

Third Year Status
on
SciDAC Center for Gyrokinetic Particle Simulation
of Turbulence Transport in Burning Plasmas

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Presented at
PSACI PAC Meeting₁
June 2007

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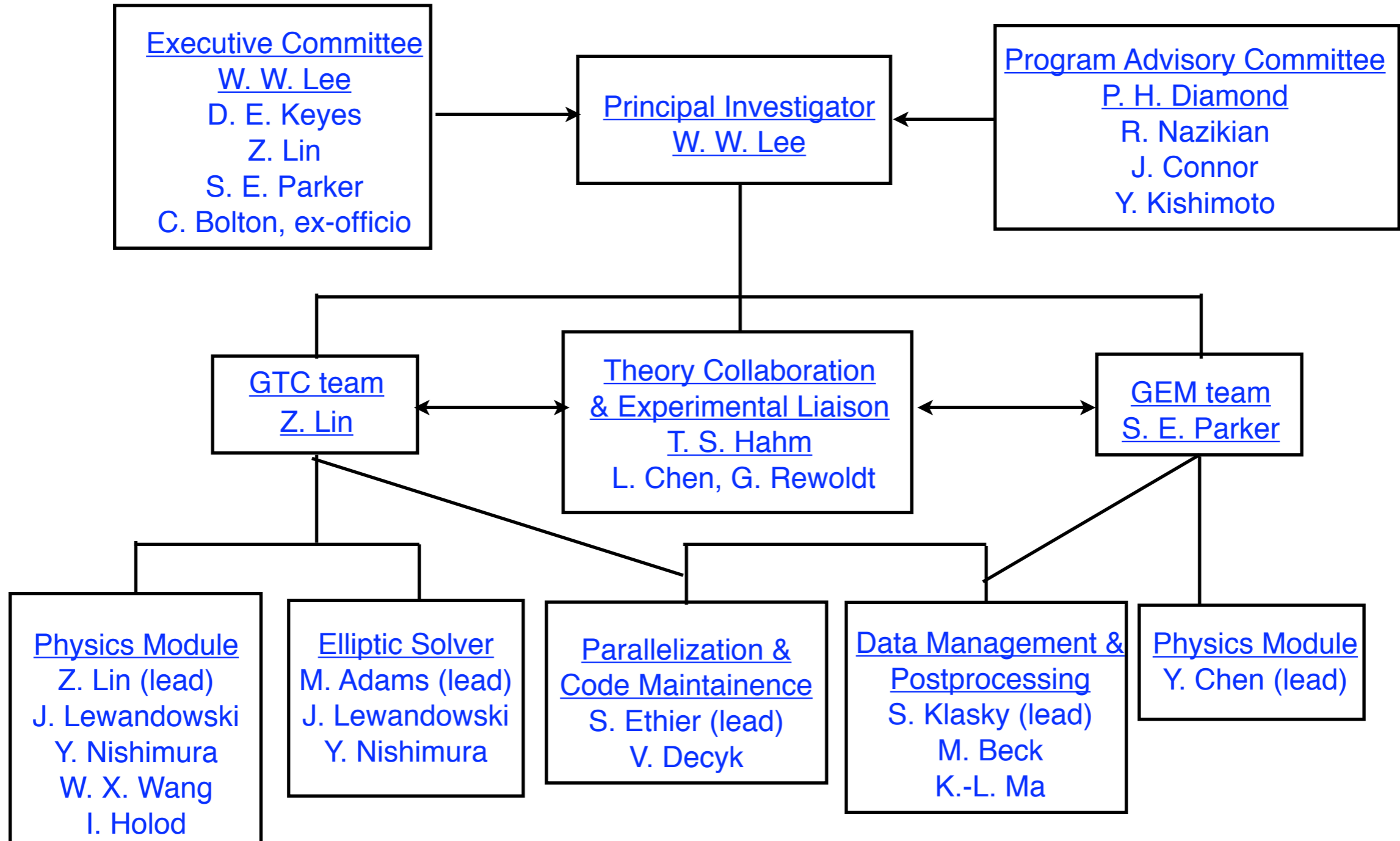
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Outline

- Center Activities
 - Fall Meeting
 - Invited talks and Publications
- Code Development
- Code Validation
- Physics Investigations
- The noise and growing weight issues
- Conclusions

Center for Gyrokinetic Particle Simulation of Turbulent Transport in Burning Plasmas



FALL MEETING

SciDAC Center for Gyrokinetic Particle Simulation
of Turbulent Transport in Burning Plasmas

Room 407-408
Philadelphia Marriot Hotel, Philadelphia, PA

November 2, 2006

- 7:00P Lee - Opening remarks
- 7:10P Wang - Shaped plasma simulations and future plans
- 7:25P Ethier - GTC performance and optimization issues
- 7:40P Lin - Status and plan in global GTC turbulence simulation
- 7:55P Nishimura - Shear Alfvén wave studies in electromagnetic global gyrokinetic simulation of tokamak plasmas
- 8:10P Parker - ETG Convergence Studies, GEM Team status
- 8:25P Y. Chen - The growing weight problem
- 8:40P Rewoldt -- Application of GEM code for experimentally-realistic tokamak cases
- 8:55P Hahm - Theory team status and plan
- 9:10P Coffee break
- 9:25P Holod - Particle noise-driven flux in GTC simulations
- 9:35P Xiao - Theory of zonal flow residual level with arbitrary wavelength and collisionality
- 9:45P Jenkins – Particle noise issues
- 9:55P Diamond – Concluding remarks
- 10:05P Klasky - Data Management, Visualization and MPP issues
- 10:15P Plans for re-competition
- 11:15P Recess

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Invited Talks and Review Papers

- Wang, W.X., T.S. Hahm, G. Rewoldt, J. Manickam and W.M. Tang, “Gyrokinetic studies of Nonlocal Properties of Turbulence-driven and Neoclassical Transport”, 21th IAEA Conference on Plasma Physics and Controlled Nuclear Fusion Research, (Chengdu, China, 2006)
- Lee, W.W., S. Ethier, T. G. Jenkins, W. X. Wang, J. L. V. Lewandowski, G. Rewoldt, W. M. Tang, S. E. Parker, Y. Chen, and Z.Lin, 21th IAEA Conference on Plasma Physics and Controlled Nuclear Fusion Research, (Chengdu, China, 2006)
- Lee, W.W., S. Ethier, W. X. Wang, W. M. Tang and S. Klasky, “Gyrokinetic particle simulation of fusion plasmas: path to petascale computing”, Presented at SciDAC 2006, Denver CO., J. of Phys.: Conference Series **46**, 73 (2006).
- Brizard, A.J., and T.S. Hahm, “Foundations of Nonlinear Gyrokinetic Theory”, Rev. Mod. Phys. **79**, 421 (2007).

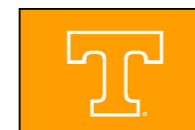
Publications

Review of Modern Physics: **1**
Physics of Plasmas: **11** published, **2** submitted
Journal of Computational Physics: **2** published
IAEA: **4** published
Other Journals: **3** published
Conference Proceedings: **11** published

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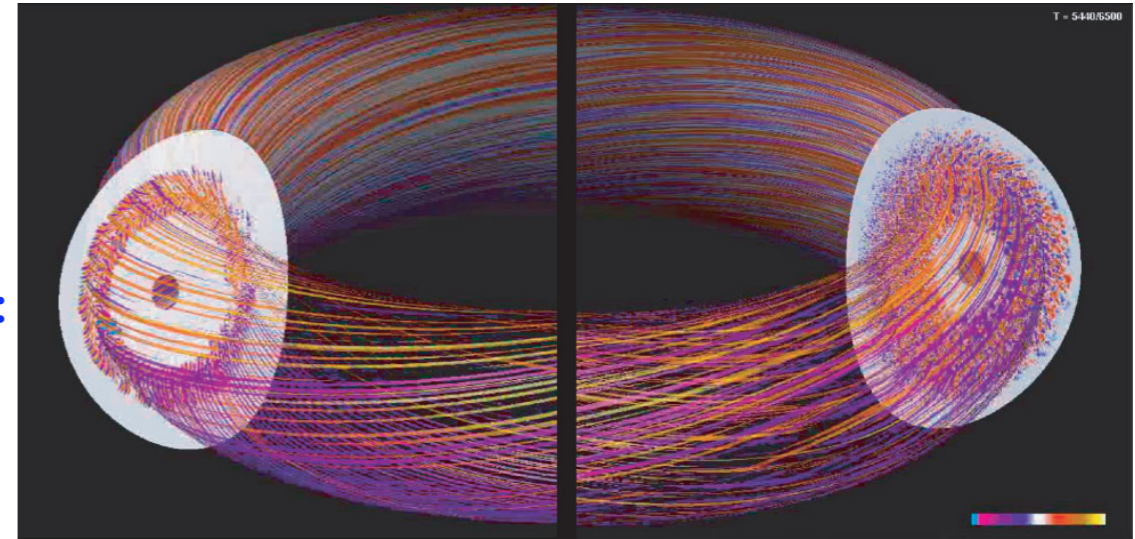
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Code Development

• GTS

- A global code for turbulence transport simulations
- Shaped plasma in general geometry interface with TRANSP, JSOLVER and ESC
- Electron dynamics based on the split-weight scheme: δh , non-adiabatic part of δf
- GK Poisson's equation is solved simultaneously for zonal flows and perturbed potentials



GTS

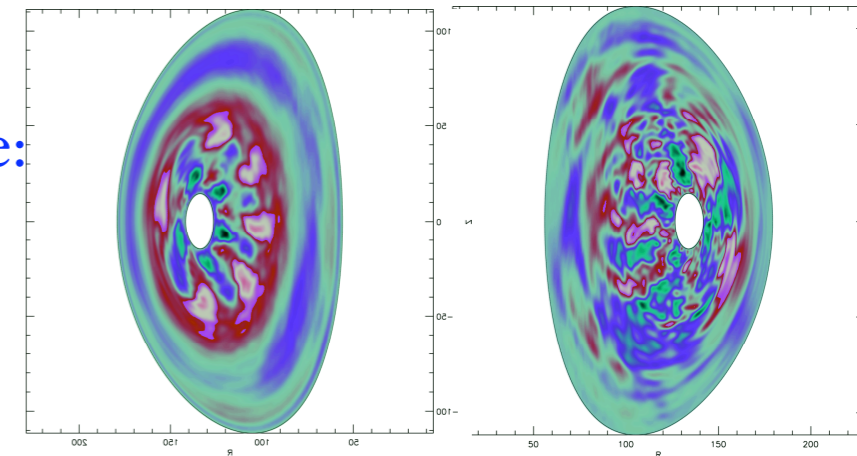
• GTC

- Adiabatic electron version for high performance computing
- Electrostatic electron dynamics based on the hybrid scheme
- Electromagnetic electron dynamics based on the hybrid scheme:

$$|\omega/k_{\parallel}v_{\parallel}| \ll 1$$

• GTC-neo -- For neoclassical transport simulations in

- General toroidal geometry
- Fully operational collision operators



GEM

• GEM

- A wedge code with multi-ion species for turbulence and gyrokinetic MHD simulations
- Shaped plasma in general geometry with interface with TRANSP and JSOLVER

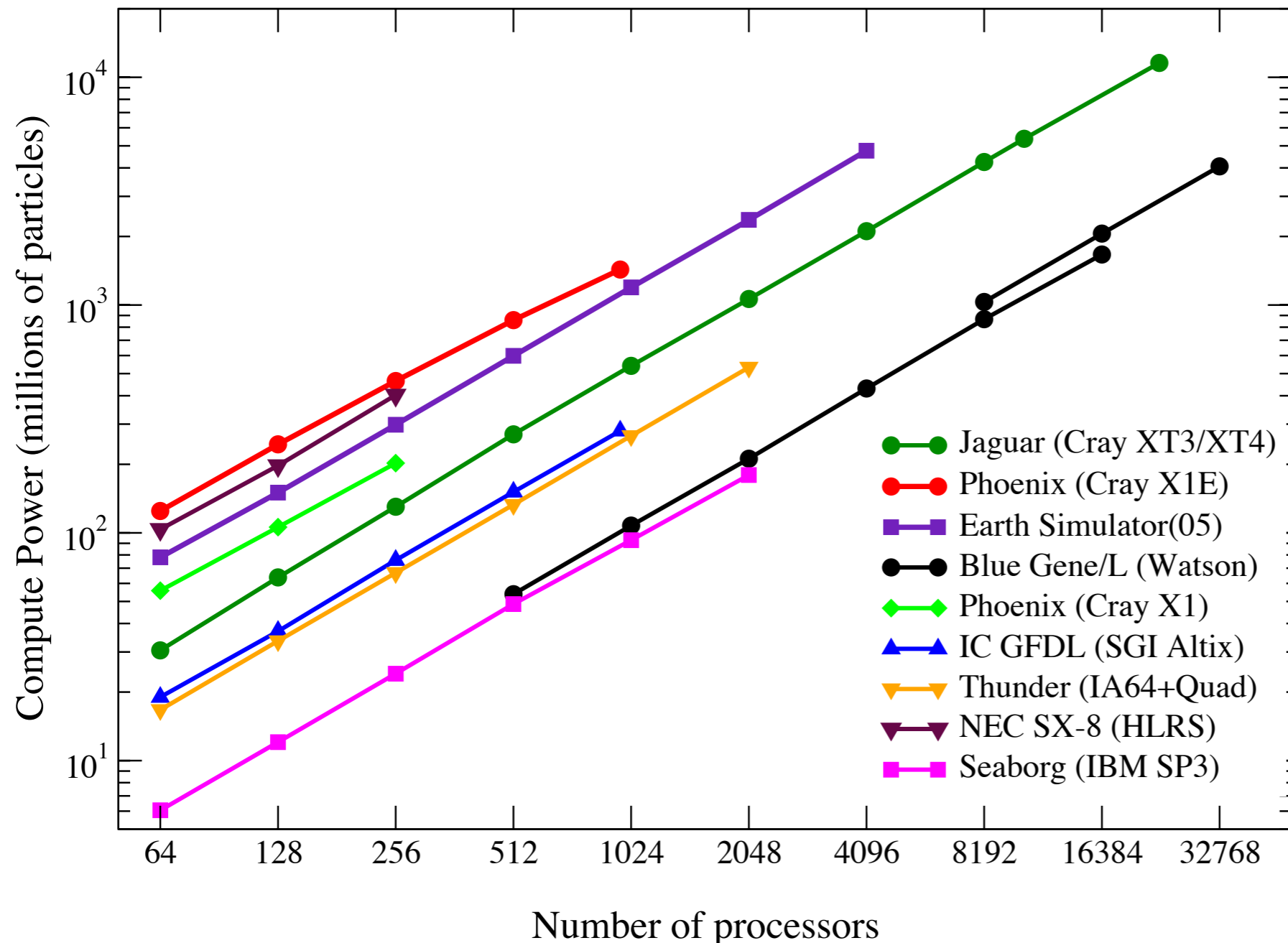
• Object Oriented GTC framework -- Based on Fortran-90 to facilitate team coding

GTC performance on MPP platforms aiming for ITER-size Plasmas

- GTC is very portable, scalable and efficient on both cache-based and vector-parallel MPPs.
- 20 TeraFlops/sec performance has been achieved with 74 billion particles on Jaguar (ORNL) with 22,976 cores and 2.8 times faster than with 32,786 BG/L cores

Compute Power of the Gyrokinetic Toroidal Code

Number of particles (in million) moved 1 step in 1 second



Joule Applications
& 250 TF campaign at
ORNL

Numerical Considerations for Gyrokinetic Simulation Codes

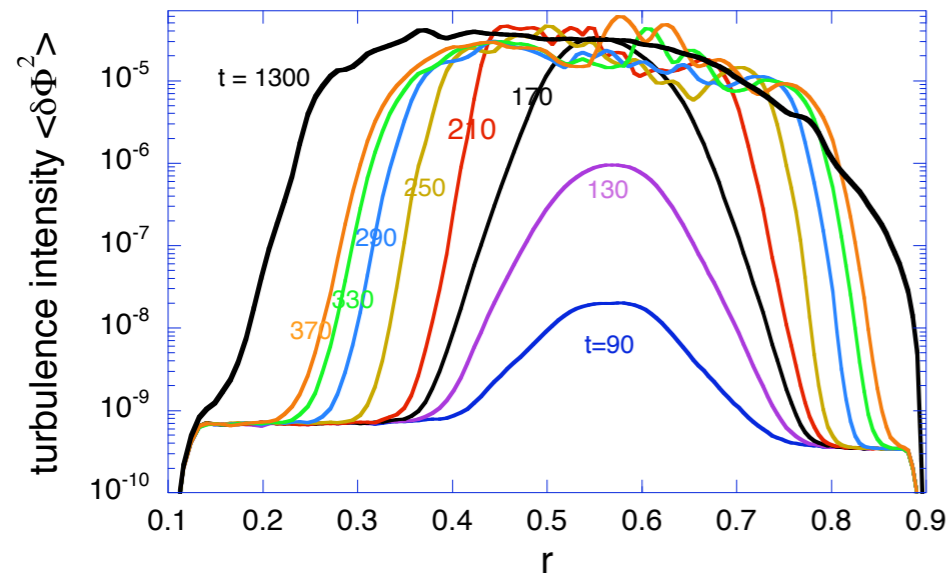
- Flux Tube codes are valid for large (m, n) modes
- Wedge codes include radial variations and some are valid only for large (m, n) modes
- Global codes are valid for any (m, n) modes and are truly five dimensional
- Physics of turbulence transport alone dictates the usefulness of these codes, i.e.,
 - are radial modes local or global?
 - does energy cascade to lower or higher (m, n) modes ?
 - how about enstrophy, to higher (m, n) modes?
 - perpendicular spatial resolution: ion gyroradius, electron skin depth or electron gyroradius?
 - parallel spatial resolution: field line following coordinates?
 - velocity space resolution?
 - ✓ trapped particle dynamics
 - ✓ wave-particle interactions
 - ✓ artificial dissipation
 - ✓ discrete particle noise
 - collisions: can neoclassical transport be simulated?

