

Electron Transport and Barriers in Reversed Shear Plasmas in TFTR

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Motivation & Outline

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- Core **Ion** thermal transport suppression demonstrated on many machines with reversed or weak magnetic shear (Internal Transport Barriers)
 - in some cases **particle** transport also suppressed
 - likely due to ExB flow-shear suppression of long-wavelength turbulence

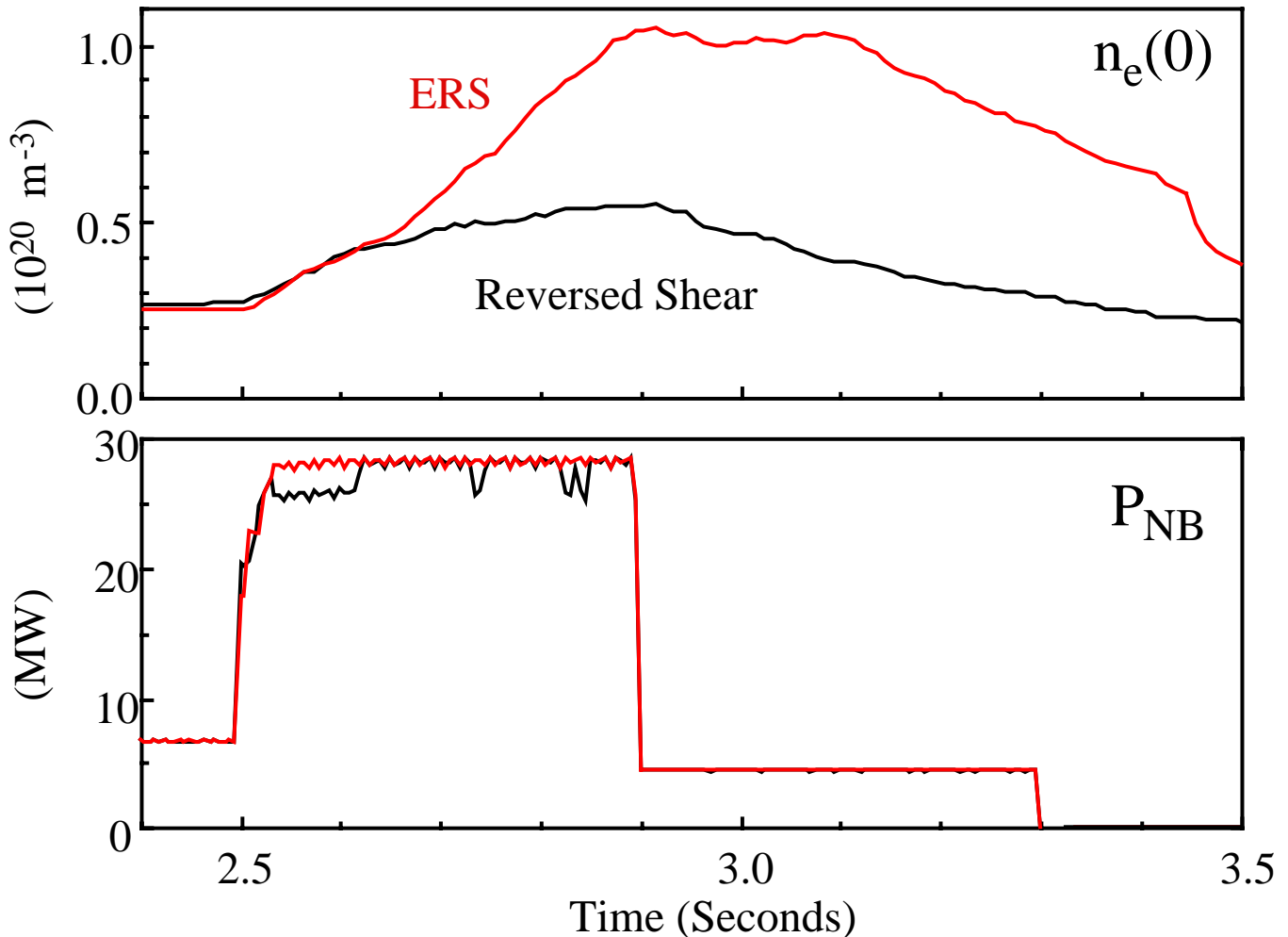
→ basis for Advanced Tokamak and Compact Stellarator Configurations

- **Electron** thermal transport is sometimes reduced, sometimes ~unchanged; always anomalous.
 - need to develop predictive understanding
 - need to understand to try to suppress

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- Brief review of Enhanced Reversed Shear (ERS) transport and profiles
 - Detailed structure of T_e profile, transport
 - Comparisons with theoretical predictions

Two Confinement Regimes Observed with Reversed Core Magnetic Shear

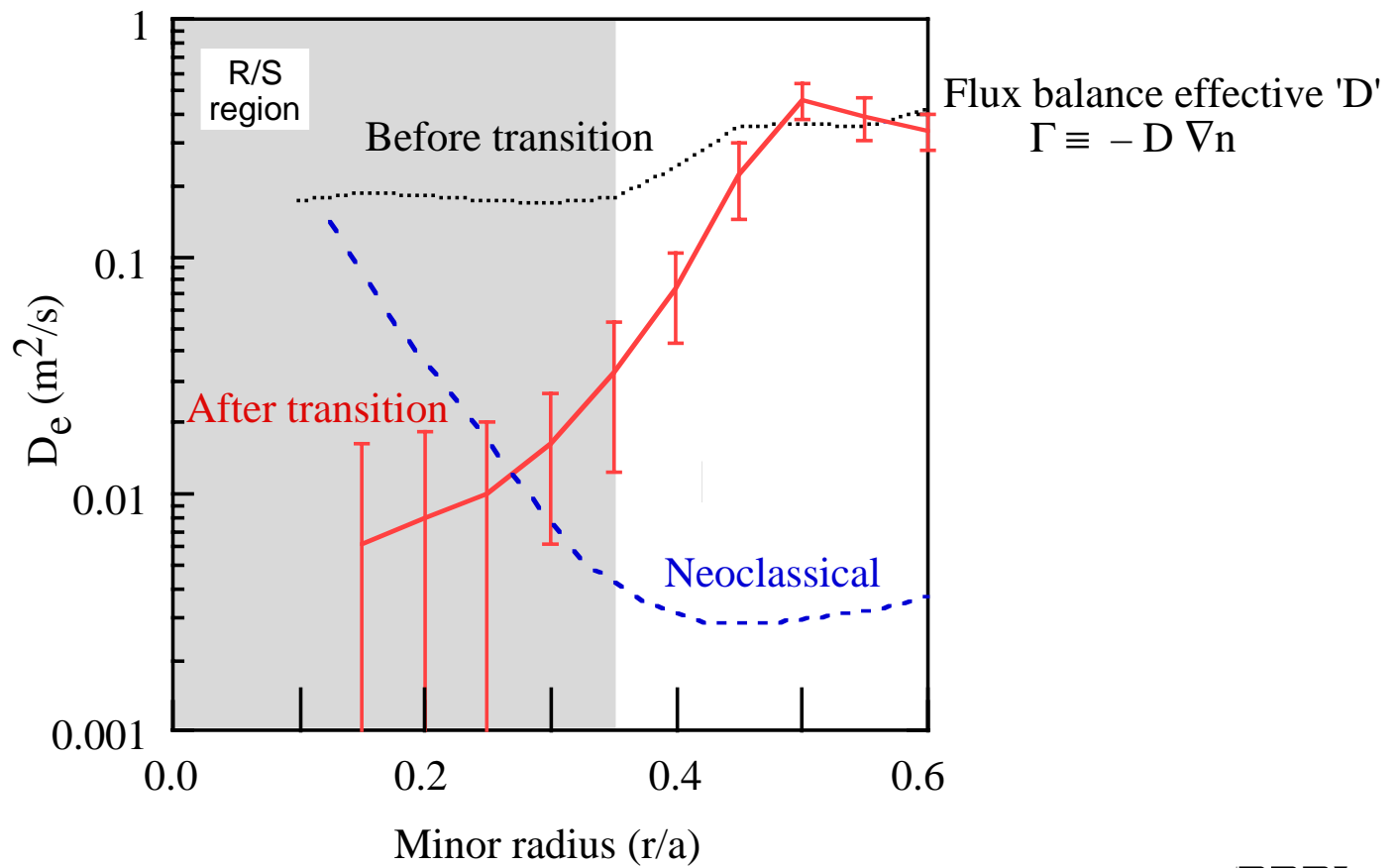
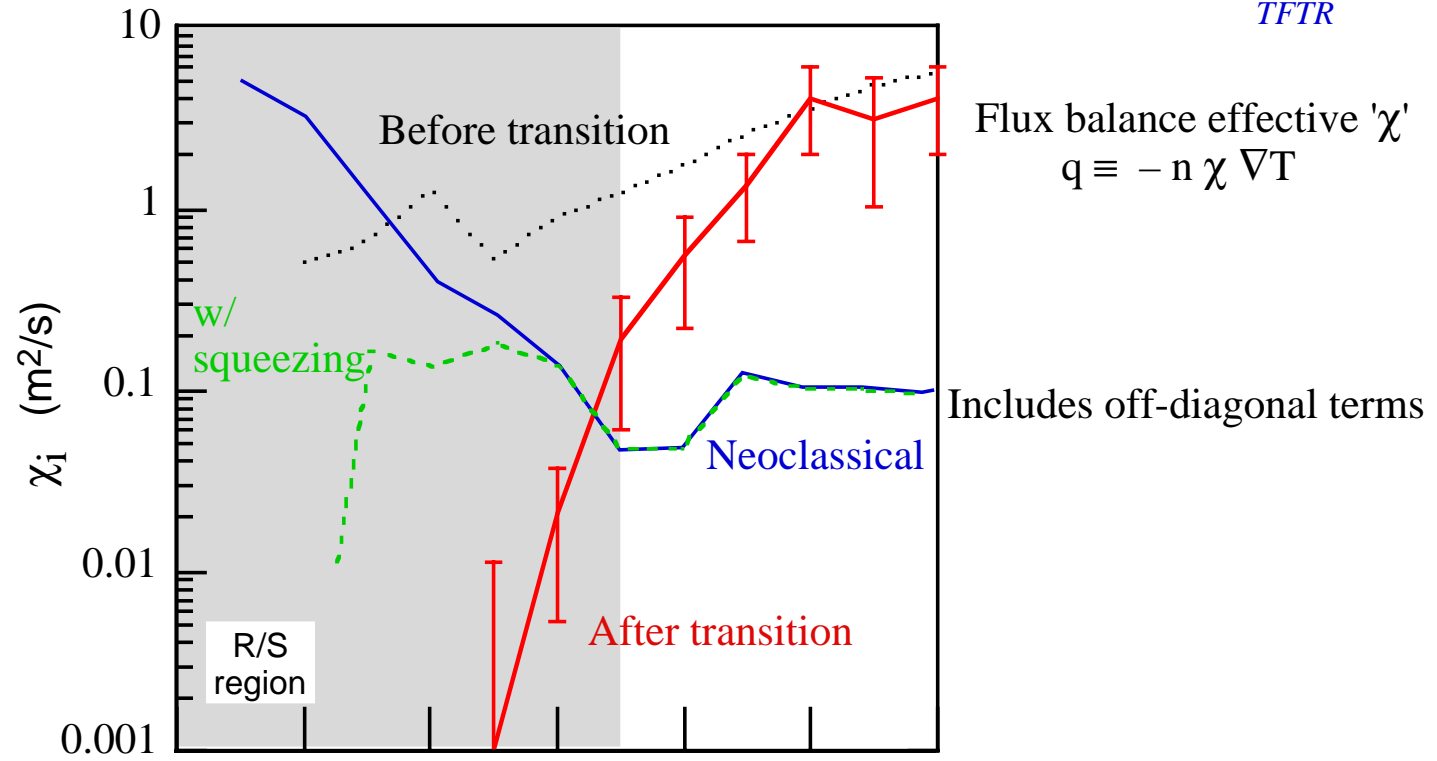
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- Plasmas with reversed or weak magnetic shear can show a sudden transition to reduced core transport of particles and energy (**ERS mode – Enhanced Reversed Shear**)
- **ERS** plasmas develop extremely peaked profiles
– $n_e(0) / \langle n_e \rangle \sim 5$, $p(0) / \langle p \rangle \sim 8$
– $p(0)$ up to ~ 6 atmospheres
- **ERS** plasmas show extreme hysteresis
High central density can be maintained with ~ 5 MW of NBI

