

# **Transport at the Magnetospheric Boundary: Magnetic Reconnection and Kelvin-Helmholtz Modes**

**A Otto**

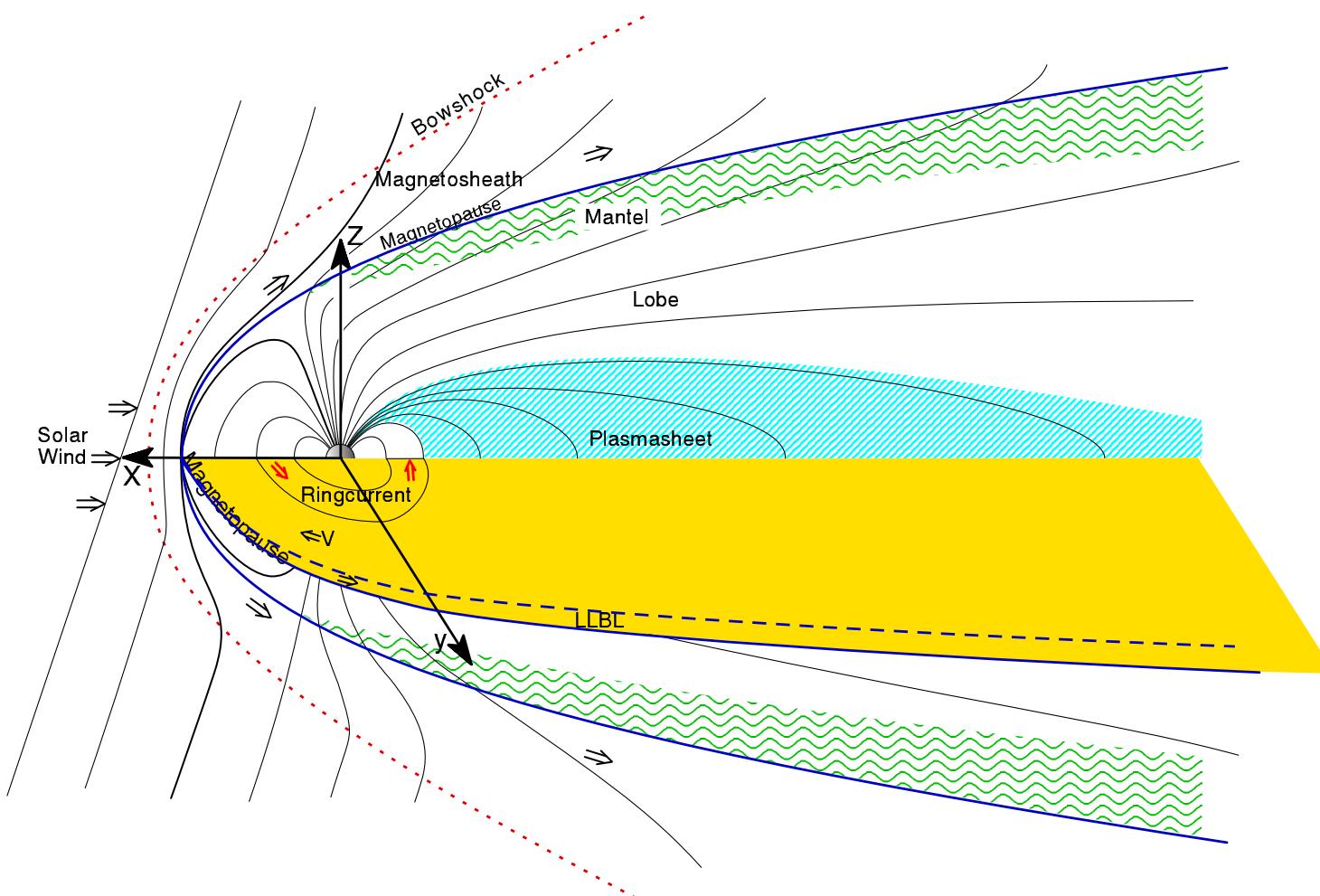
Geophys. Inst., Univ. Alaska, Fairbanks

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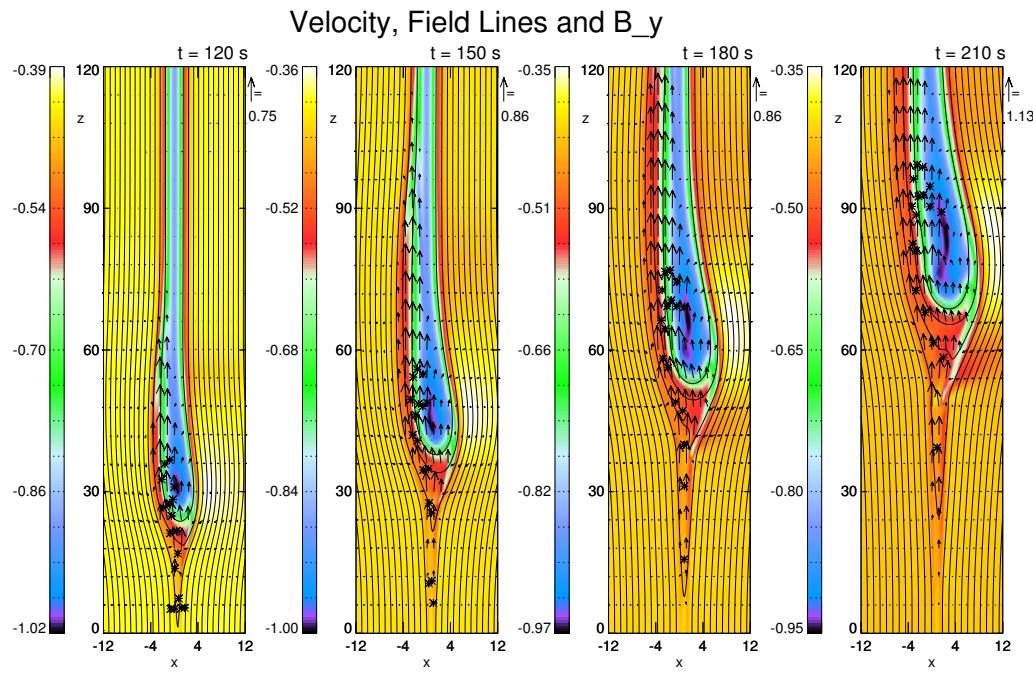
- **Motivation and Introduction**
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**Acknowledgements:** K. Nykyri, Q. Chen, F. Hall, J. Buechner, D. Fairfield

## Magnetospheric Boundaries:



## Dayside Reconnection for Southward Interplanetary Magnetic Field (IMF)



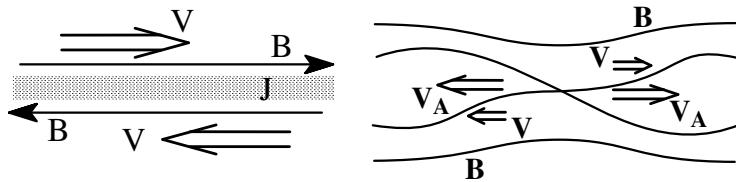
# Solar Wind - Magnetosphere Interaction

## Diffusion

- Thin boundaries  $r_{ci}, c/\omega_{pi}$

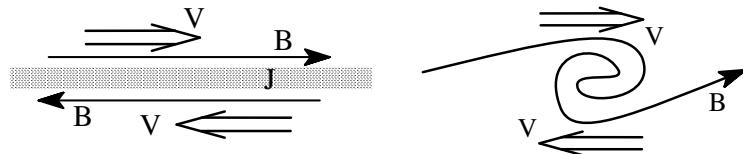
## Magnetic Reconnection:

- Transport of plasma, momentum, energy, and magnetic flux
- Requires large anti-parallel magnetic field components  $\Delta V_A > \Delta V$



## Kelvin-Helmholtz mode:

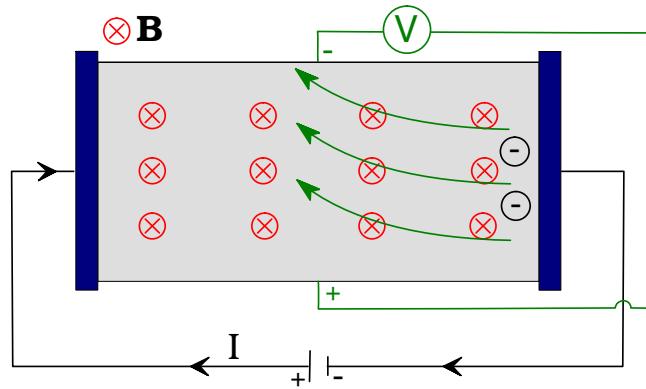
- Ideal instability => Transport of momentum and energy (viscous coupling)
- Requires  $\Delta v > V_{A,typ}$  along the k vector of the instability



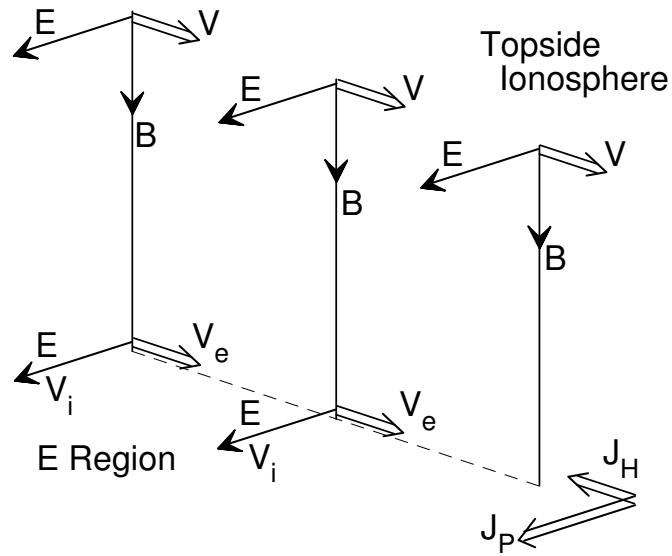
## Combination of tearing/reconnection with KH (Vortex induced reconnection) in 2D??