General Geometry GTS

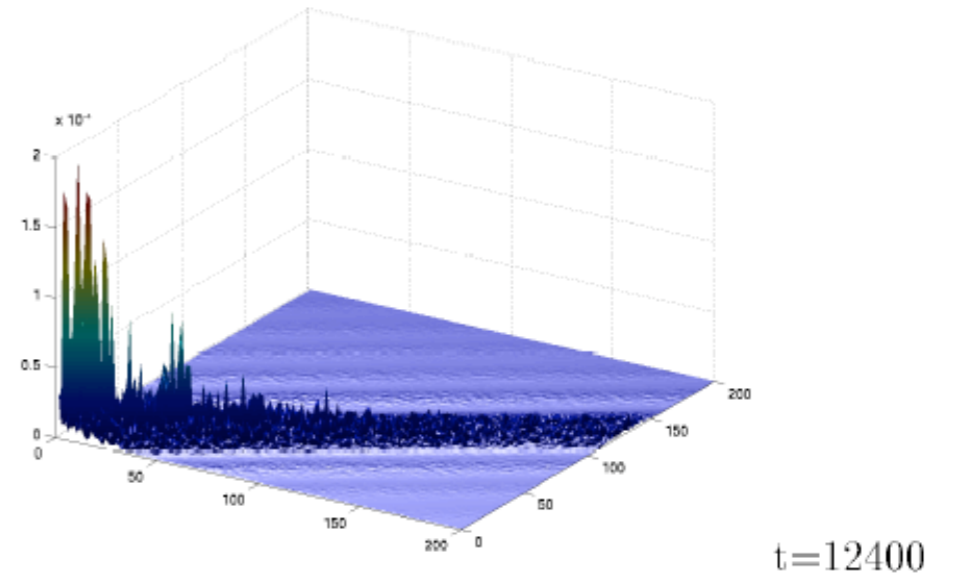
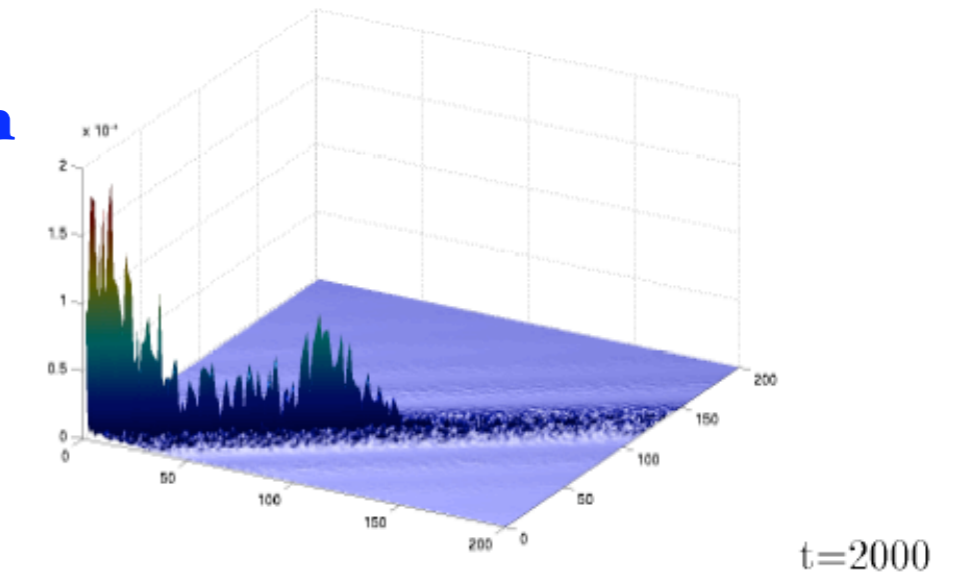
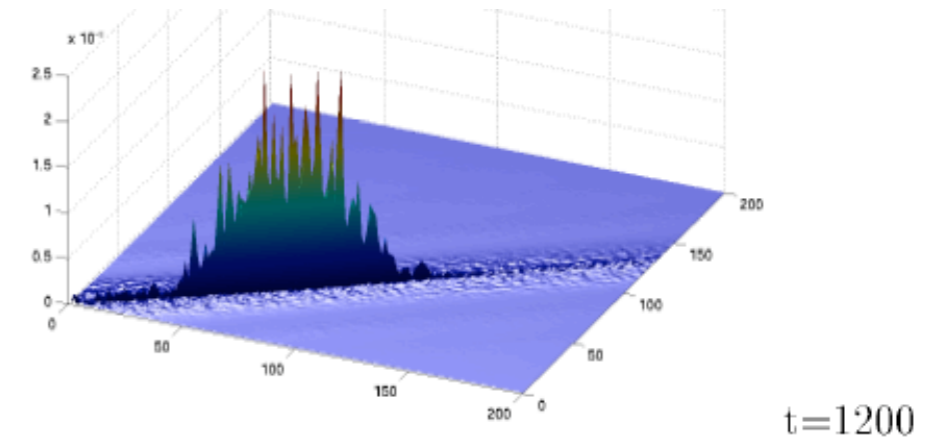
W. X. Wang [PoP '06]

- Global Turbulence Dynamics in Shaped Plasmas
- Interfaced with TRANSP and JSOLVER and ESC
- Re-Write of GTC

ITG Turbulence Spreading

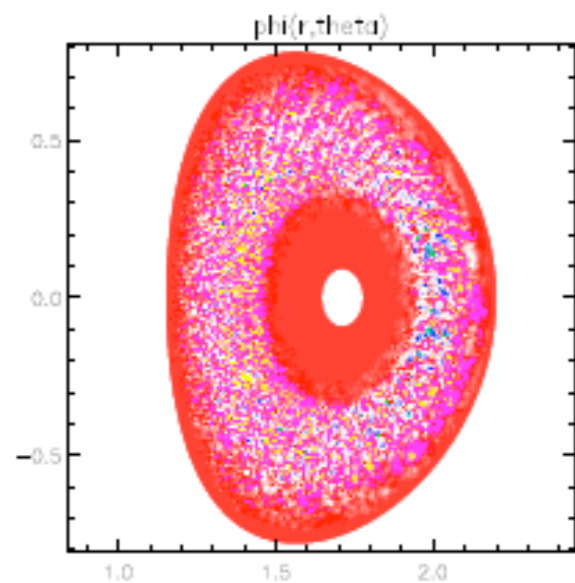
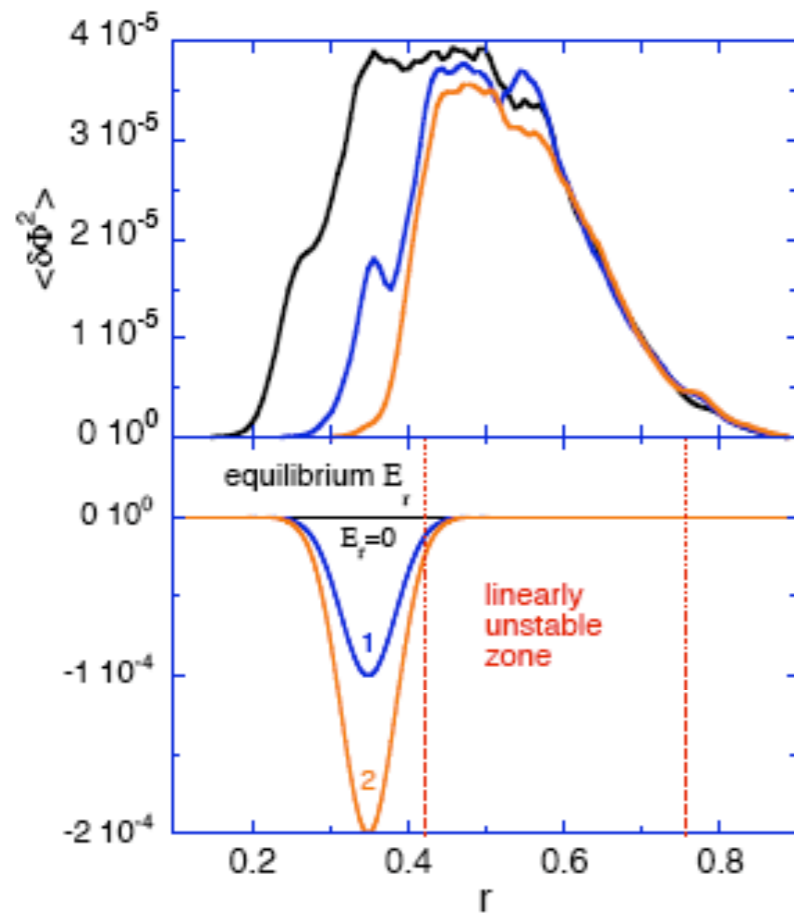


ITG
Energy
Cascade to
longer wavelength
modes



Turbulent transport is a global phenomena

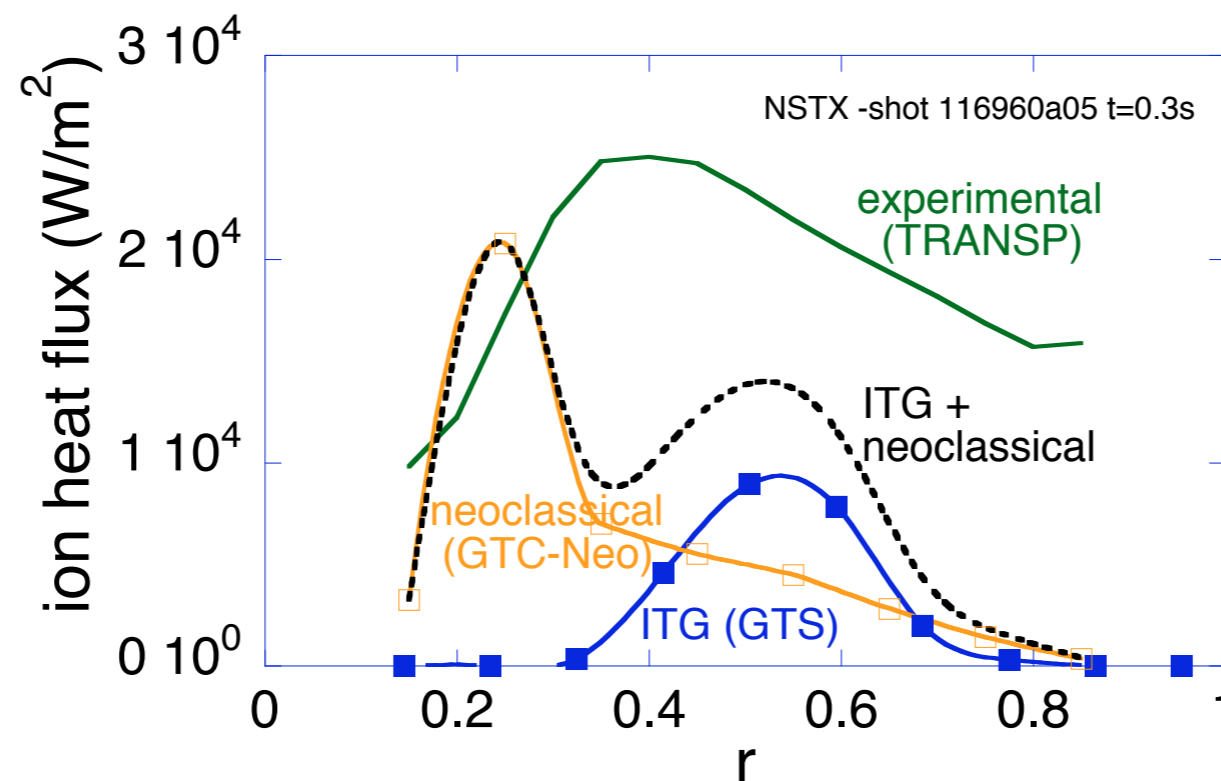
$E \times B$ Shear Layer Blocks Turbulence Spreading



- $\omega_{E \times B}^{max} = 0$: turbulence widely spreads to fill up big area in both directions
- $\omega_{E \times B}^{max} = 0.13 \frac{c_s}{a}$: inward spreading partially blocked
- $\omega_{E \times B}^{max} = 0.26 \frac{c_s}{a}$: almost completely blocked
- Shear layer not only reduces turbulence spreading extension but also slows down the spreading
- Turbulence level not increased in source region as spreading blocked
- Outward spreading is not affected

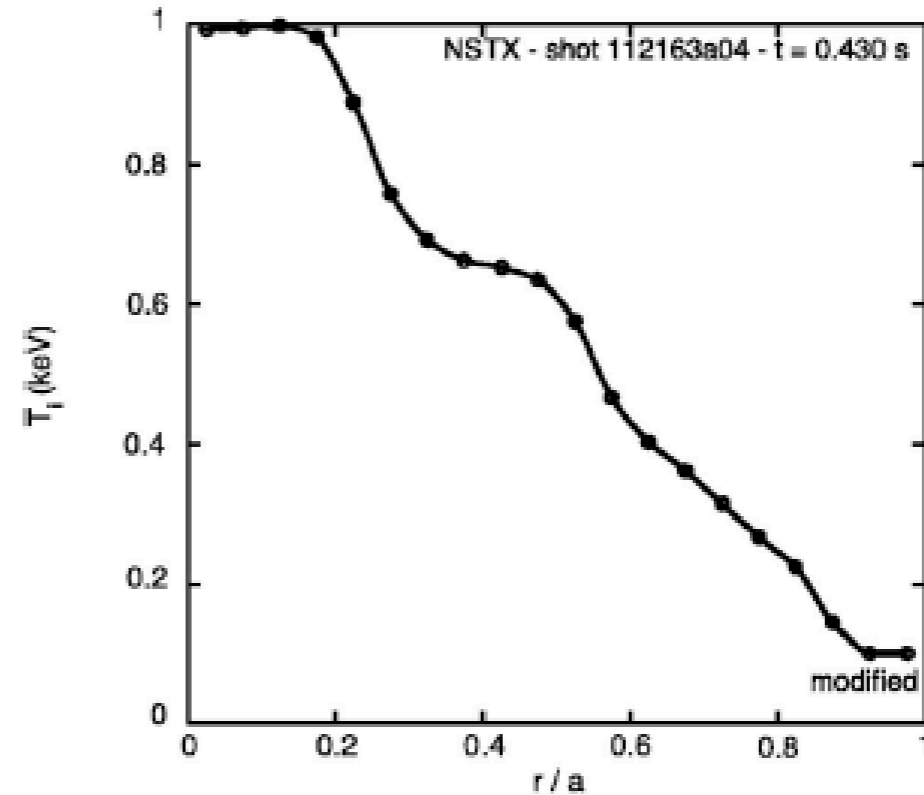
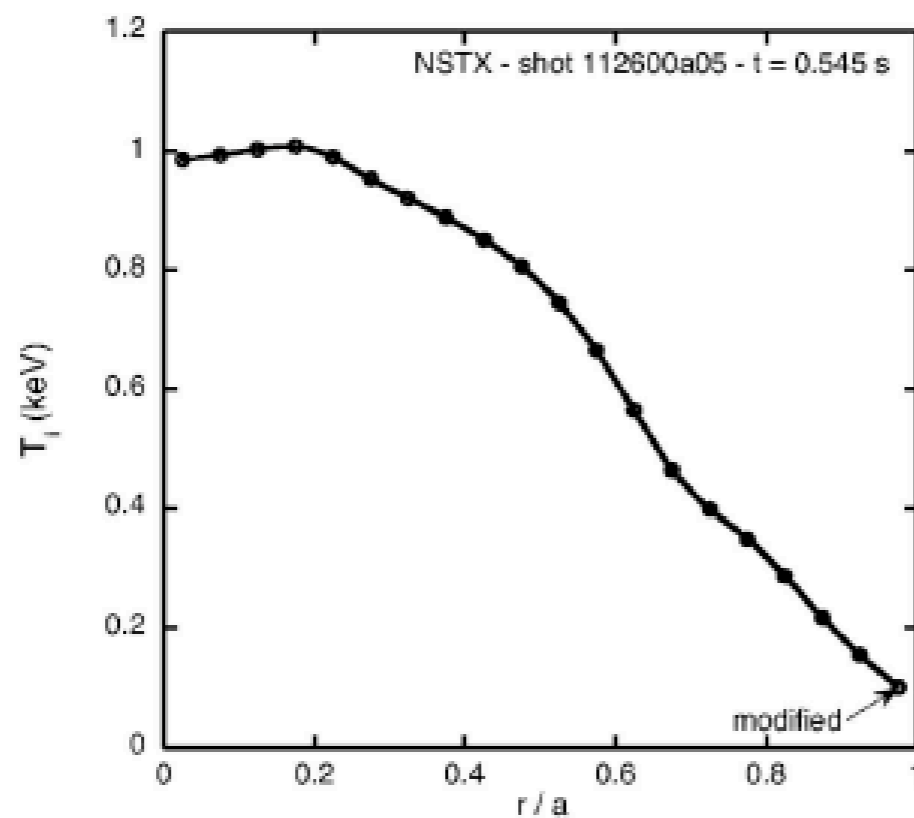
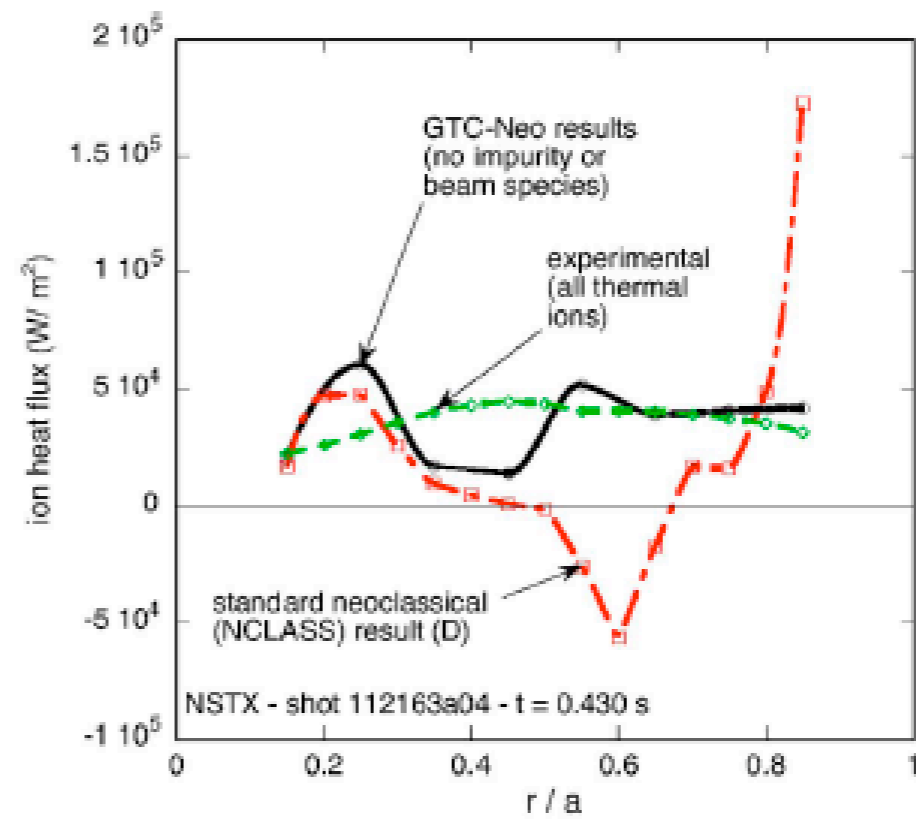
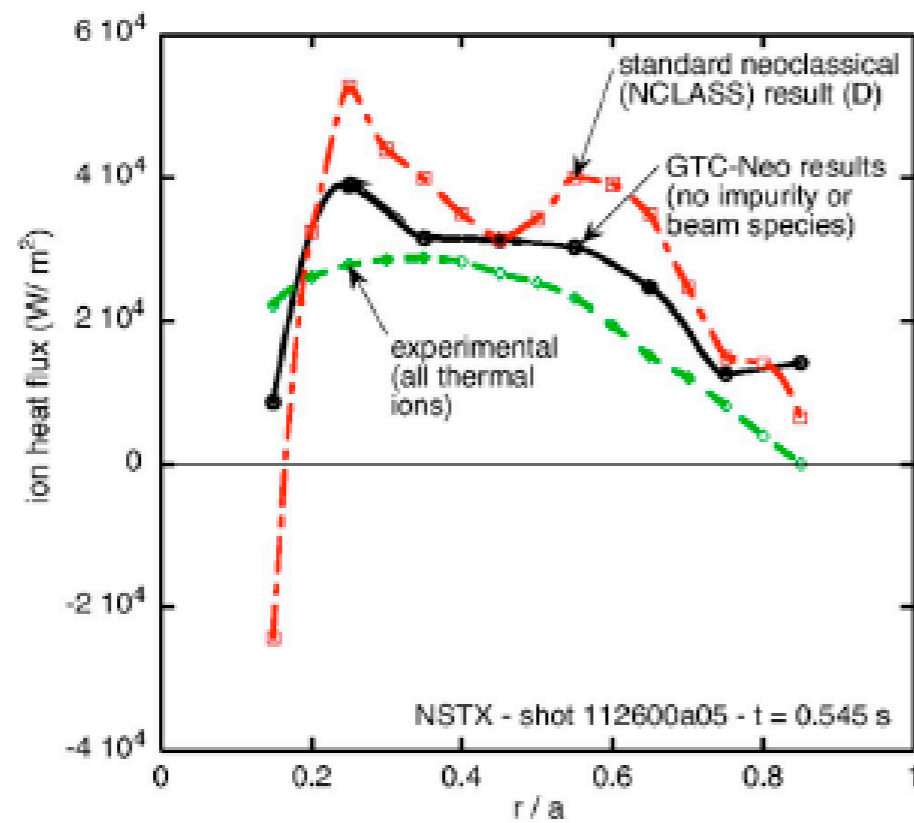
Recent Development of GTC-S and NSTX-physics-oriented Algorithm

