

# Exchange and polarization effects on elastic electron-atom/ion scattering

*(aka. a small part of a bigger project)*

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\*Supported by the US National Science Foundation

## **The objectives:**

- understand the physics of weakly-ionized gases
- develop tools for optimizing the efficiency of Plasma Display Panels (PDPs)

## **Tools used:**

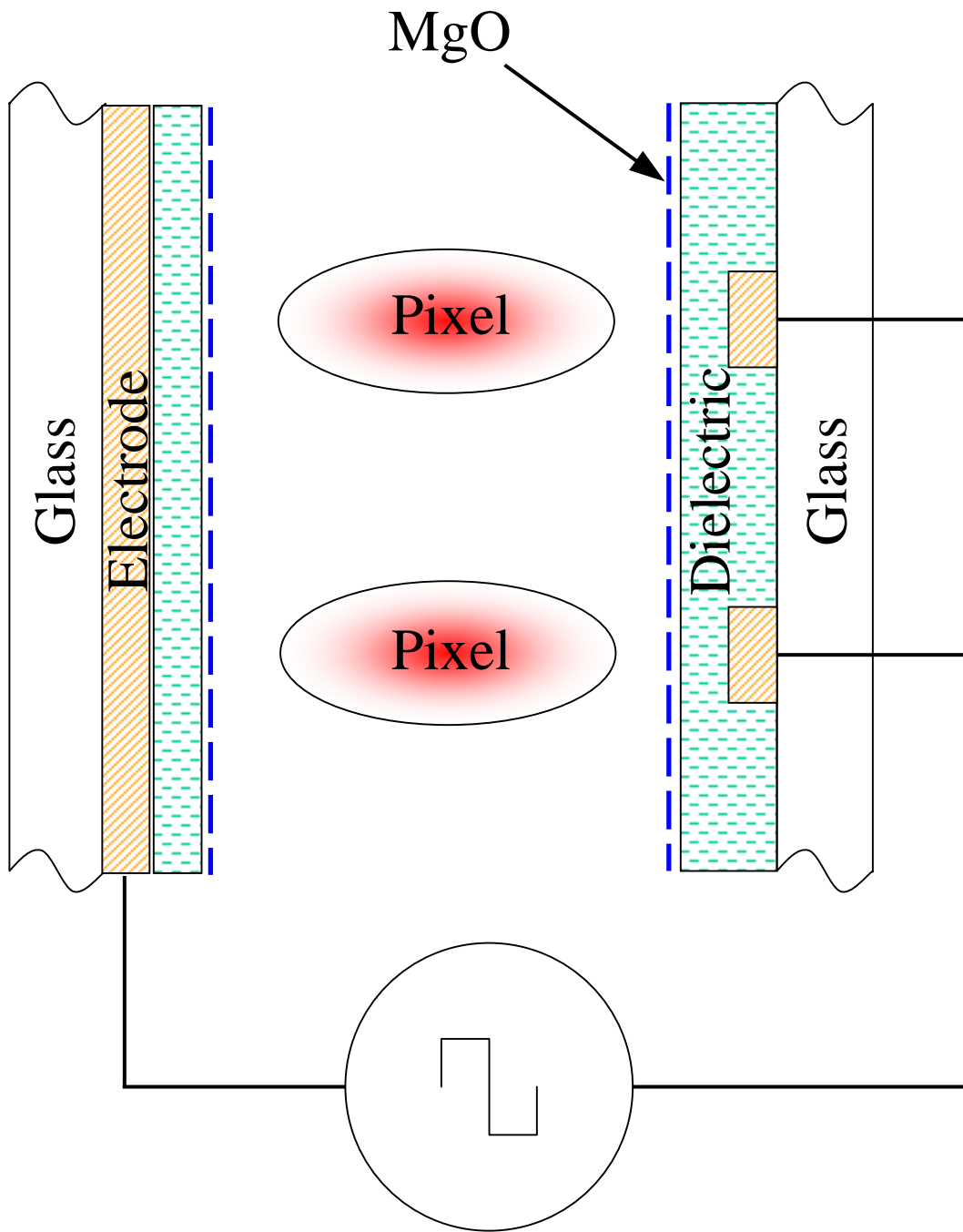
- A combination of experimental and theoretical approaches

## **Needed for modeling and simulations :**

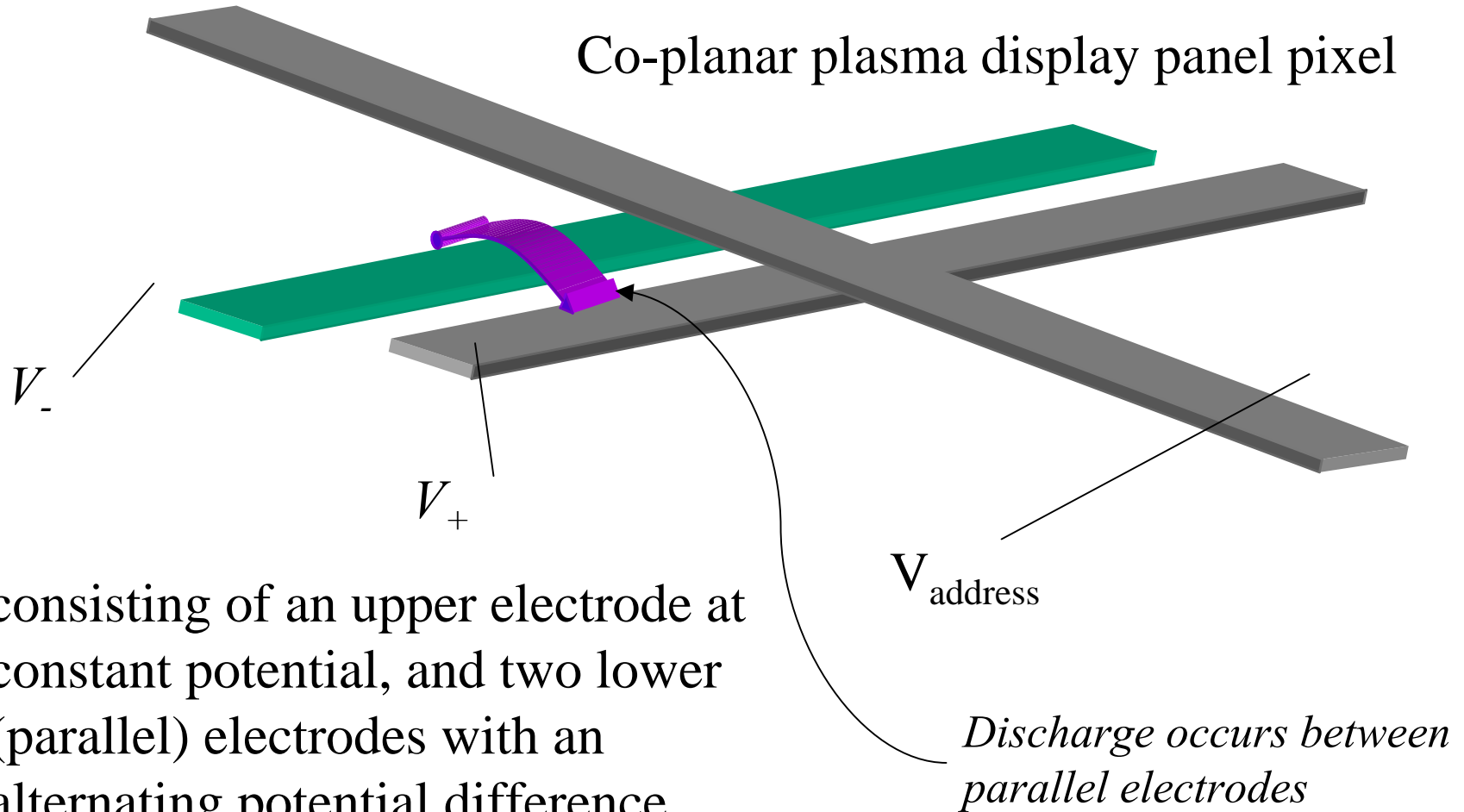
- accurate atomic and molecular data ( $\sigma$ ,  $d\sigma/d\Omega$ ,  $\tau$ , rates)
- accurate beam-surface interaction characteristics data ( $\gamma$ )
- accurate statistical description of the collisions in a non-thermal equilibrium environment ( $f(v)$ ).