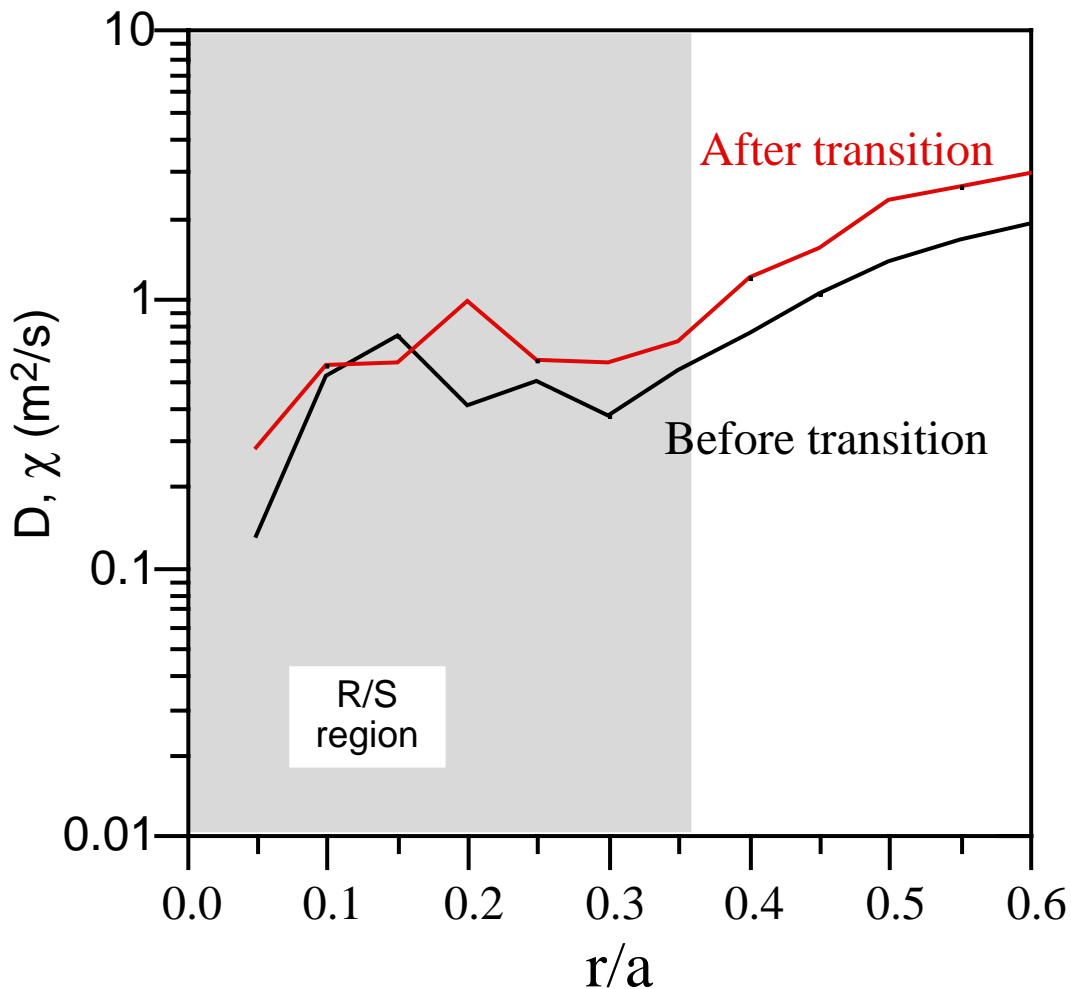
χ_i and D_e are Sharply Reduced after Transition to below neoclassical level

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χ_e is Not Reduced after Transition

