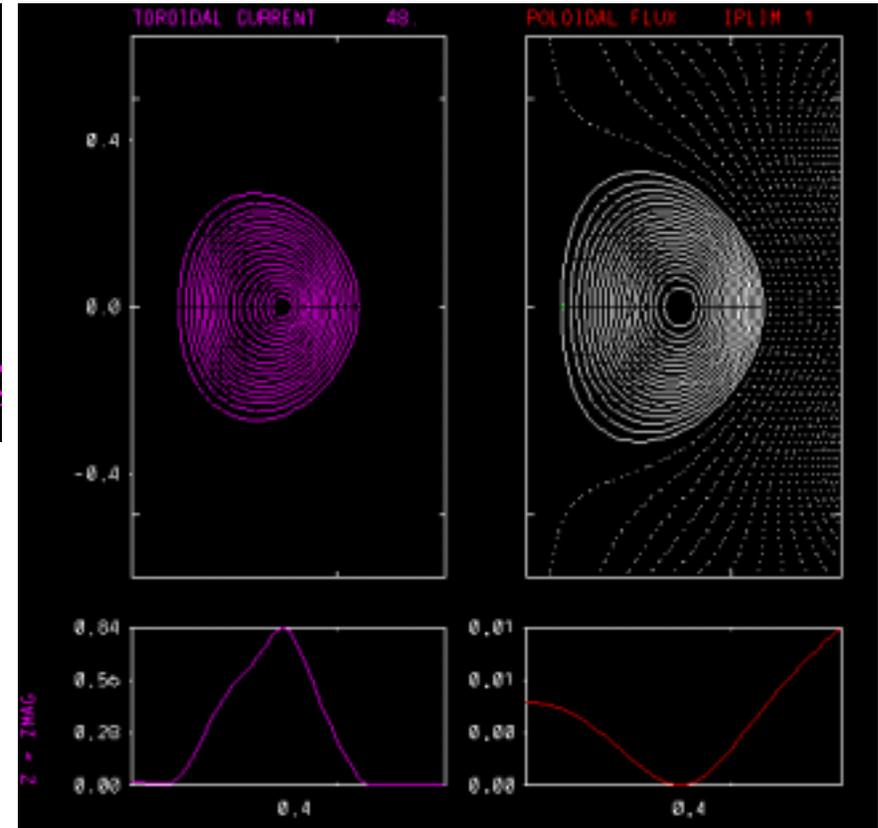
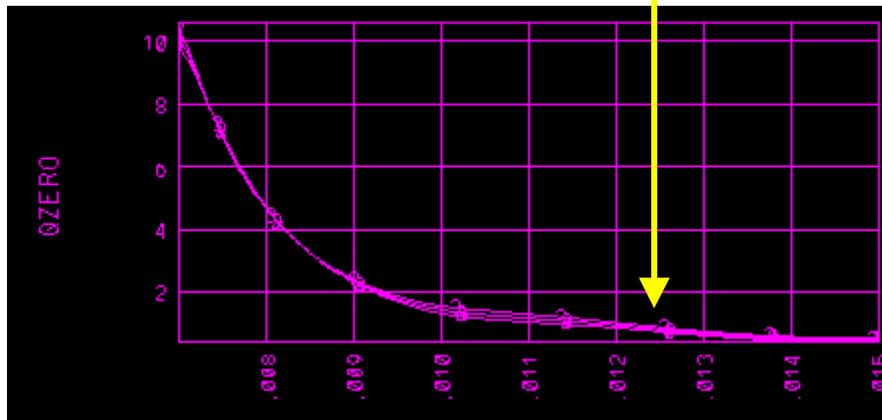
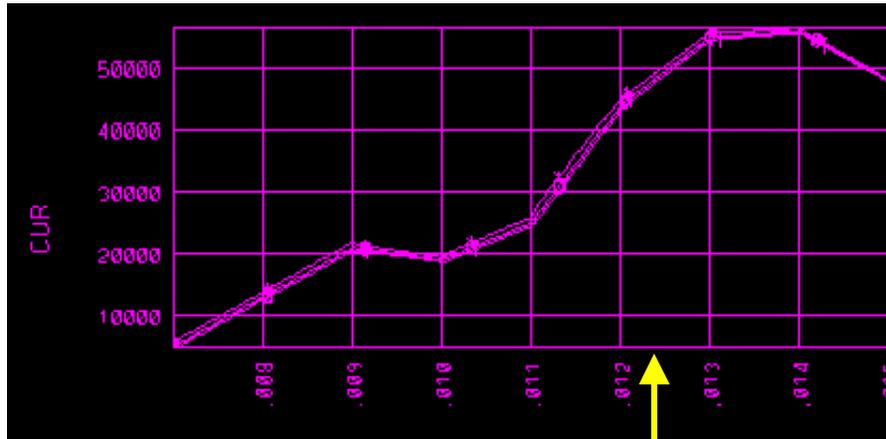


- Low aspect ratio tokamak ($R_0/a = 1.4 - 1.5$)
- Small ($R_0 = 33.5$ cm)
- Elongation $\kappa \sim 1.6$
- $B_T \sim 2300$ gauss
- $n_e \sim 4 \times 10^{13}$ cm⁻³
- $T_e \sim 100$ eV
- $I_p \sim 70$ kA

- Soft X-ray signals from typical discharges indicate two predominant types of low- n MHD activity:
 - sawteeth
 - “snakes”

The TSC Sequence

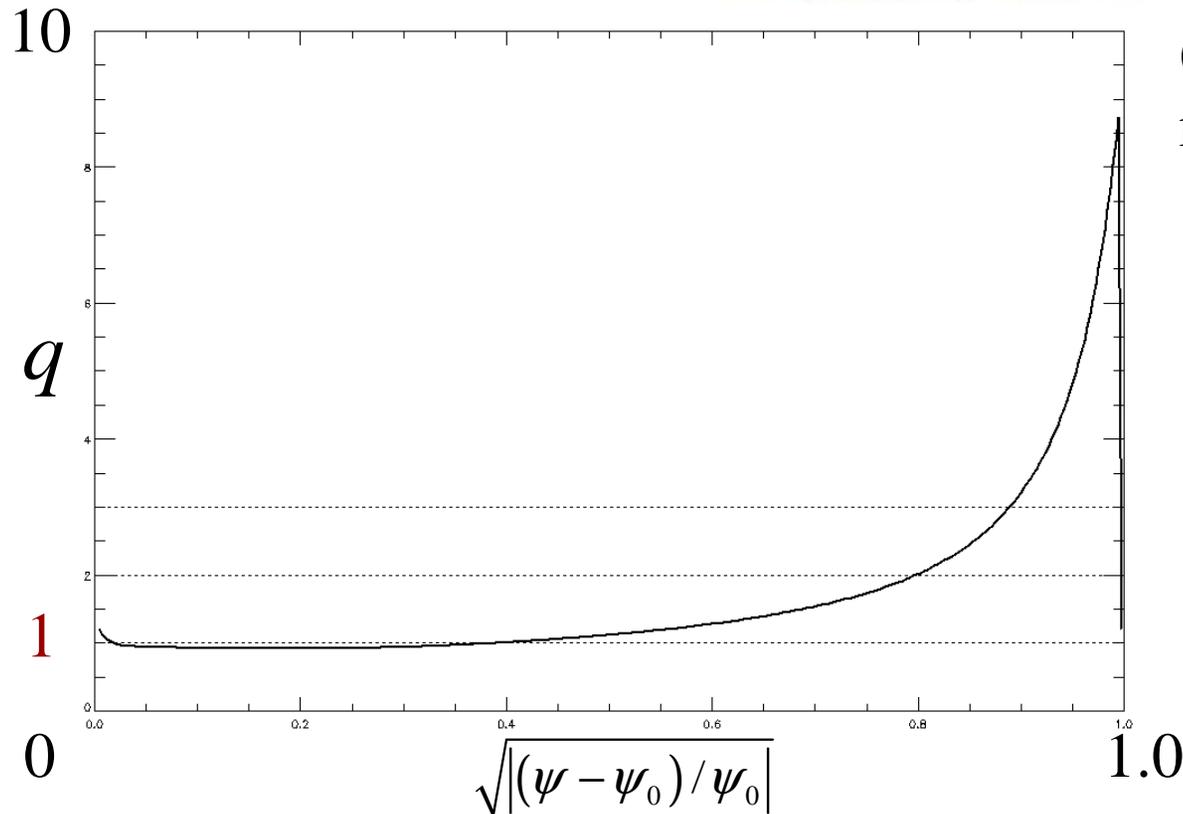
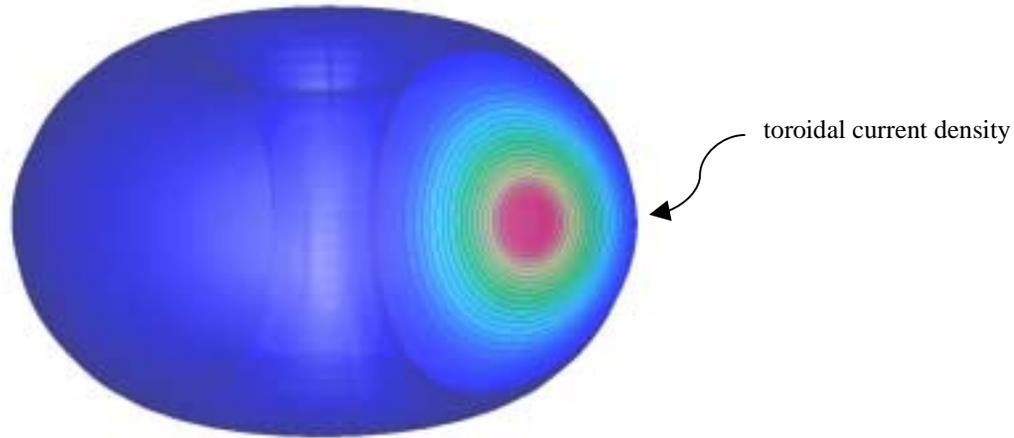
TSC follows 2D (axisymmetric) evolution of typical CDX-U discharge



Equilibrium at $t=12.40\text{ms}$
(as q_0 drops to 0.92) is used
to initialize 3D runs

Device Aspect Ratio (Low-A) Case

- Equilibrium taken from a TSC sequence (Jsolver file).
- Shaped cross-section
- $R/a = 1.4$
- $q_{\min} \approx 0.922$
- $q(a) \sim 9$