EPI prototype color AC PDP - 30 inch (diagonal)



## Co-planar plasma display panel pixel

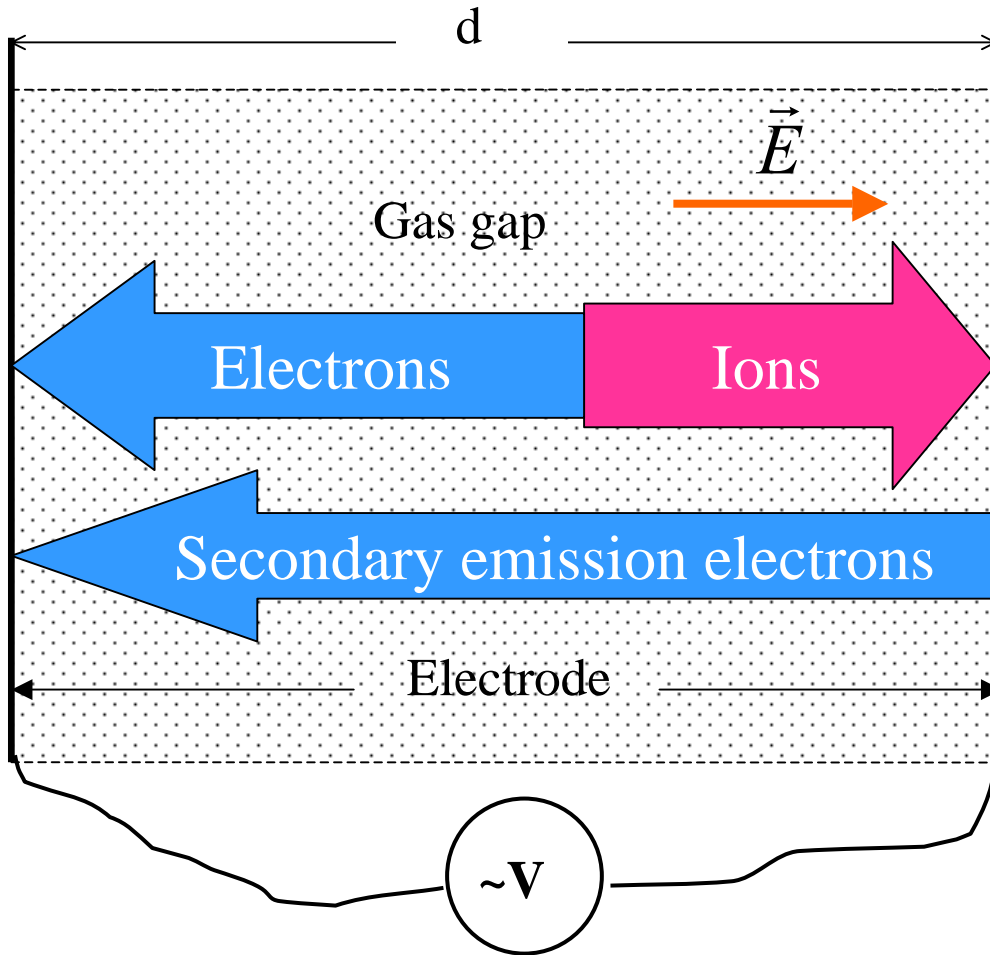


# Setup of our general problem

plasma physics is like astrophysics:

you need to know

- hydrodynamics
- statistical physics
- atomic physics
- molecular physics
- beam-surface interactions



### Boundary conditions

- Densities of all species = 0 at the boundaries
- For secondary electrons,  $J_e = -\sum_i \gamma_i J_i$

### Equations

$$\frac{\partial n_e}{\partial t} + \frac{\partial J_e}{\partial x} = S_e$$

$$\frac{\partial n_i}{\partial t} + \frac{\partial J_i}{\partial x} = S_i$$

$$\frac{\partial n_n}{\partial t} + \frac{\partial J_n}{\partial x} = S_n$$

$$J_e = n_e \mu_e E - \frac{\partial (D_e n_e)}{\partial x}$$

$$J_i = n_i \mu_i E - \frac{\partial (D_i n_i)}{\partial x}$$

$$J_n = - \frac{\partial (D_n n_n)}{\partial x}$$

$$\frac{dE}{dx} = \frac{e}{\epsilon_0} (n_e - n_p)$$

11 equation for species coupled with 1 for electric field

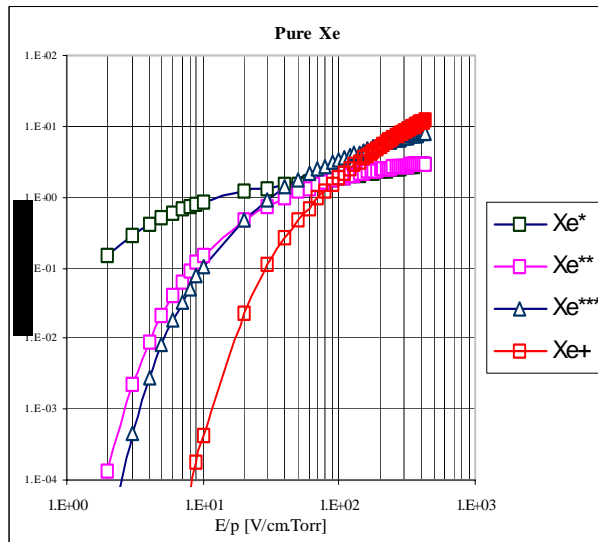
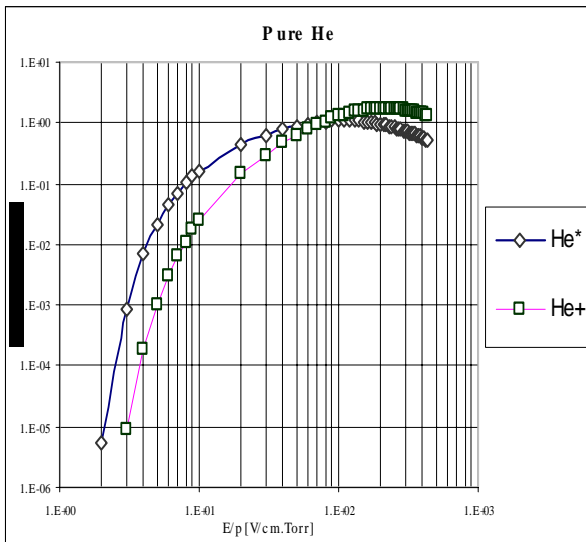
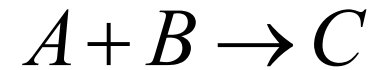
# Information needed and used as input to the solution

- **Electron collision rates**

found by solving the Boltzmann equation in stationary electric field for two-term expansion

$$S = \alpha(E/p)J_e$$

- **Chemical reaction rates**



$$rate = k[A][B]$$

$$S_A = -rate,$$

$$S_B = -rate,$$

$$S_C = rate$$

- **Secondary electron coefficients**  $\gamma_i, \gamma_n$ .