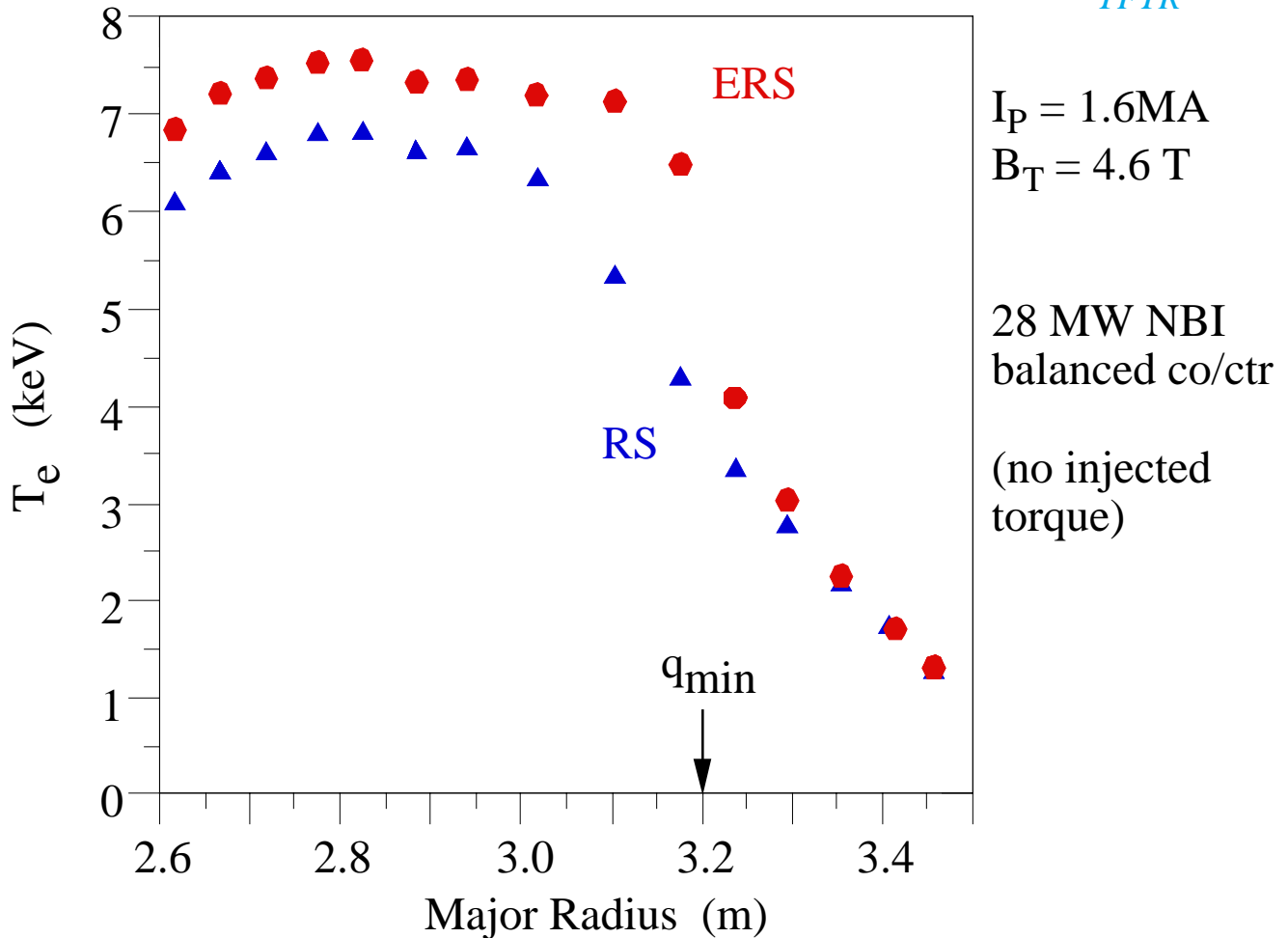
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- Is χ_e larger in ERS?
or \sim equal within the uncertainties
 - uncertainties due to large time derivatives and profile measurements of ∇T
- Why?

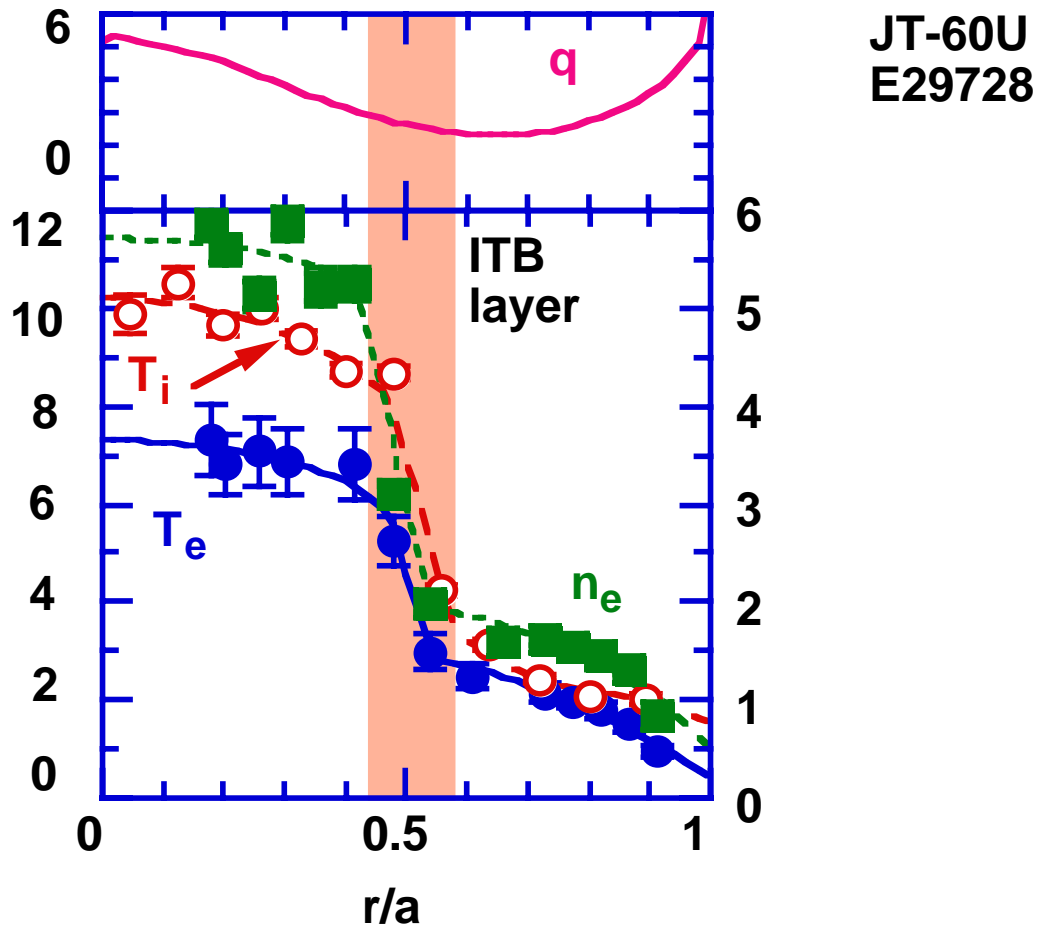
ERS Te Profiles Are Broader, Squarer than RS

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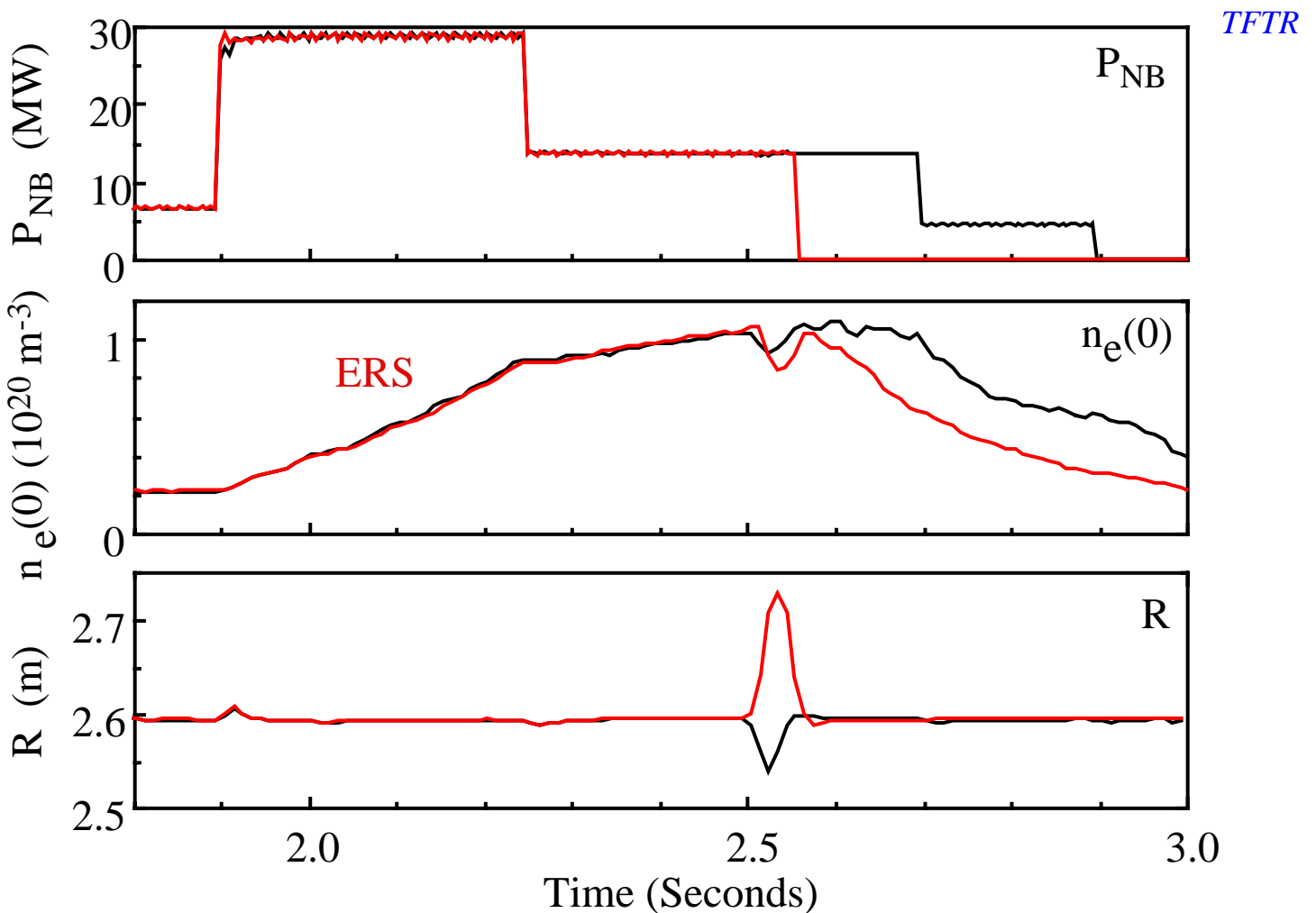
- T_e measured by ECE grating polychromator, each channel cross calibrated to Michelson interferometer
- Adjacent identical shots, except one transitions to ERS
- ERS T_e profile shape develops ~ 0.1 sec after transition
- Observed on all ERS, Type I transitions
Not observed on Type II transitions
- Gradient in reversed shear region is within systematic uncertainty of diagnostic calibration