- Requires  $\Delta v > V_{A,typ}$  for KH and  $\Delta v < \Delta V_A$  for reconnection
- For strongly northward IMF:  $\Delta B$  is small along the flanks of the magnetosphere

## Hall Effect (Edwin Hall, 1879)



## Hall Currents in the Ionosphere



## Generalized Ohm's Law

$$\mathbf{E} + \mathbf{u} \times \mathbf{B} = \frac{m_e}{e^2 n} \left[ \frac{\partial \mathbf{j}}{\partial t} + \nabla \cdot (\mathbf{u} \mathbf{j} + \mathbf{j} \mathbf{u}) \right] - \frac{1}{en} \nabla p_e + \frac{1}{en} \mathbf{j} \times \mathbf{B} + \eta \mathbf{j}$$

with  $\eta = m_e \nu_c / ne^2$ .

## Normalized form ( $L_0, B_0, n_0, u_0, \dots$ )

$$\mathbf{E} + \mathbf{u} \times \mathbf{B} = \frac{c^2}{\omega_{pe}^2 L_0^2} \left[ \frac{\partial \mathbf{j}}{\partial t} + \nabla \cdot (\mathbf{u} \mathbf{j} + \mathbf{j} \mathbf{u}) \right] - \frac{c}{\omega_{pi} L_0} \nabla p_e + \frac{c}{\omega_{pi} L_0} \mathbf{j} \times \mathbf{B} + \eta \mathbf{j}$$

### Electron and ion inertia scale:

$$\lambda_e = \frac{c}{\omega_{pe}} = \left[ \frac{\epsilon_0 m_e c^2}{n_e e^2} \right]^{1/2} \approx 5 \text{ km (MSP)}, 1 \text{ to } 15 \text{ cm (SUN)}$$

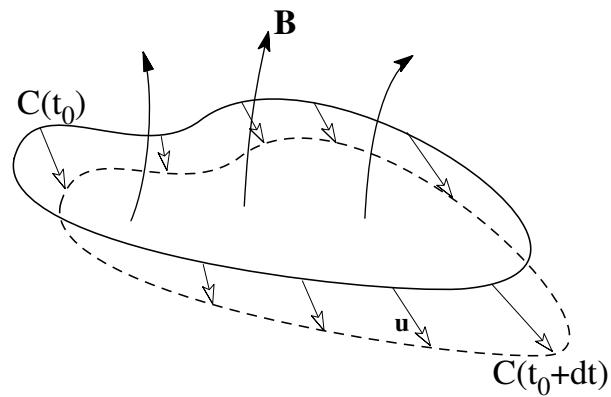
$$\lambda_i = \frac{c}{\omega_{pi}} = \left[ \frac{m_i}{m_e} \right]^{1/2} \frac{c}{\omega_{pe}} \approx 200 \text{ km (MSP)}, 3 \text{ to } 100 \text{ m (SUN)}$$

Magnetic diffusion time scale:  $\tau_{diff} = L_0^2 / (\lambda_i^2 \nu_c)$

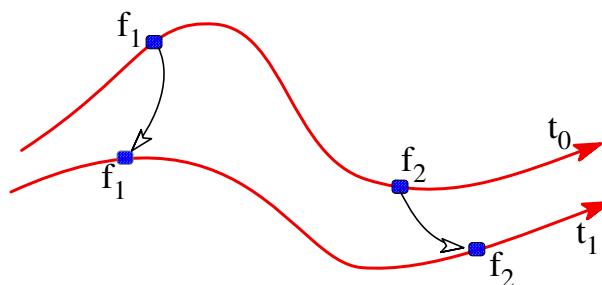
## Frozen-in Condition

$$\mathbf{E} + \mathbf{u} \times \mathbf{B} = 0$$

**(a) Magnetic flux through any cross section  $C(t)$  moving with velocity  $\mathbf{u}$  remains constant**



**(b) Two fluid elements which are on the same field line remain connected by this field line for all times.**



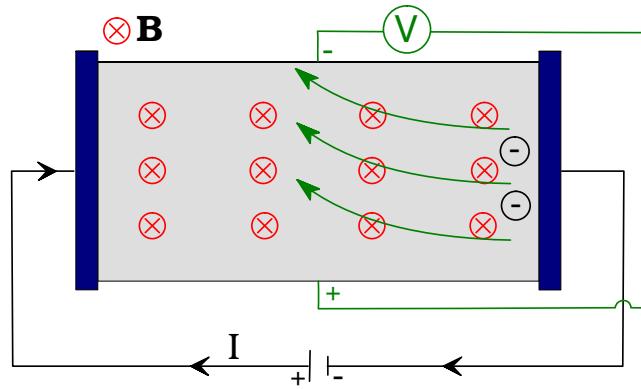
**Including the Hall term:**

$$\mathbf{E} + \mathbf{u} \times \mathbf{B} = \frac{1}{en} \mathbf{j} \times \mathbf{B}$$

$$\mathbf{u}_e = \mathbf{u} - \frac{1}{en} \mathbf{j} \quad \Rightarrow$$

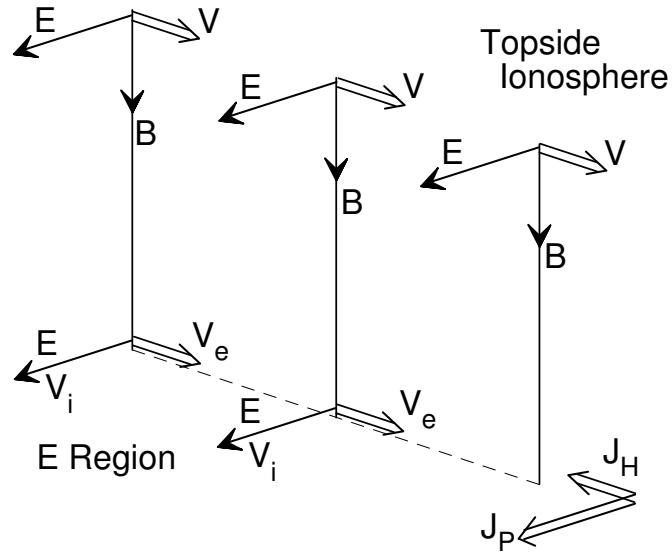
$$\mathbf{E} + \mathbf{u}_e \times \mathbf{B} = 0$$

## Hall Effect



$$\mathbf{E} + \mathbf{u}_e \times \mathbf{B} = \eta \mathbf{j}$$

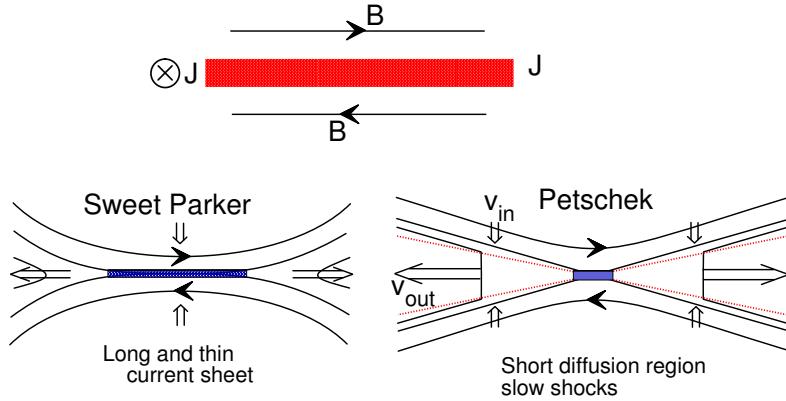
## Hall Currents in the Ionosphere



$$\omega_{gi} < \nu_{ic} \quad \text{and} \quad \omega_{ge} > \nu_{ec}$$

# Magnetic Reconnection

$$\mathbf{E} + \mathbf{u} \times \mathbf{B} = \eta \mathbf{j}$$



**Reconnection rate:** Rate with which magnetic flux is transported from the inflow to the outflow region

$$R \approx \frac{v_{in}}{v_{out}} \approx \frac{d}{l}$$

- **Sweet and Parker:**

$$R_{SP} = R_m^{-1/2} = [L\sigma v_{A0}\mu_0]^{-1/2} = \text{slow!!!!}$$

- **Petschek:**

$$R_{SP} = \frac{\pi}{8 \ln R_m} \approx 0.1 = \text{fast!!!!}$$

# Hall Physics

## Inclusion of Hall term

$$\mathbf{E} + \mathbf{u} \times \mathbf{B} = \frac{1}{en} \mathbf{j} \times \mathbf{B}$$