# Species and reactions for He Xe mixture

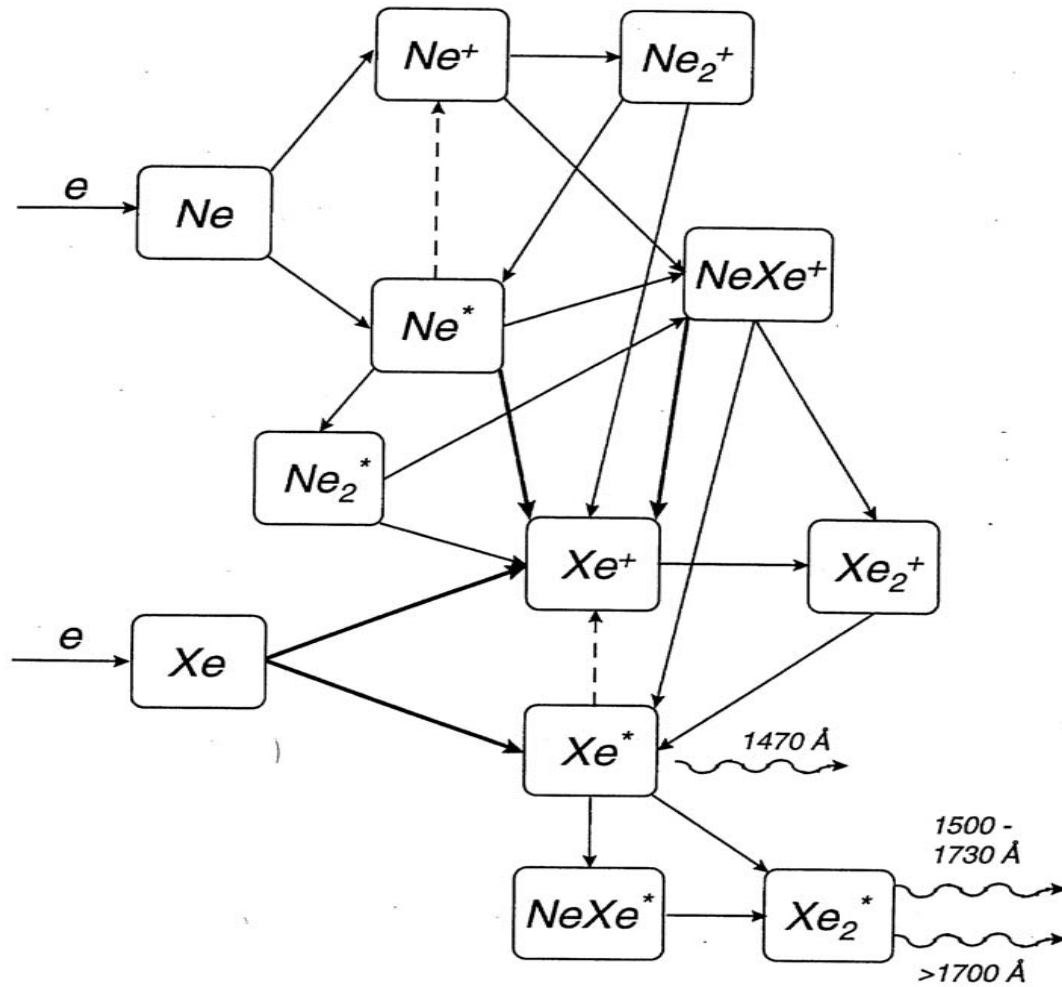
## Electrons,

**Ions:** He<sup>+</sup>, Xe<sup>+</sup>, He<sub>2</sub><sup>+</sup>, Xe<sub>2</sub><sup>+</sup>

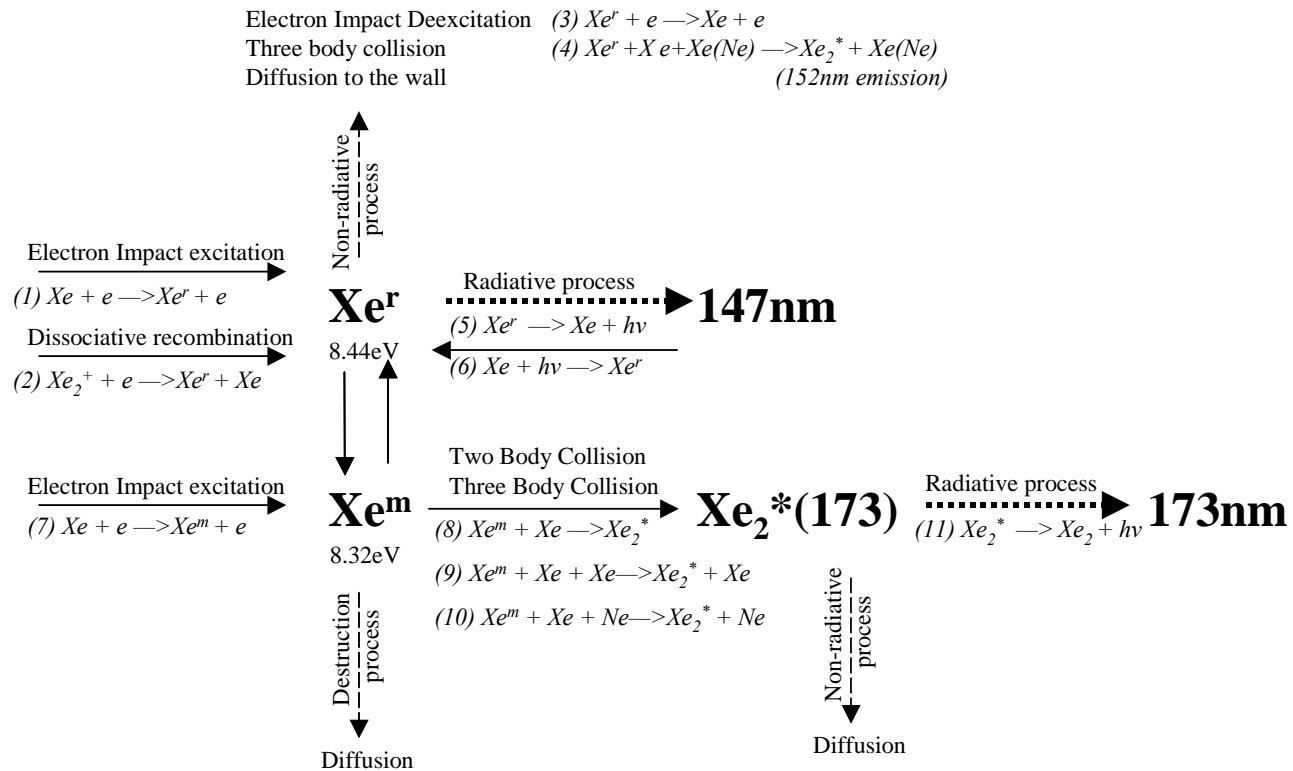
**Neutrals:** He\*, Xe\*, Xe\*\*, Xe\*\*\*, He<sub>2</sub>, Xe<sub>2</sub>