Similar flattening is Observed on JT-60U & AUG

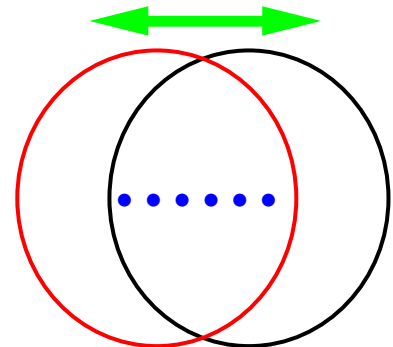


- see H. Shirai et al, Phys. Plasmas 5, 1712 (1998) and 1998 IAEA, Yokohama, Japan.
- also observed on T_e only in Asdex-Upgrade, see paper by R. Wolf et al, 1998 IAEA, Yokohama, Japan.

Plasma 'Jog' Used to Improve Radial Resolution

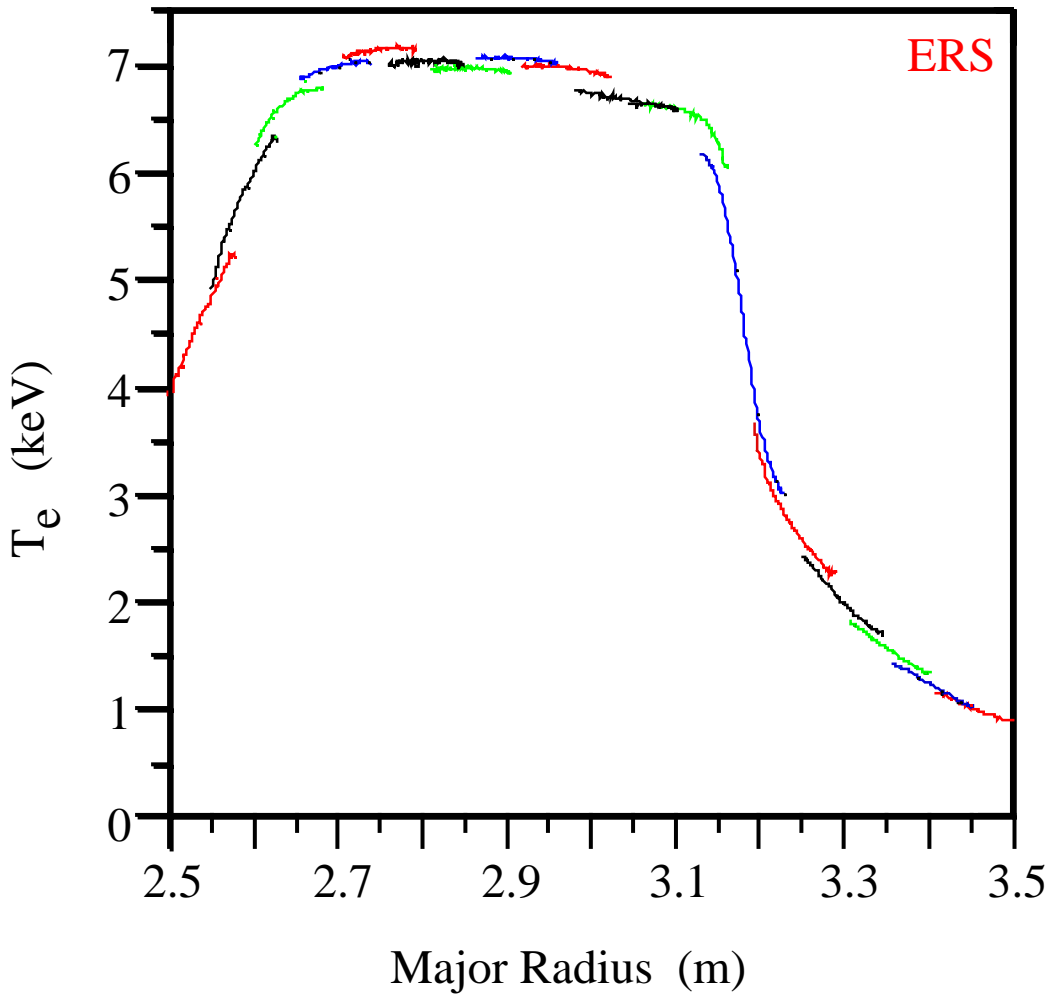


- Sweep plasma past fixed detectors to improve radial resolution both inward and outward motions used
- Maximum velocity ~ 3 m/sec; sample rate up to 500 kHz
- Jog during reduced power 'postlude' plasma near steady state
- Gives single detector measurement of gradients \Rightarrow reduced systematic uncertainty



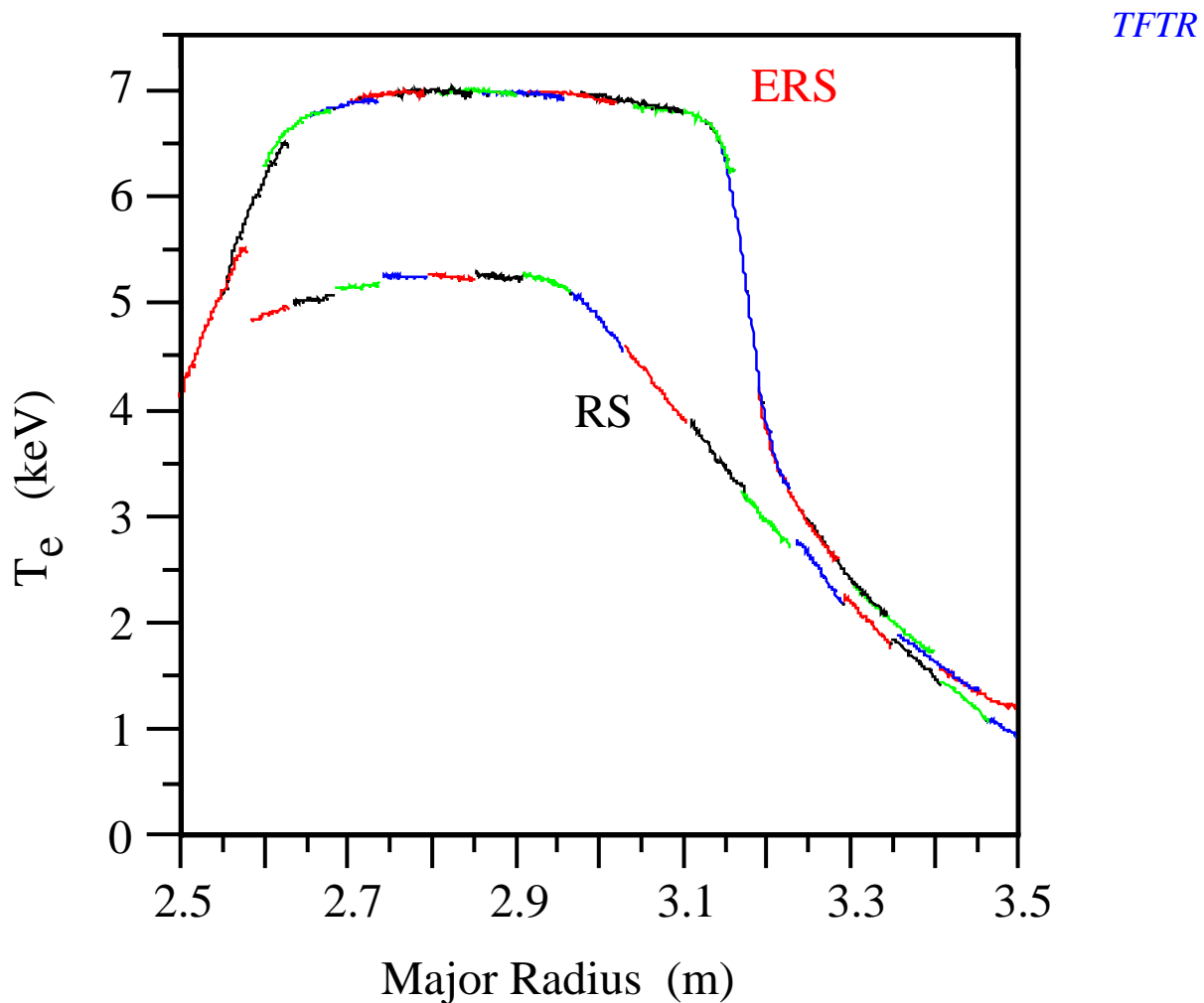
Jog'd ECE Shows Core T_e is Flat in ERS !