- Generalized Poisson Solver to solve integral equation for total potential $\Phi = \delta\Phi + \langle\Phi\rangle$ using superLU/PETSc
- previous solver solves $\delta\Phi$ and $\langle\Phi\rangle$ separately using approximations:
- i) Pade approximation $\Gamma_0(b) \equiv I_0(b)e^{-b} \approx 1/(1+b)$ and
 - ii) $\langle\tilde{\Phi}\rangle \approx \langle\tilde{\Phi}\rangle$ – **not justified for NSTX geometry!**



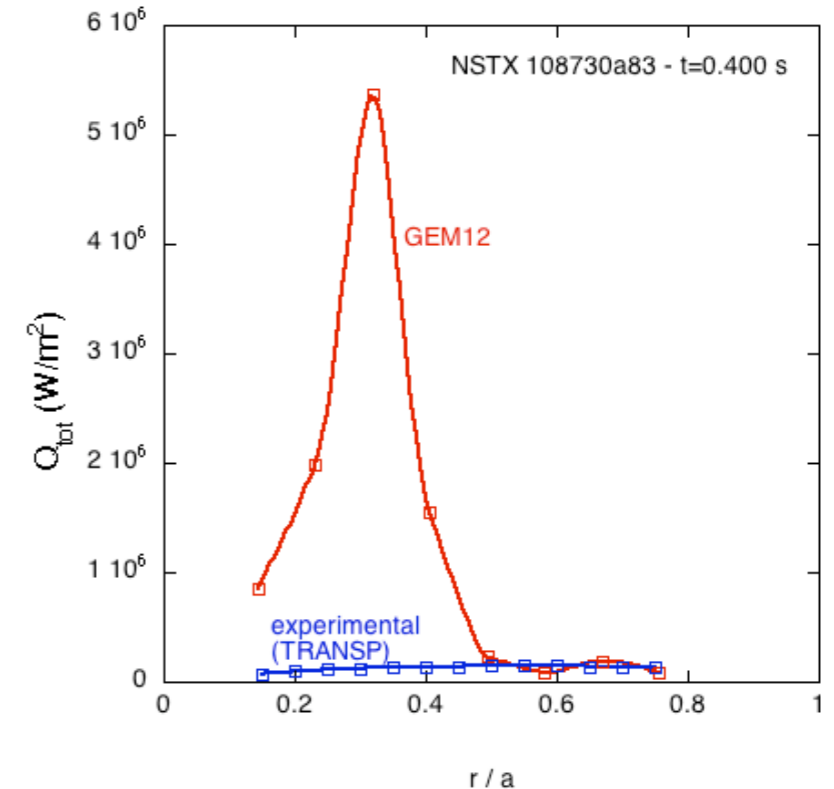
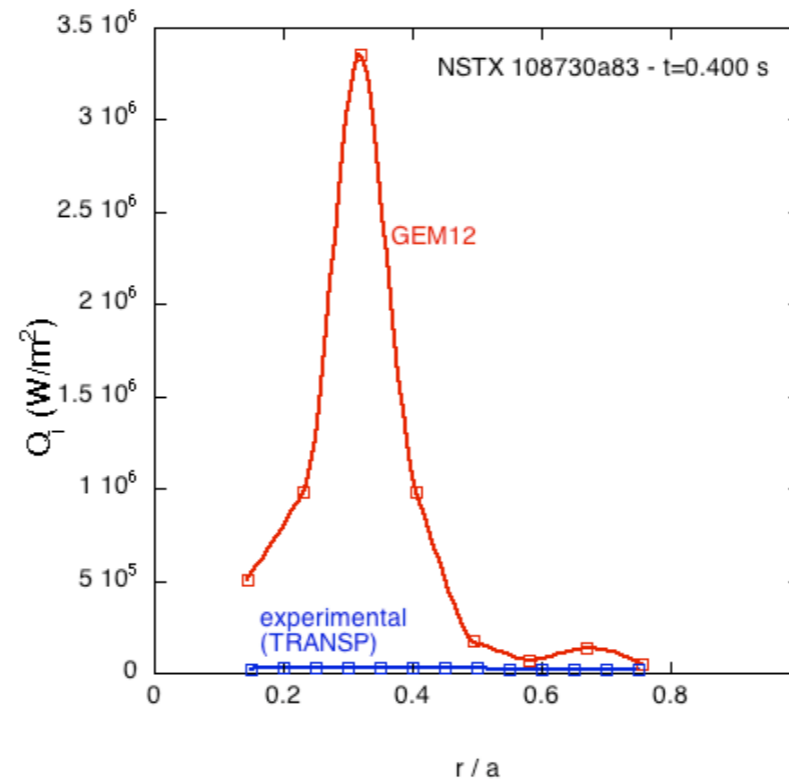
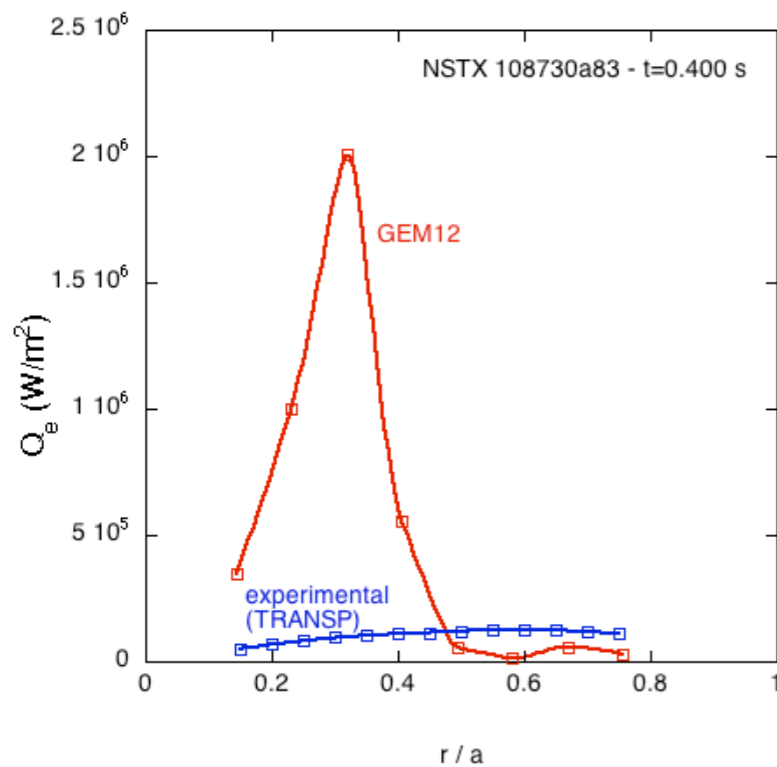
General Geometry GTC-neo

W. X. Wang [PoP, '06]



Comparisons of GEM with NSTX: Energy Flux Measurements

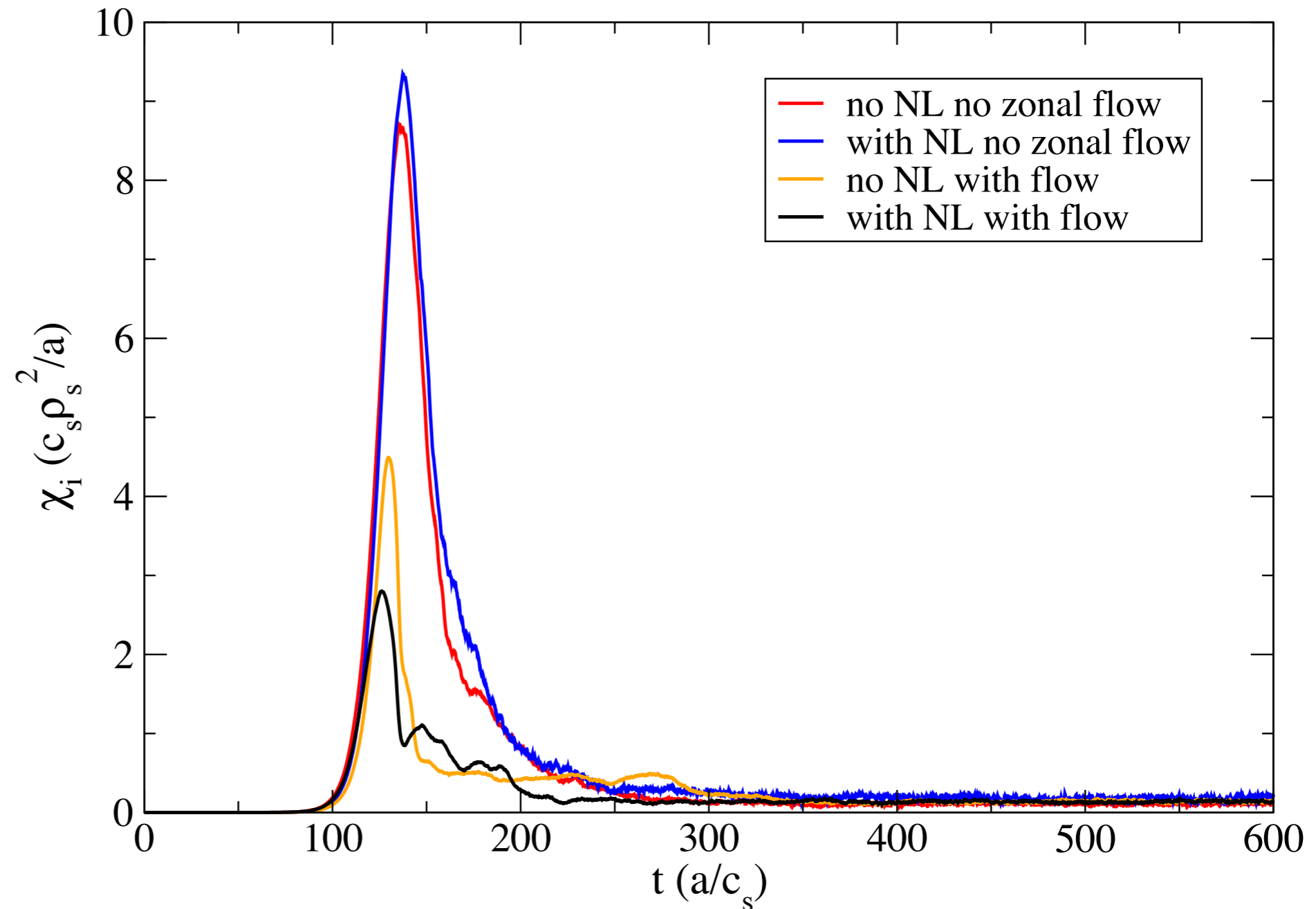
- First results - does not include parallel ion equilibrium flows! (which are transonic or supersonic)
- 128 particles / species / grid cell, 3 ion species, experimental β



ITG simulations with for adiabatic electrons based on Cyclone-based parameters using GTC

Peak χ_i (bin 3) - $a/\rho_i=125$

(part/cell=10, mzetamax=32, cyclone case profile, parallel non-linearity study)

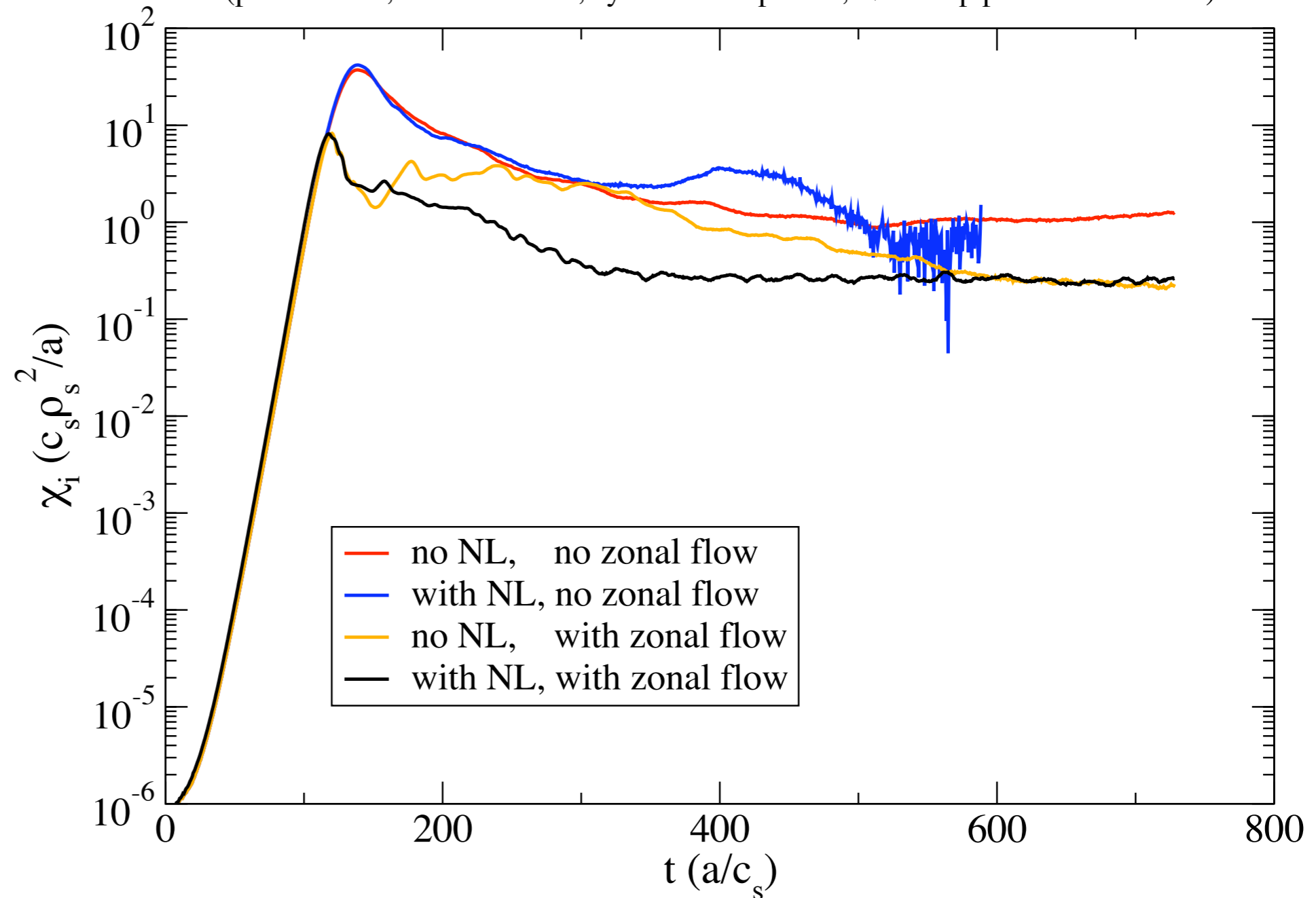


- Steady state fluxes remain essentially the same with or without zonal flows and with or without parallel velocity space nonlinearity for a small simulation volume.
- But, these nonlinearities become progressively important for larger systems

ITG simulations using GTC (cont.)