Process	Process	Process
<b>Electron impact excitation</b>	<b>Electron impact ionisation</b>	<b>Charge exchange in</b>
<i>of atoms in the ground state</i>	<i>of atoms in excited states</i>	two-body heavy particle collisions
e + He → He* + e	e + He* → He + 2e	<i>of atoms and molecules</i>
e + Xe → Xe* + e	e + Xe* → Xe + 2e	He <sub>2</sub> <sup>+</sup> + Xe → Xe <sup>+</sup> + 2He
e + Xe → Xe** + e	e + Xe** → Xe + 2e	
e + Xe → Xe*** + e	e + Xe*** → Xe + 2e	Excited molecule formation in
		three-body heavy particle collisions
<b>Electron impact de-excitation</b>	<b>Electron impact ionisation</b>	<i>of atoms</i>
<i>of excited atoms to the ground state</i>	<i>of excited molecules</i>	He* + 2He → He <sub>2</sub> <sup>+</sup> + He
e + He* → He + e	e + He <sub>2</sub> <sup>+</sup> → He <sub>2</sub> + 2e	He* + He + Xe → He <sub>2</sub> <sup>+</sup> + Xe
e + Xe* → Xe + e	e + Xe <sub>2</sub> <sup>+</sup> → Xe <sub>2</sub> + 2e	Xe* + 2Xe → Xe <sub>2</sub> <sup>+</sup> + Xe
e + Xe** → Xe + e		Xe** + 2Xe → Xe <sub>2</sub> <sup>+</sup> + Xe
e + Xe*** → Xe + e	<b>Electron impact recombination</b>	Xe*** + 2Xe → Xe <sub>2</sub> <sup>+</sup> + Xe
	<i>of molecules</i>	Xe* + He + Xe → Xe <sub>2</sub> <sup>+</sup> + He
<b>Electron impact excitation</b>	e + He <sub>2</sub> <sup>+</sup> → He* + He	Xe** + He + Xe → Xe <sub>2</sub> <sup>+</sup> + He
<i>of atoms in excited states to a higher excited state</i>	e + Xe <sub>2</sub> <sup>+</sup> → Xe* + Xe	Xe*** + He + Xe → Xe <sub>2</sub> <sup>+</sup> + He
e + Xe* → Xe** + e	e + Xe <sub>2</sub> <sup>+</sup> → Xe** + Xe	
e + Xe* → Xe*** + e	e + Xe <sub>2</sub> <sup>+</sup> → Xe*** + Xe	Ionised molecule formation in
e + Xe** → Xe*** + e		three-body heavy particle collisions
	Ions formation in	<i>of atoms and ions</i>
<b>Electron impact de-excitation</b>	two-body heavy particle collisions	He <sup>+</sup> + 2He → He <sub>2</sub> <sup>+</sup> + He
<i>of atoms in excited states to a lower excited state</i>	<i>of atoms</i>	Xe <sup>+</sup> + 2Xe → Xe <sub>2</sub> <sup>+</sup> + Xe
e + Xe** → Xe* + e	2He* → He <sup>+</sup> + He + e	Xe <sup>+</sup> + He + Xe → Xe <sub>2</sub> <sup>+</sup> + He
e + Xe*** → Xe* + e	2Xe* → Xe <sup>+</sup> + Xe + e	
e + Xe*** → Xe** + e	2Xe** → Xe <sup>+</sup> + Xe + e	UV Radiation
	2Xe*** → Xe <sup>+</sup> + Xe + e	Xe <sub>2</sub> <sup>+</sup> → 2Xe + hv
<b>Electron impact ionisation</b>	He* + Xe → Xe <sup>+</sup> + He + e	Xe <sup>+</sup> → Xe + hv
<i>of atoms in the ground state</i>	<i>of atoms and molecules</i>	Xe** → Xe + hv
e + He → He <sup>+</sup> + 2e	He <sub>2</sub> <sup>+</sup> + Xe → Xe <sup>+</sup> + 2He + e	
e + Xe → Xe <sup>+</sup> + 2e		

# Ne-Xe reaction scheme



# Xe - vuv emission process



$$\frac{\partial f}{\partial t} + \frac{\partial}{\partial x_i} (f v_i) + \frac{\partial}{\partial v_i} (f \dot{v}_i) = \left( \frac{df}{dt} \right)_{\text{collisional}}$$