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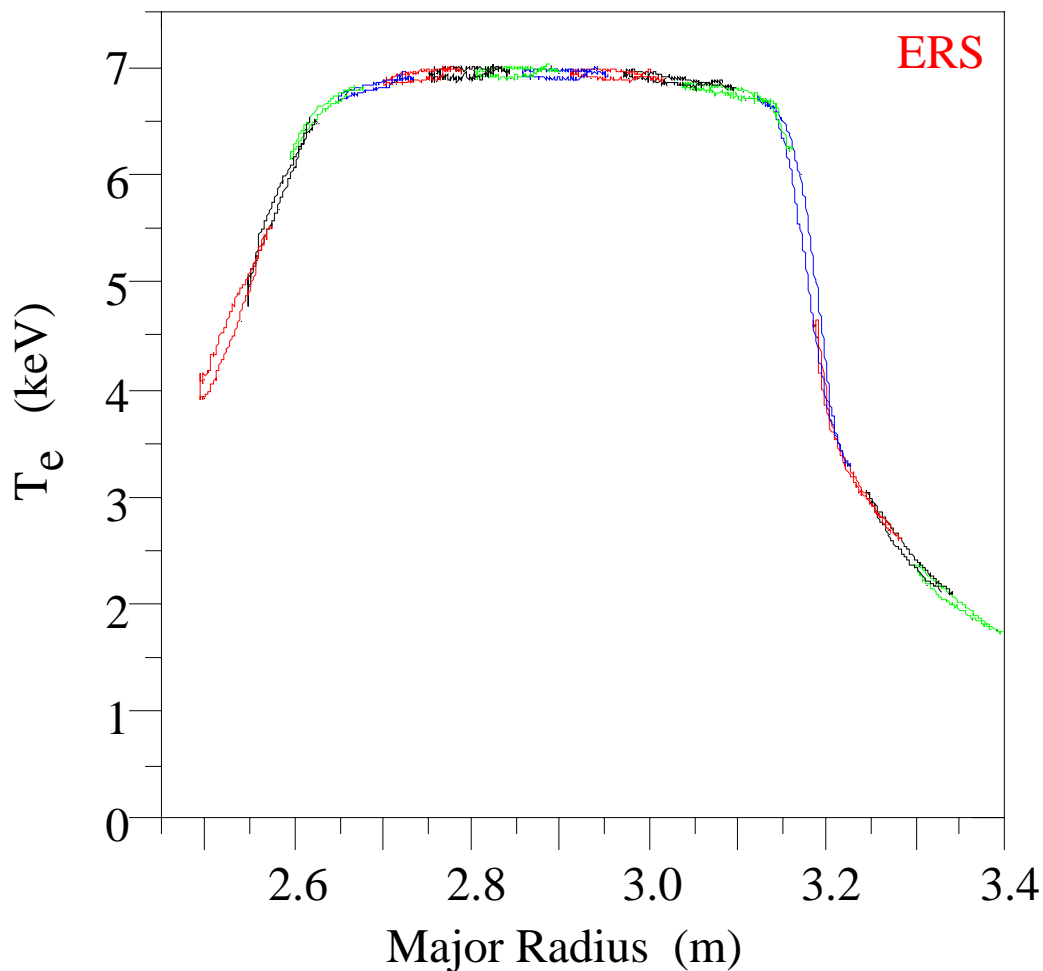
- Each colored segment shows the trajectory of a single detector during the plasma motion, mapped to the pre-motion position
- Similar profiles obtained at low and high B ERS shots
- Note corner in T_e profile near shear-reversal point

Jog'd ECE Shows Core T_e is Flat in ERS !



- Detector calibrations corrected
correction factors averaged over 4 shots
in core-region < 4 % corrections, typ. ~ 2%
- ERS Not profile consistent !!
- Most of the ERS core ∇T_e in the standard analysis is
from systematic errors in detector calibration
- RS profile similar to jog-measured T_e profile for supershots
profile shape ~ similar to other regimes
L-mode, supershot, ohmic,...

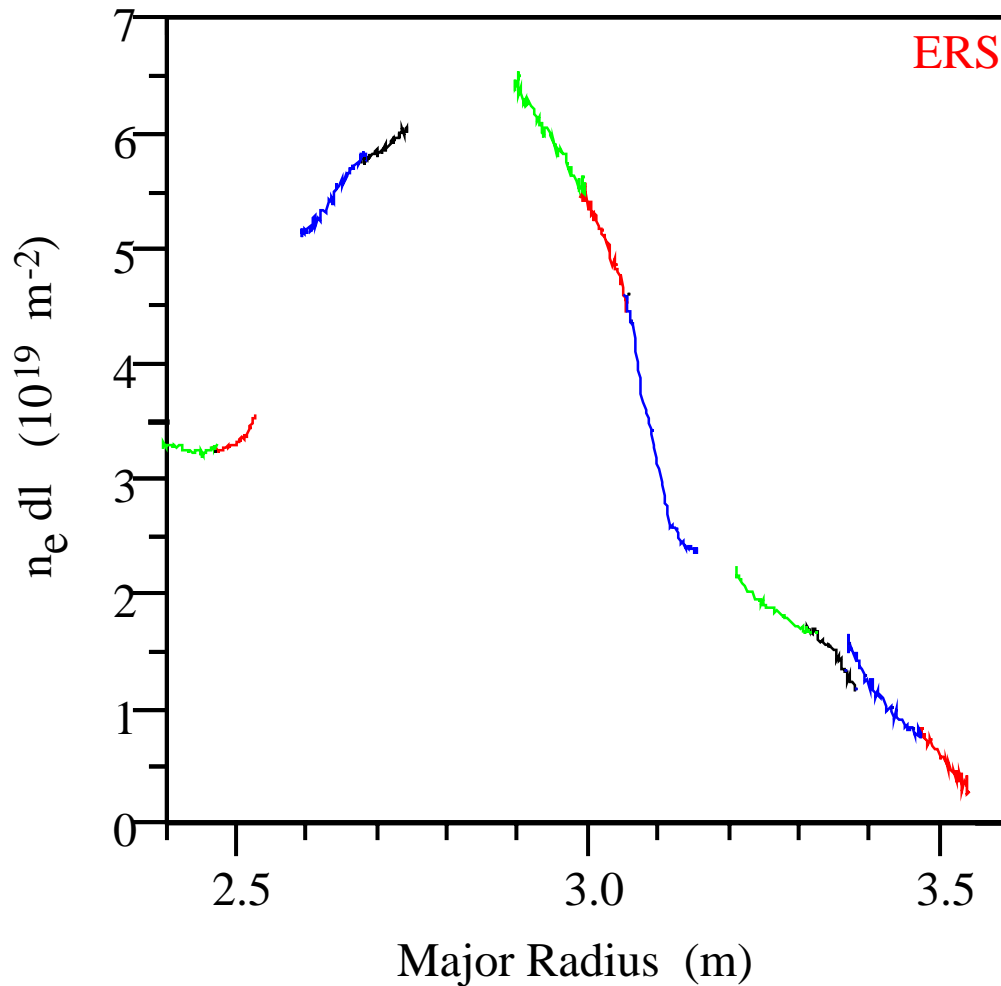
Forward/Back Jog Shows Plasma is not Damaged