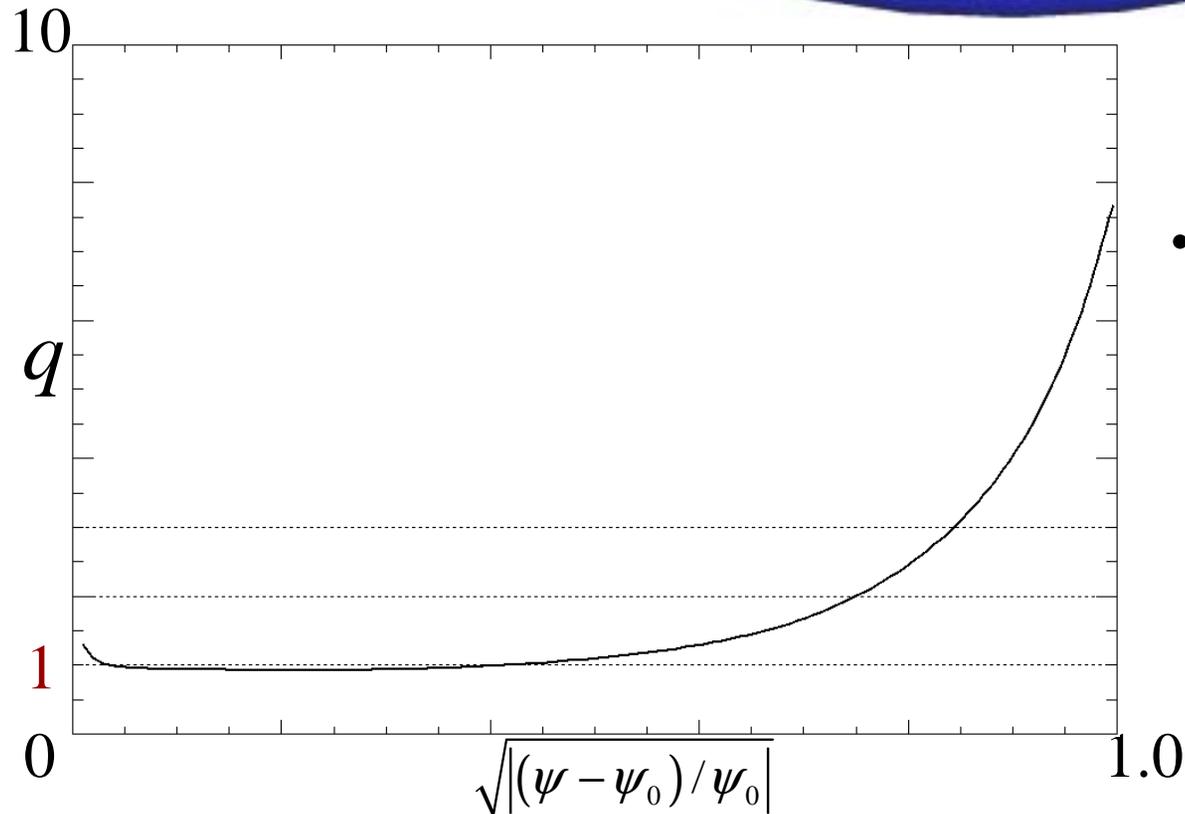
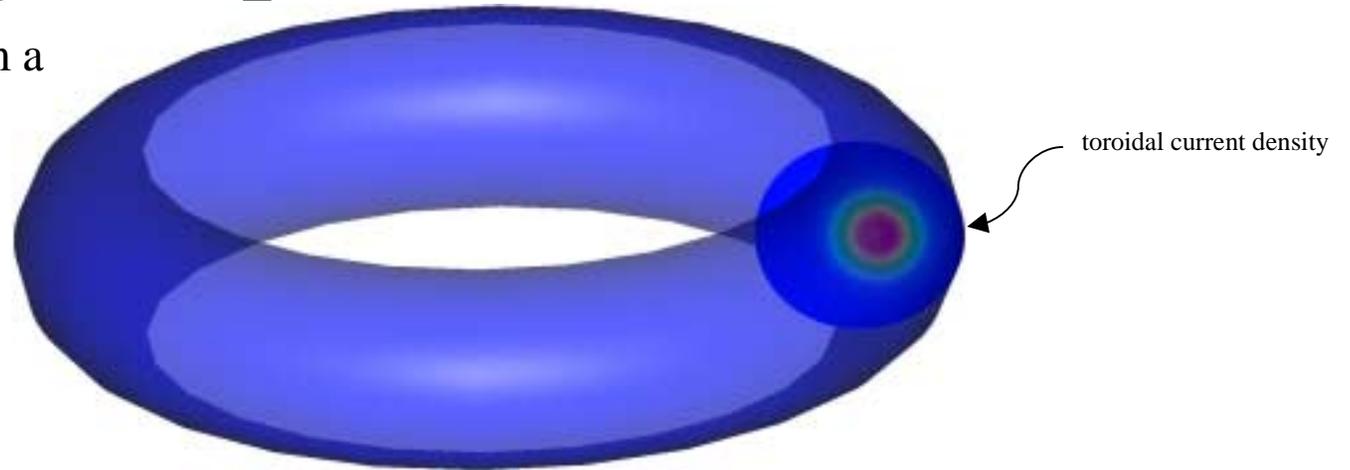


Questions to investigate:

- Linear growth rate and eigenfunctions
- Nonlinear evolution
 - disruption?
 - stagnation?
 - repeated reconnections?

Large Aspect Ratio Case

- Equilibrium taken from a TSC sequence (Jsolver file).
- Circular cross-section
- $R/a = 3.75$
- $q_{\min} \approx 0.927$
- $q(a) \sim 8$



- Note: this does not correspond to any actual device. Chosen for ease of cross-code comparison.

Baseline Parameters for Cross-code Benchmark

Lundquist Number S	$\sim 2 \times 10^4$ on axis.
Resistivity η	Spitzer profile $\propto T_{\text{eq}}^{-3/2}$, cut off at $100 \times \eta_0$
*Prandtl Number Pr	10 on axis.
*Viscosity μ	Constant in space and time.
*Perpendicular thermal conduction κ_{\perp}	0
*Parallel thermal conduction κ_{\parallel}	0
Peak Plasma β	$\sim 3 \times 10^{-2}$ (low-beta).
Density Evolution	Turned on for nonlinear phase.
Nonlinear initialization	Pure $n=1$ perturbation such that $\frac{\max(B_{\text{pol}}^1)}{\max(B_{\phi}^0)} = 10^{-4}$

*Non-physical values chosen to aid in cross-code benchmark.

Comparing the Codes

M3D and NIMROD are both parallel 3D nonlinear extended MHD codes in toroidal geometry maintained by multi-institutional collaborations, and comprise the two members of the Center for Extended MHD Modeling (CEMM) SciDAC.

M3D

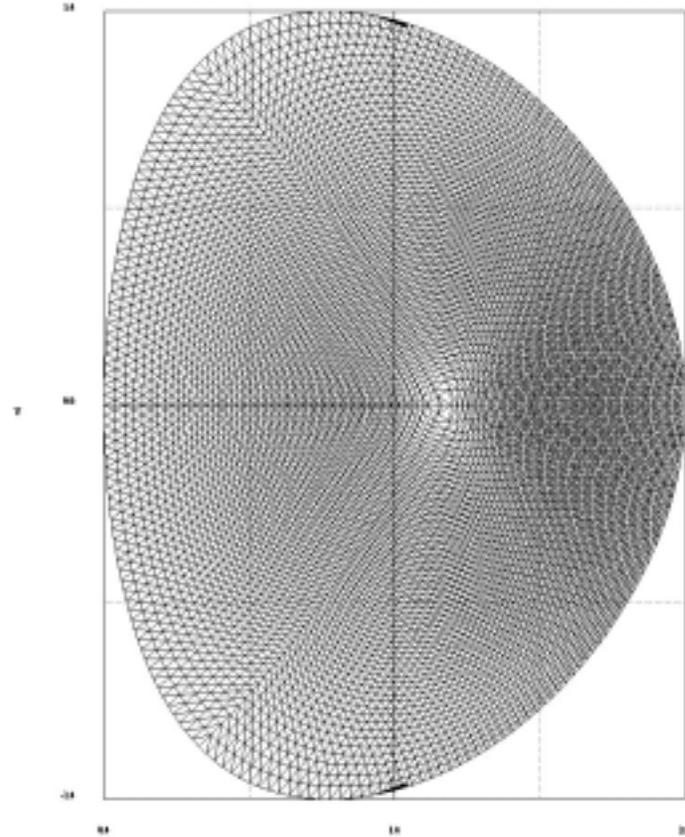
- Uses linear finite elements in-plane.
- Uses finite differences between planes or pseudo-spectral derivatives.
- Partially implicit treatment allows efficient time advance but requires small time steps.
- Linear operation: full nonlinear + filtering, active equilibrium maintenance.
- Nonlinear operation: all components of all quantities evolve nonlinearly.

NIMROD

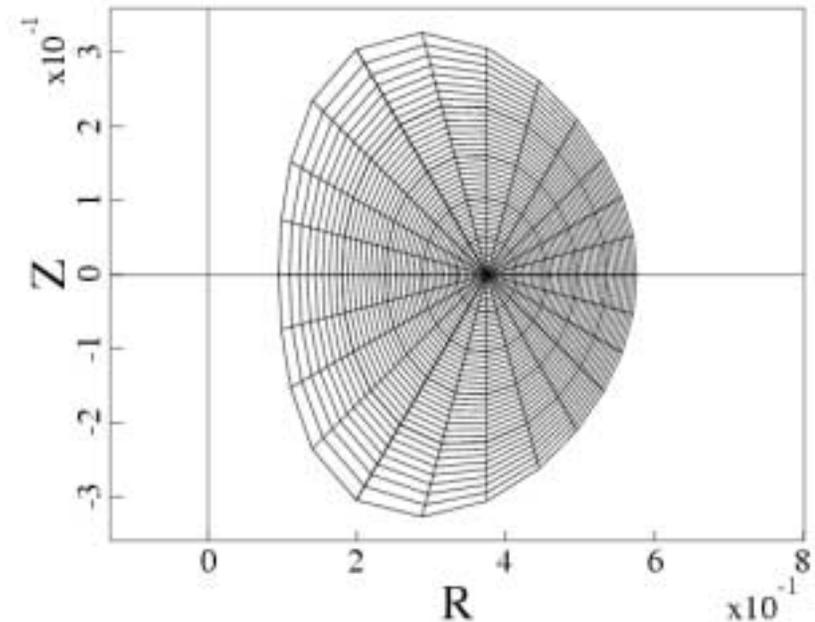
- Uses high-order finite elements in-plane.
- Uses Fourier decomposition in toroidal direction.
- Fully implicit treatment requires costly matrix inversions but allows large time steps.
- Linear operation: evolve perturbations to particular modes only.
- Nonlinear operation: perturbations to fixed equilibrium are evolved, with nonlinear couplings between modes.