Requires **accurate knowledge** of many processes in

- atomic physics
- molecular physics
- beam-surface interactions

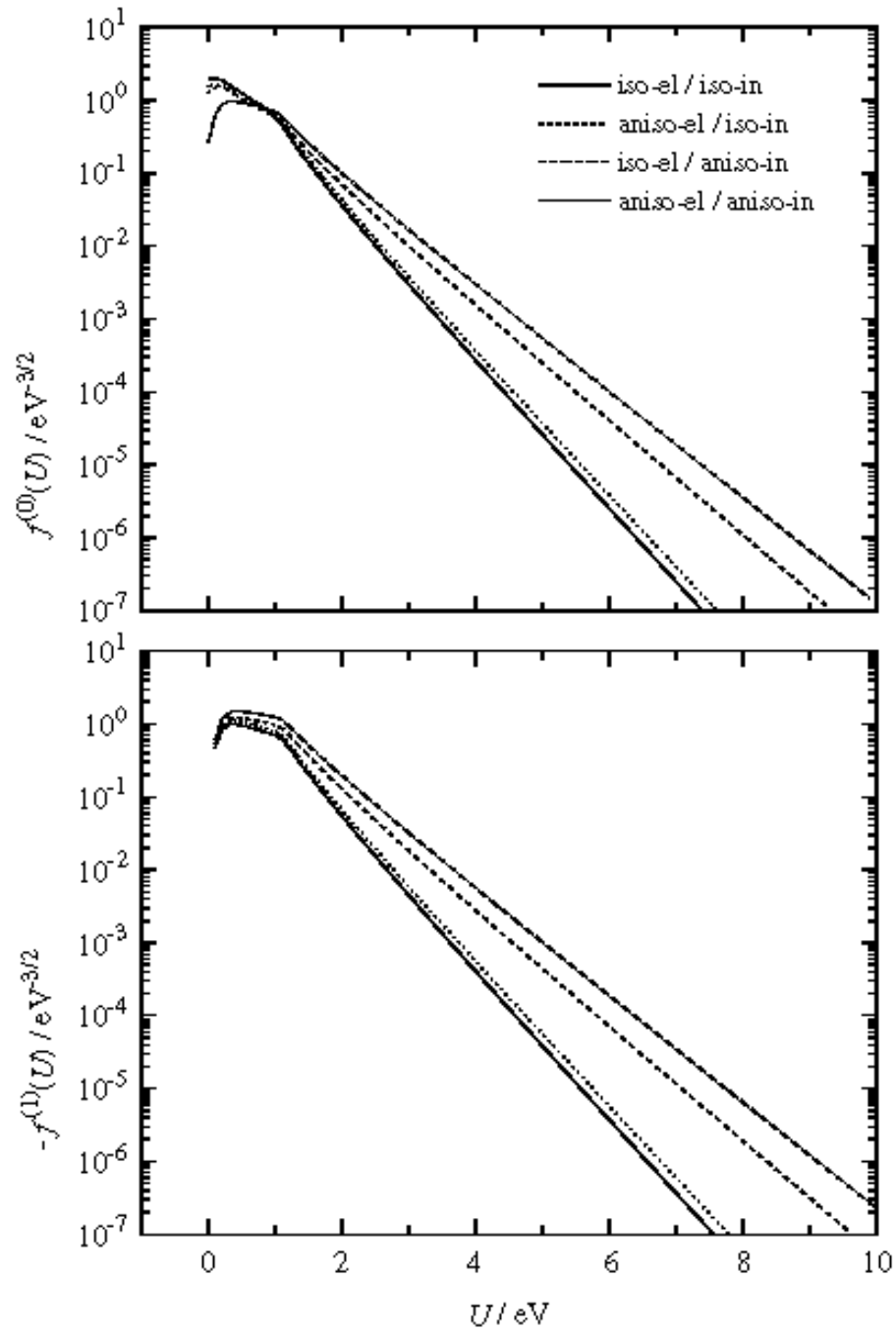
# Elastic cross sections

are among the most important

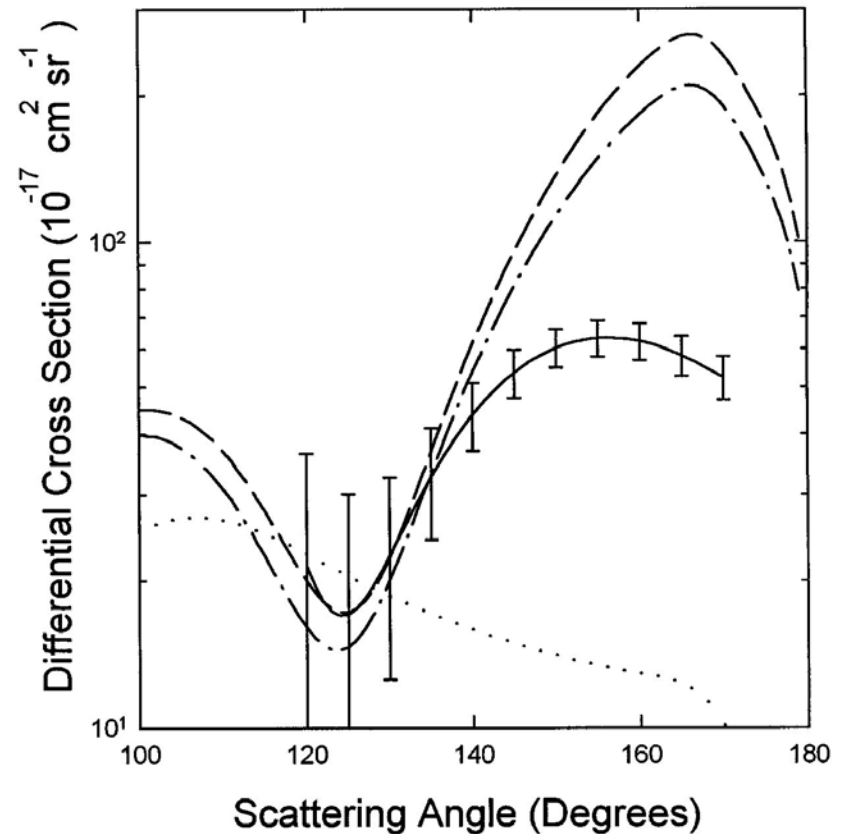
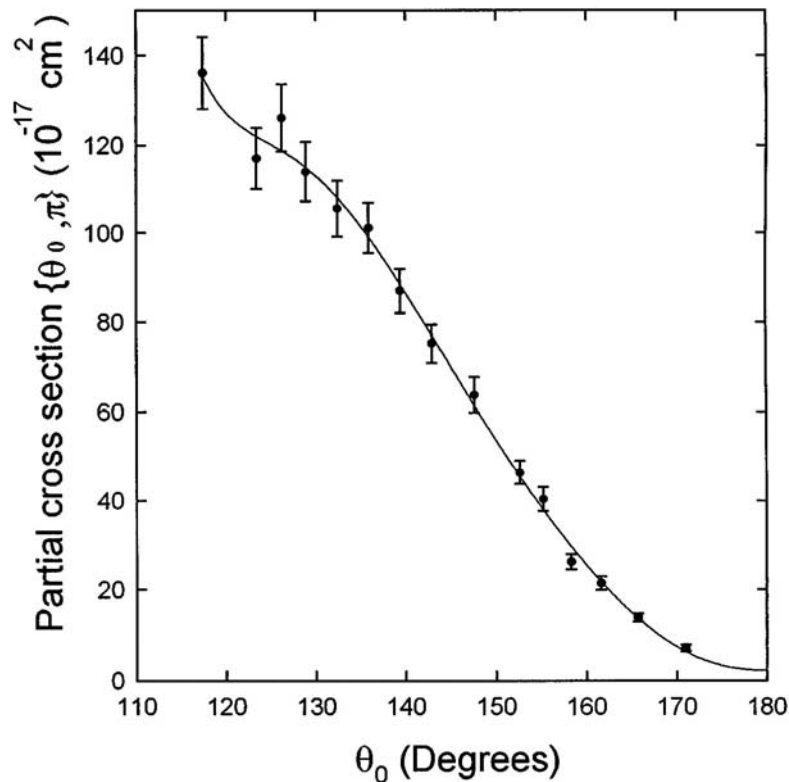
Work by the Greifswald Plasma Group has shown [e.g., Leyh et al. *Comput. Phys. Rev. Commun.* **113**, 33 (1998)] that:

- **anisotropic vs. isotropic** elastic (and inelastic) collisions can make a big difference in the electron velocity distribution of electrons in weakly ionized plasmas