=> Electron whistler wave

$$\omega_w = k^2 c^2 \omega_{ge} \omega_{pe}^{-2}$$

- Phase and group velocities  $\sim k$ . => Numerical limitation

## Courant Condition for explicit simulation methods:

$$\Delta t \leq \frac{v_{max}}{\Delta x}$$

but including Whistler dynamics:

$$\Delta t \leq const \frac{v_{Alfven}}{(\Delta x)^2}$$

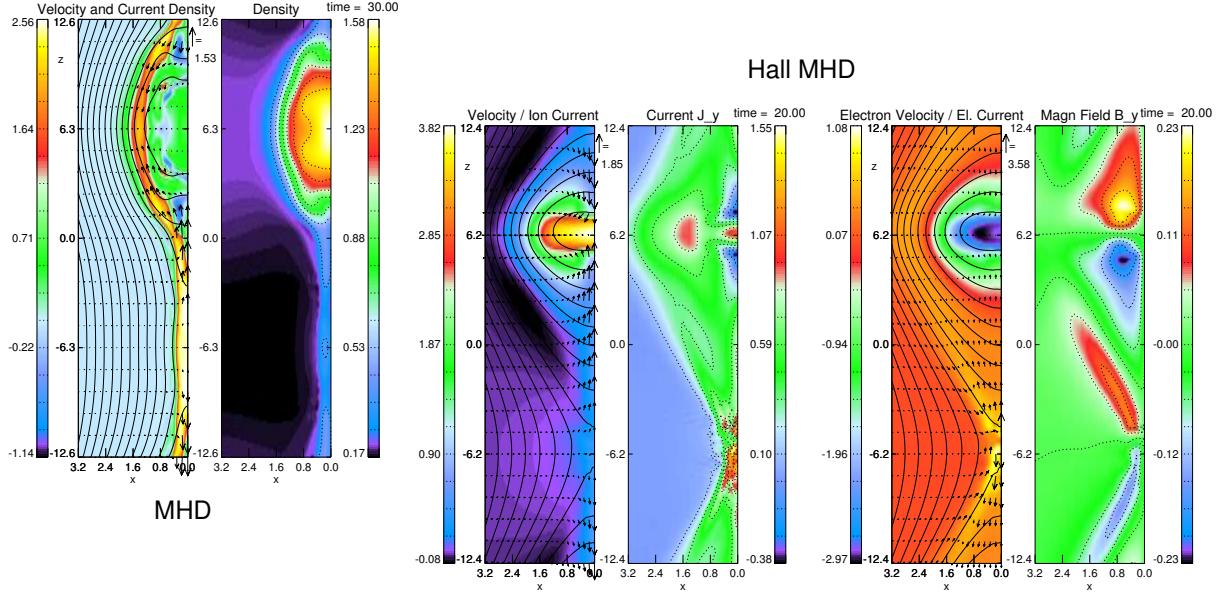
Numerical effort in two dimensions  $\sim \frac{1}{(\Delta x)^4}$

## GEM Reconnection Challenge

- Identical (*simple*) initial and boundary conditions
- Two dimensions
- Comparison of different plasma approximations
  - Full particle electromagnetic simulation
  - Hybrid (ions kinetic and electrons fluid) simulation
  - Two-fluid simulation
  - Hall MHD simulation
  - MHD simulation

## Results: Hall MHD

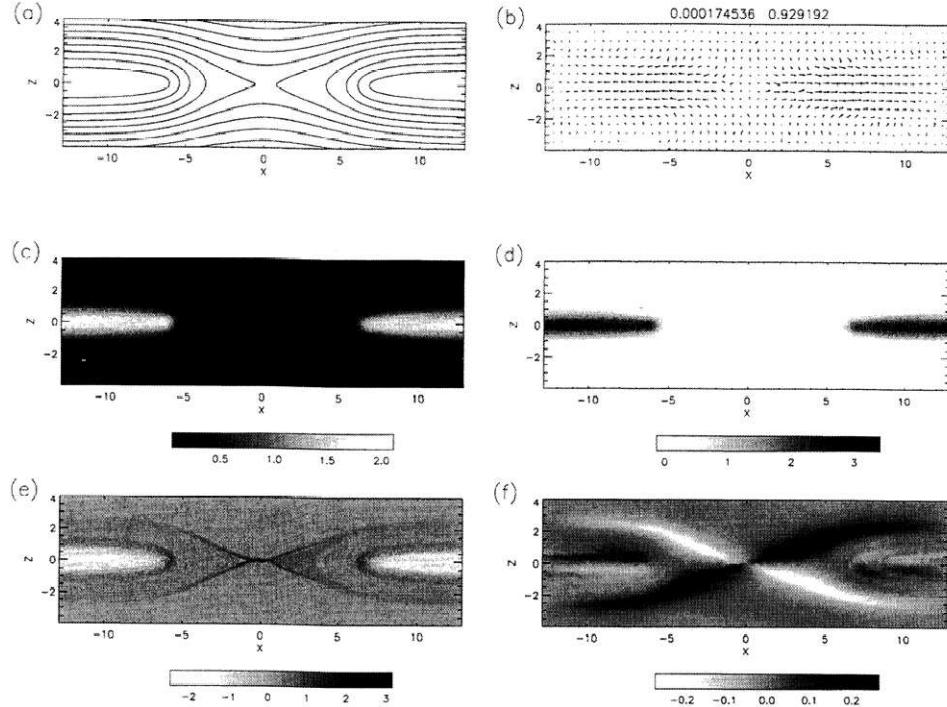
Unit distance = Ion inertia scale ( $c/\omega_{pi}$ )



- Electron current concentrated in the diffusion region
- Electron & ion dynamics decouple on scale  $< c/\omega_{pi}$
- ‘Quadrupolar’  $B_y$  signature
- Fast reconnection

## Results: Hybrid model (ions kinetic + electrons fluid)

Shay et al. (2001)



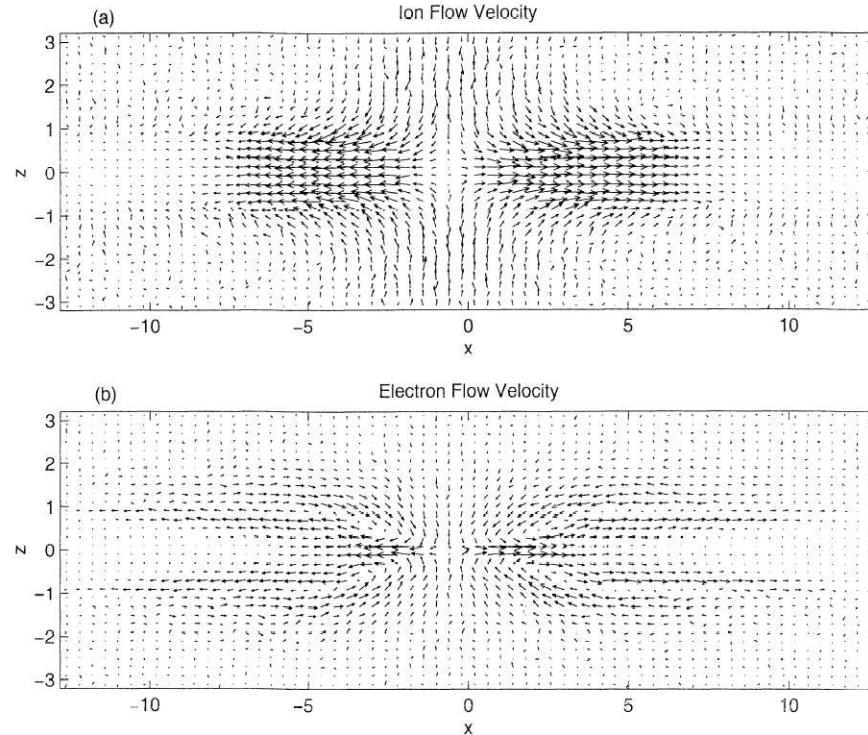
**Figure 2.** Results from a hybrid run with  $m_e/m_i = 1/100$ : (a) Magnetic field lines, (b) ion in-plane flows, (c) density, (d) ion out-of-plane current, (e) electron out-of-plane current, and (f) out-of-plane magnetic field.

## Results: Full particle (electromagnetic) model

Pritchett (2001)

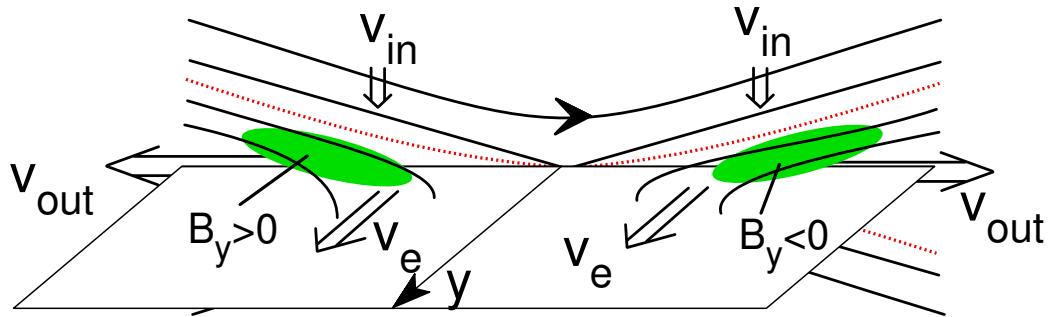
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PRITCHETT: COLLISIONLESS RECONNECTION



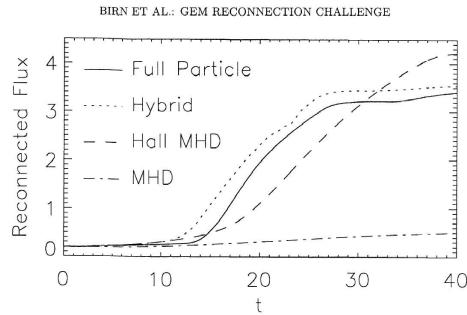
**Figure 3.** The in-plane ( $x, z$ ) (a) ion and (b) electron flow velocities at the same time as in Plate 1.