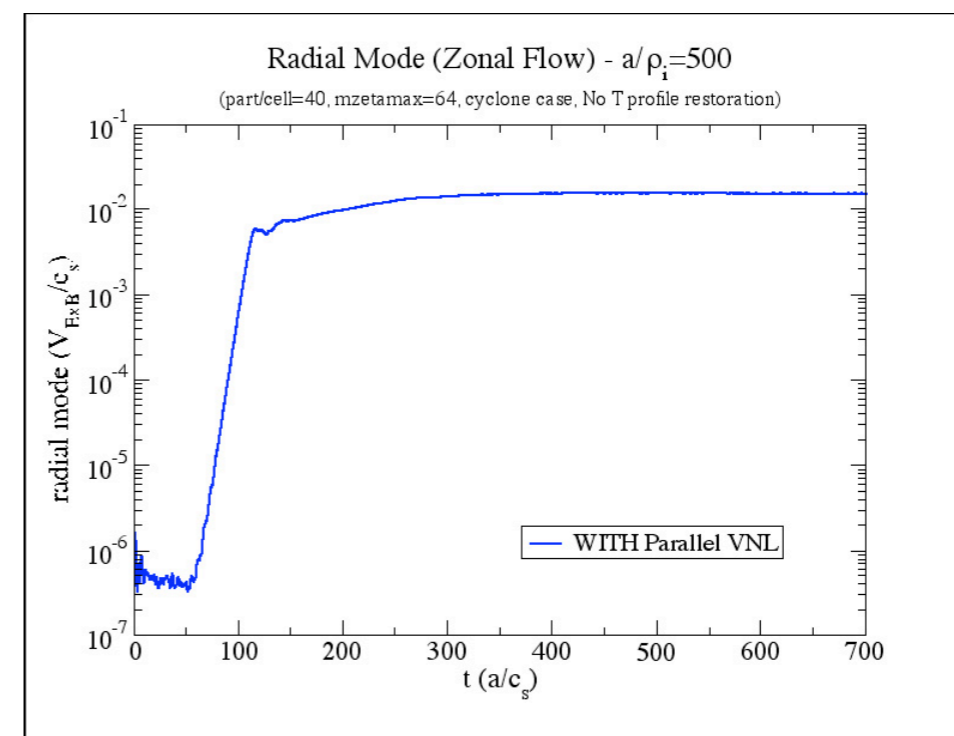
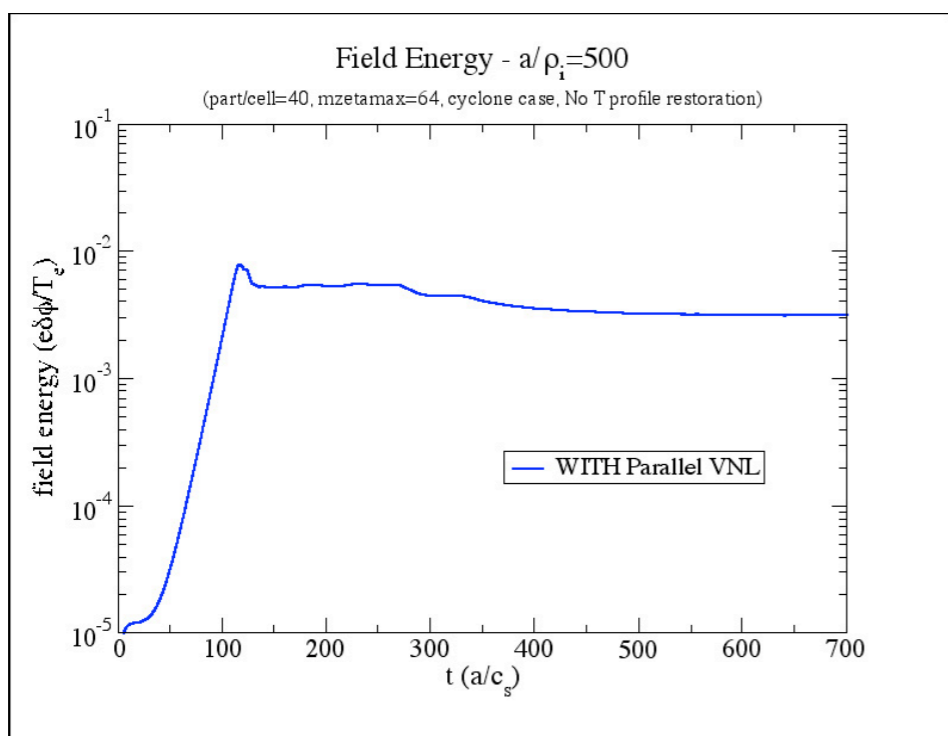
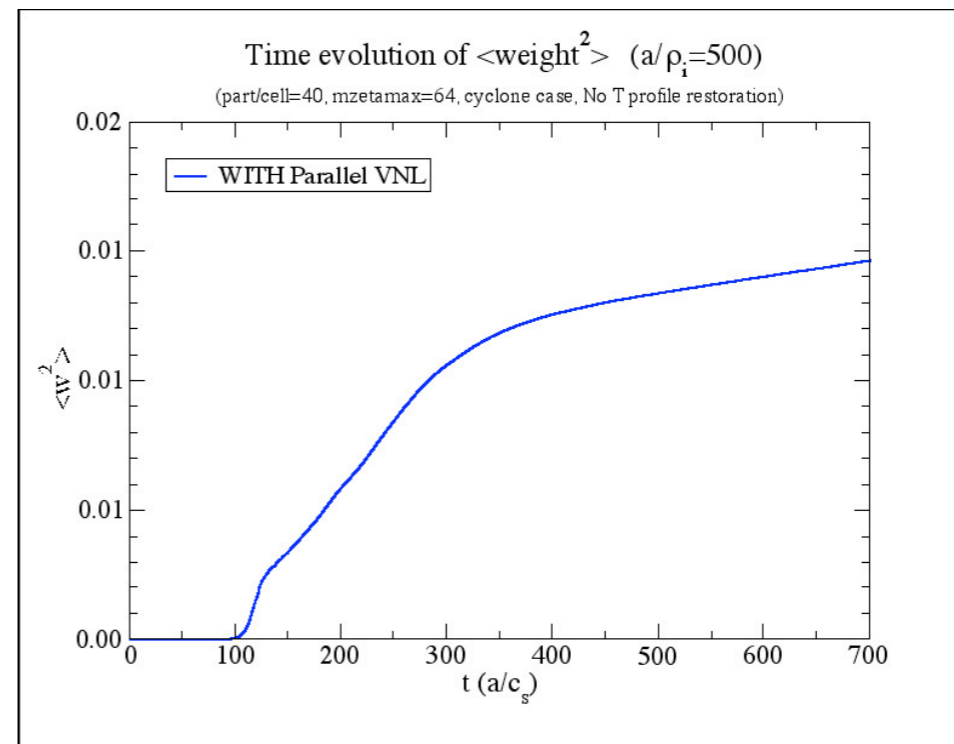
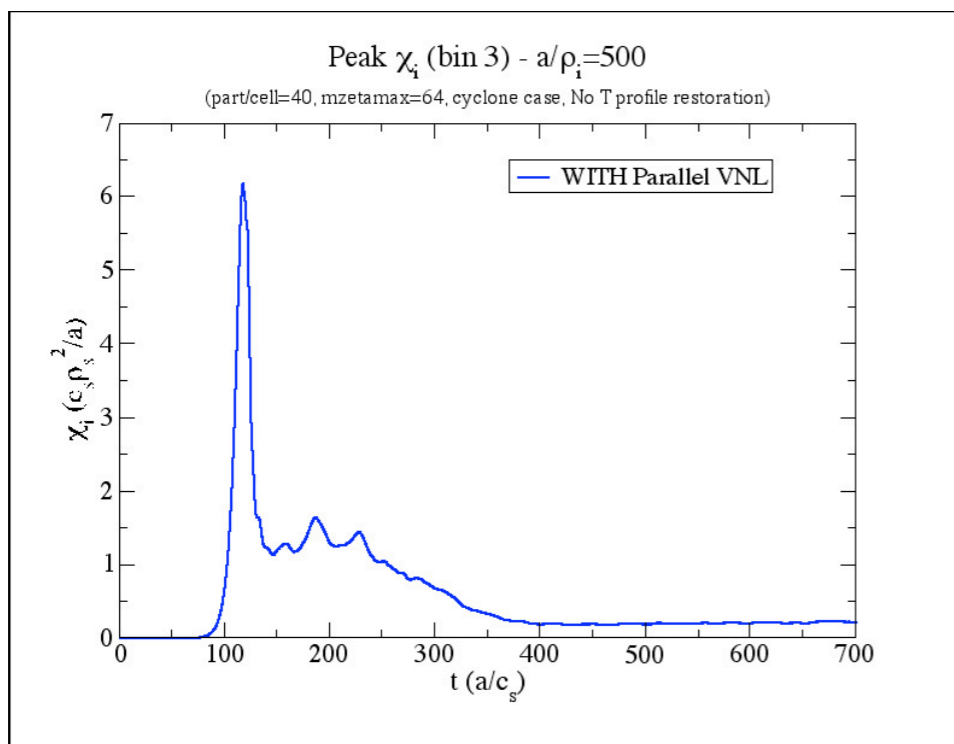
Peak χ_i (bin 3) - $a/\rho_i=250$

(part/cell=40, mzetamax=64, cyclone case profile, No Temp profile restoration)



- For example, zonal flows are essential for maintaining steady state flux for a larger simulation volume.

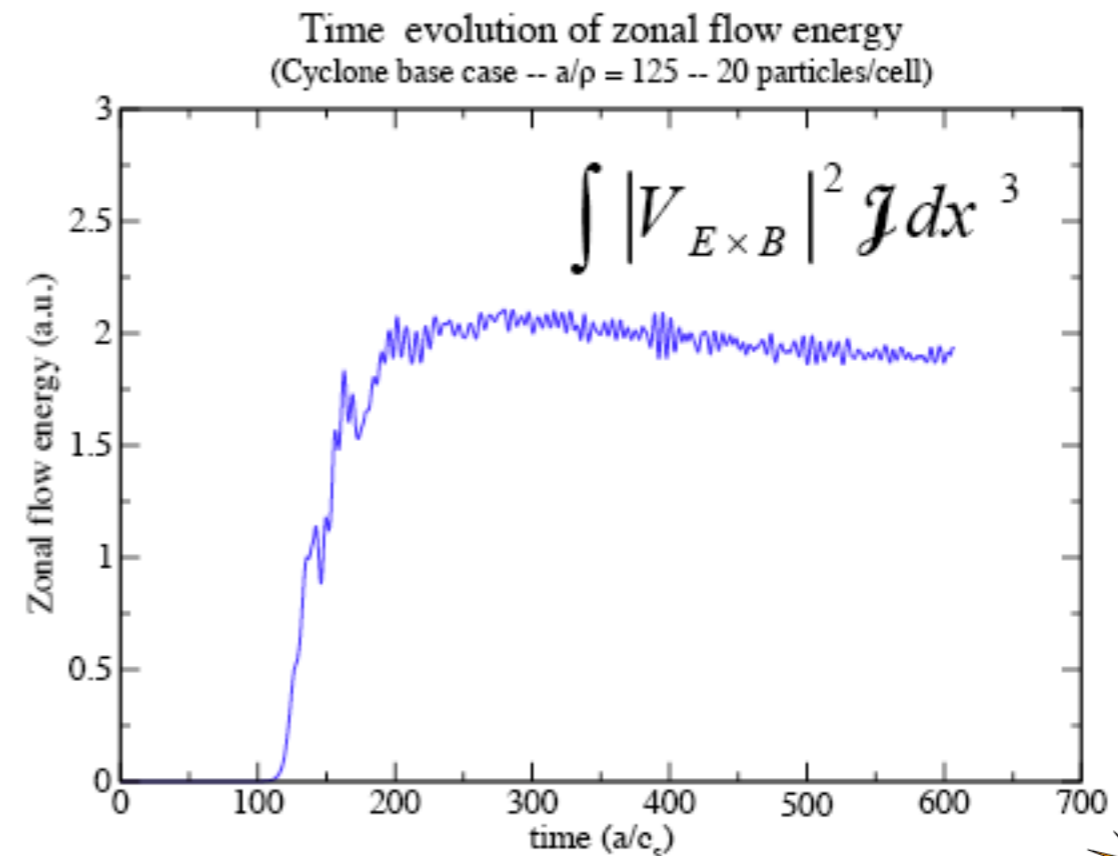
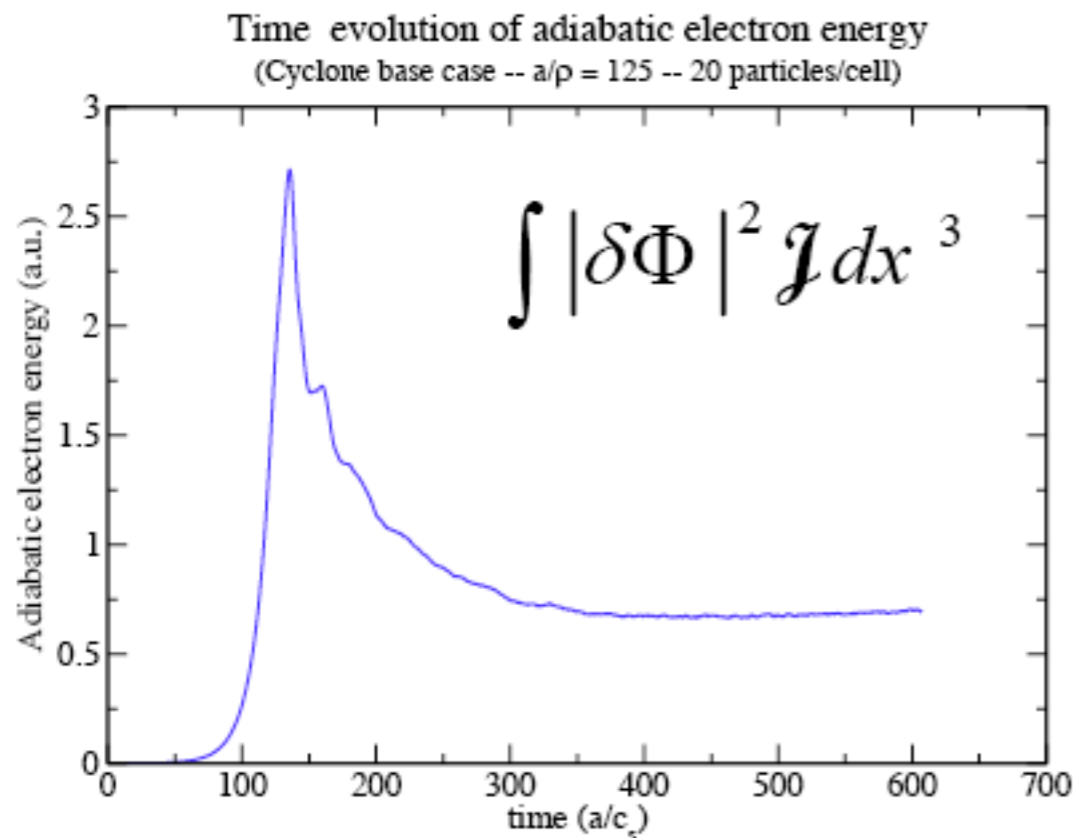
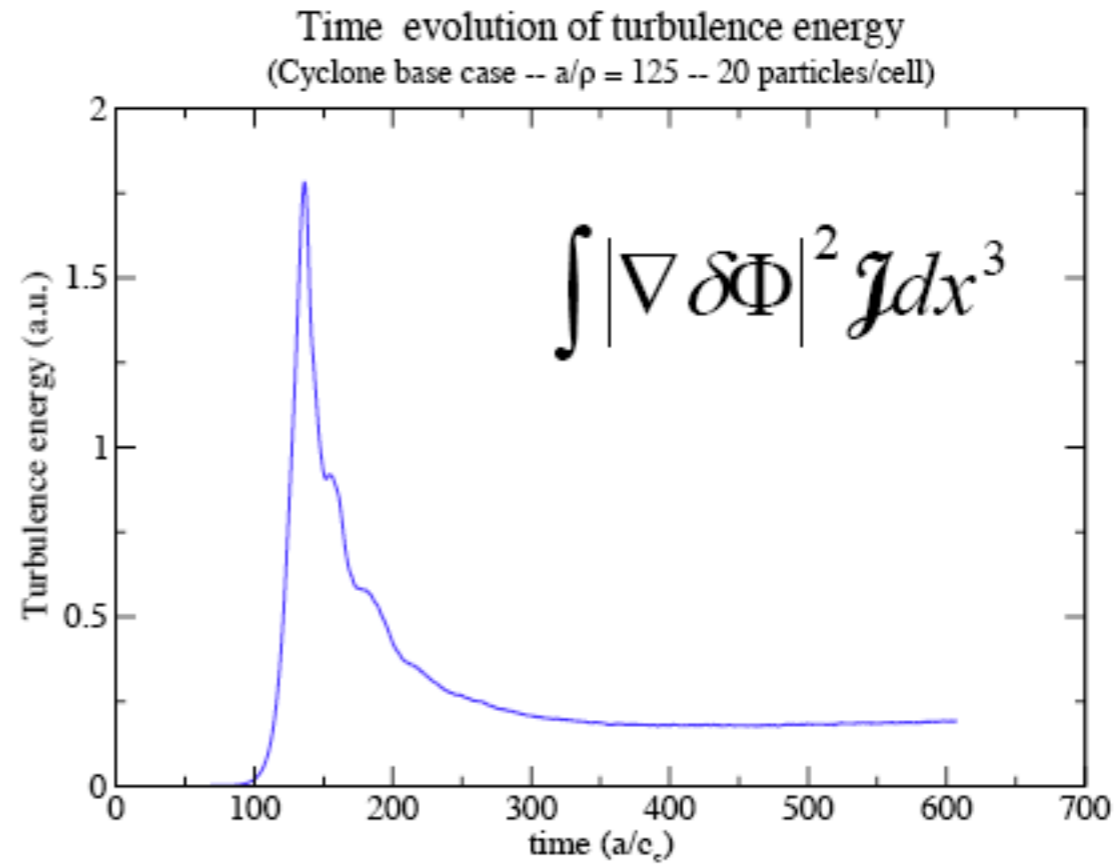
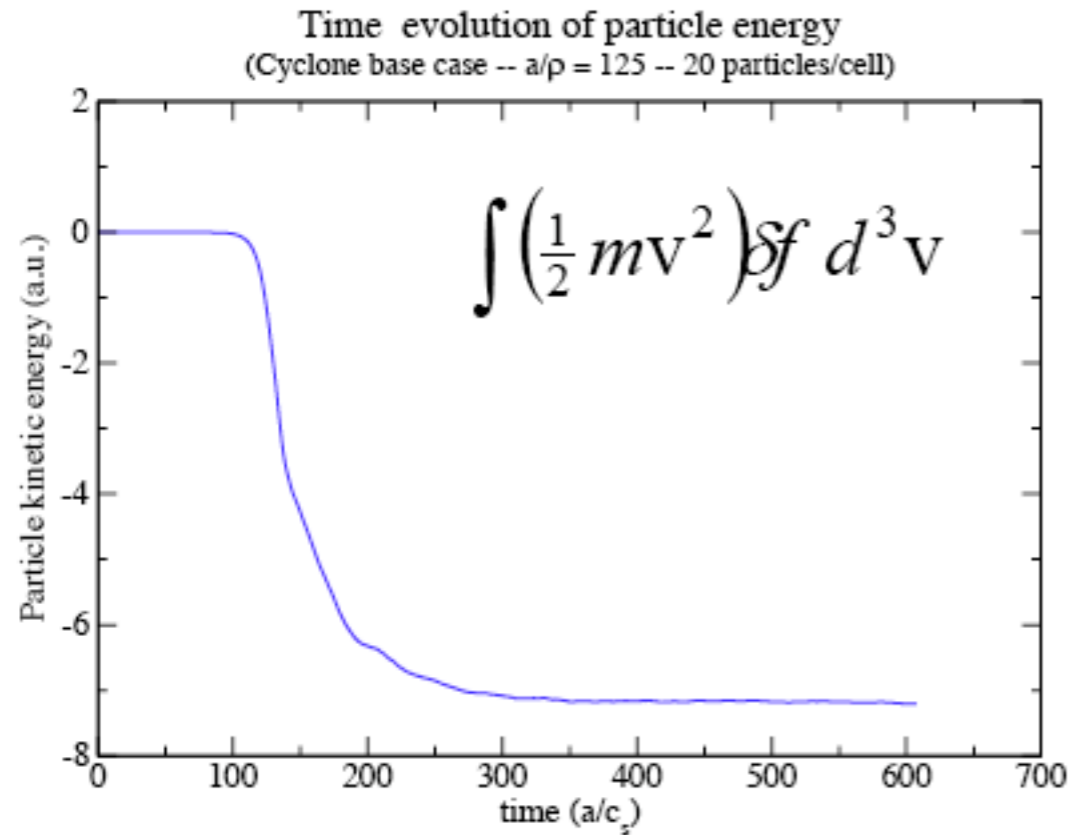
Both zonal flows and velocity-space nonlinearity are essential for maintaining steady state flux for an even larger system ($a/\rho_i = 500$).