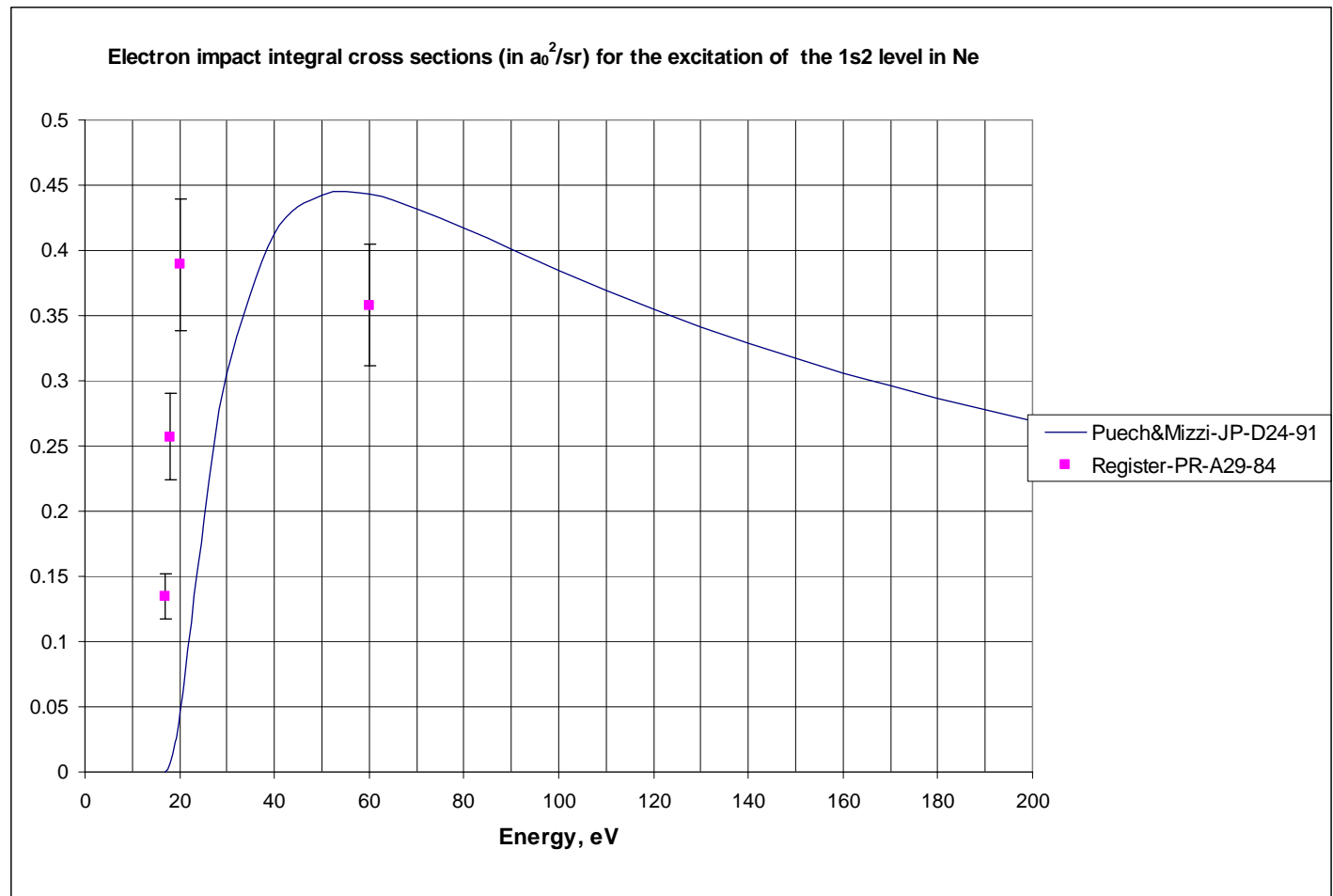
# simulations



Greenwood et al., [Phys. Rev. Lett. **75**, 1062 (1995)] claimed **backward** elastic scattering of **3.3 eV electrons on Ar<sup>+</sup>** measurements in disagreement with theory