- Shows full jog, moving out and then back to original position
elapsed time: 70 msec
- Individual detectors trace loops, height ~ 50 eV in core
- Loops close
 \Rightarrow likely residual problems with plasma position measurement

Jog'd Interferometry: Core n_e is Peaked

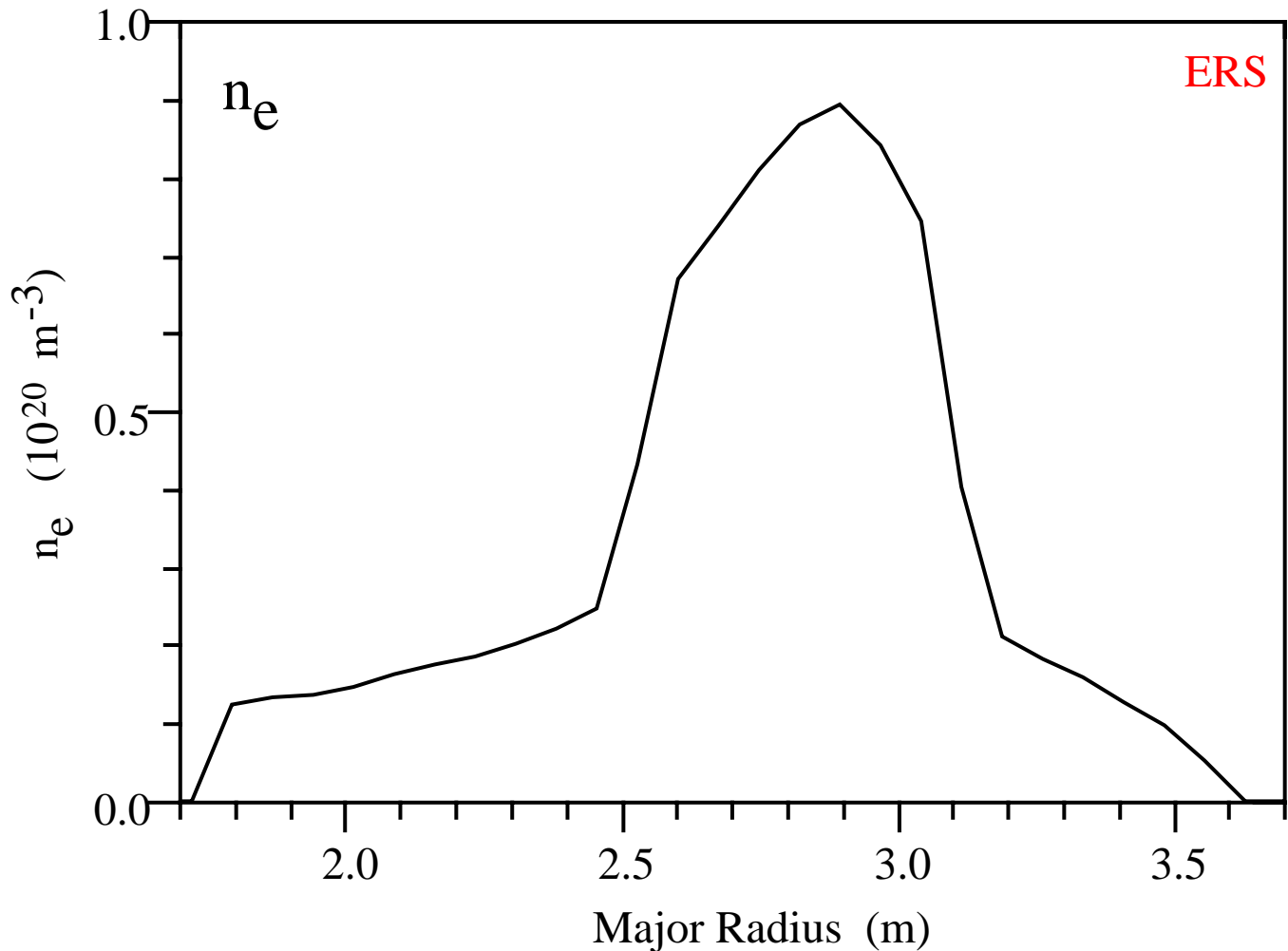
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- Combination of two identical shots, one jog'd inward, one jog'd outwards
- Also: jog'd core Visible Bremsstrahlung emission $\propto n_e^2 T_e^{1/2} Z_{\text{eff}}$ very peaked, consistent with peaked n_e

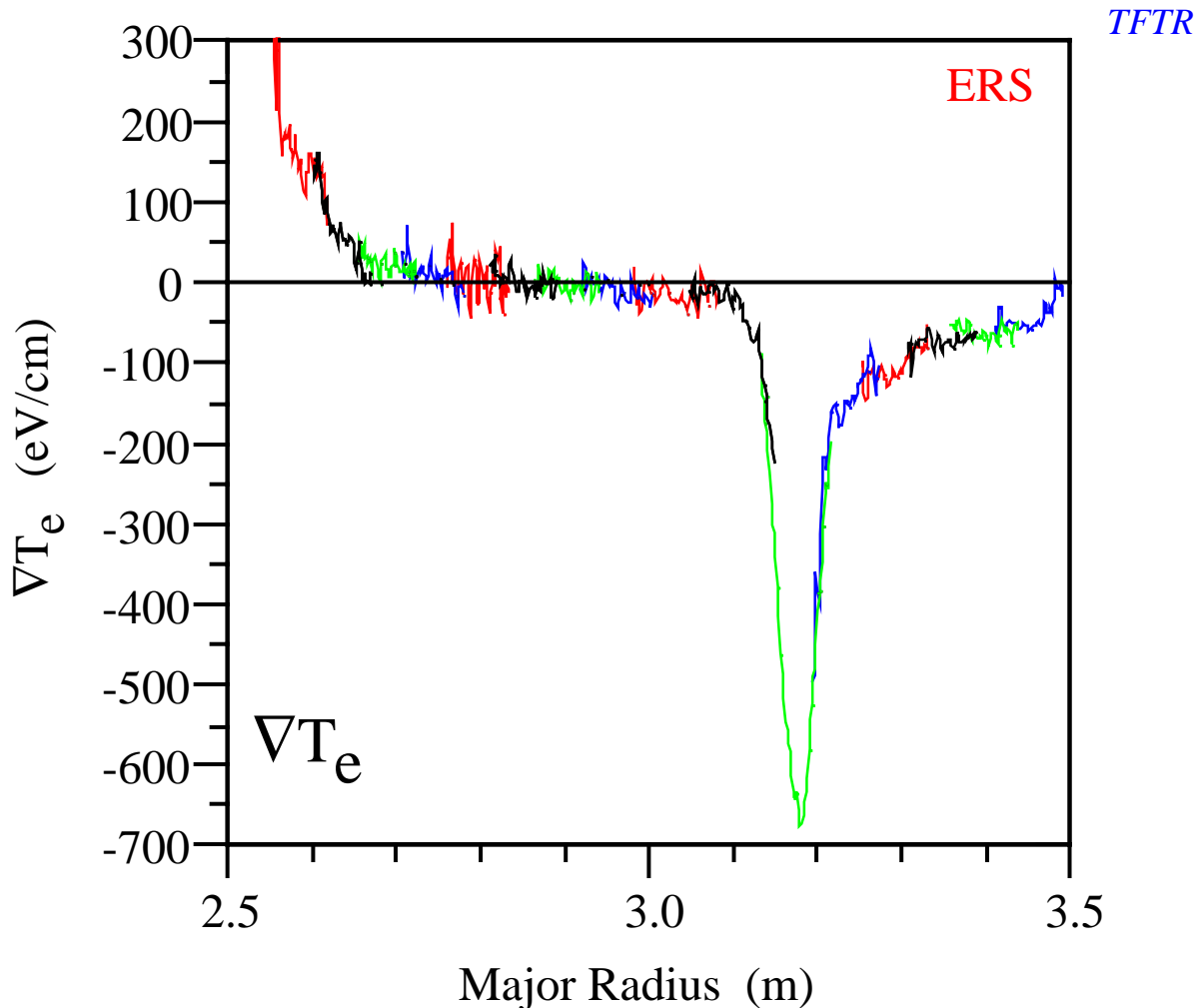
Jog'd Interferometry: Core n_e is Peaked