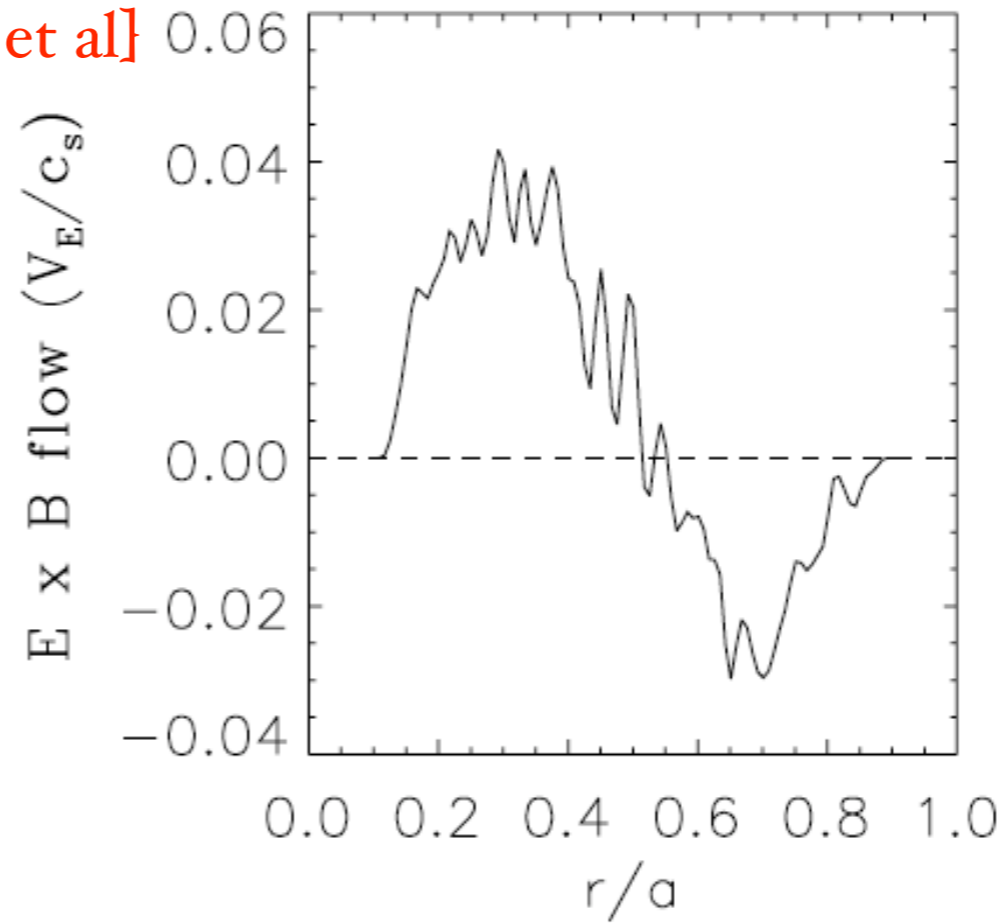
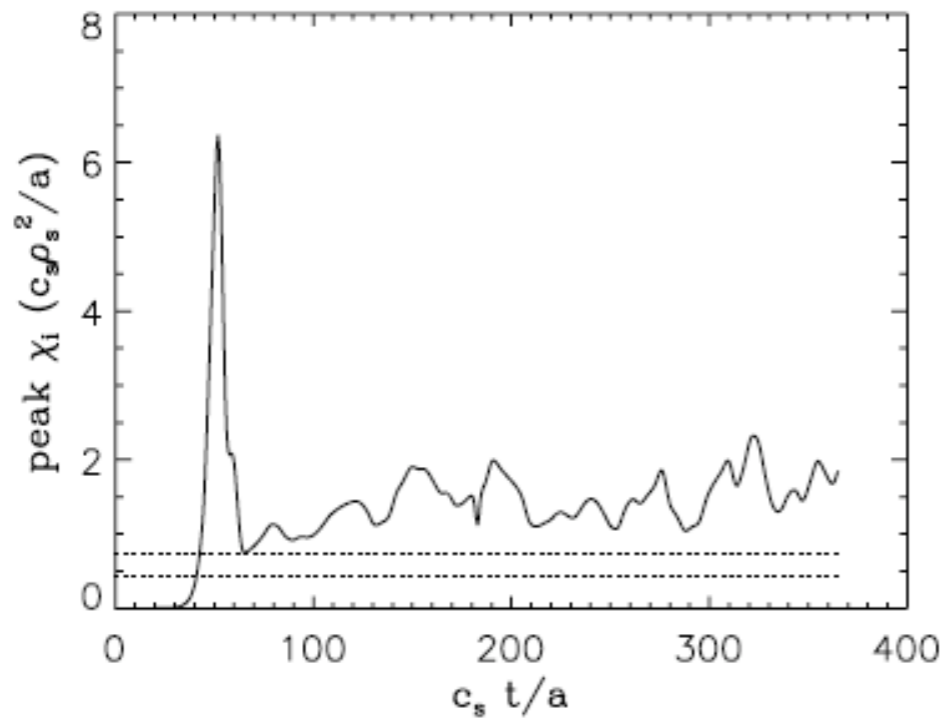
- There is no steady state without these nonlinearities in large scale global simulations.

Lee, Ethier and Kolesnikov

Conservation properties of ITG simulation (20 particles/cell)



- ITG simulation with kinetic electrons using GTC with the split-weight scheme [Lewandowski et al]

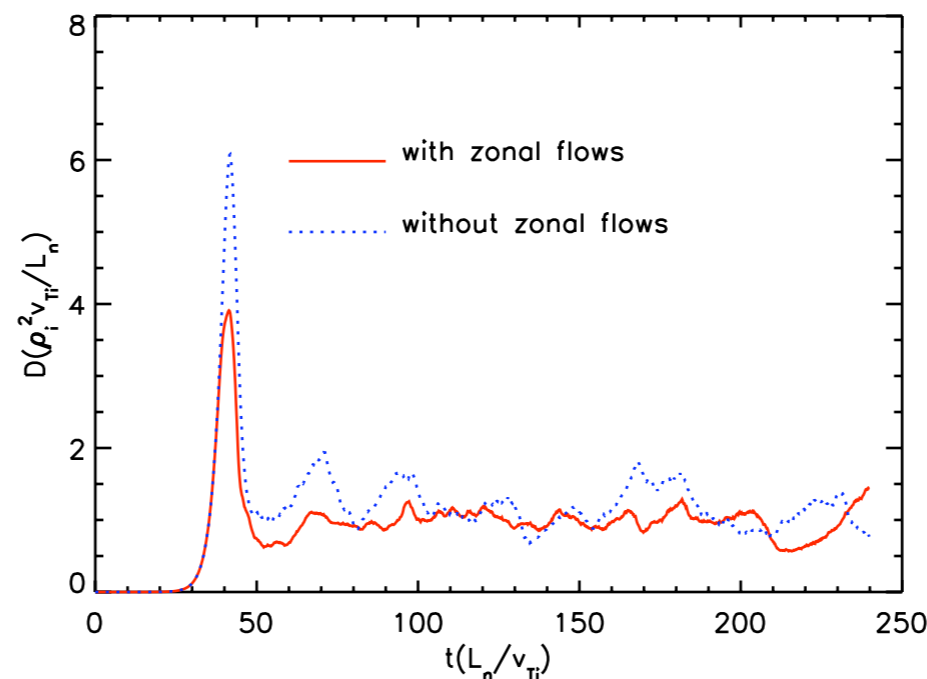


Zonal flows develop finer structures

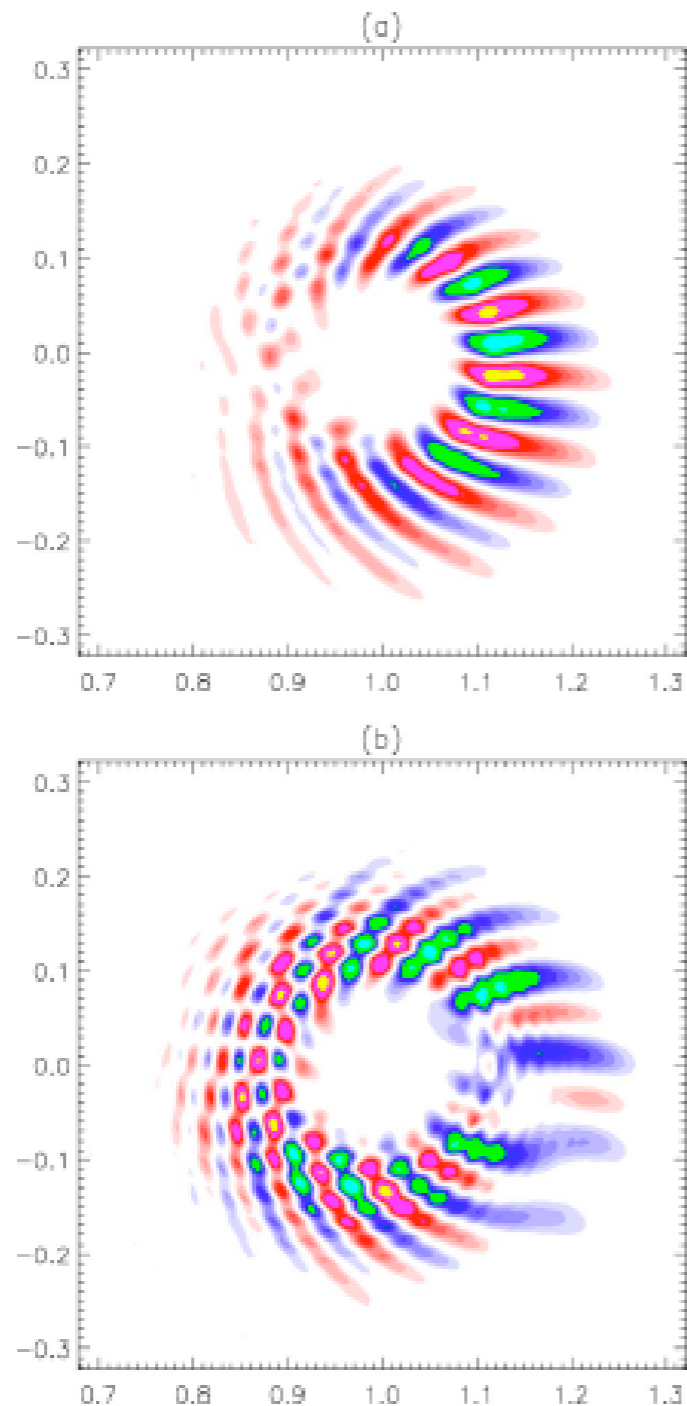
χ_i is enhanced above the adiabatic electron level (with NLV)

- TEM simulation with collisionless electrons using GEM with the split-weight scheme [Lang, Parker and Chen]

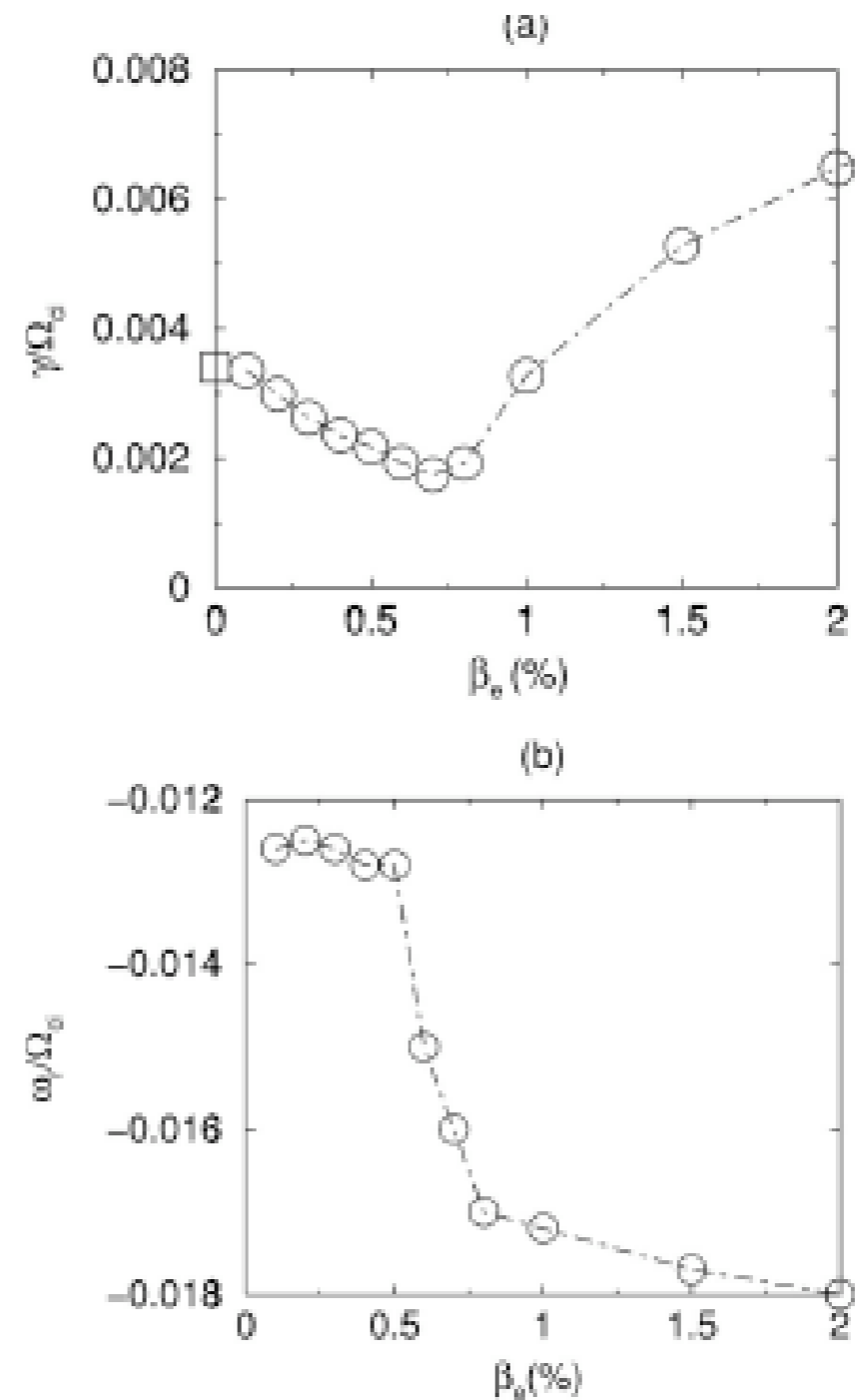
Enhanced electron particle flux



GTC-EM : fluid electron model



Mode structures for (a) the electrostatic potential and (b) the vector potential



Growth rates and real frequencies for finite-beta modified ITG modes

Momentum Pinch from $E \times B$ shear and Magnetic Curvature

Two different mechanisms for non-diffusive momentum flux

$$m_i B^* dv_{\parallel} / dt = -(eB + m_i c v_{\parallel} \nabla \times \hat{b}) \cdot \nabla \delta\phi$$