## Reconnection with Hall Physics

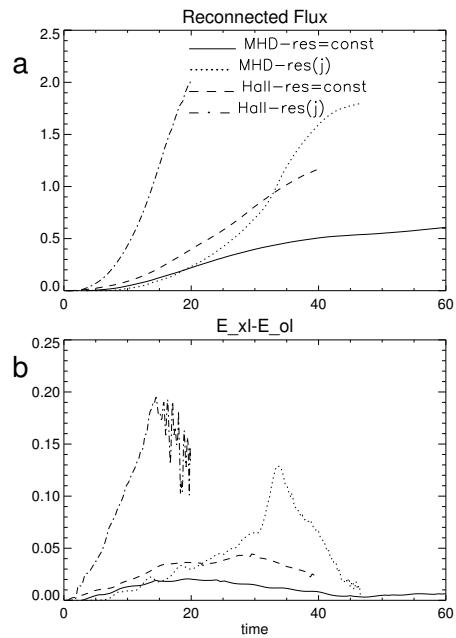


- Simple Initial geometry with magnetic field in the  $x, z$  plane ( $B_y(t = 0) = 0$ )
- Reconnection generates  $B_z$
- Current in diffusion region = electron current  $\Rightarrow v_e$  in the  $y$  direction  
 $\Rightarrow$  motion of magnetic foot points in symmetry plane along  $y$  (frozen-in condition)  
 $\Rightarrow$  generation of  $B_y$  with correct polarity

## Reconnection rates:



Birn et al. (2001)



- Fast reconnection with  $R=0.2$
- Reconnection rates very similar for different plasma approximations
- Exception: simple MHD
- All models with fast reconnection include Hall physics

## Fast Reconnection by including Hall physics?

- Bulk velocity for current carriers in a current sheet with  $L_0 = c/\omega_{pi}$ :

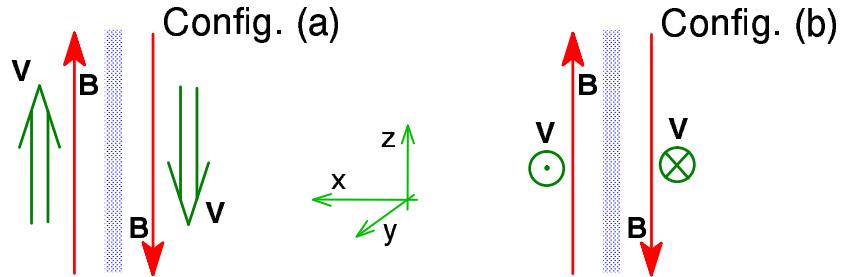
$$v_e = \frac{B_0 \omega_{pi}}{en\mu_0 c} = v_A$$

- Length of current sheet =  $O(c/\omega_{pi})$
- Width limited because of the generation of  $B_y$

=> limits aspect ratio of the diffusion region  $d/l$  => fast reconnection

## Reconnection determined by Hall physics with a unique reconnection rate?

Simple configurations with plasma flow:

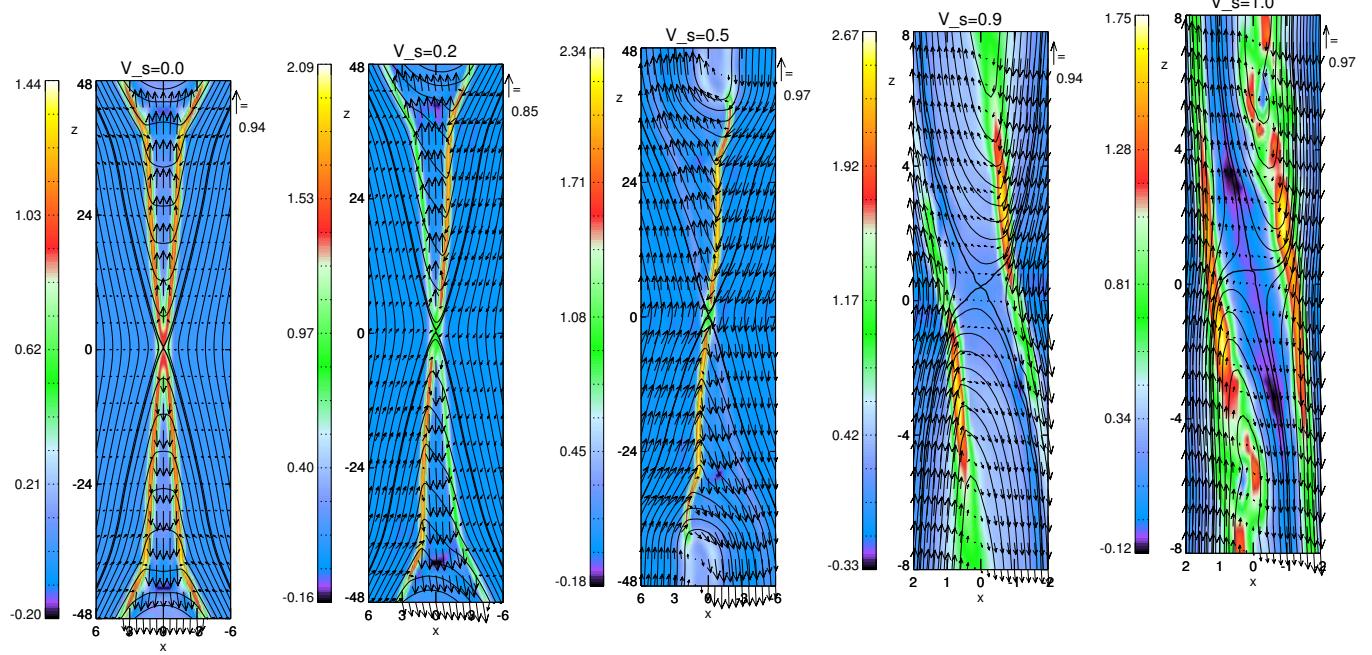


- Configuration (a): Plasma flow aligned with the magnetic field
- Configuration (b): Plasma flow in the invariant direction

## Configuration (a): Flow shear along antiparallel magnetic field

MHD Case:

Magn. Field, Velocity, and Current Density  $t = 110$



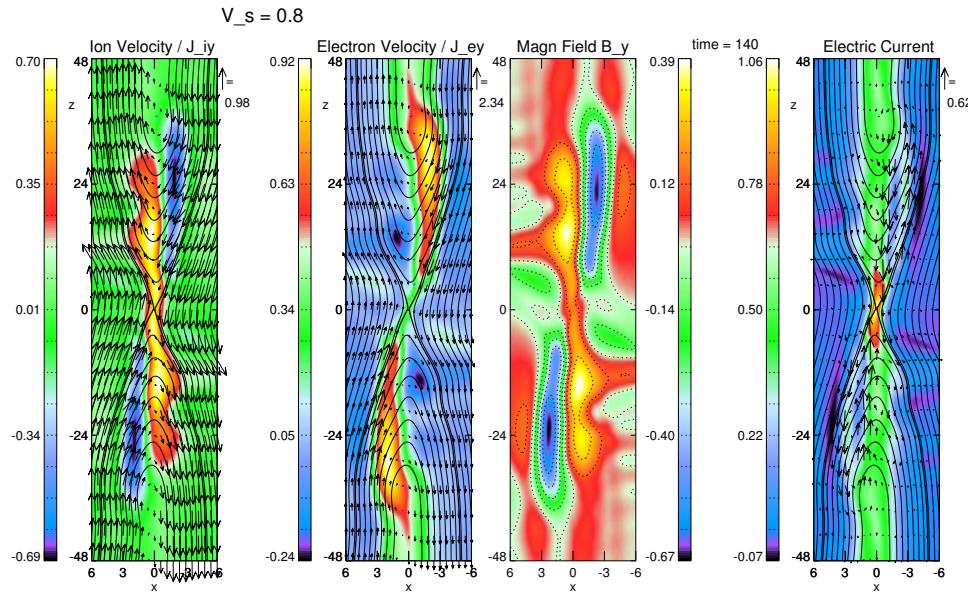
Main results:

- Asymmetric current layers
- Reconnection is switched off for  $v_{shear} = v_A$

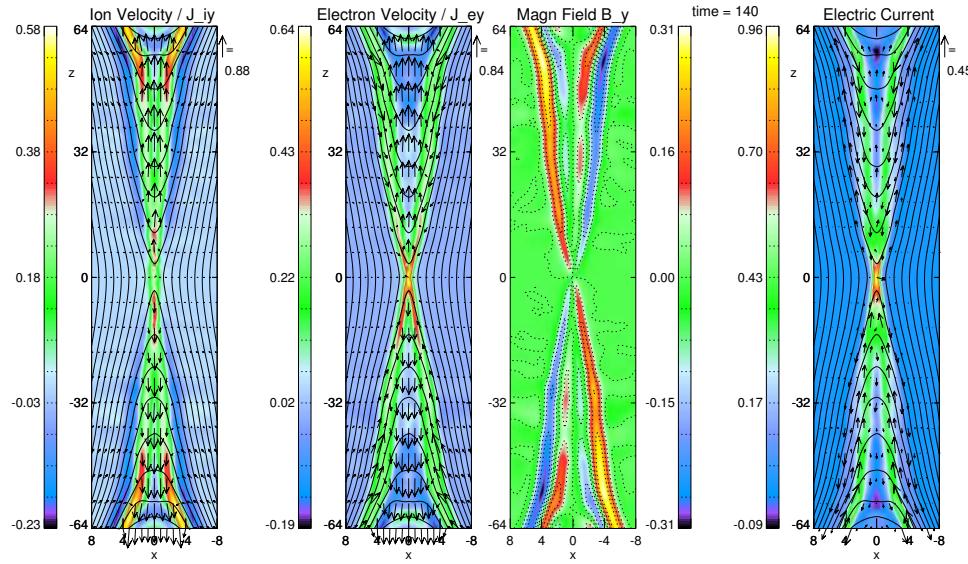
Mechanisms:

- Acceleration into outflow region is caused  $\mathbf{j} \times \mathbf{B}$  force  
=> larger acceleration required for side with shear flow opposite to outflow  
=> larger current density required on this side
- Shear flow faster than Alfvén speed does not permit information to travel upstream from  $x$  line  
=> reconnection is switched off

**Hall MHD for  $v_{sh} = 0.8$ :**



$v_{sh} = 0$ :