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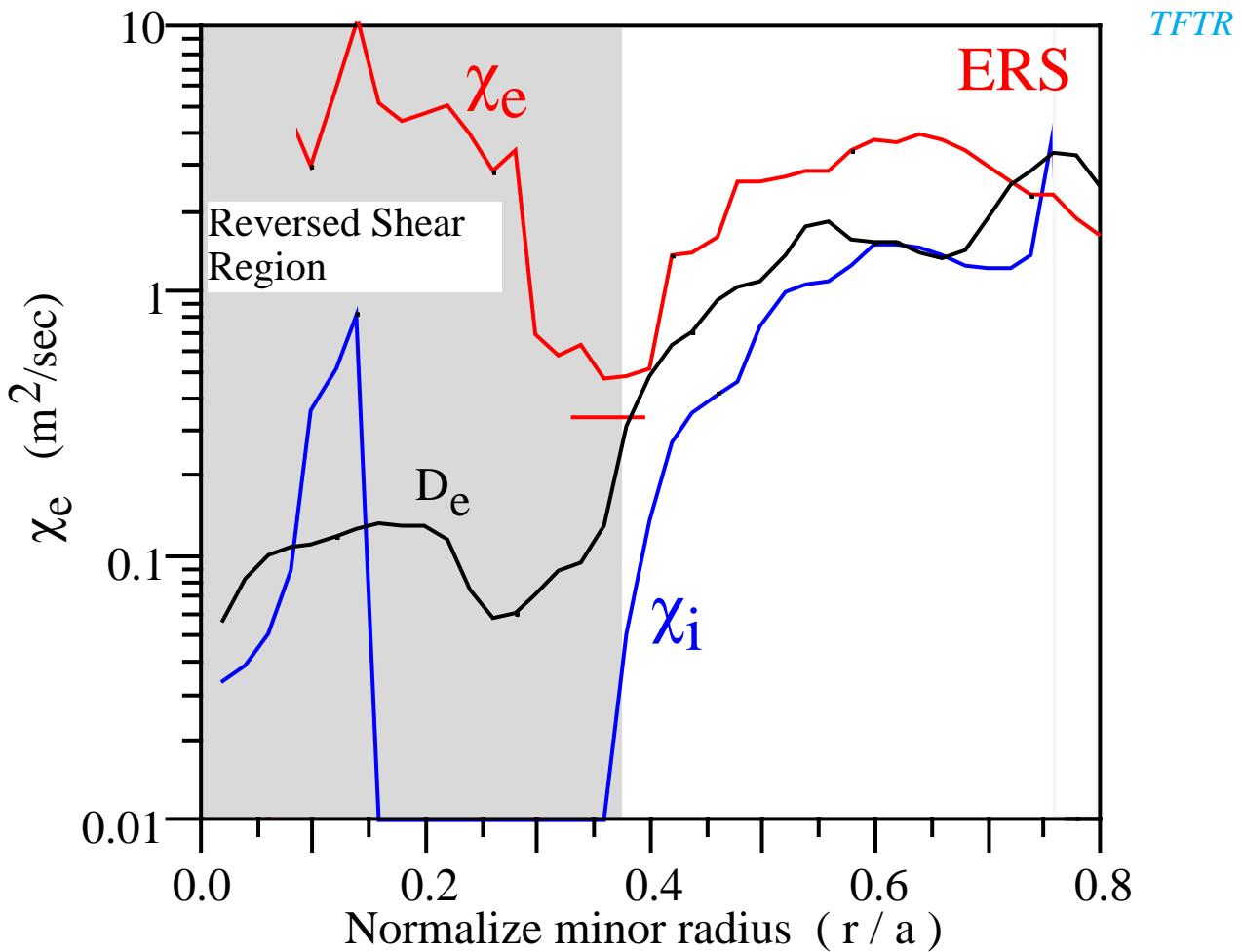
- From Abel Inversion of jog'd $n_e l$ profile
- Peak n_e is $\sim 10\%$ lower than non-jog standard analysis
- Outer edge of steep gradient is at same location as for T_e
- T_i and v_ϕ measurement averaging time too long for jog-technique
 - however, those profiles are peaked outside the error bars

Core ∇T_e is Extremely Low in ERS



- ∇T_e measured by a single detector in each spatial region, from change in T_e during plasma motion
– minimize systematic uncertainty
- $\nabla T_e \sim 15$ eV/cm in core, averaged over one cm.
- Very high ∇T_e near shear-reversal surface
- transition from ~ 60 eV/cm to > 300 eV/cm
with < 2 cm separation !
At the limit of instrumental spatial (frequency) resolution

ERS χ_e is Increased in Core, Reduced in Barrier



- Power balance analysis of jog'd profiles for T_e and n_e
- Analysis during near-steady-state 'postlude'
Reduced uncertainties relative to earlier analysis with large time-derivatives. Largest uncertainty now, for thermal transport, is due to T_i and the ripple modeling.
- $\chi_e / D_e \sim 50$ in the core !!
~ 4 in the barrier
- $\chi_e / \chi_i \sim 100$ in the core

Most Plausible Explanation:

Stochastic B !

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In core: $\chi_e / D_e \sim 50$, $\chi_e / \chi_i > 100$

- Difficult to understand how electrostatic modes could give such large χ_e / D_e
 - Rechester-Rosenbluth: $\chi_e / D_e \sim m^{1/2} = 67$
would require $\tilde{B}_r / B \sim 3 \times 10^{-4}$
- Must be high $k_{\perp} \rho_i$ turbulence to avoid ion transport via ion orbit averaging.
 - $\rho_i = 0.3 - 0.5$ cm

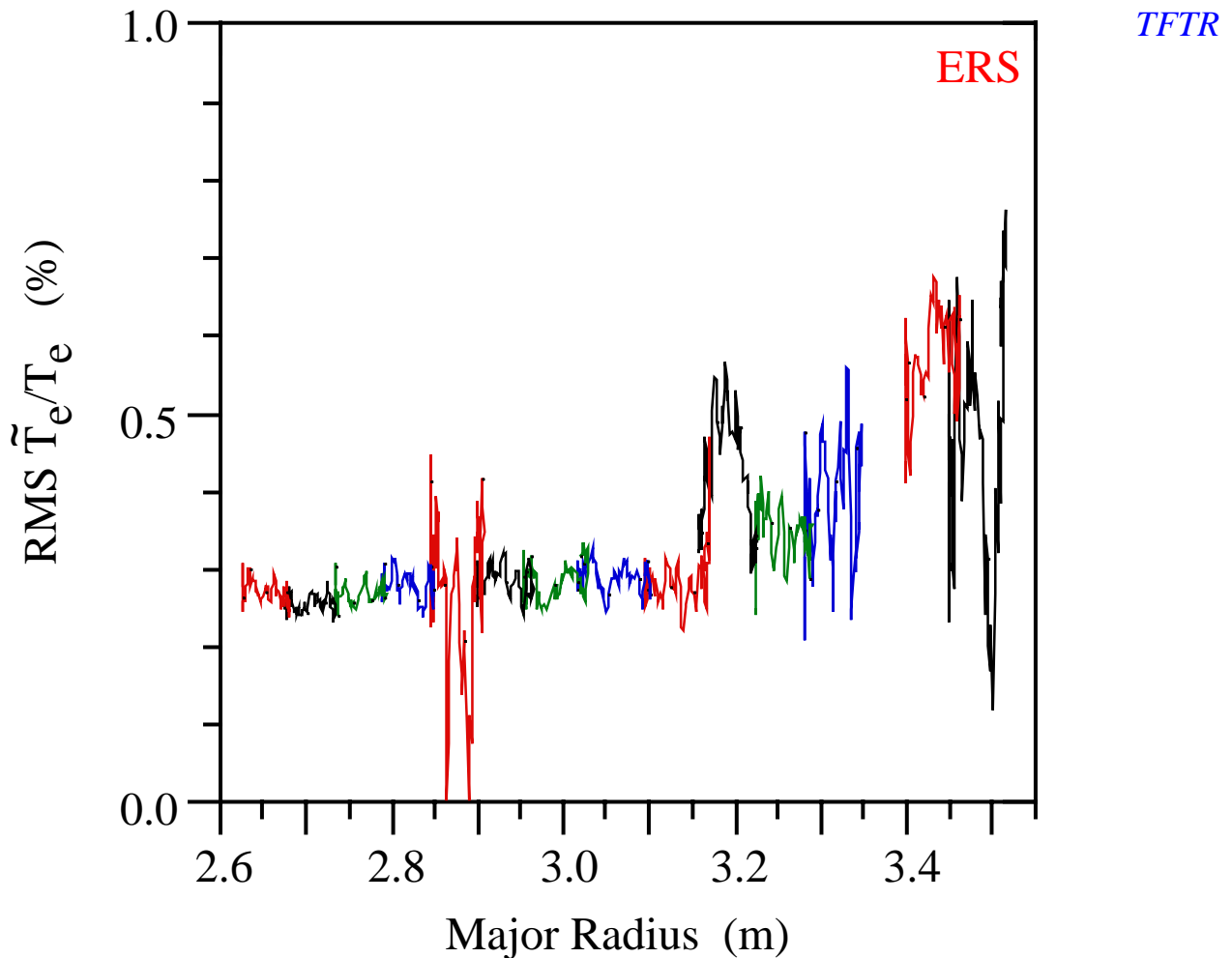
What Instability?