Poloidal Meshes for the low-A Case

M3D



NIMROD

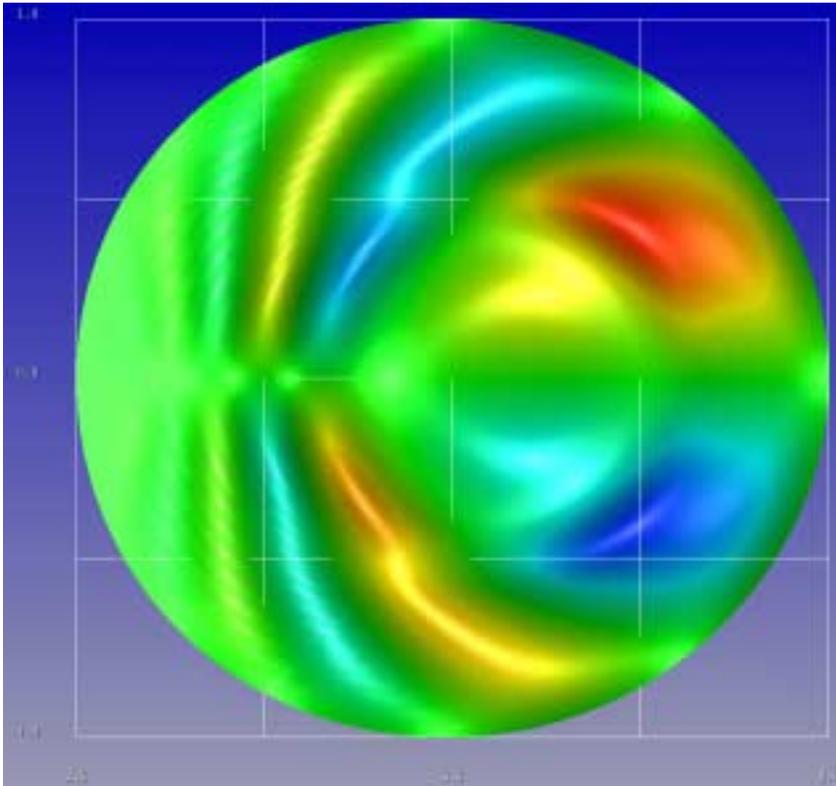


- 89 radial zones, up to 267 in θ in unstructured mesh
- Linear basis functions on triangular elements
- Conducting wall
- Finite differences toroidally; 24 planes

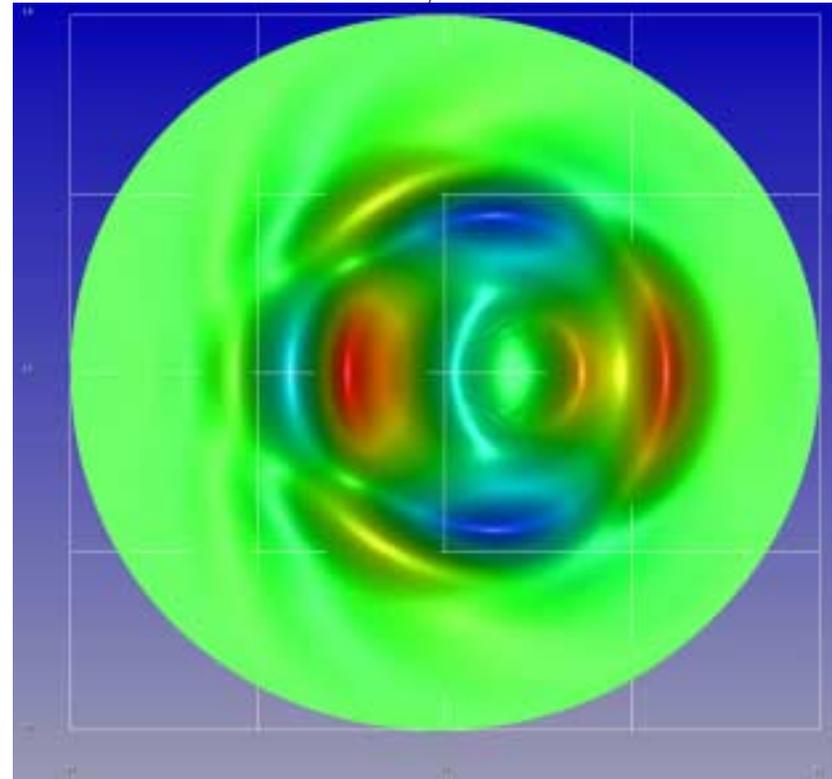
- 40×24 structured grid
- 4th order basis functions on quadrilateral elements
- Conducting wall
- Fourier decomposition toroidally; 10 or more modes retained

Large Aspect Ratio: $n=1$ Eigenmode

Incompressible velocity
stream function U



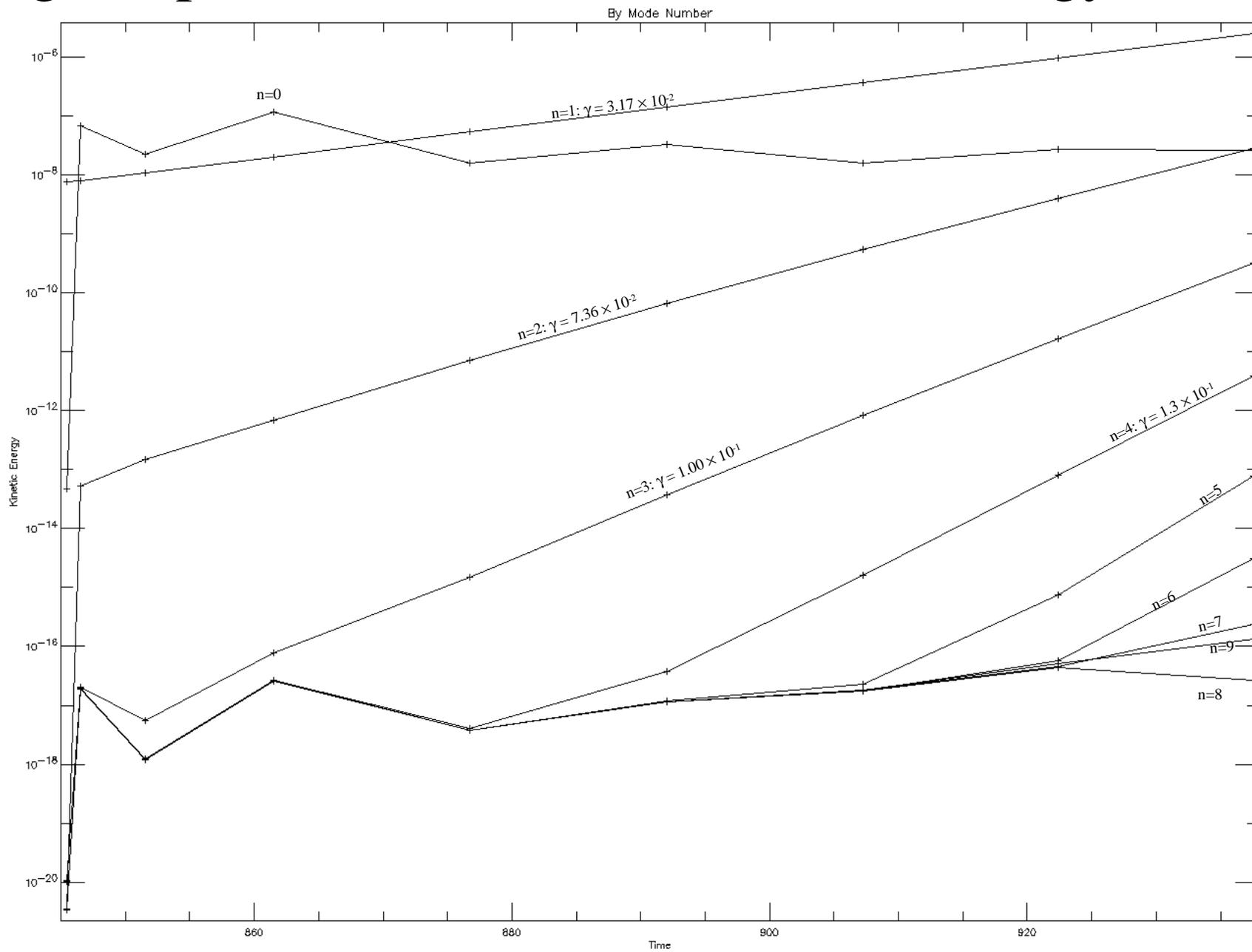
Toroidal current density
 J_ϕ



$$\gamma\tau_A = 3.28 \times 10^{-2} \rightarrow \text{growth time} = 30.5 \tau_A$$

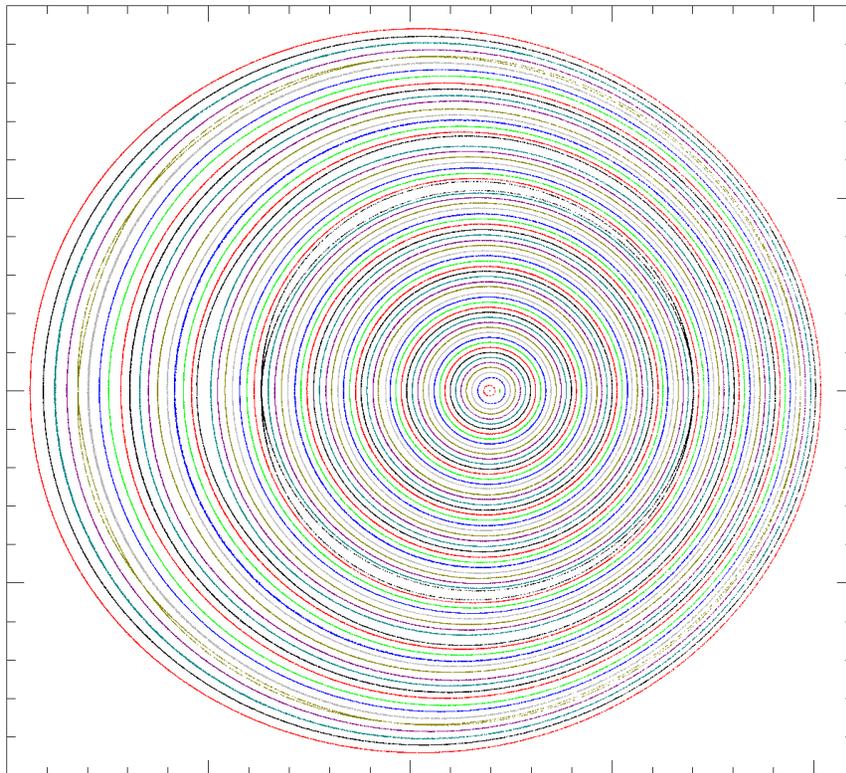
Agrees with NIMROD value to within 7%.