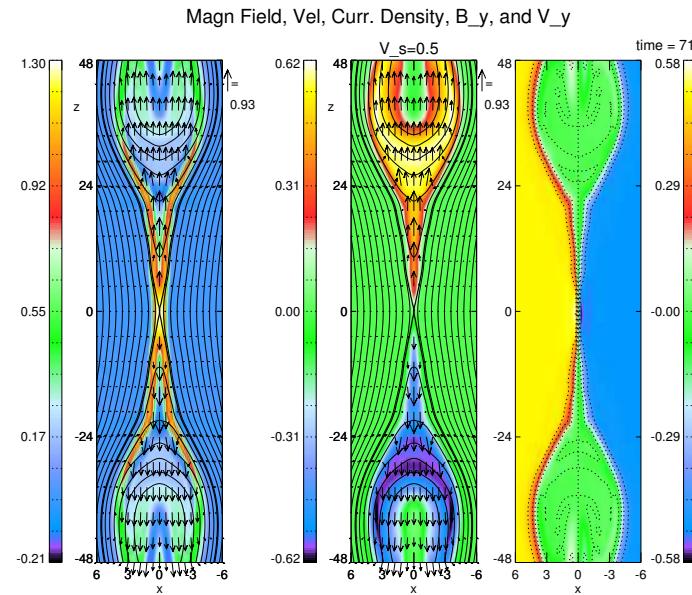


**Many similarities between MHD cases and the cases with Hall term included!**

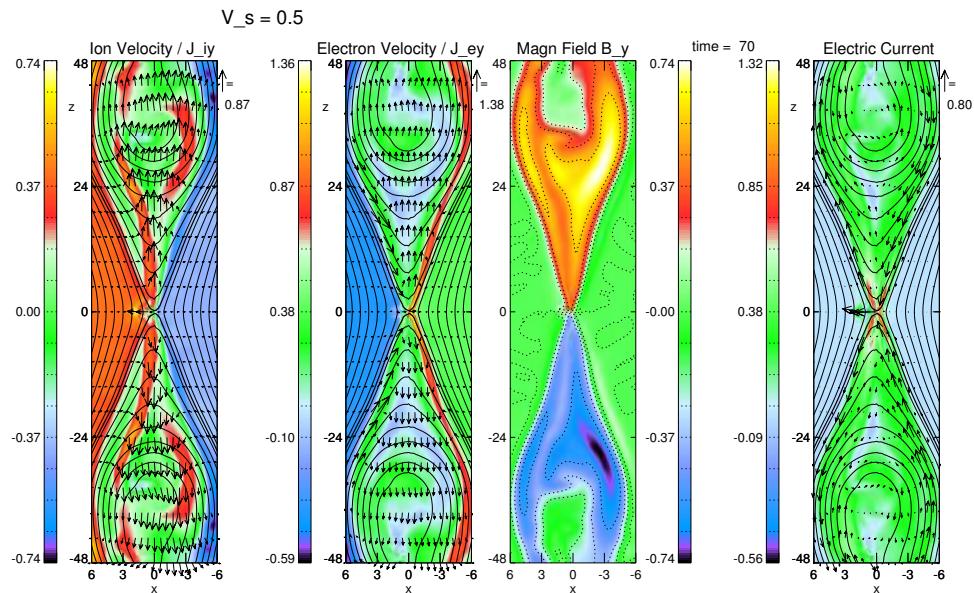
- **Same asymmetry:** Main current layer shifted toward the side with flow antiparallel to the respective outflow region.
- **Same stabilization mechanism and limits ( $v_{sh} \simeq v_A$ )**

## Configuration (b): Flow shear along the current direction

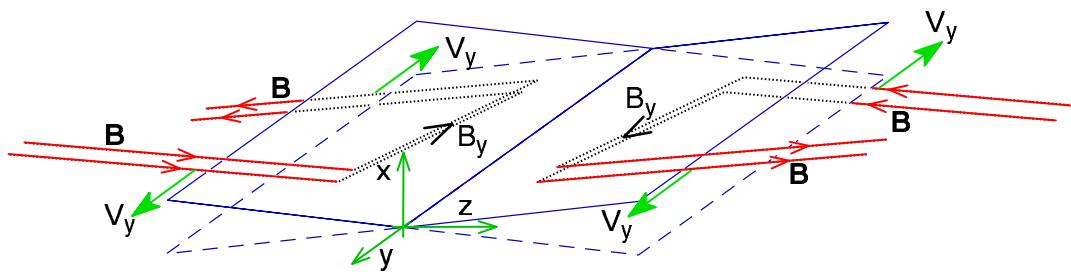
**MHD Case:**



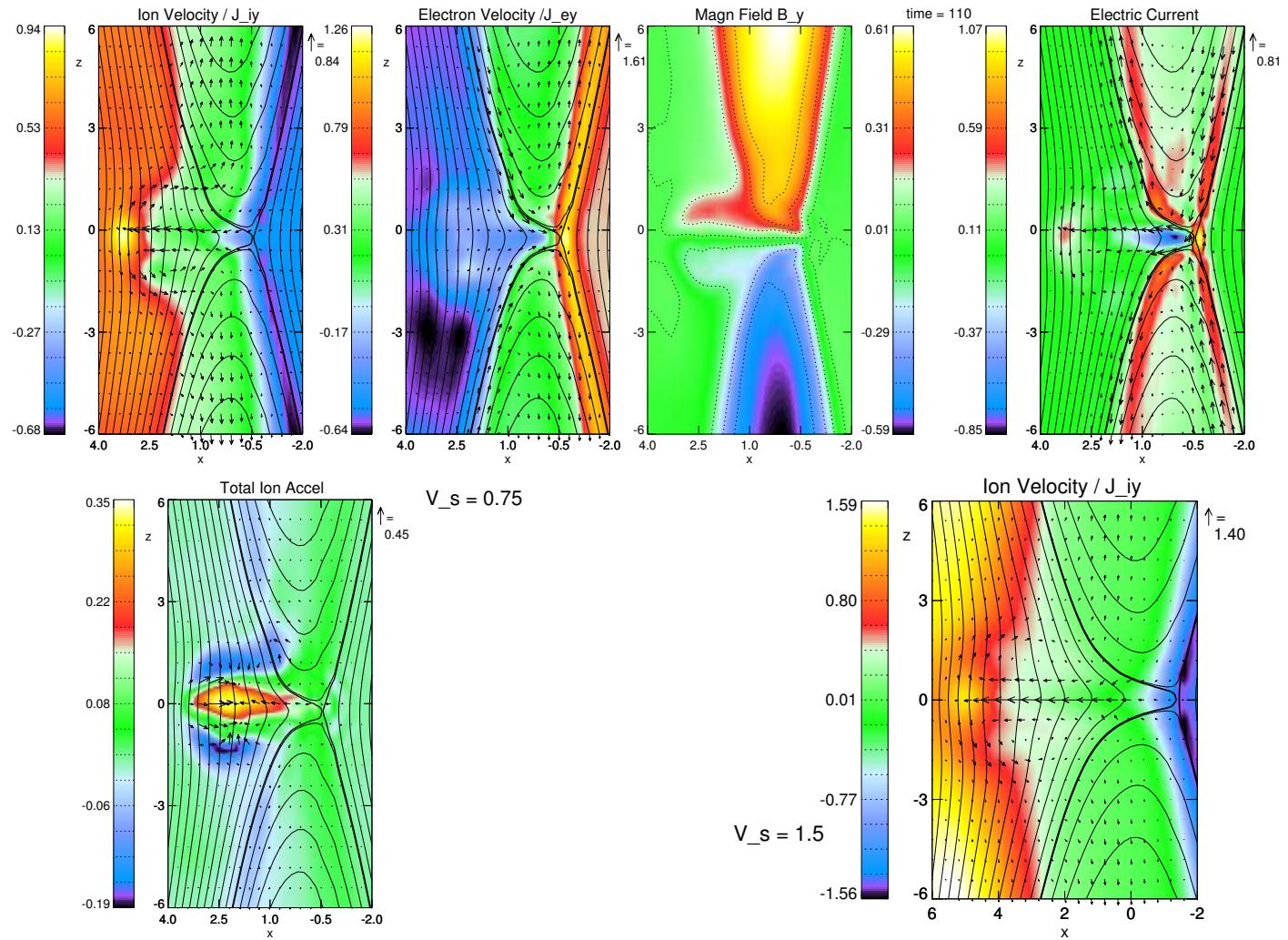
**Hall Case:**



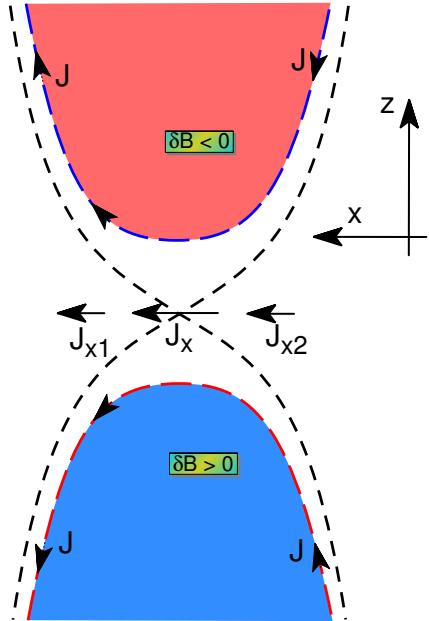
## Physical mechanism for the generation of $B_y$ :



## “Ion fountain” reconnection

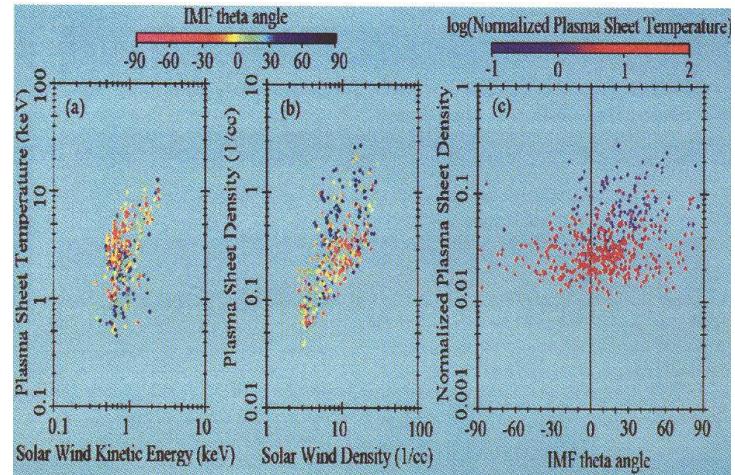
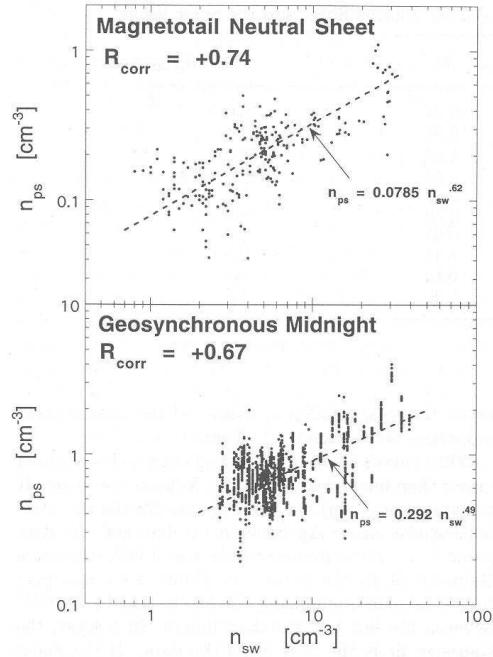