Fundamental problem: What gradient provides drive?
 T_e is being transported, but $\nabla T_e \sim 0$.

Substantial ∇T_i and ∇n_e , but they are not being dissipated.

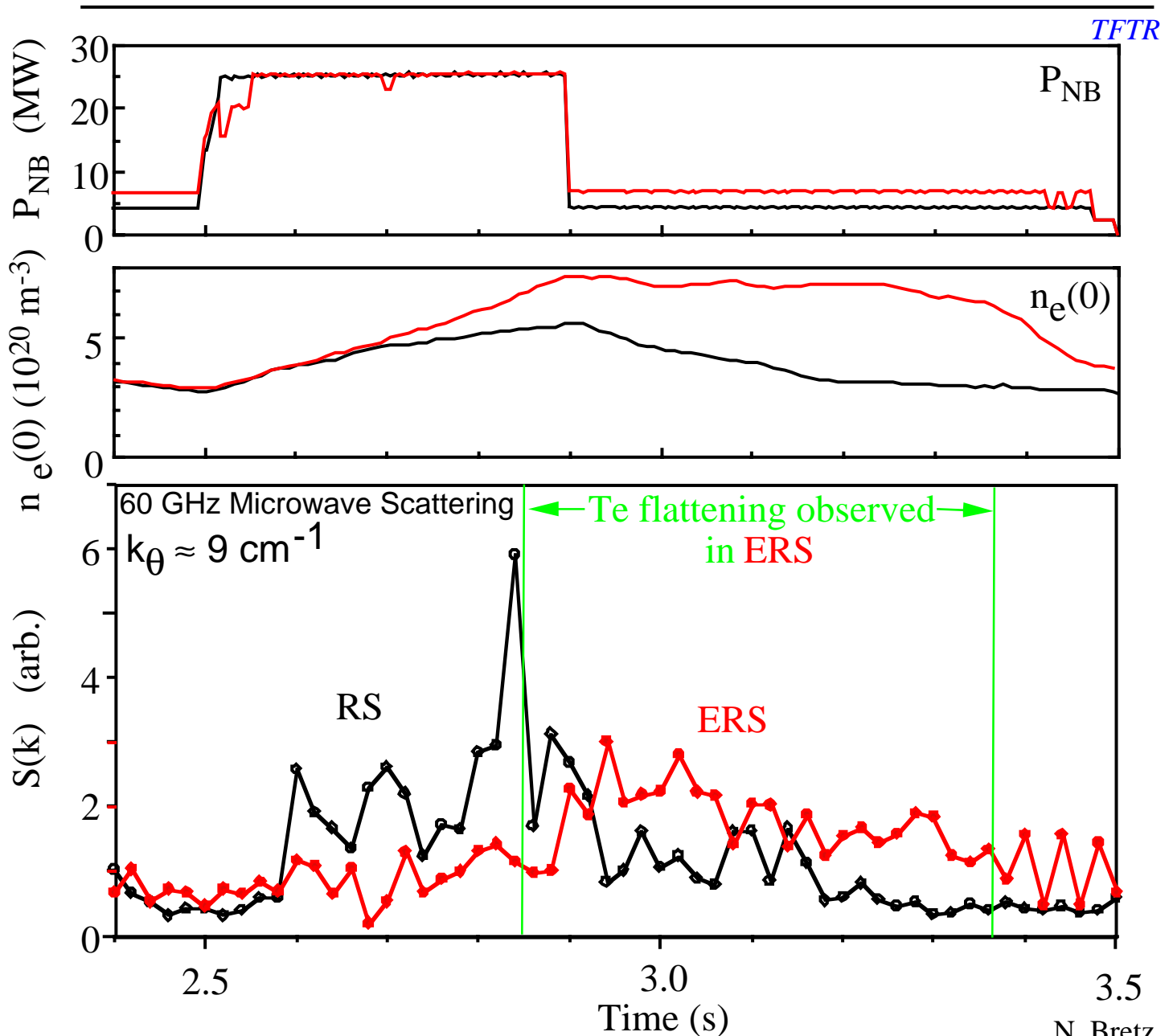
- η_e or Electron Temperature Gradient-Mode (ETG)
has no drive
 $L_{T_e} \sim 3.5$ m $> R = 2.6$ m
 $\eta_e < 0.1$

No Apparent Core T_e Fluctuations



- RMS T_e fluctuation in 1ms intervals, mapped to pre-motion position
- Amplifier noise subtracted. Channel at $R \sim 2.87\text{m}$ very noisy
- In core, measured T_e fluctuations consistent with expected ECE blackbody noise.
 - As expected for convective modes with $\nabla T_e \sim 0$.
- Very small 60 ± 10 kHz fluctuations near shear reversal also observed on reflectometer, Mirnov array $n = -3$, $m = 5 - 7$ from Mirnov displacement ≤ 1 mm

No Strong High- k_θ Density Fluctuations



- Integrating shifted scattered spectrum, ignoring unshifted line center
- Core T_e flattening observed in **ERS** case from ~ 2.85 s thru ~ 3.37 s
- Scattered signal is proportional to density \Rightarrow fluctuations similar in postlude phase for **ERS** and RS
- For $k_\theta \sim 2$ cm^{-1} , RS fluctuation level higher than in ERS, similar to reflectometer measurements

Core χ_e : Possible Models

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- No unstable ideal MHD modes found (J. Manickam)
- Electromagnetic skin-depth
$$\frac{c^2}{\omega_{pe}^2} \frac{v_{Te}}{R q} \sim 1 \text{ m}^2/\text{s} \ll \chi_e$$
- Resistive Interchange: $D_R \sim 1$ both ERS and RS (M. Hughes)
 - mode is barely unstable, easily stabilized kinetically
 - at higher β values, $D_R \sim 70$ has been calculated
 - appears to be uncorrelated with transport
- Resistive pressure gradient turbulence [e.g. Carreras-Diamond, Phys. Fluids B1, 1017 (1989)]
 - predicts large enough transport,
 - Gives same level for ERS and RS: no discrimination.
 - Should be re-examined, including flow shear
- GS code by M. Kotschenreuther [CPC 88, 128 (1995)]
 - Comprehensive linear stability analysis of full gyrokinetic equations (Antonsen and Lane) in ballooning representation.
 - Should address all short wavelength electrostatic and electromagnetic modes, including resistive interchange, resistive ballooning, micro-tearing, ETG, ...
 - finds no linearly-unstable modes !!

Need to look at non-linear instabilities
and (possibly) tearing modes (below neoclassical stabilization)

Transport in the Barrier: **ETG?**

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The GS code has also analyzed the microstability of the electron transport barrier region, where χ_e is suppressed in ERS.

Preliminary results:

(some cross checking in progress)

- unstable mode with $k_{\perp}\rho_i \sim 100$, $k_{\perp}\rho_e \sim 1$
- ∇T_e is $\sim 30\%$ higher than the critical gradient for ETG

ETG is analogous to ITG (with $i \leftrightarrow e$), but with strictly adiabatic ions due to the very high $k_{\perp}\rho_i$

\Rightarrow strong gyro-averaging.

From the analogous ITG calculation, can estimate

- $\chi_e(\text{ETG}) > 30 D_{\text{mix}} \sim 3 \chi_e(\text{Exp}) \sim 100 \chi_i(\text{ETG})$

where D_{mix} is the Kadomtsev mixing length estimate

$$D_{\text{mix}} \sim \gamma / k_{\perp}^2$$

\Rightarrow ETG may be strong enough to enforce marginality
uncertainties need to be investigated

Summary

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- T_e is very flat in the core of ERS plasmas (inside the shear reversal surface)
- χ_e in the ERS core is much larger than in RS
 - in contrast to D_e , χ_i , χ_ϕ
- $\chi_e \sim 50 D_e$, $\sim 100 \chi_i$ in core!
 - may imply that the core magnetic field is stochastic, on a very fine scale
 - dissipating ∇T_e but no ∇T_e drive!
what is driving the turbulence?
 - no instabilities found by comprehensive code
- ∇T_e is locally very large (~ 680 eV/cm) in ~ 5 cm layer near reversal surface
- χ_e is 4 times lower in ERS than RS in layer near reversal surface
 - clear electron thermal transport
 - may be limited by ETG