Large Aspect Ratio: Nonlinear Kinetic Energy History

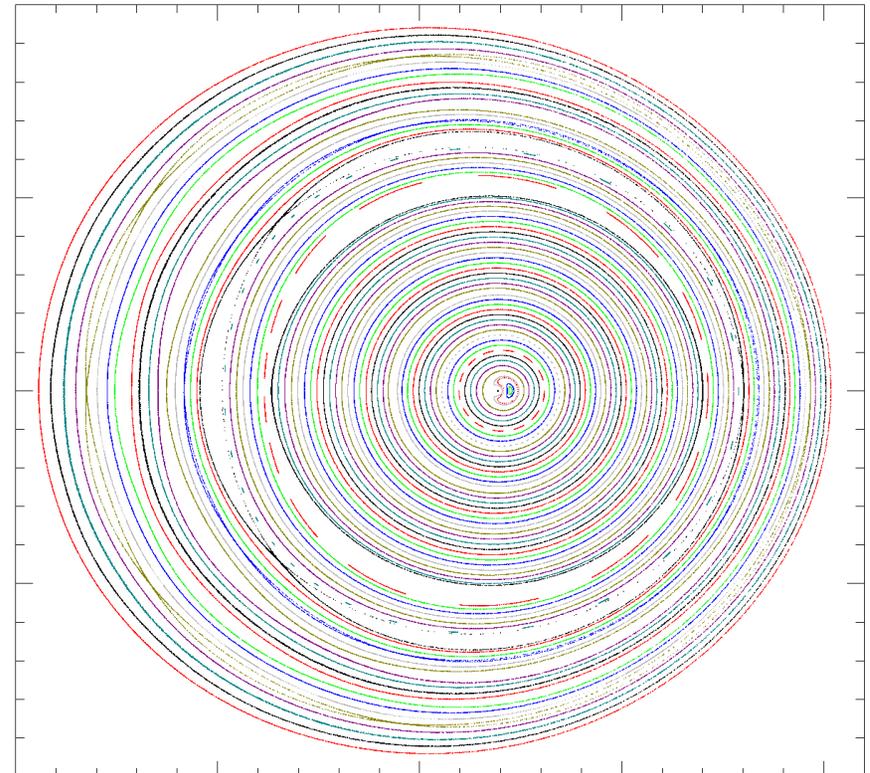


Large Aspect Ratio: Nonlinear Time Series

Poincaré Plots



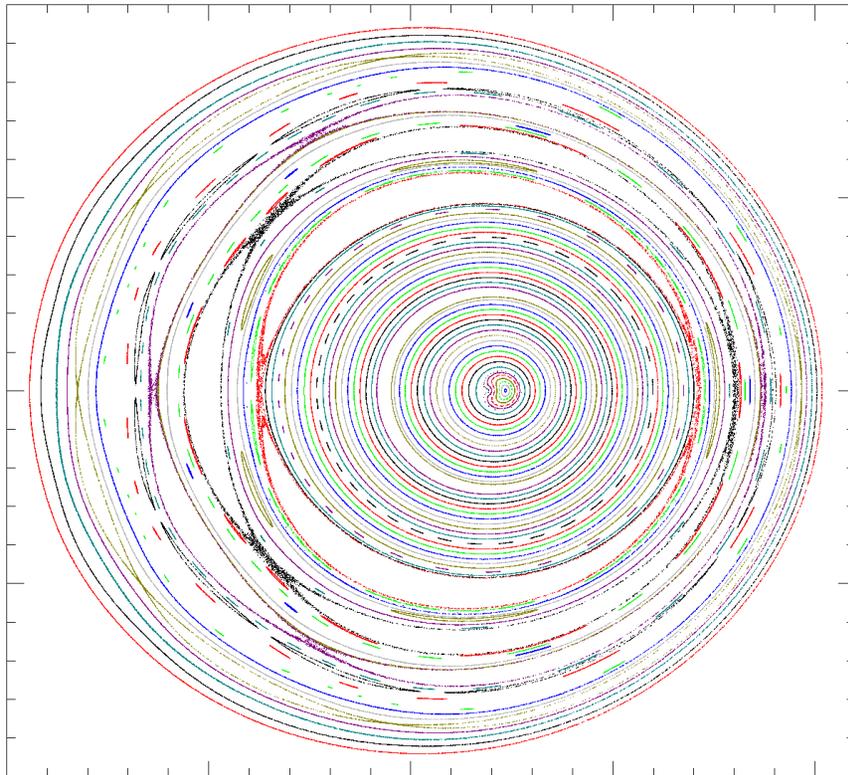
$t = 845.45$; chopped $\times 6.498543$



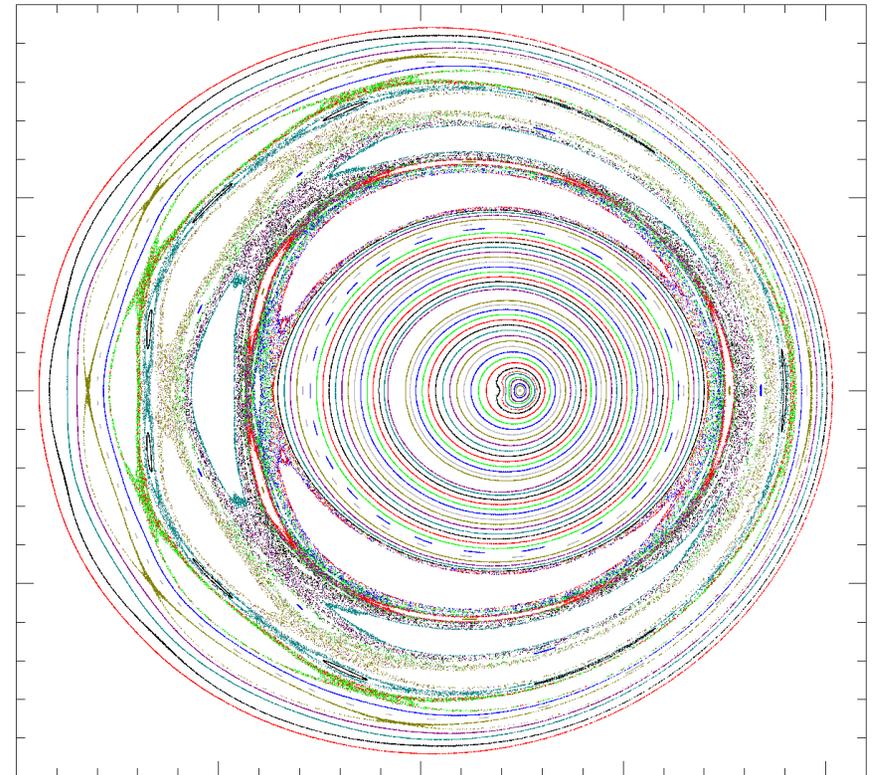
$t = 892.00$

Large Aspect Ratio: Nonlinear Time Series

Poincaré Plots, Continued



$t = 922.37$



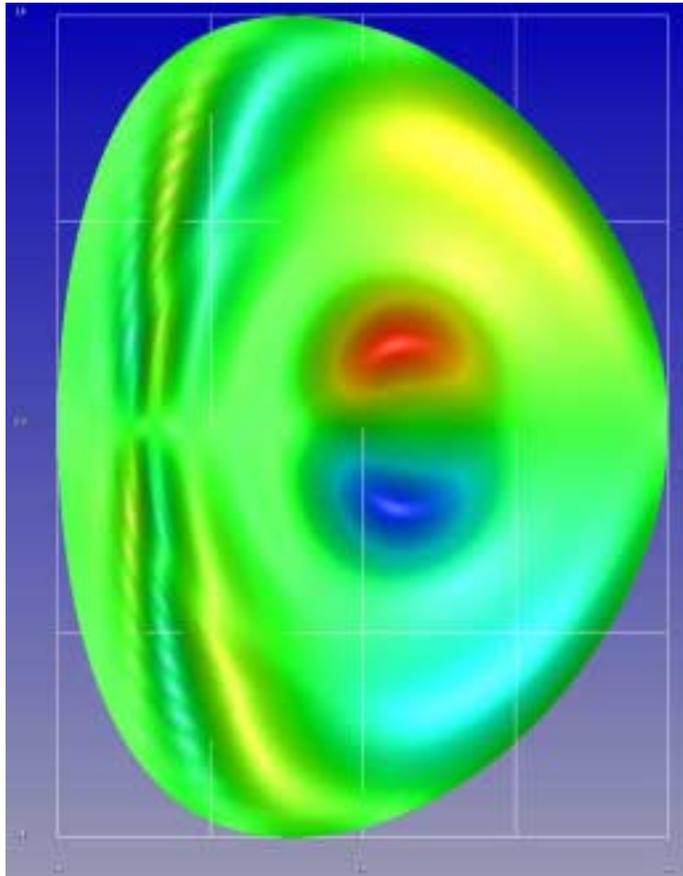
$t = 937.56$

All islands visible in these plots have toroidal mode number $n=1$.