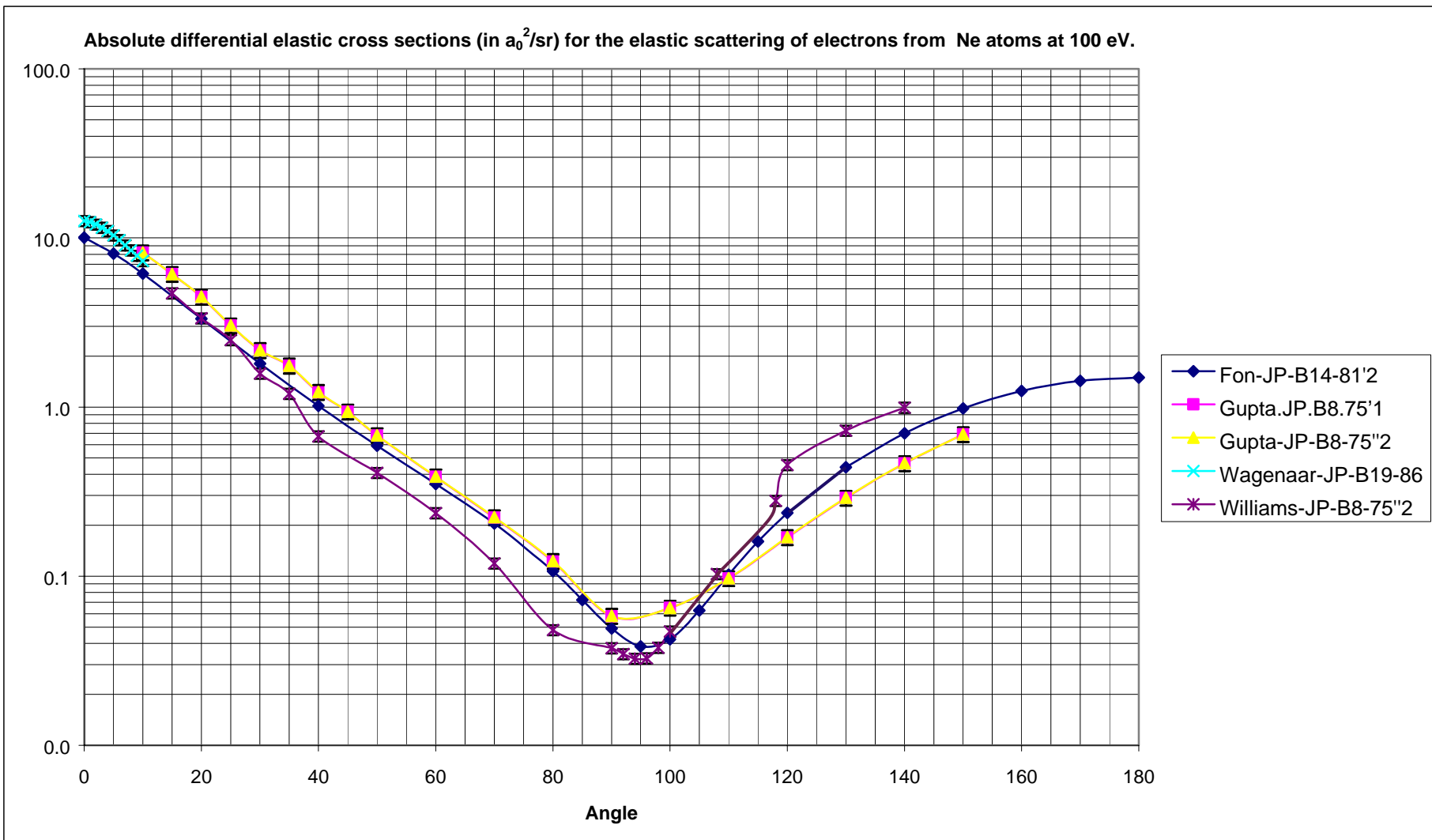
## Problems – 2



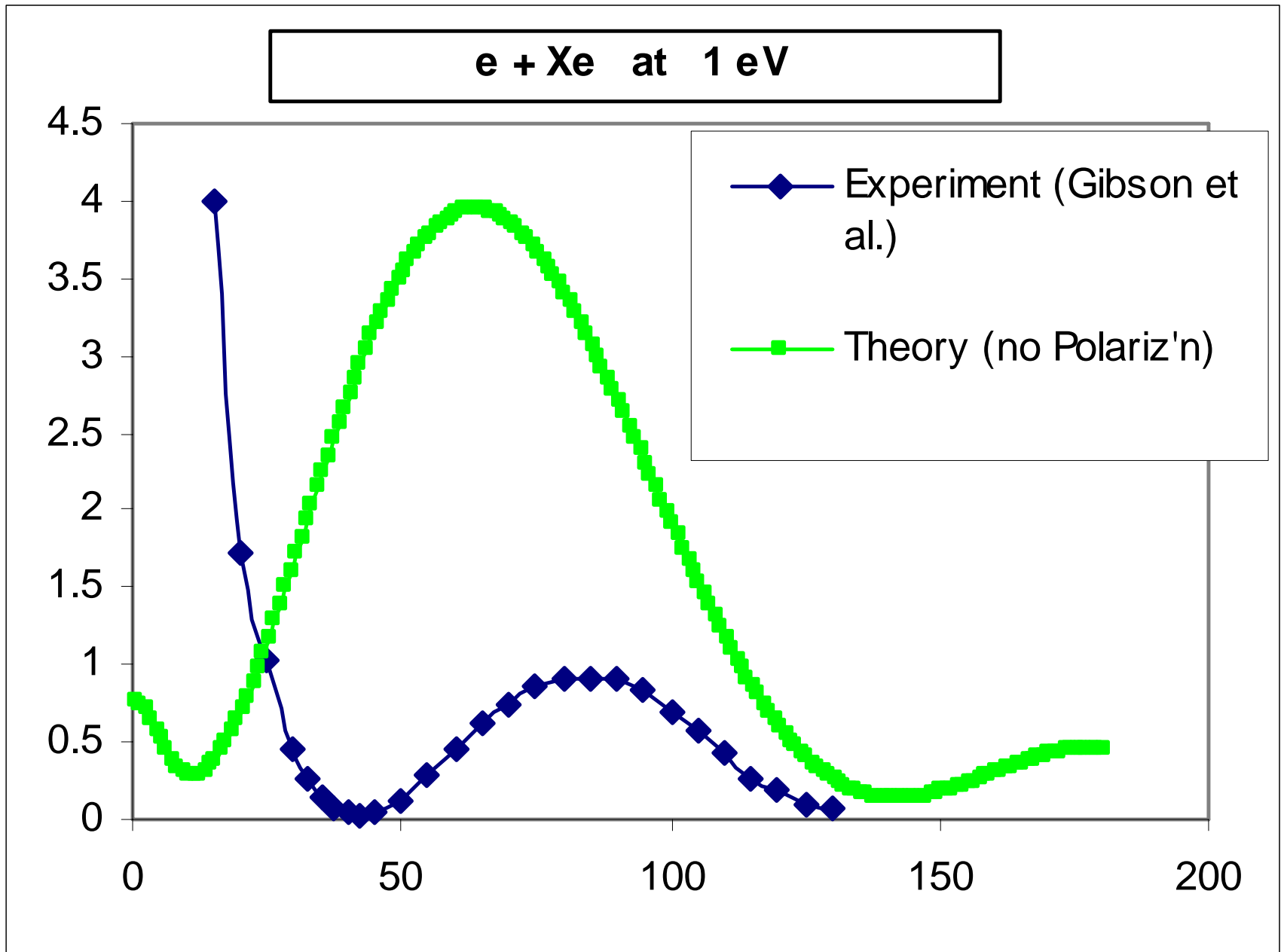
*Demonstrates magnitude of disparity possible among similar data sets*



# Problems - 3



*A chart demonstrating a typical comparison of similar data sets from different sources*



# Fits of differential cross sections

- The experimental cross sections were fitted to the form:

$$\frac{d\sigma}{d\Omega}(\theta) = A \exp(-\theta / B) + C \exp[-(\pi - \theta)^4 / D] \\ + \sum_{l=0}^{4-10} a_l P_l(\cos \theta)$$

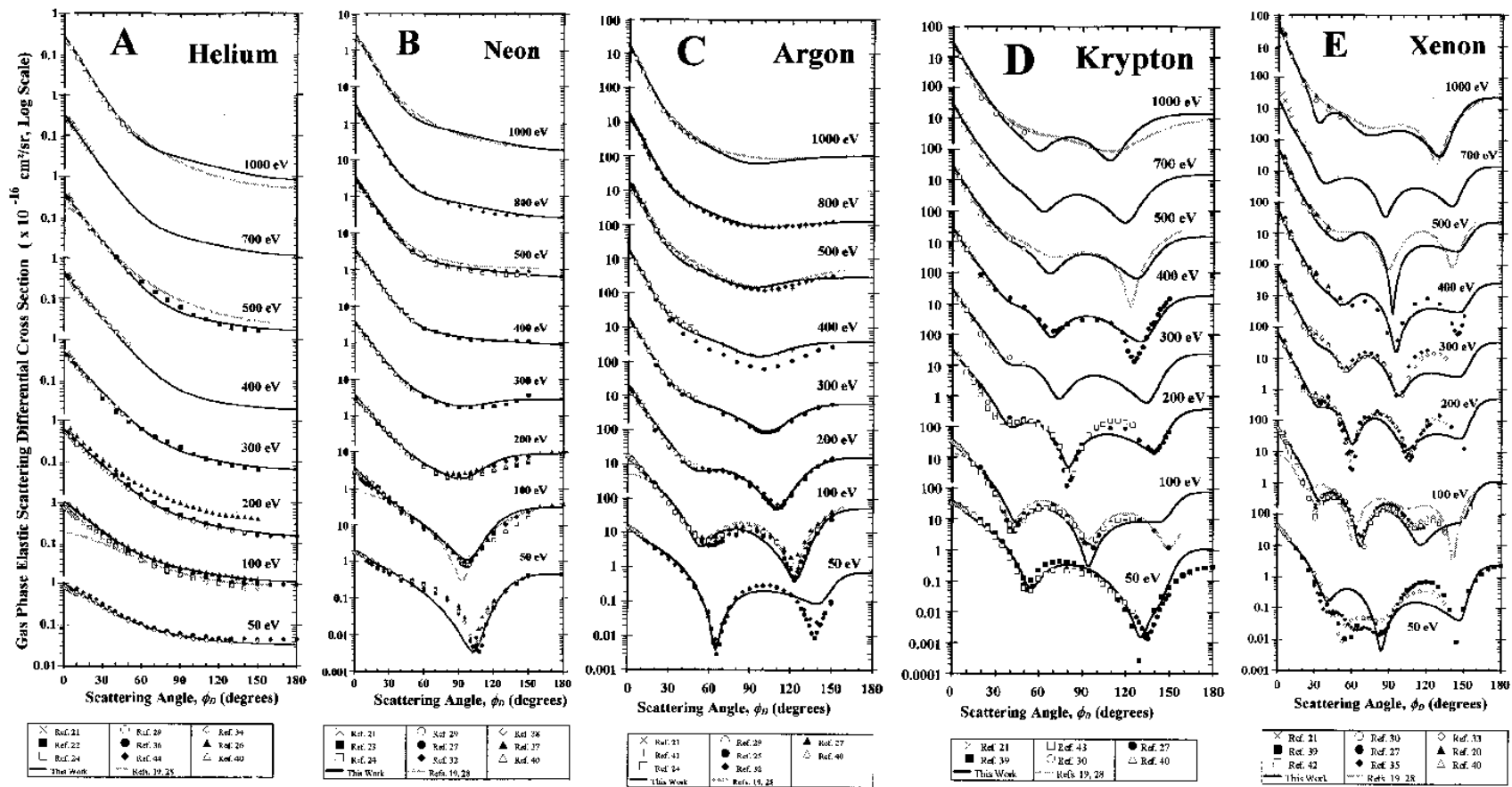


Figure 5. Logarithmic plots of experimental elastic differential cross sections for the rare gases. Gas phase measurements (individual points) are compared with the present description (heavy solid lines) and previous theoretical treatments (fine lines). References are shown in the legend below each plot. The logarithmic format allows the large-angle portions of the distributions to be seen more clearly by exaggerating the low intensities.