## Physical mechanism for the unusual ion flow:



- Opposite polarities of  $B_y$  require current in the  $x$  direction:  $j_x \simeq \frac{\delta B_y \omega_{pi}}{\mu_0 c}$  with a drift velocity of  $v_{xd} = v_A \frac{\delta B_y}{B_0}$
- Reconnection rate:  $\epsilon_r \Rightarrow$  electron inflow velocity of  $v_{ex} = \epsilon_r v_A$
- Electron inflow not sufficient to carry the required current for  $v_{ex} < v_{xd}$  or  $\epsilon_r < O(1) \frac{\delta B_y}{B_0}$   
 $\Rightarrow$  ion jet needed to carry current to support the bipolar  $B_y$

# Plasma Entry into the Magnetotail (for Northward IMF)

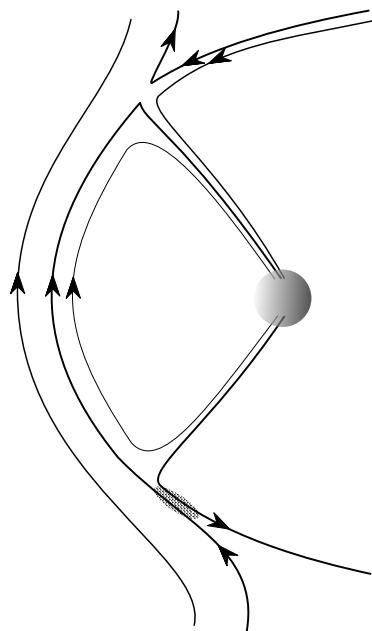


**Borovsky et al. (1998):** Neutral sheet density versus solar wind density.

**Fujimoto et al. (1998):** Neutral sheet temperature and density versus solar wind conditions.

## Properties:

- Correlation between plasma sheet density and SW density
- Cold dense plasma sheet for northward IMF
- Penetration times:
  - Midtail: 1 hour
  - Near Earth region: several hours
- High density and low temperature filaments on closed field lines (slow convection)
- Energetic (magnetospheric) and cold ion (magnetosheath) populations
- Possible Mechanisms for entry:
  - Re-reconnection at high latitudes (Song and Russell, 1992)
  - Diffusion

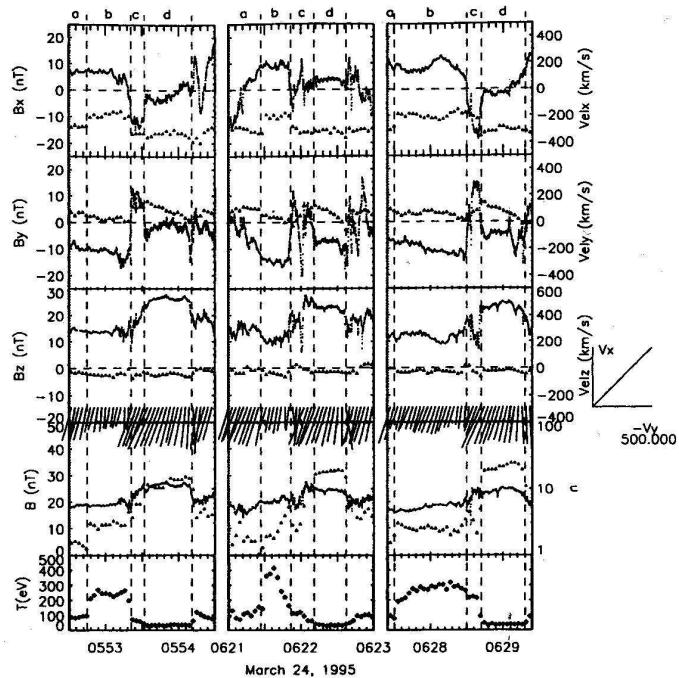


# Kelvin-Helmholtz Mode at the Magnetospheric Boundary

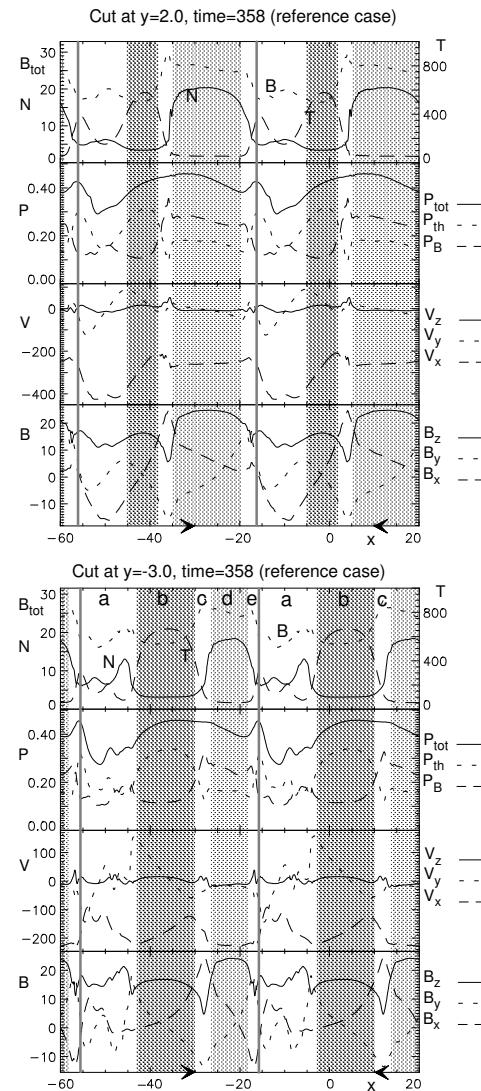
Many Observations of quasi-periodic boundary oscillations

(Schopke et al., 1981; Ogilvie and Fitzenreiter, 1989; Chen and Kivelson, 1993, ..)

Fairfield et al. (2000)

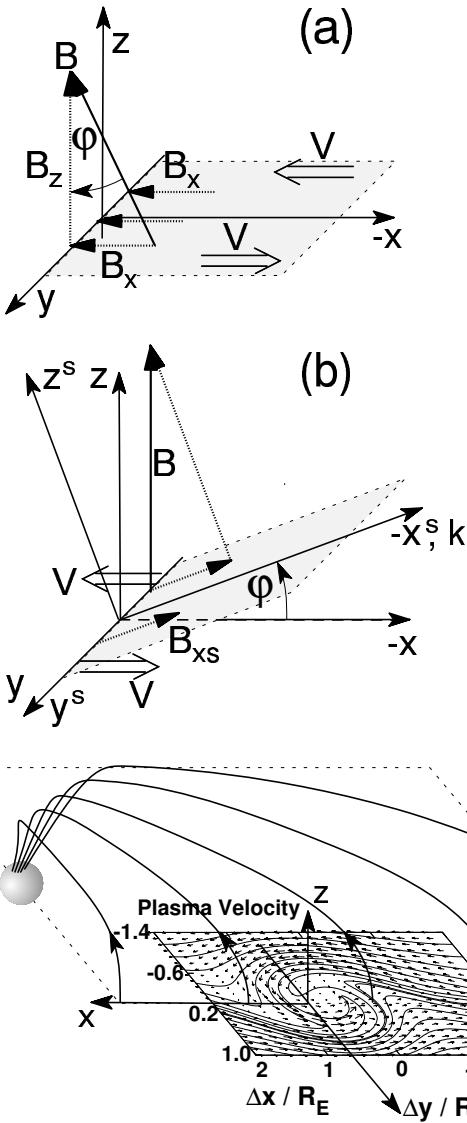


Otto and Fairfield (2000)



## Initial Geometry:

No magnetic field deformation if  $\mathbf{B} \perp \mathbf{k}_{KH}$ !!