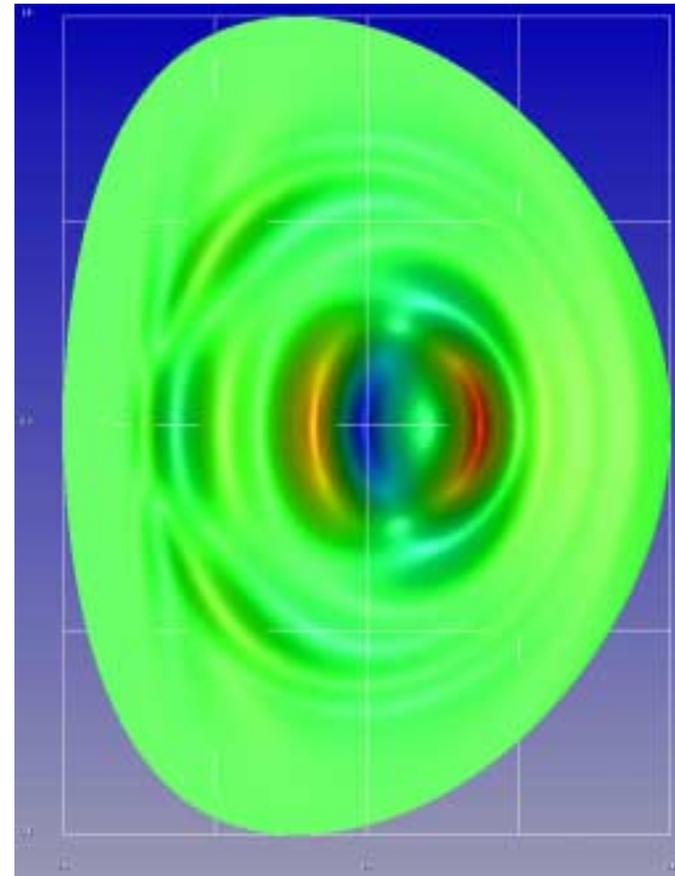
Note: Run halted because of onset of stochasticity.

Low Aspect Ratio: $n=1$ Eigenmode

Incompressible velocity
stream function U

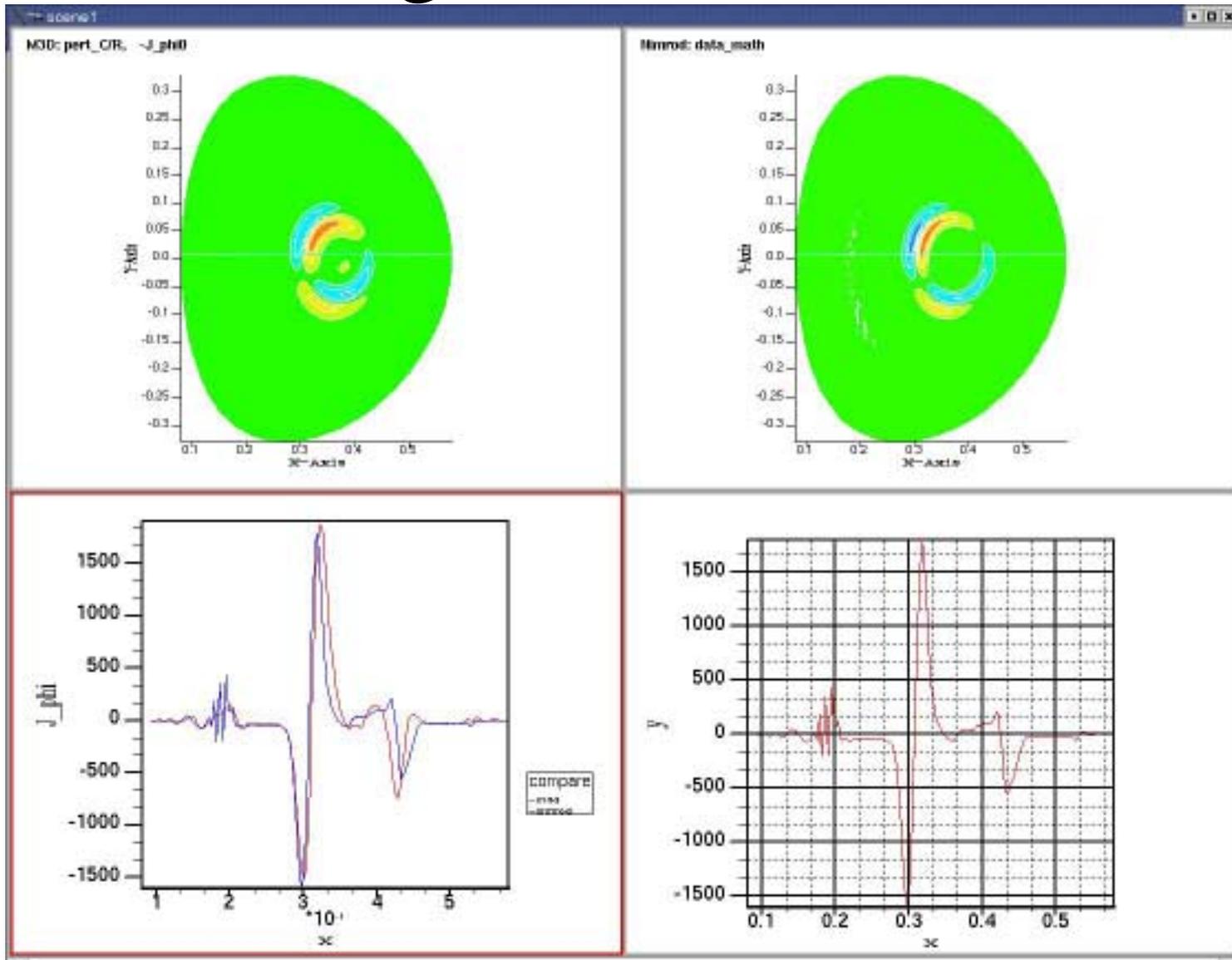


Toroidal current density
 J_ϕ



$$\gamma\tau_A = 8.61 \times 10^{-3} \rightarrow \text{growth time} = 116 \tau_A$$

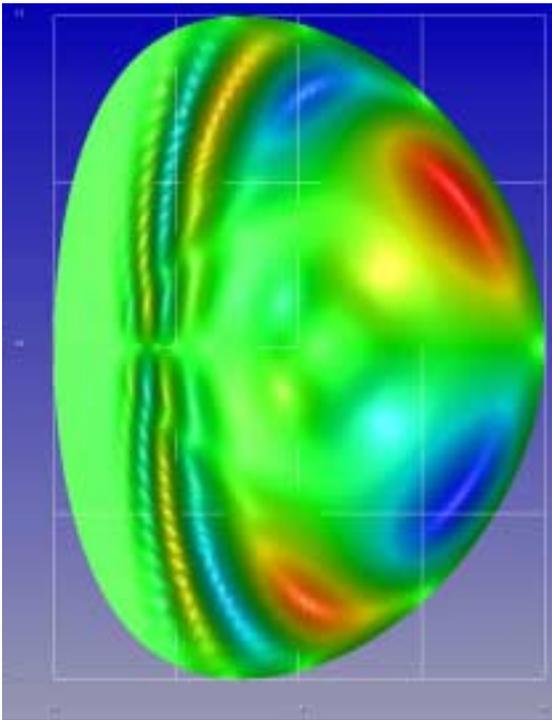
Predicted Eigenmode Shows Qualitative Agreement with NIMROD



Low Aspect Ratio: Higher n Eigenmodes

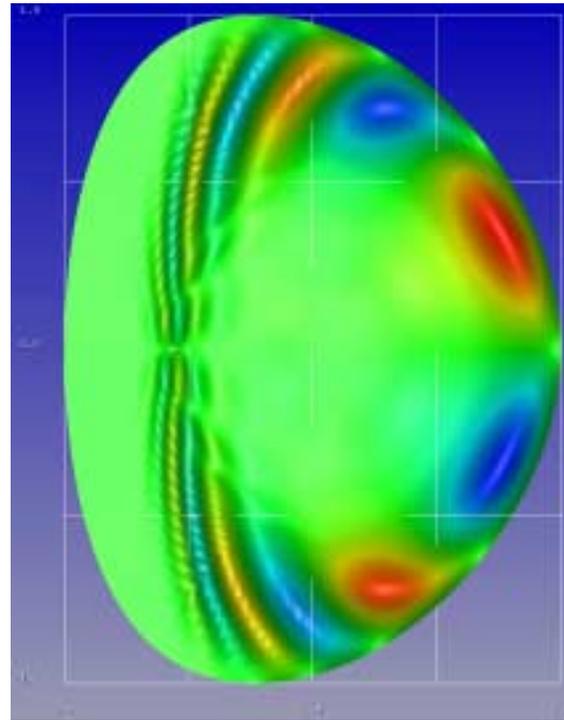
Incompressible velocity
stream function U

$n = 2$



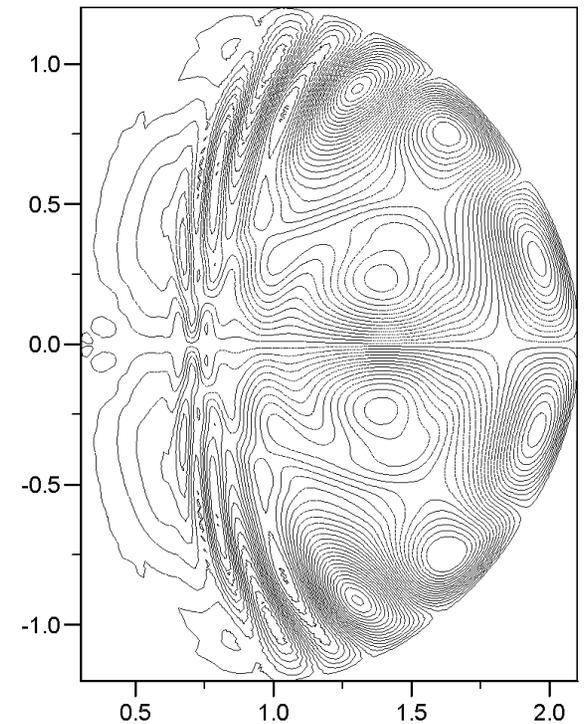
$m \geq 5$
 $\gamma\tau_A = 1.28 \times 10^{-2}$