	Gurcan, Diamond, Hahm, Singh [Phys. Plasmas 14, 042306 '07]	Hahm, Diamond, Gurcan, Rewoldt [Phys. Plasmas, June '07]
Net acceleration of parallel flow:	$-e_i B \nabla_{\parallel} \delta\phi$	$-m_i c v_{\parallel} \nabla \times \hat{b} \cdot \nabla \delta\phi$
Symmetry-breaking:	k_{\parallel} over the spectral width	curvature drift $\sim \hat{b} \times (\hat{b} \cdot \nabla) \hat{b}$ over the flux surface
Provided by:	mean $E \times B$ shear shifting fluctuations radially	ballooning mode structure causing finite net parallel acceleration over the flux surface
Main consequence:	residual stress driven by $E \times B$ shear (or $\nabla P_i / n_i$ and velocity shear via radial force balance)	convective pinch-like term (the TEP-like piece is insensitive to mode details)
Most likely to be relevant for:	plasmas with strong $E \times B$ shear, incl. H-mode, ITB's	pinch is likely to be inward for OH and electron-heated plasmas

Discrete particle noise in particle-in-cell simulations of plasma microturbulence

[Nevins, Hammett, Dimits, Dorland, and Shumaker, PoP 12, 122305 (2005)]

$$\{1 + [1 - \Gamma_0(k_{\perp}^2 \rho_e^2)]\} \frac{e\phi_{\mathbf{k}}}{T} = \frac{S_G(\mathbf{k})}{N_p} \sum_{\mathbf{p}} S(\mathbf{k}_p) \sum_i w_i J_0(k_{\perp} \rho_i) \times \exp(-i\mathbf{k}_p \cdot \mathbf{x}_i), \quad (1)$$

$$\left\langle \left| \frac{e\phi}{T} \right|^2 \right\rangle = \sum_{\mathbf{k}} \left\langle \left| \frac{e\phi_{\mathbf{k}}}{T} \right|^2 \right\rangle = \frac{\langle w^2 \rangle}{n_p V_{\text{shield}}}, \quad (8)$$

$$V_{\text{shield}}^{(N)} \equiv \left\{ \frac{1}{(2\pi)^3} \int d^3\mathbf{k} \frac{S_G^2 \Gamma_0(k_{\perp}^2 \rho_e^2)}{[2 - \Gamma_0(k_{\perp}^2 \rho_e^2)]^2} \right\}^{-1}, \quad (9)$$

$$V_{\text{shield}}^{(H)} \equiv \left\{ \frac{1}{(2\pi)^3} \int d^3\mathbf{k} \frac{S_G^2 \Gamma_0(k_{\perp}^2 \rho_e^2)}{(2 - \Gamma_0)[2 - (1 - S_G d_{\parallel}) \Gamma_0]} \right\}^{-1}. \quad (10)$$

- **Since dynamic plasma response is not included in the calculation, it is difficult to assess the effect of the shielding volume noise on long wavelength modes.**

Discrete Particle Noise for Equilibrium Plasmas

- Fluctuation-Dissipation Theorem

$$L|E(k, \omega)|^2/8\pi = -(T/\omega)Im(1/\epsilon)$$

- Fluctuations per k -mode

$$L|E(k)|^2/8\pi = \int (d\omega/2\pi) L|E(k, \omega)|^2/8\pi = (T/2)[1/\epsilon(k, \omega = \infty) - 1/\epsilon(k, \omega = 0)].$$

$$L|E(k)|^2/8\pi = (T/2)(\lambda_D^2/\rho_s^2)$$

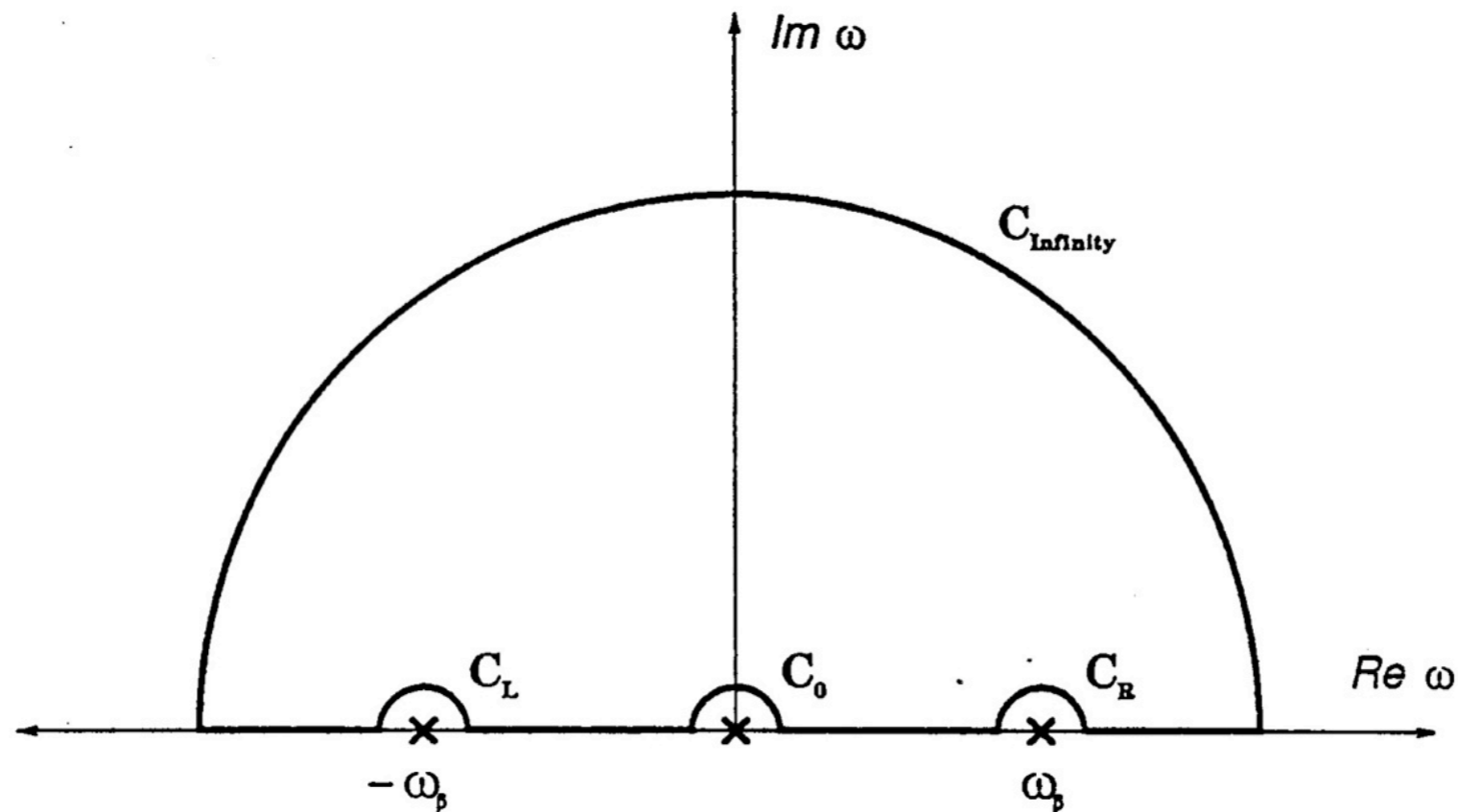


Figure 4.1: Contour integration

(John Reynders, PhD thesis, Princeton University, 1992)

Fluctuation Dissipation Theorem applied to a nonlinearly saturated driven system

$$1 + k_{\perp}^2 - i\sqrt{\frac{\pi}{2}} \frac{\omega_* - \omega}{k_{\parallel} v_{te}} \left[1 - \frac{k_{\perp}^4}{\gamma_l^2} |\phi|^2 \right] - \frac{\omega_*}{\omega} = 0$$

- Noise level for high frequency modes

$$\left| \frac{e\Phi}{T_e} \right|_{HF-noise}^2 = \frac{\langle w^2 \rangle}{N k_{\perp}^2 \rho_s^2},$$

- Noise level for low frequency modes

$$\left| \frac{e\Phi}{T_e} \right|_{LF-noise}^2 = \frac{\langle w^2 \rangle}{N(1 + k_{\perp}^2 \rho_s^2)}$$

- Nonlinear saturation amplitude

$$\left| \frac{e\Phi}{T_e} \right|_{NL}^2 = \frac{\gamma_L / \Omega_i}{2k_x k_y \rho_s^2}$$

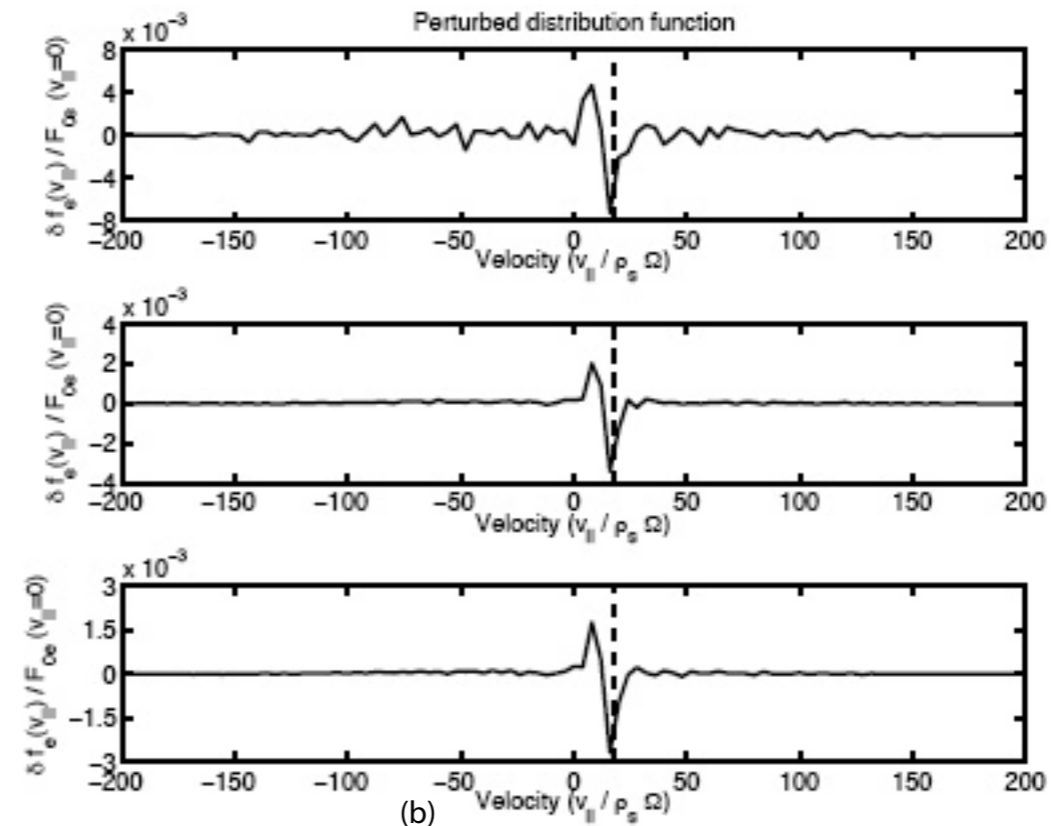
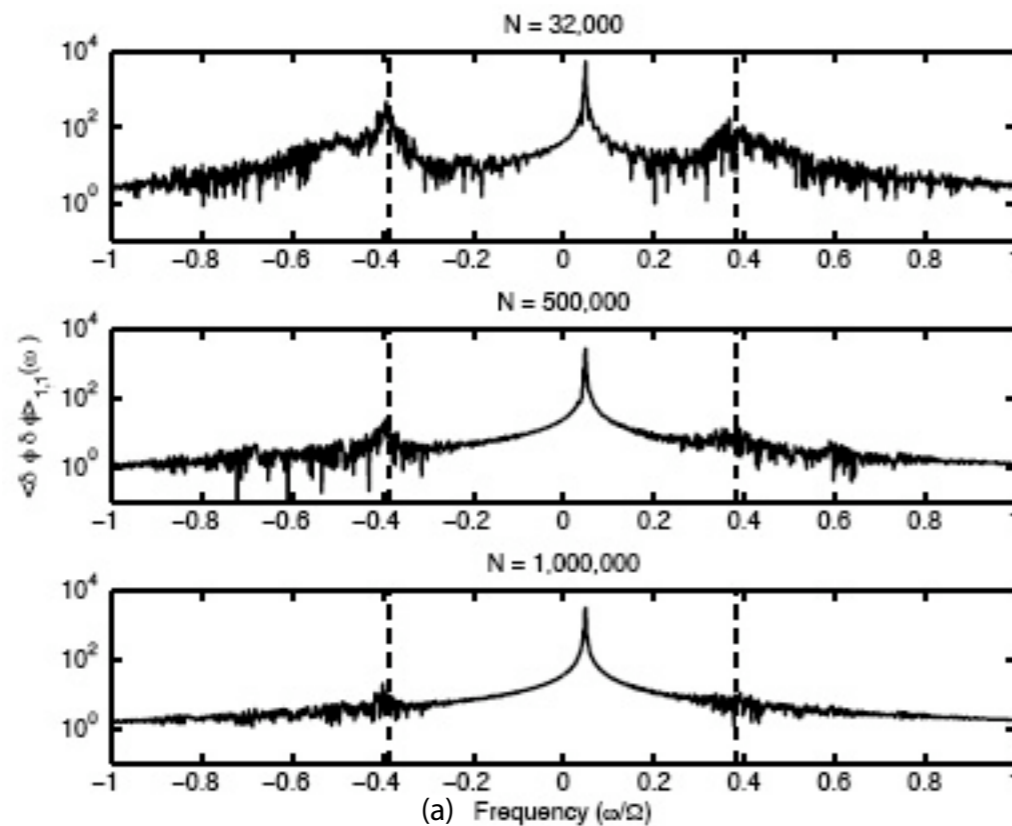
[Jenkins and Lee, PoP '07]

Discrete Particle Noise in Nonlinearly Saturated Plasmas

2D drift wave simulations

with

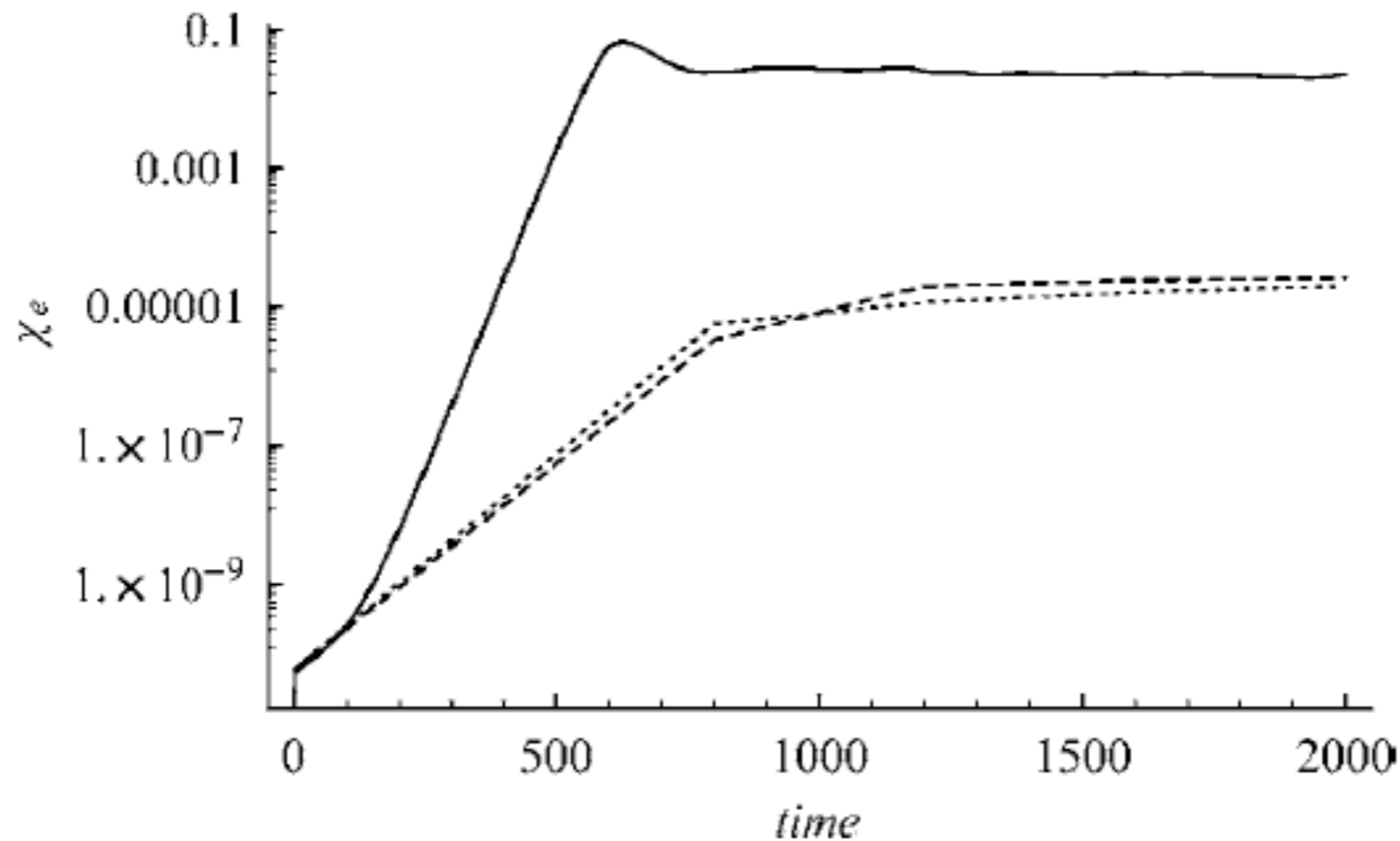
$N = 32K, 500K, 1M$



- high frequency noise decreases with particle number,
- saturation level is independent of particle number,
- background change is small.