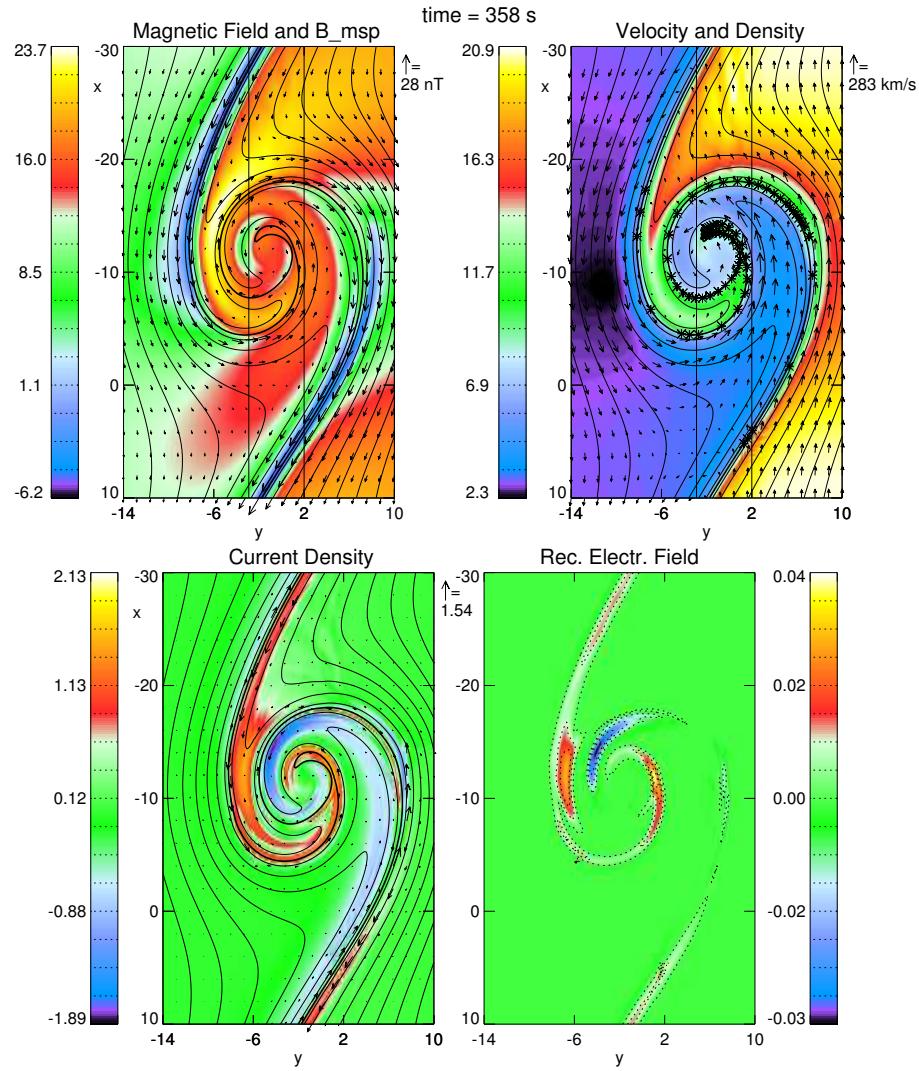
## Initial Parameters:

Location	$B$ , nT	$V$ , km s <sup>-1</sup>	$n$ , cm <sup>-3</sup>	$\beta$
M'sphere	16	0	2.8	2.8
M'sheath	24	315	19.2	0.7

## Normalization:

	Units
Magnetic field $B_0$	16 nT
Number density $n_0$	11 cm <sup>-3</sup>
Current density $j_0$	22 nA m <sup>-3</sup>
Length scale $L_0$	600 km
Velocity $v_A$	105 km s <sup>-1</sup>
Time $\tau_A$	5.7 s

# Structure of the KH Vortex



## **Signatures:**

### **Region (a):**

- **Core of the KH vortex**
- **Strong fluctuations in all plasma and field parameters**
- **high density spikes**

### **Region (b):**

- **Magnetospheric plasma trailing the core of the vortex**
- **Magnetospheric temperatures and densities**
- **Small fluctuations Steady decrease in  $v_y$**

### **Region (c):**

- **Outbound transition into sheath-like material**
- **Decrease in temperature and increase in density**
- **Minimum in  $B_z$**
- **Spikes in  $B_x$  and  $B_y$  with opposite polarity**

### **Region (d):**

- **Magnetosheath-like interval**
- **Steady plasma and field properties**

### **Region (e):**

- **Inbound transition into vortex**
- **Minimum in  $B_z$**
- **Spikes in  $B_x$  and  $B_y$**

# Stability

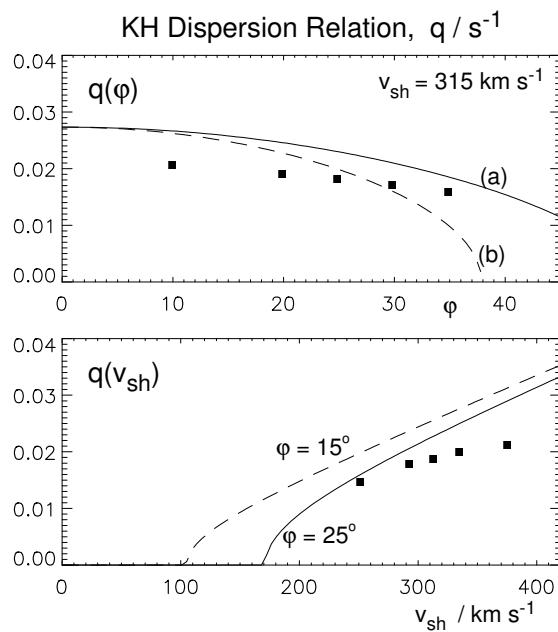
**Growth rate:**

$$q = [\alpha_1 \alpha_2 [(\mathbf{V}_1 - \mathbf{V}_2) \cdot \mathbf{k}]^2 - \alpha_1 (\mathbf{V}_{A1} \cdot \mathbf{k})^2 - \alpha_2 (\mathbf{V}_{A2} \cdot \mathbf{k})^2]^{1/2}$$

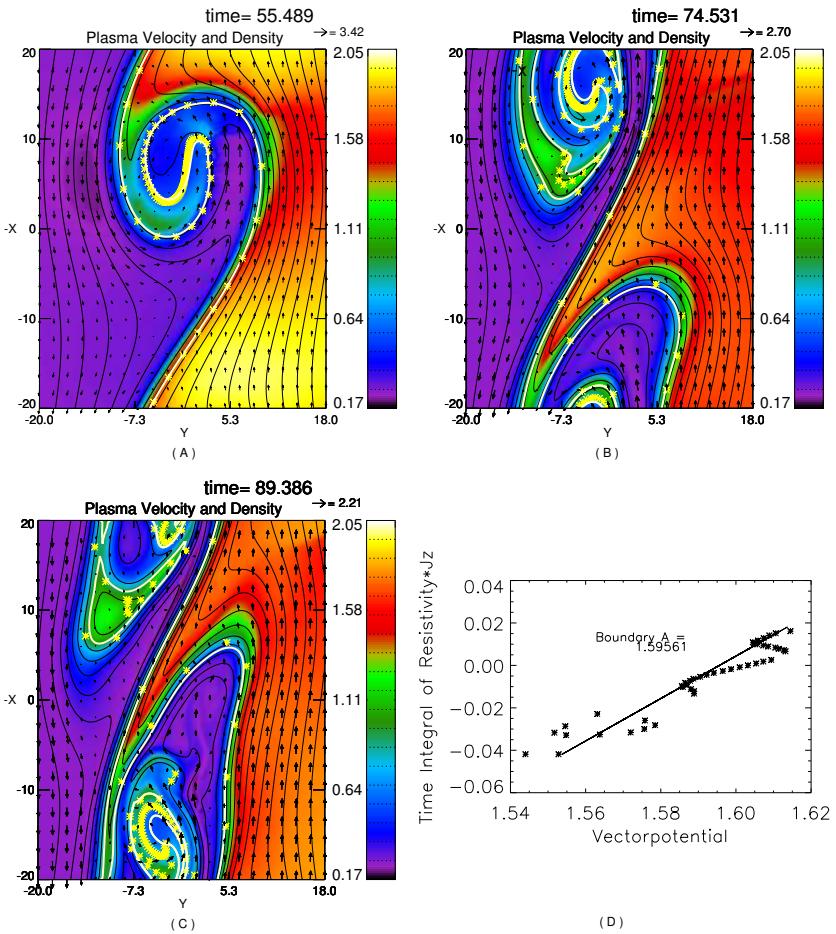
with  $\alpha_i = n_i / (n_1 + n_2)$

**Onset condition:**

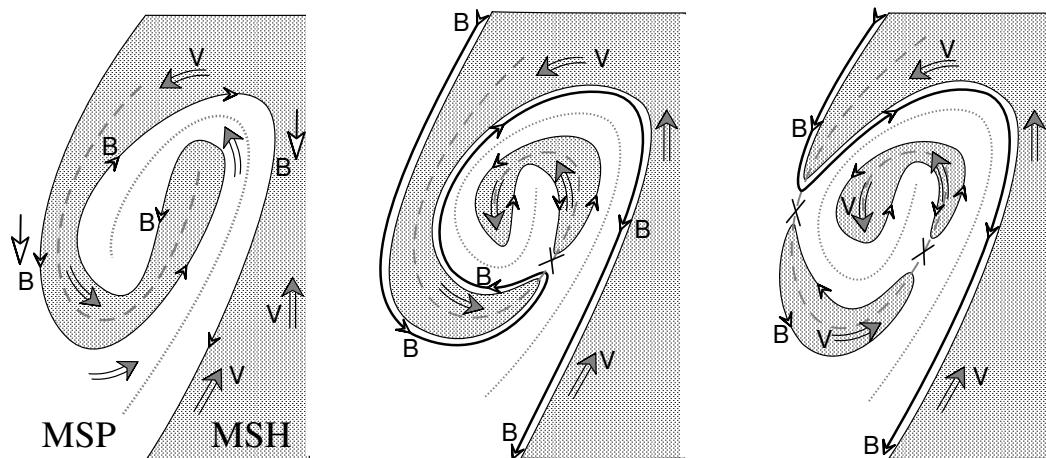
$$[(\mathbf{V}_1 - \mathbf{V}_2) \cdot \mathbf{k}]^2 > \frac{n_1 + n_2}{4\pi m_0 n_1 n_2} [(\mathbf{B}_1 \cdot \mathbf{k})^2 - (\mathbf{B}_2 \cdot \mathbf{k})^2]$$



# Magnetic Reconnection in KH Vortices



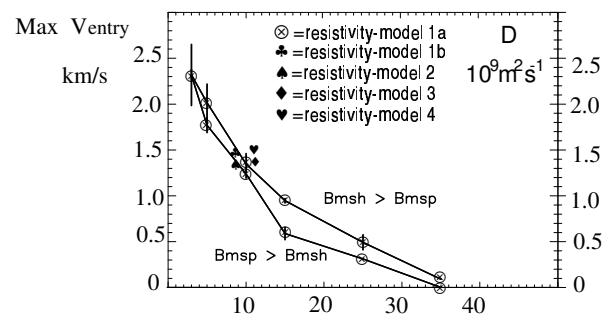
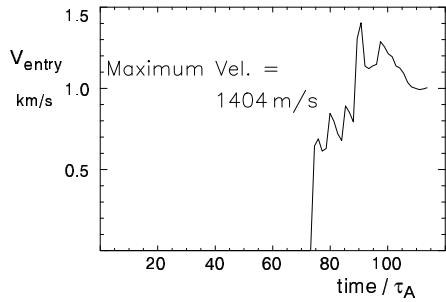
Nykyri and Otto (2001)