$n = 3$



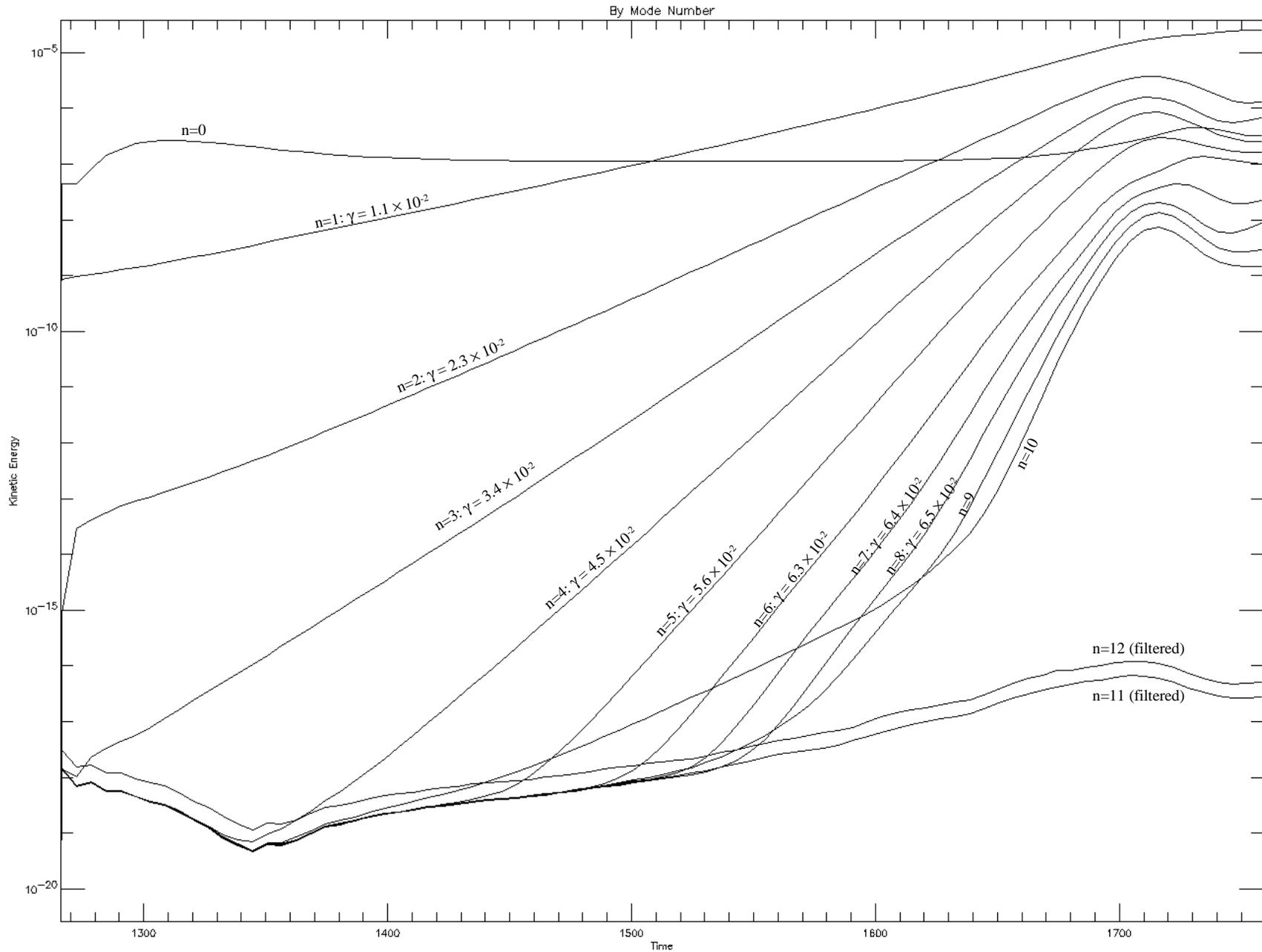
$m \geq 7$
 $\gamma\tau_A = 1.71 \times 10^{-2}$

$n = 4$ (projected)



$m \geq 8$
not converged

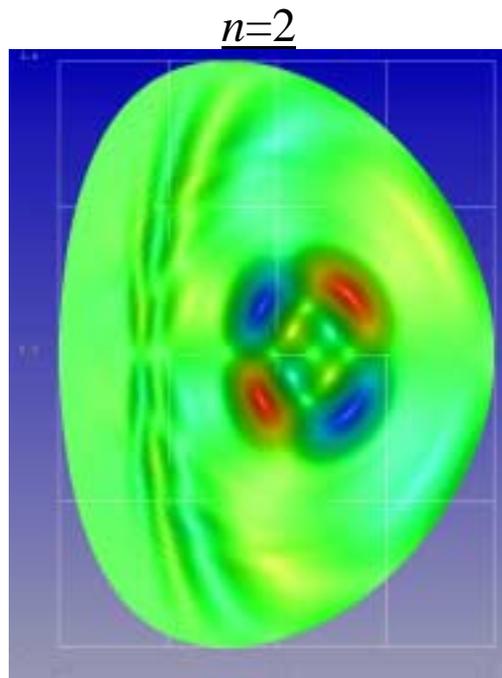
Low Aspect Ratio: Nonlinear Kinetic Energy History



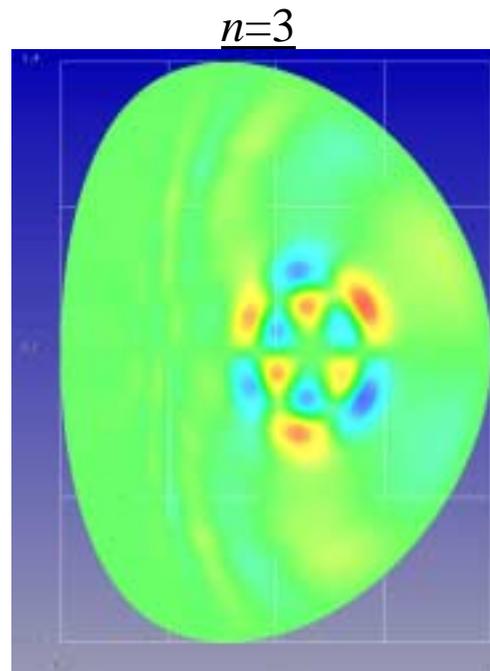
“Linear” high- n modes are driven, not eigenmodes

Incompressible velocity stream function U

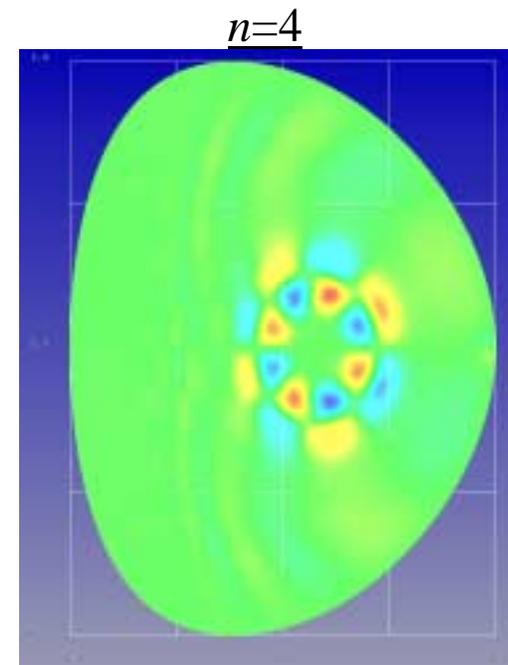
Component of “linear” mode
in nonlinear run