[Jenkins and Lee, PoP '07]

ETG simulation using GTC: Noise-driven transport vs. fluctuation-driven transport



Electron transport in ETG simulation: total (solid line), noise driven contribution estimated by scramble test (dashed line) and estimated from δf weight (dotted line).

[Holod and Lin, PoP '07]

Entropy conservation in ITG turbulence:

velocity-space nonlinearity, collisions and numerical noise & dissipation in steady state

$$\frac{\partial}{\partial t} \left\langle \int \frac{\delta f_i^2}{F_{0i}} dv_{\parallel} + \tau \phi^2 + \tau |\nabla_{\perp} \phi|^2 \right\rangle + \left\langle \tau \frac{\partial \phi}{\partial x_{\parallel}} \int v_{\parallel} \frac{\delta f_i^2}{F_{0i}} dv_{\parallel} + 2\tau \nu \int \frac{dv_{\parallel}}{F_{0i}} \left(\frac{\partial \delta f_i}{\partial v_{\parallel} / v_{ti}} + \frac{v_{\parallel}}{v_{ti}} \delta f_i \right)^2 \right\rangle = \kappa_{Ti} \langle Q_{ix} \rangle$$

Monotonic
increase in
time

Vanishes in
steady state

Velocity Space
Nonlinearity

Collisional and/or numerical
dissipation

Entropy
production

[Lee and Tang, '88]

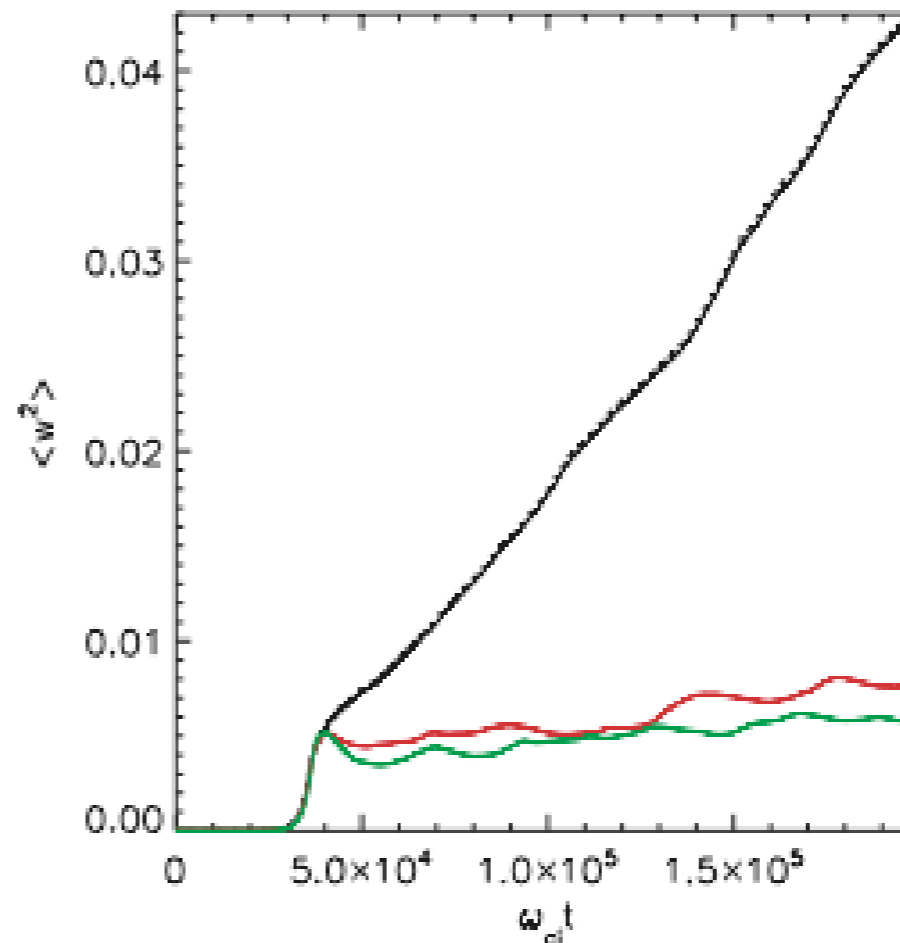
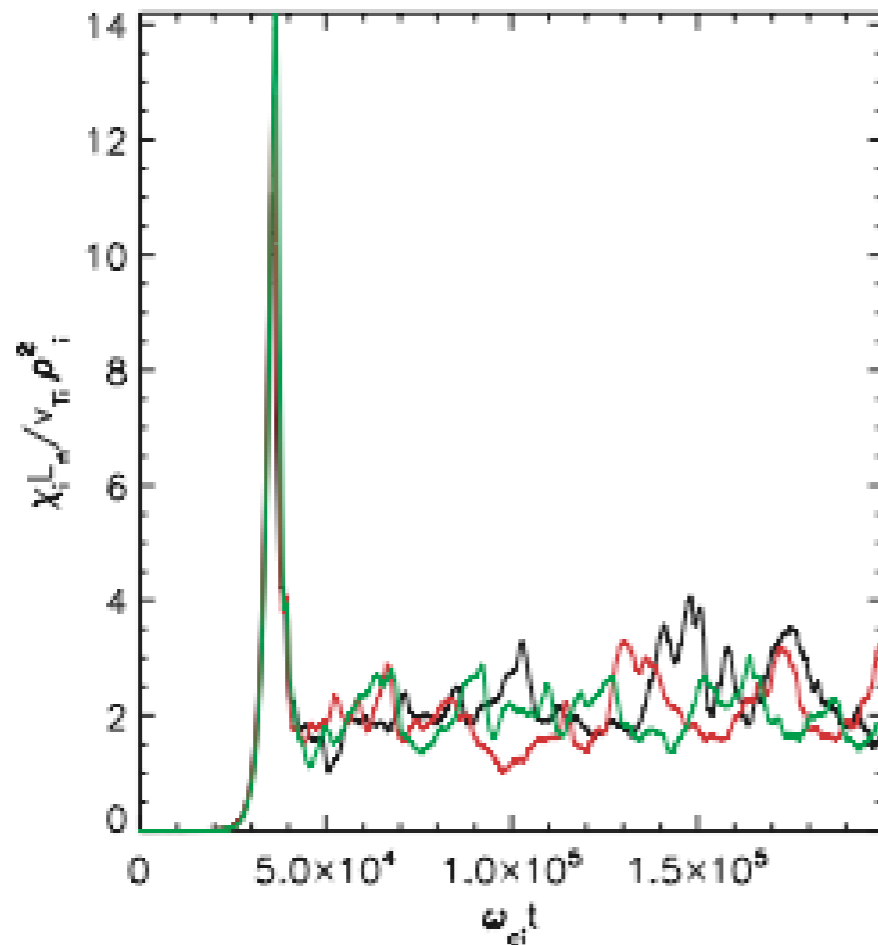
$$\tau \equiv T_e / T_i$$

$$\kappa_{Ti} \equiv -d \ln T_{0i} / dx$$

$$\langle \dots \rangle \equiv \int d\mathbf{x} / V$$

- Coarse graining in velocity space has to be taken with great care [Parker and Chen '06, Watanabe and Sugama '06]

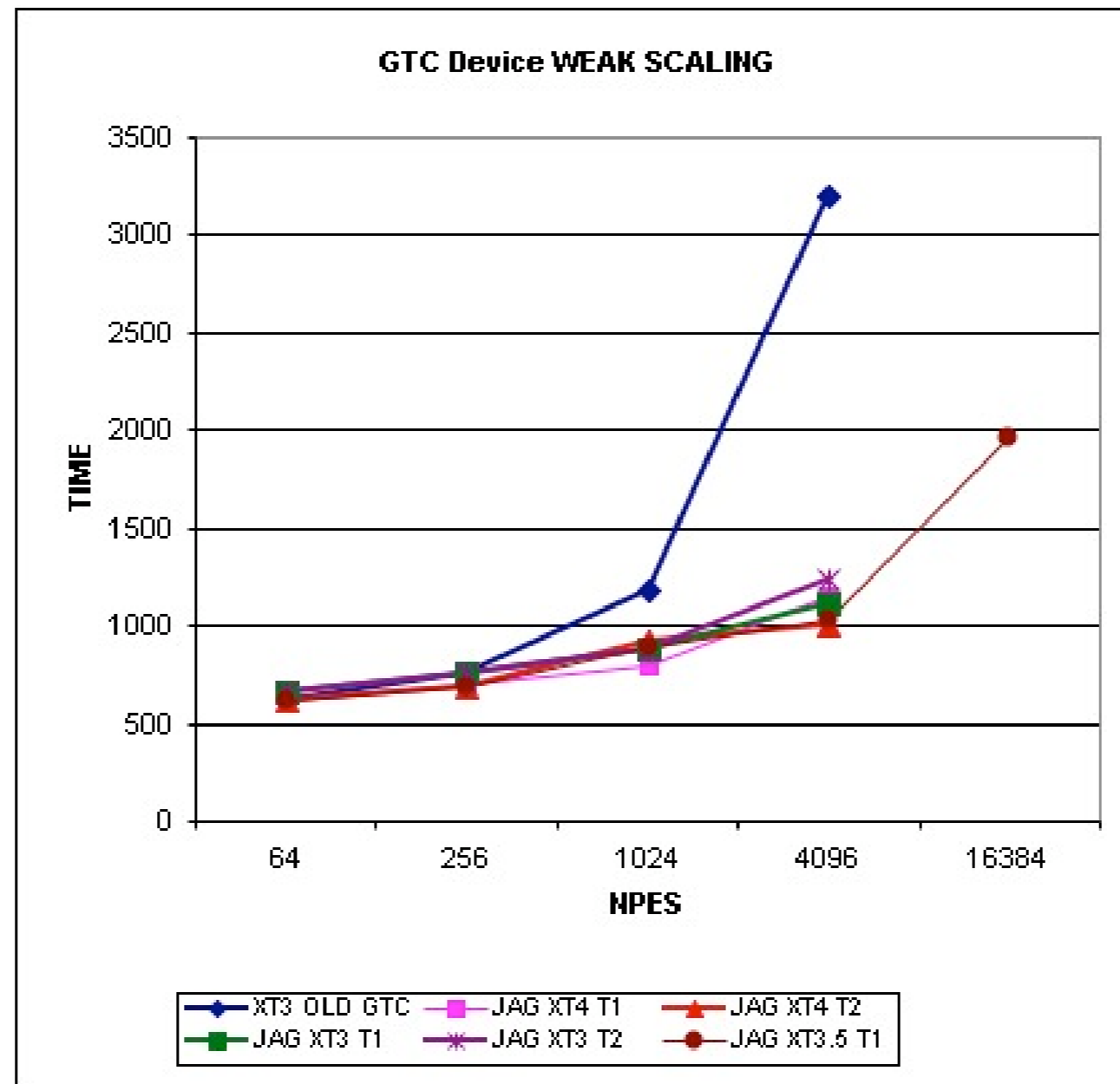
The Growing Weight Problem? and The Particle-Continuum Method



Resetting particle weights on a phase space grid periodically solves the so-called growing weight problem: no re-setting (black) vs. resetting (green and red)

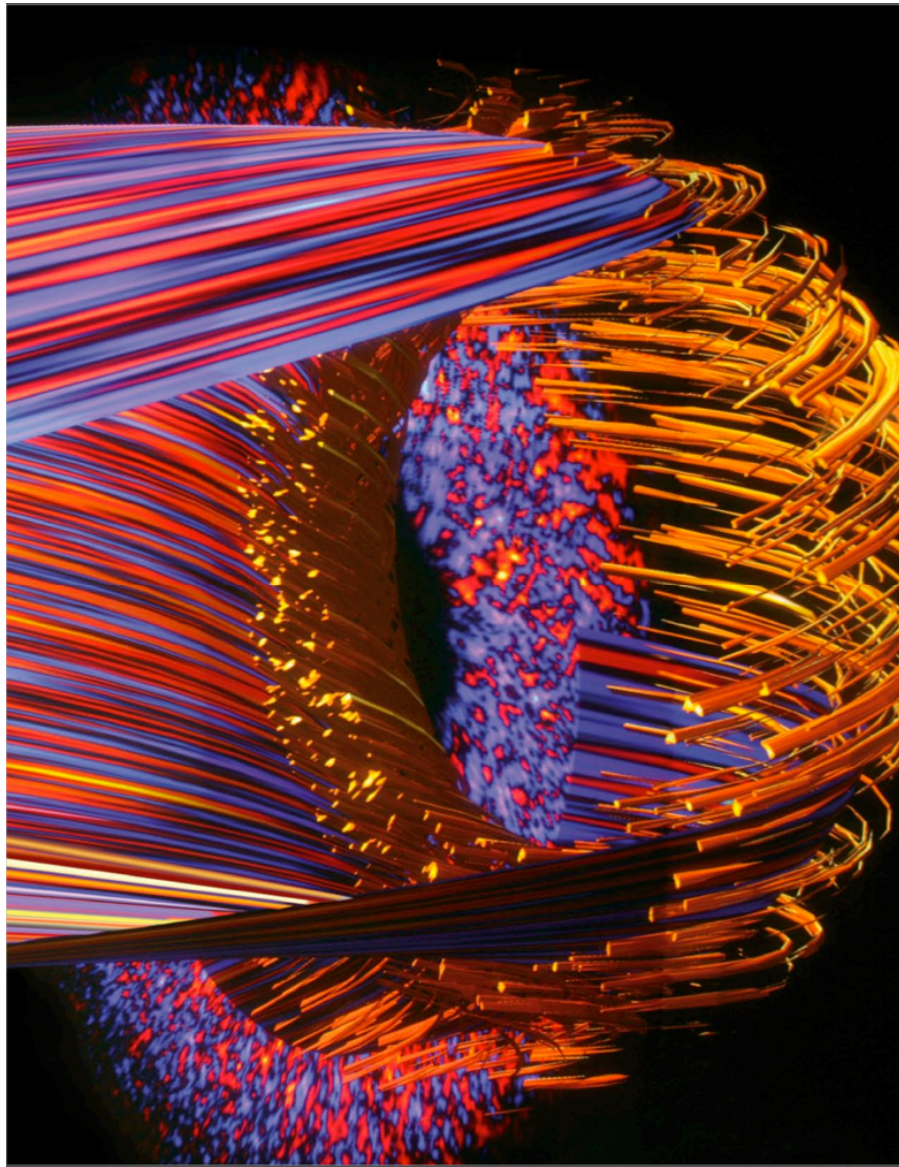
[Chen and Parker]

Two dimensional grid domain decomposition on GTC [Adams, Ethier, Wichmann]

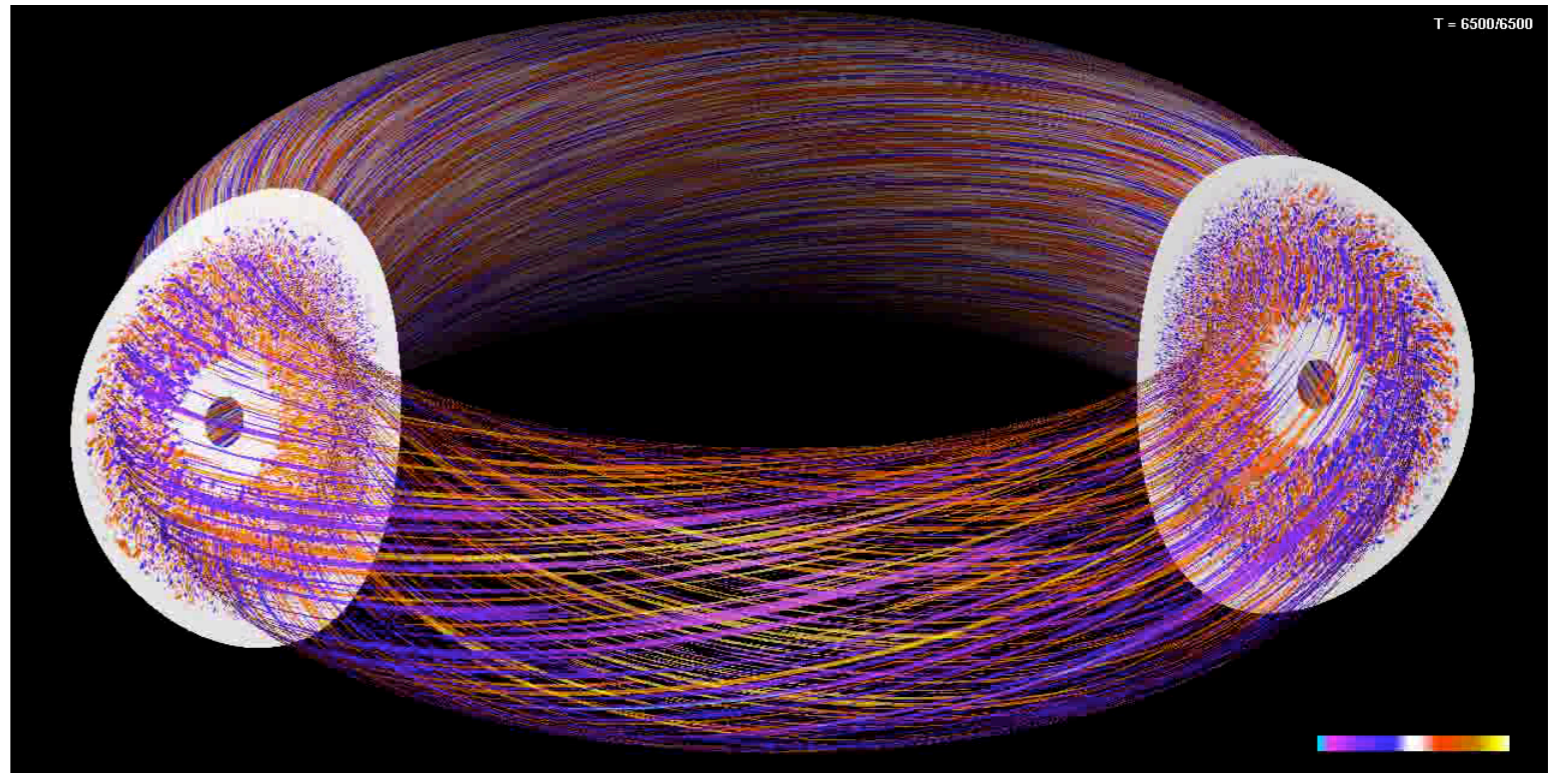


- It is important for simulating ITER-size devices
- Scaling inefficiencies point to large numbers of Translation Lookaside Buffer (TLB) misses on some processes on XT4/3 with larger grids.