# Plasma Transport

Evolution of the magnetic boundary:

$$\frac{dA_z}{dt} \Big|_{Fluid element} = \frac{\partial A_z}{\partial t} + v \cdot \nabla A_z = -\eta J_z \Big|_{Fluid element}$$



## Simulation cases:

Case	$\phi$	$B_{msp}$	$B_{msh}$
1	3	16	24
2	3	24	16
3	5	16	24
4	5	24	16
5	10	16	24
6	10	24	16
7	15	16	24
8	15	24	16
9	25	16	24
10	25	24	16
11	35	16	24
12	35	24	16

$$v_{entry} = \frac{\Delta M}{\Delta t} \frac{1}{\rho_{msh} L_x}$$

## **Summary of Properties:**

- Mass transport from magnetosheath to magnetosphere
- Entry velocities of 0.5 to 2 km/s
  - => 45min to 3 hours to fill the plasma sheet ( $1 \text{ cm}^{-3}$ )
- Limits:
  - Stabilization of the KH (nonlinear) by large magnetic shear  $|\phi| < 25^\circ$
  - $\phi \rightarrow 0$  no identification of magnetic boundary

**KH is an ideal mode:**

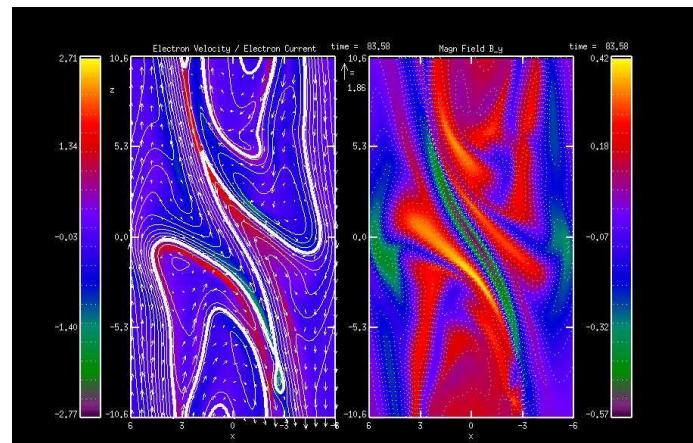
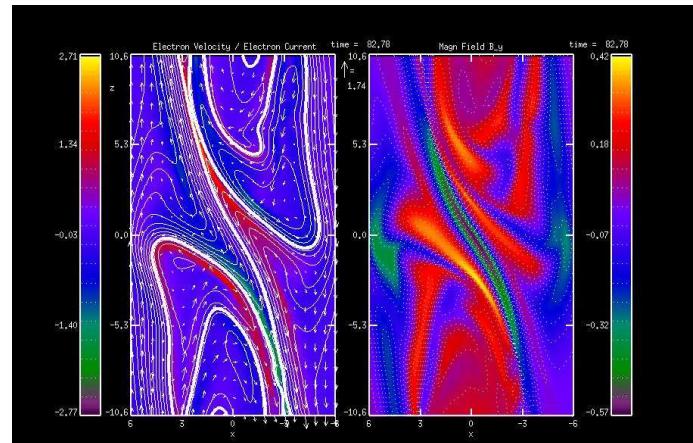
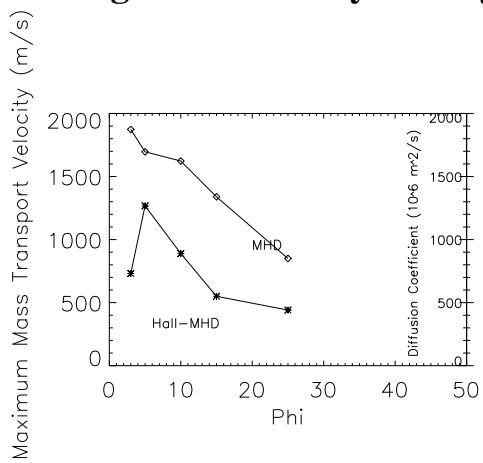
$$\Delta t \sim L_{KH} \text{ (for } L_{KH} \gg L_{grad})$$

$$\Delta M \sim L_{KH}^2 \Rightarrow v_{entry} \text{ independent of wavelength to lowest order!}$$

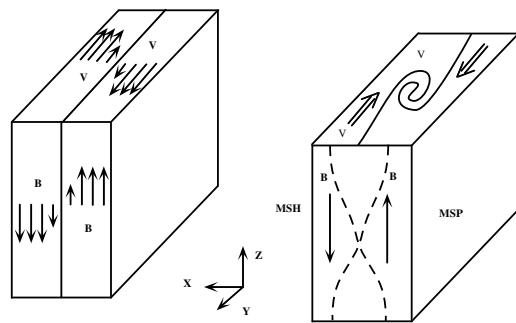
# KH Mode and Hall Dynamics

**Initial configuration:**

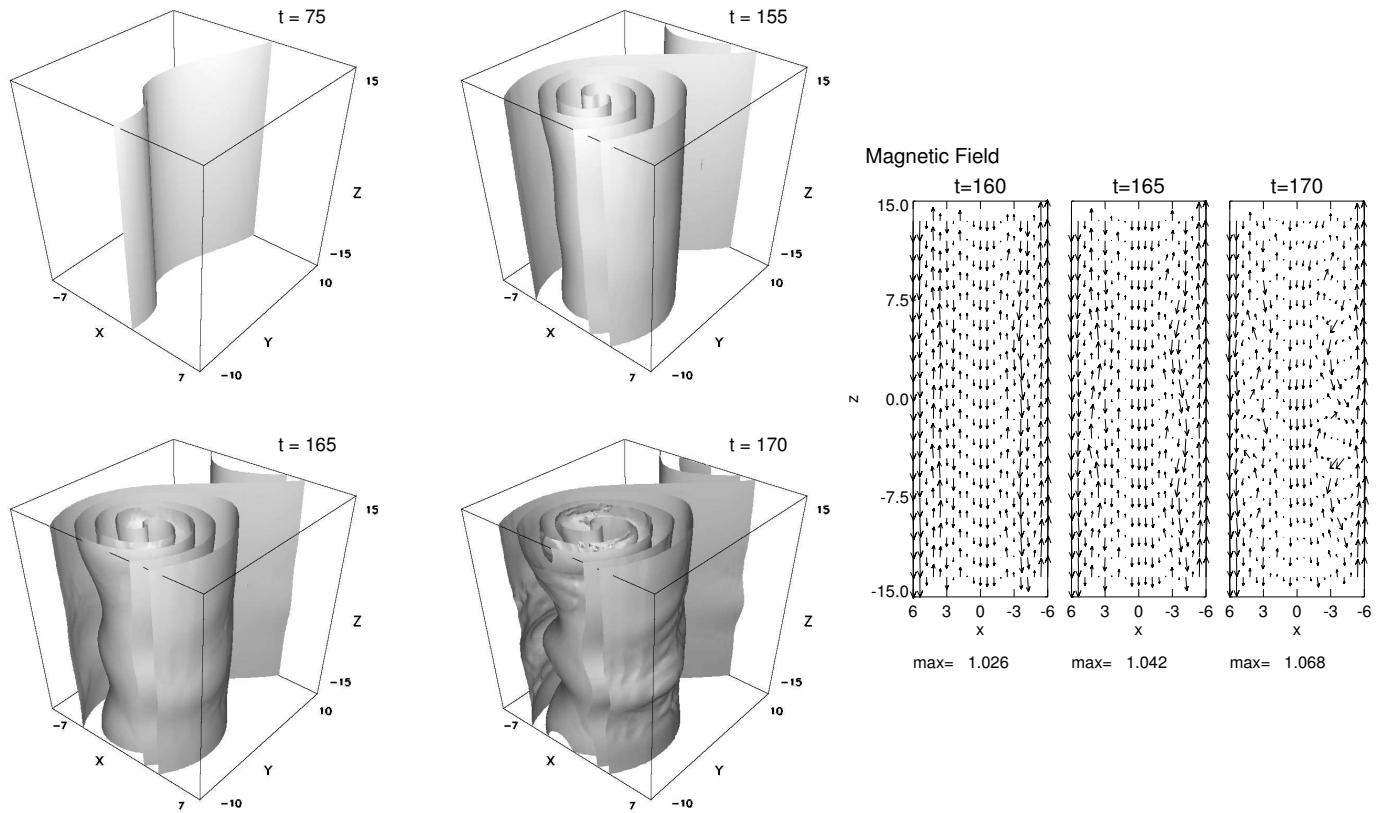
- **Density Asymmetry**
- **Magnetic Field Asymmetry**



## Example of three-dimensional KH and reconnection



## Evolution of the $B=0$ surface:



## **Summary and Conclusions**

### **Importance of Hall physics for magnetic reconnection**

- Enhanced reconnection rate
- Influence on reconnection geometry but
- Geometrical properties (boundary conditions, shear flow, asymmetries) equally important
- Hall ion fountain reconnection

### **Kelvin-Helmholtz mode (ideal instability):**

- Amplification and distortion of the magnetic field in KH Vortices
- Plasma entry through
  - Diffusion over thin boundaries (Fujimoto and Terasawa, 1994,1995)
  - Reconnection in KH vortices
- Plasma transport sufficient to explain the formation of a cold dense plasma sheet for northward IMF
- Preference for strongly northward IMF ( $\pm 50^\circ$ ) and/or small  $V_{A,MSP}$
- Mass transport not strongly dependent on Hall physics