$m = 2$
 $\gamma = 2.3 \times 10^{-2}$



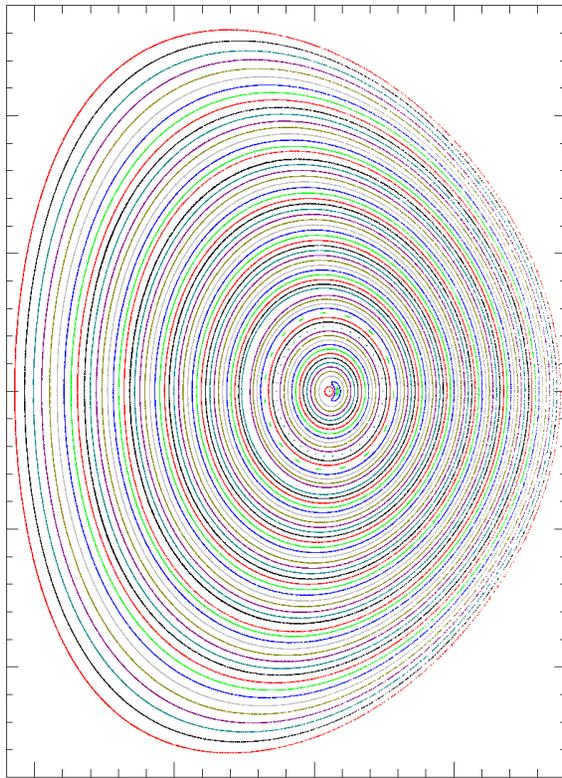
$m = 3$
 $\gamma = 3.4 \times 10^{-2}$



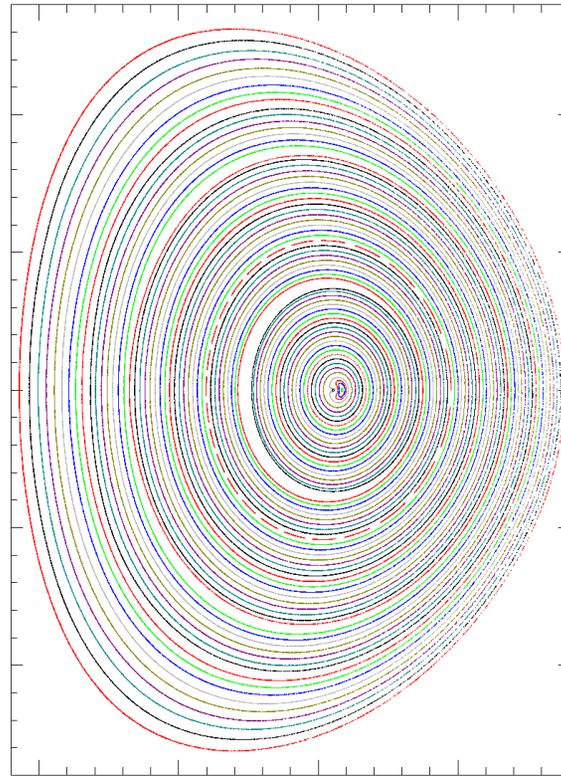
$m = 4$
 $\gamma = 4.5 \times 10^{-2}$

Low Aspect Ratio: Nonlinear Time Series

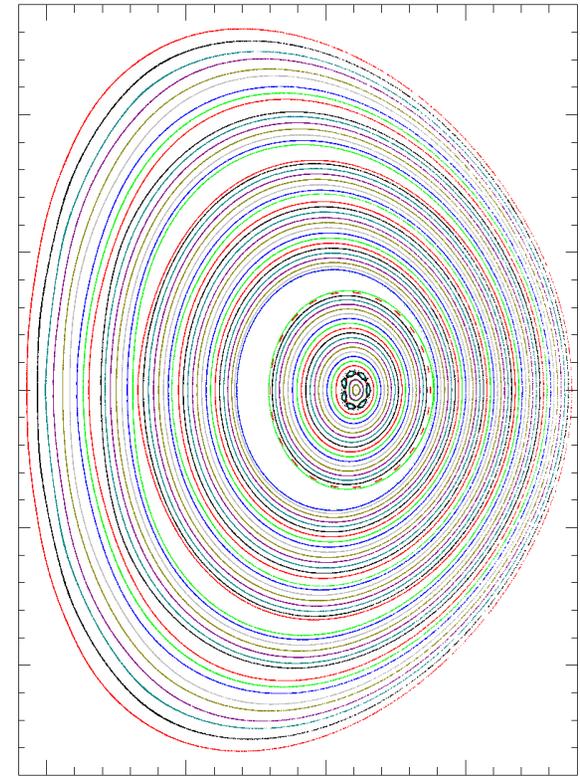
Poincaré Plots



$t = 1266.17$, chopped $\times 4.470358$



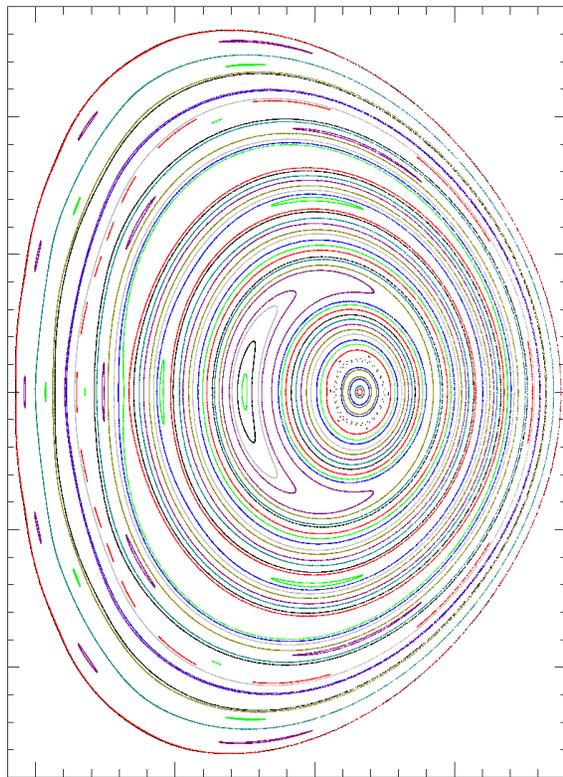
$t = 1404.57$



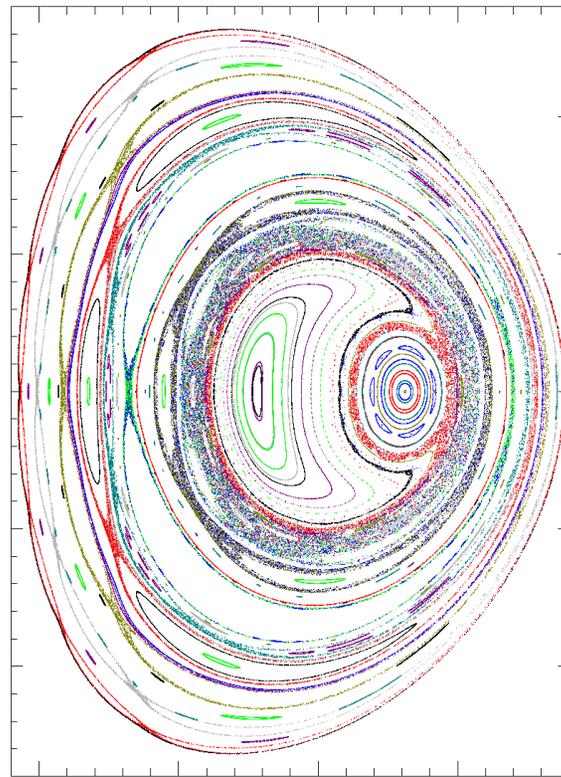
$t = 1548.68$

Low Aspect Ratio: Nonlinear Time Series

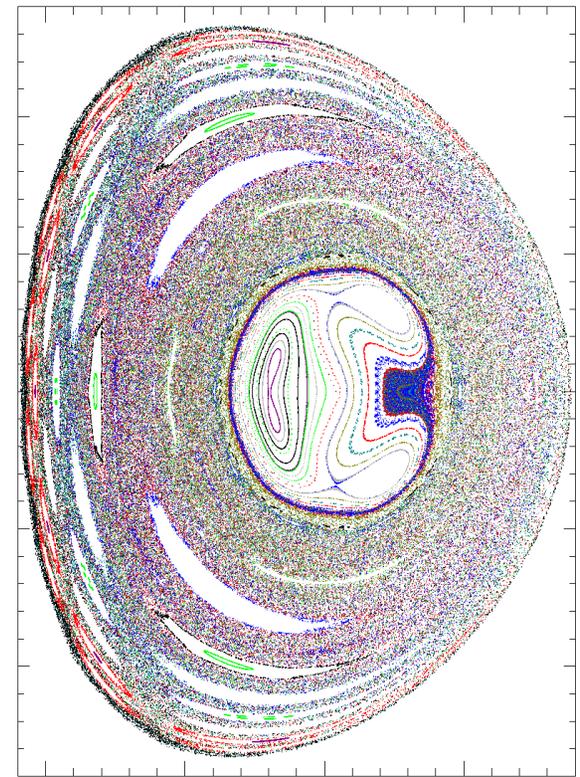
Poincaré Plots, Continued



$t = 1620.62$



$t = 1686.41$



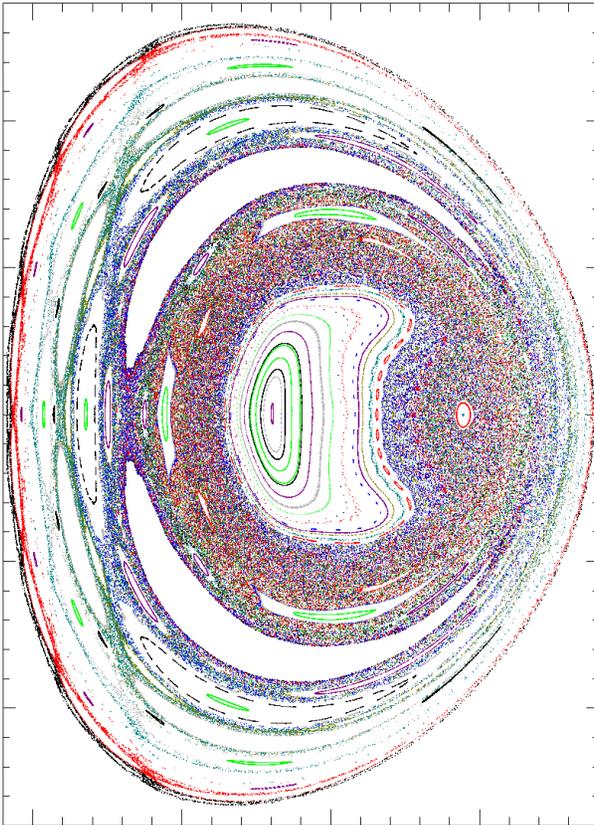
$t = 1758.34$

Disruption occurs before completion of sawtooth crash.

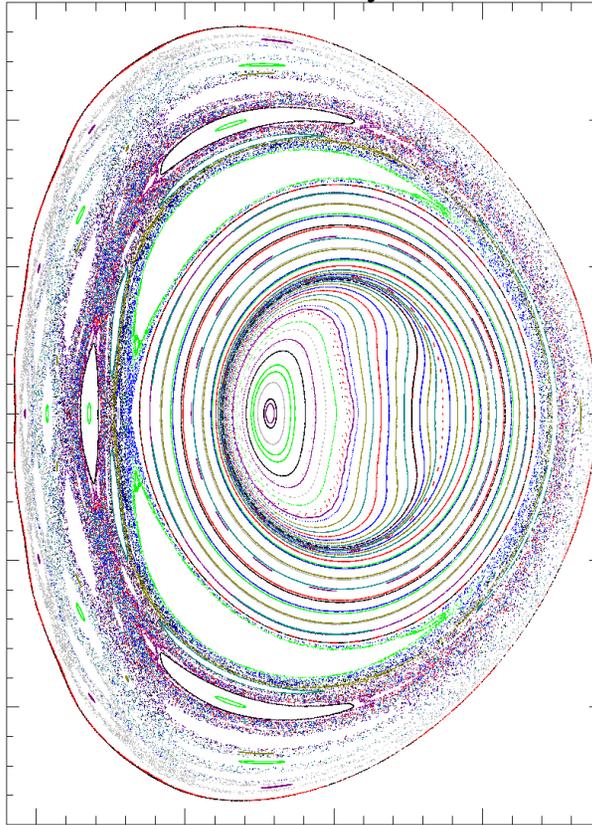
Focus on Saturation Time $t=1710.38$

(Time when $n=2$ kinetic energy peaks)

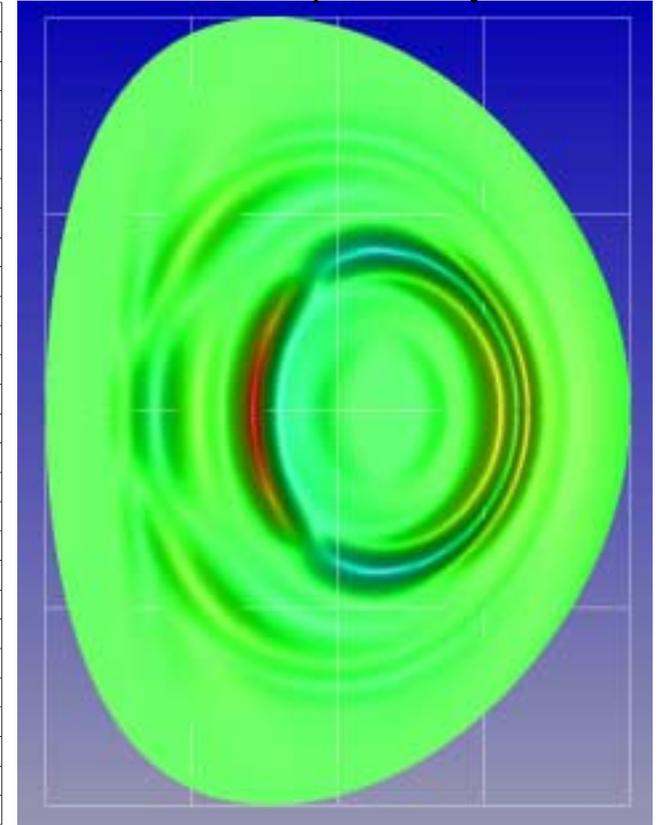
Poincaré plot,
all modes



Poincaré plot,
 $n=0,1$ only



J_ϕ
 $n=1$ component only



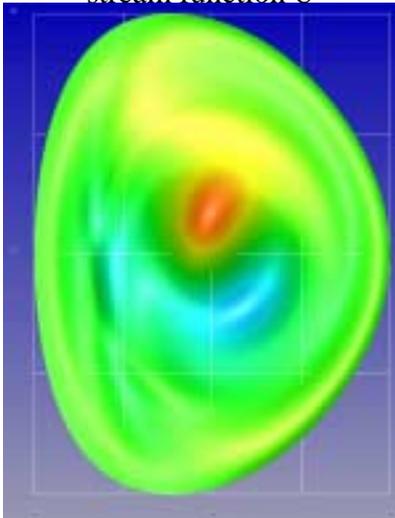
Flow Stabilization of the $n=1$ Mode in the low- A MHD Case