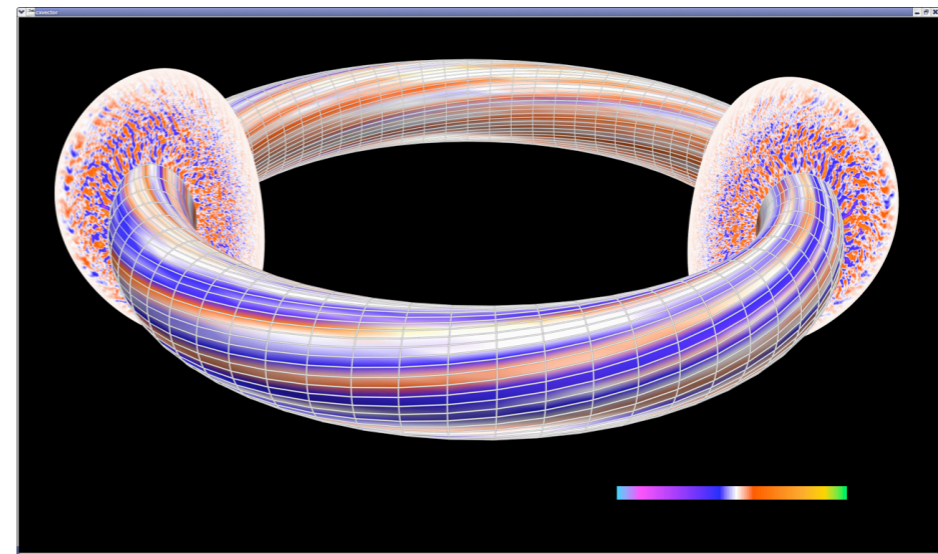
Visualization



[Klasky, ORNL; Ethier, Wang, PPPL]



[Klasky, ORNL; Ethier, Wang, PPPL]



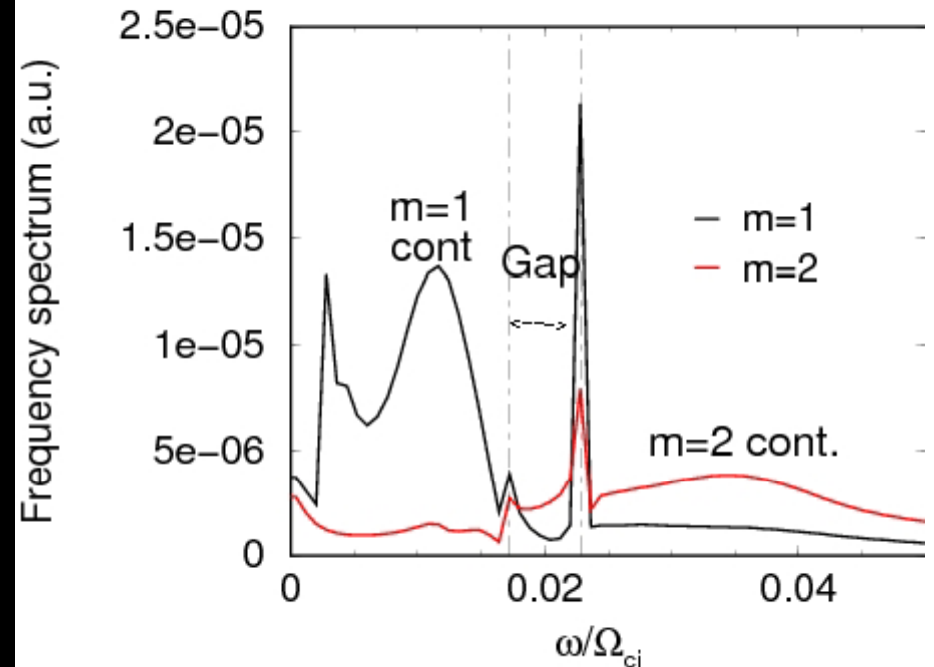
[Ma, UC-Davis]

Electromagnetic GTC development at UCI UCI

- **Kinetic electrons and electromagnetic fluctuations implemented using fluid-kinetic hybrid electron model**
- Enabled by a FEM field solver [*Nishimura, Lin, Lewandowski, Ethier, JCP2006*]
- FEM solver extension for edge geometry in collaboration with CPES [*Nishimura and Lin, CPP 2006*]
- A shaped version with 3D equilibrium is being applied for DIII-D resonant magnetic perturbation (RMP) experiments in collaboration with Evans and Joseph of GA
- Decyk of UCLA, in collaboration with Nishimura and Ethier, developed OO version for supporting multiple field solvers and multiple species (impurity and energetic particle), and for integrating with shaped version
- GTC selected for early application of 250TF ORNL computer

GTC electromagnetic simulation UCI

- Electromagnetic fluctuation via toroidal fluid-kinetic hybrid electron model
- Demonstrated finite- β stabilization of ITG mode
- Demonstrated Alfvén wave propagation in tokamak, damping via phase mixing, and existence of toroidal frequency gap
- [Nishimura, Lin, and Wang, *PoP*2007]



GTC nonlinear convergence in ETG simulation

UCI

Convergence from 400 to 2000 particles per cell

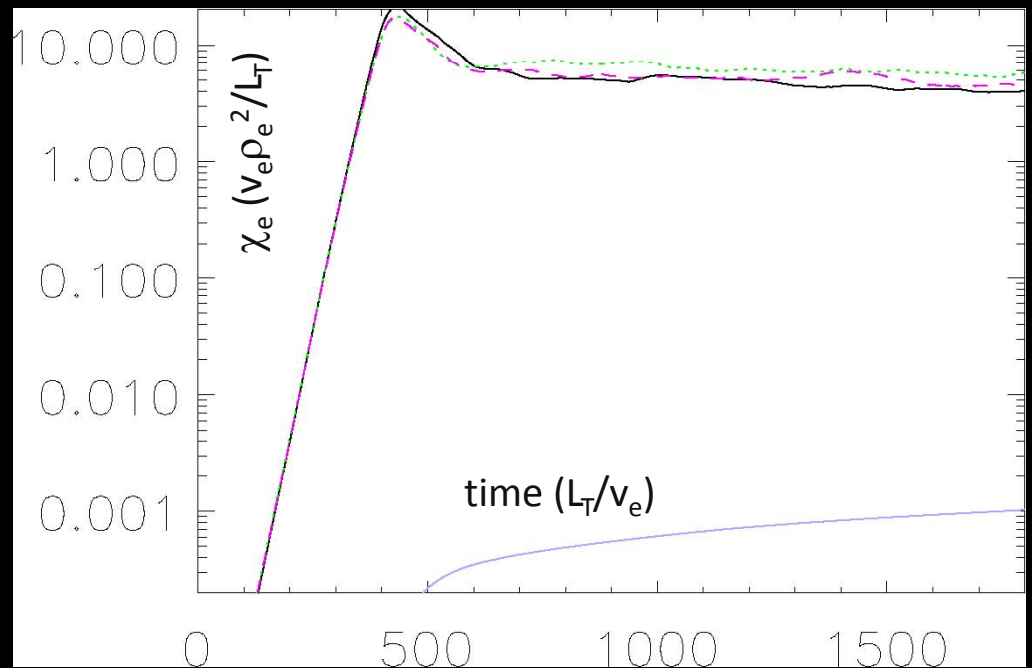
- ▶ Weak Cyclone case: $R/L_T=5.3$, $s=0.78$, $q=1.4$, $a/\rho_e=500$, $\gamma\sim\omega_r/4$

Noise driven flux is 4000 times smaller than ETG driven flux

Noise spectrum in ETG simulation measured. Noise driven flux calculated & measured, both in good agreement

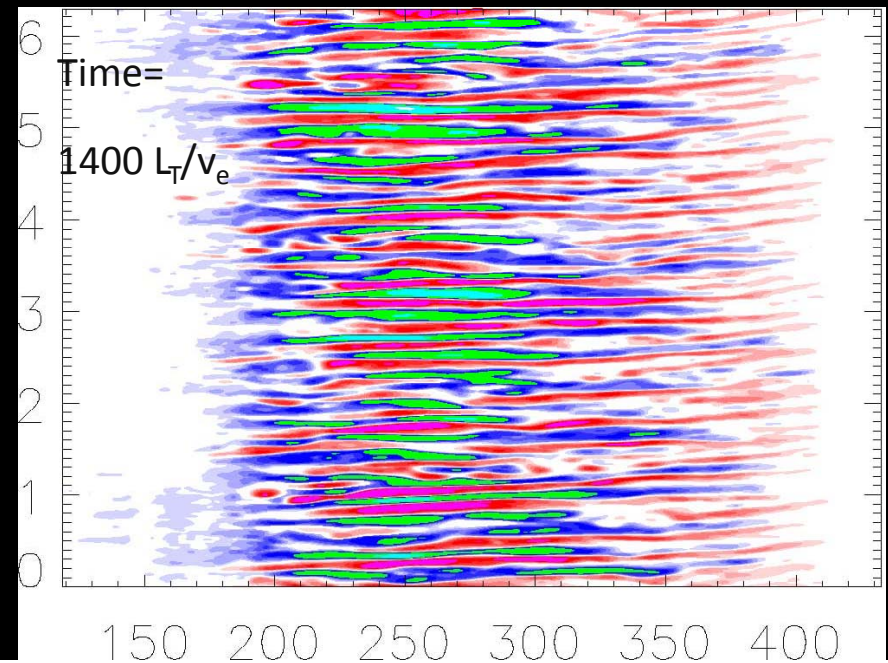
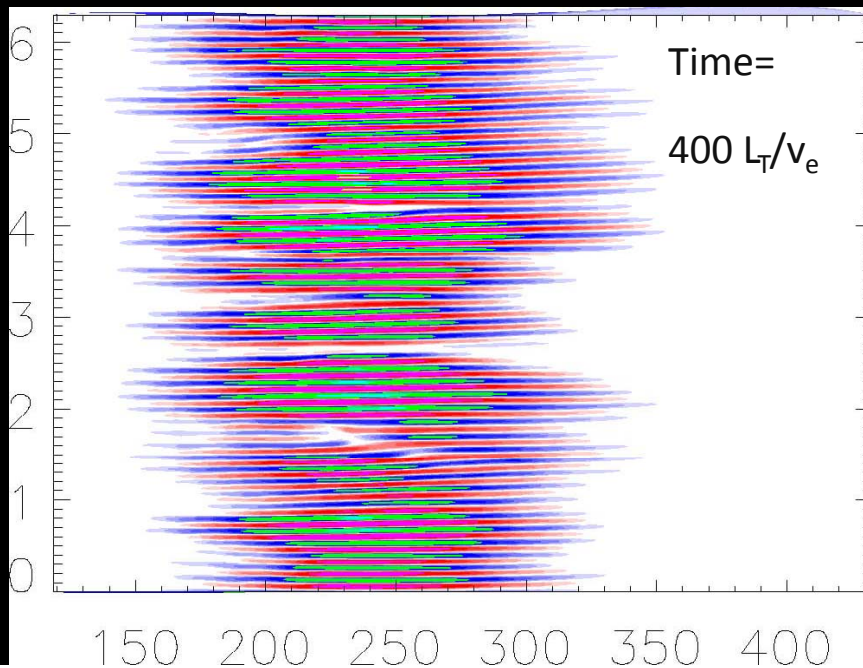
ORNL Cray XT3, 6400 PE, 4×10^{10} particles

[Holod and Lin, PoP2007]



Electron transport via wave-particle decorrelation UCI

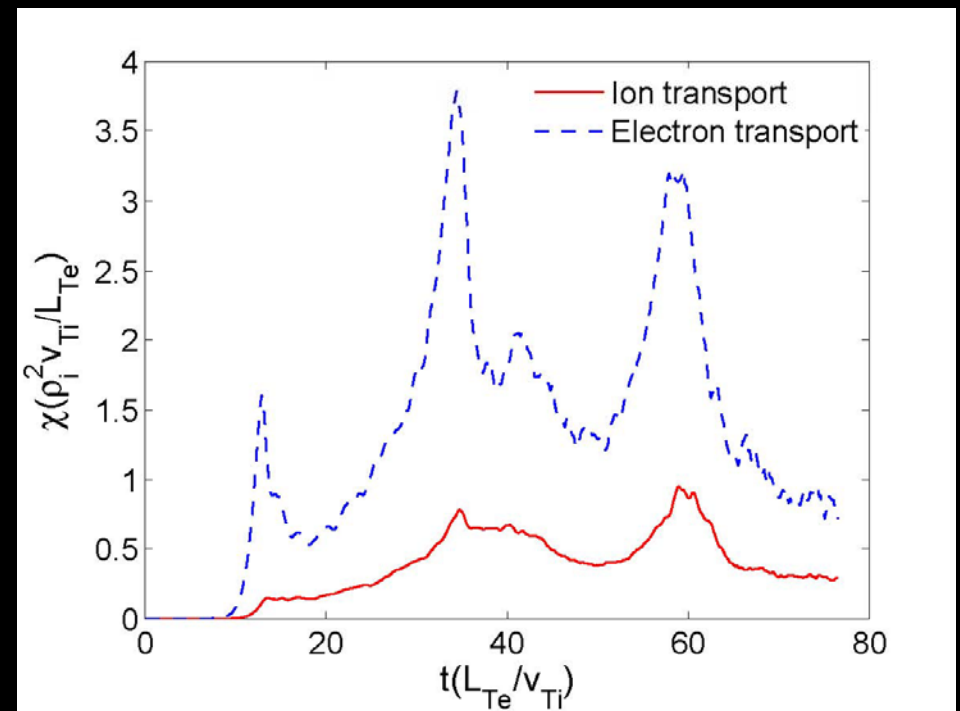
- Time scale separation between kinetic and fluid processes
 $\tau_{auto} \gg 1/\gamma \gg \tau_{wp} \sim 1/\Delta k_{\parallel} v_e$
- Electron heat transport dominated by wave-particle decorrelation of parallel resonance $\delta(\omega - k_{\parallel} v_{\parallel})$
- Quasilinear calculation agrees well with GTC simulation
- *[Lin et al, EPS2007 invited]*



GTC Simulation of CTEM turbulence UCI

- CTEM turbulence: electron transport dominant & bursty, zonal flows dominate saturation mechanism but not regulate the bursting.
- In collaboration with CPES [*Xiao and Lin, CPES spring meeting, 2007*]
- A kinetic theory was developed for radial propagation and damping of kinetic GAM mode [*Chen and Zonca, TTF2007*]

CTEM: $R/L_n=R/L_{Ti}=2.2$,
 $R/L_{Te}=6.9$, $s=0.78$, $q=1.4$



Conclusions

- It has been an exciting three years
- The PPPL team's work on
 - GTC performance on MPP
 - GTS and GTC-Neo and their V&V work
 - PNL and noise
 - Theory
- The UCI team's work on
 - ETG, TEM and noise using GTC
 - EM capability for GTC
- The Colorado team's work on
 - GEM for TEM
 - Particle continuum method
 - EM capability in GEM and the V&V work
- The SAPP team on
 - Solvers
 - Visualization
 - Data management

Future Directions

- Verification and & Validation
- Electromagnetic physics in GTS
- ITER simulation capabilities
- Integrated simulation: Heating, Turbulence, MHD, Transport