- Introduce steady-state 2D toroidal flow profile of the form

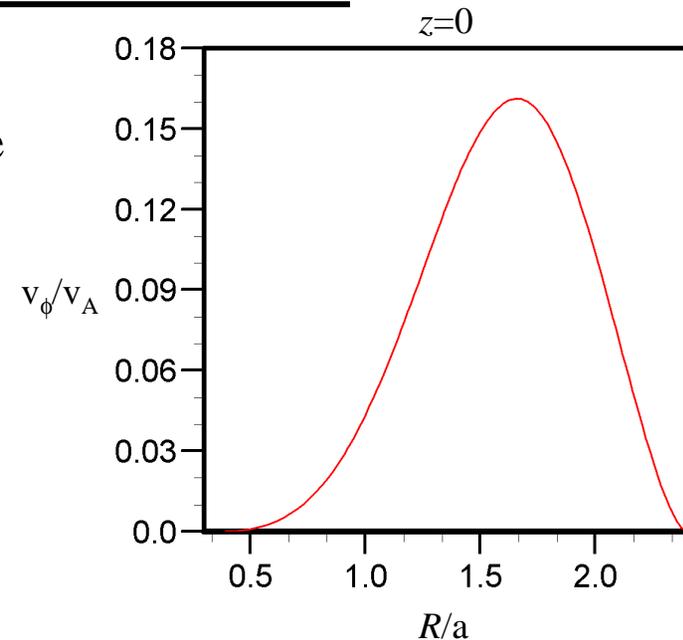
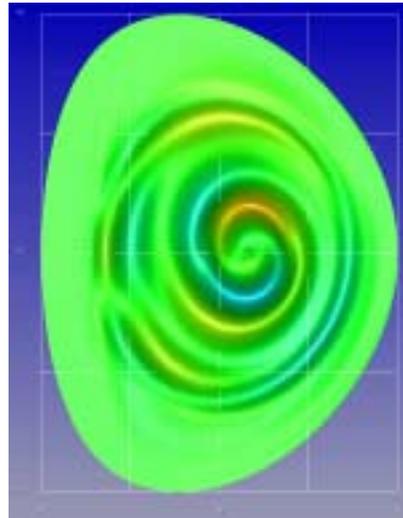
$$\frac{v_\phi(R, z)}{v_A} = 0.1 \frac{R}{a} \left| \frac{\psi(R, z)}{\psi_0} \right|^{1.5} :$$

New $n=1$ eigenmode:

Incompressible velocity stream function U



Toroidal current density J_ϕ



Growth rate $\gamma\tau_A$ decreases from 8.61×10^{-3} to 2.4×10^{-3} .

Extended MHD Effects

(Initial density profile uniform)

	MHD		MHD + ω_i^* term		
Evolving Density		$v_{\parallel}/v_A = 0$ (χ_{\parallel} off)	$v_{\parallel}/v_A = 2.5$ (χ_{\parallel} on)	$v_{\parallel}/v_A = 0$ (χ_{\parallel} off, H=0.10)	$v_{\parallel}/v_A = 2.5$ (χ_{\parallel} on, H=0.15)
	$n=1$	0.016	0.014	0.0165	0.0137
	$n=4$	0.0307	0.0075	0.0304	0.0095
Constant Density		$v_{\parallel}/v_A = 0$ (χ_{\parallel} off)	$v_{\parallel}/v_A = 2.5$ (χ_{\parallel} on)	$v_{\parallel}/v_A = 0$ (χ_{\parallel} off, H=0.15)	$v_{\parallel}/v_A = 2.5$ (χ_{\parallel} on, H=0.15)
	$n=1$	0.017	0.021	0.021	0.021
	$n=4$	0.0349	0.0132	0.0354	0.0123

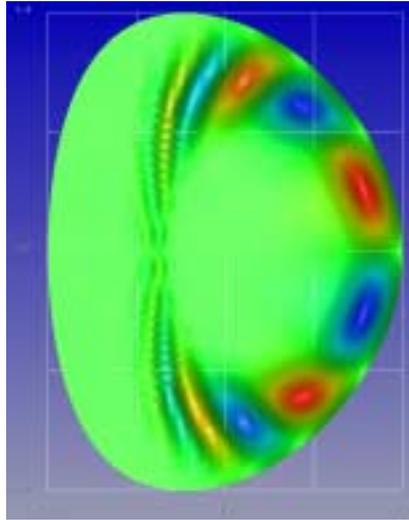
- Large parallel thermal conductivity reduces growth rate of higher n modes.
- ω_i^* term alone does not have stabilizing effect on either mode.

$n=4$ Eigenmodes, Evolving Density

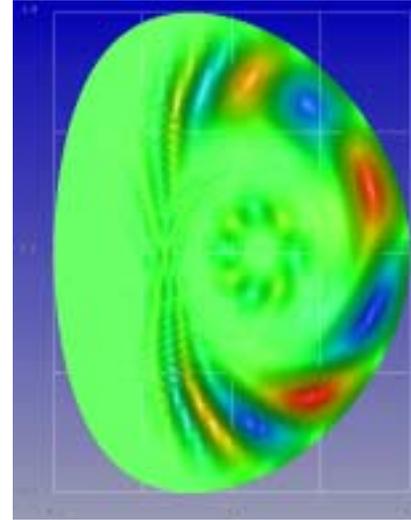
MHD

MHD + ω^*

$\chi_{||} = 0$

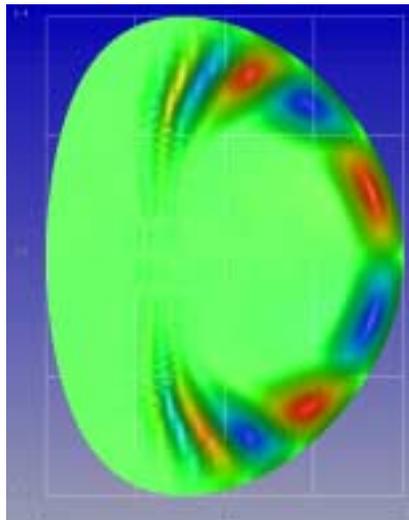


$\gamma\tau_A = 0.0307$

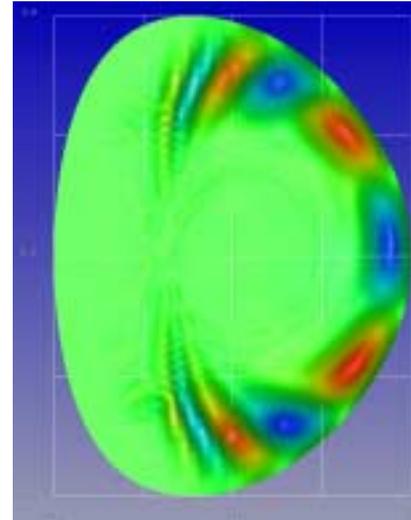


$\gamma\tau_A = 0.0304$

$\chi_{||} \neq 0$



$\gamma\tau_A = 0.0075$



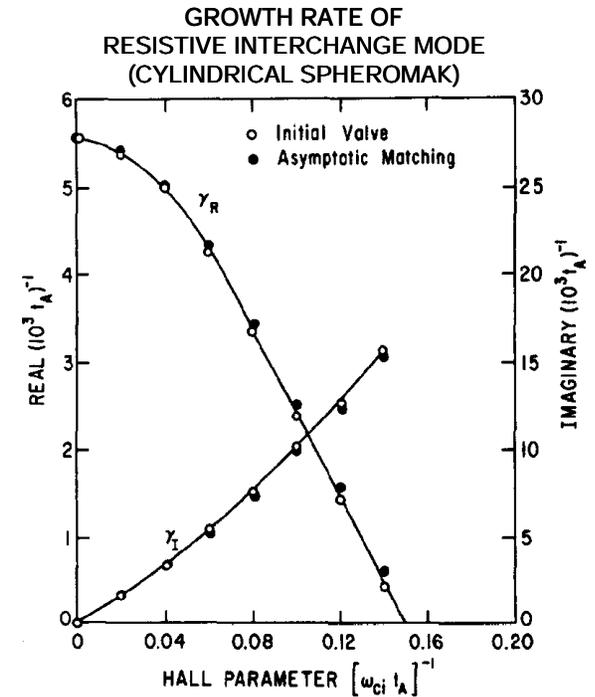
$\gamma\tau_A = 0.0095$

Summary

- All toroidal modes of the $q_{min} = 0.92$ CDX equilibrium are linearly MHD-unstable.
 - $n = 1$ is an internal kink mode
 - $n > 1$ are ballooning instabilities
 - Higher n modes have higher growth rates.
- Nonlinear MHD evolution beginning with just an $n=1$ perturbation disrupts within a sawtooth crash time.
 - High poloidal mode number m components of the $n=1$ mode interact to create islands, stochasticity in outer region.
 - $n = 1$ mode couples to and drives higher n modes at $q=1$ rational surface to create stochasticity in inner region.
- Adding toroidal flow to the model reduces the growth rate of the $n=1$ mode.
- Adding large parallel thermal conductivity (via artificial sound wave) has a stabilizing effect on higher n modes, but not on $n=1$.
- Adding the ω^* term to the MHD equations does not appreciably alter the growth rates of either the $n = 1$ or the $n > 1$ modes.

Future Work

Previous studies¹ indicate that resistive interchange modes can be stabilized with a sufficiently large Hall term:



- Run linear and nonlinear cases with Hall term on.
- Run linear and nonlinear cases with heat conduction on.
- Do convergence studies.

¹J. DeLucia, S.C. Jardin, and A.H. Glasser, *Phys. Fluids* **27** (6), 1470 (1984).