**Our Objective:** A **simple** and **dependable** calculation for **elastic scattering cross sections** of electrons from atoms and ions for gaseous discharge studies.

**Reasons:**

Realistic calculations of gas discharges require

- **total**
- **momentum transfer**
- **angle-differential cross sections**

(to account for **anisotropic scattering events**)

The vast majority of such calculations assume isotropic elastic scattering

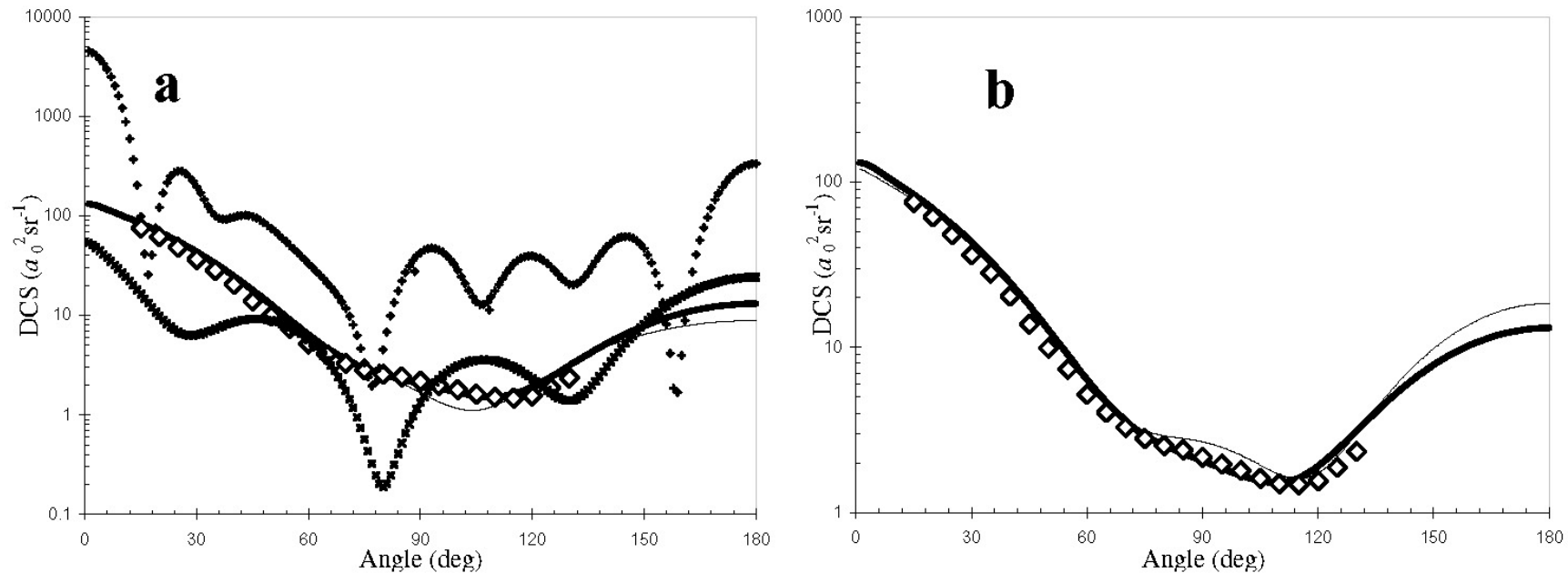


FIG. 2: Studying the effect of different choices of central atomic and exchange potentials on present DCS values for xenon at 10 eV: **a)** Comparison with experimental values while the central atomic potential varies, but Exchange (SC) and polarization (Buckingham-type II) potential remain unchanged:  $\diamond$ , Gibson *et al.* [40]; thick solid line, DS potential [8]; thin solid line, Salvat *et al.* potential [7]; +, HF potential;  $\times$ , Green *et al.* potential [5, 6]. **b)** Comparison with experimental values while the exchange potential varies, but Central atomic (DS) and polarization (Buckingham-type II) potential remain unchanged:  $\diamond$ , Gibson *et al.* [40]; thick solid line, SC exchange potential; thin solid line, FEG exchange potential.

# General theoretical expression

$$\frac{d\sigma}{d\Omega}(\theta) = |f_c(\theta) + f_{nc}(\theta)|^2$$

$$f_c = \frac{-\alpha \exp(2i\sigma_0)}{2\mu v \sin^2(\theta/2)} \exp[-i\alpha \log(\sin^2(\theta/2))]$$

$$\alpha = -q/v \quad \sigma_l = \arg(\Gamma(l+1+i\alpha))$$

$$f_{nc} = \frac{1}{2ik} \sum_{l=0}^{l_{\max}} (2l+1) \exp(2i\sigma_l) (\exp(2i\delta_l) - 1) P_l(\cos \theta)$$

# General theoretical expressions - neutral targets

## Scattering amplitudes

$$f(\theta) = \frac{1}{k} \sum_{l=0}^{\infty} \left[ (l+1) e^{i\delta_l^+(k)} \sin \delta_l^+(k) + l e^{i\delta_l^-(k)} \sin \delta_l^-(k) \right] P_l(\cos \theta)$$

$$g(\theta) = \frac{1}{k} \sum_{l=0}^{\infty} \left[ e^{i\delta_l^-(k)} \sin \delta_l^-(k) - e^{i\delta_l^+(k)} \sin \delta_l^+(k) \right] P_l^1(\cos \theta)$$

## Differential elastic cross section

$$\frac{d\sigma}{d\Omega}(\theta, k) = |f(\theta)|^2 + |g(\theta)|^2$$

## Total elastic cross section

$$\sigma(k) = \frac{4\pi}{k^2} \sum_l \left[ (l+1) \sin^2 \delta_l^+(k) + l \sin^2 \delta_l^-(k) \right]$$



## Momentum transfer cross section

$$\sigma_m(k) = \frac{4\pi}{k^2} \sum_{l=0}^{\infty} \left\{ \begin{array}{l} \frac{(l+1)(l+2)}{2l+3} \sin^2 \left( \delta_l^+(k) - \delta_{l+1}^+(k) \right) \\ + \frac{l(l+1)}{2l+1} \sin^2 \left( \delta_l^-(k) - \delta_{l+1}^-(k) \right) \\ + \frac{(l+1)}{(2l+1)(2l+3)} \sin^2 \left( \delta_l^+(k) - \delta_{l+1}^-(k) \right) \end{array} \right\}$$

## Sherman function

$$S(\theta, k) = i \frac{f(\theta)g^*(\theta) - f^*(\theta)g(\theta)}{|f(\theta)|^2 + |g(\theta)|^2}$$

# Approximations

## Atomic potential

## Central Dirac-Slater potential

## Electron exchange

**Local approximation of Furness and McCarthy's** [JPB 6, 2280 (1973)]

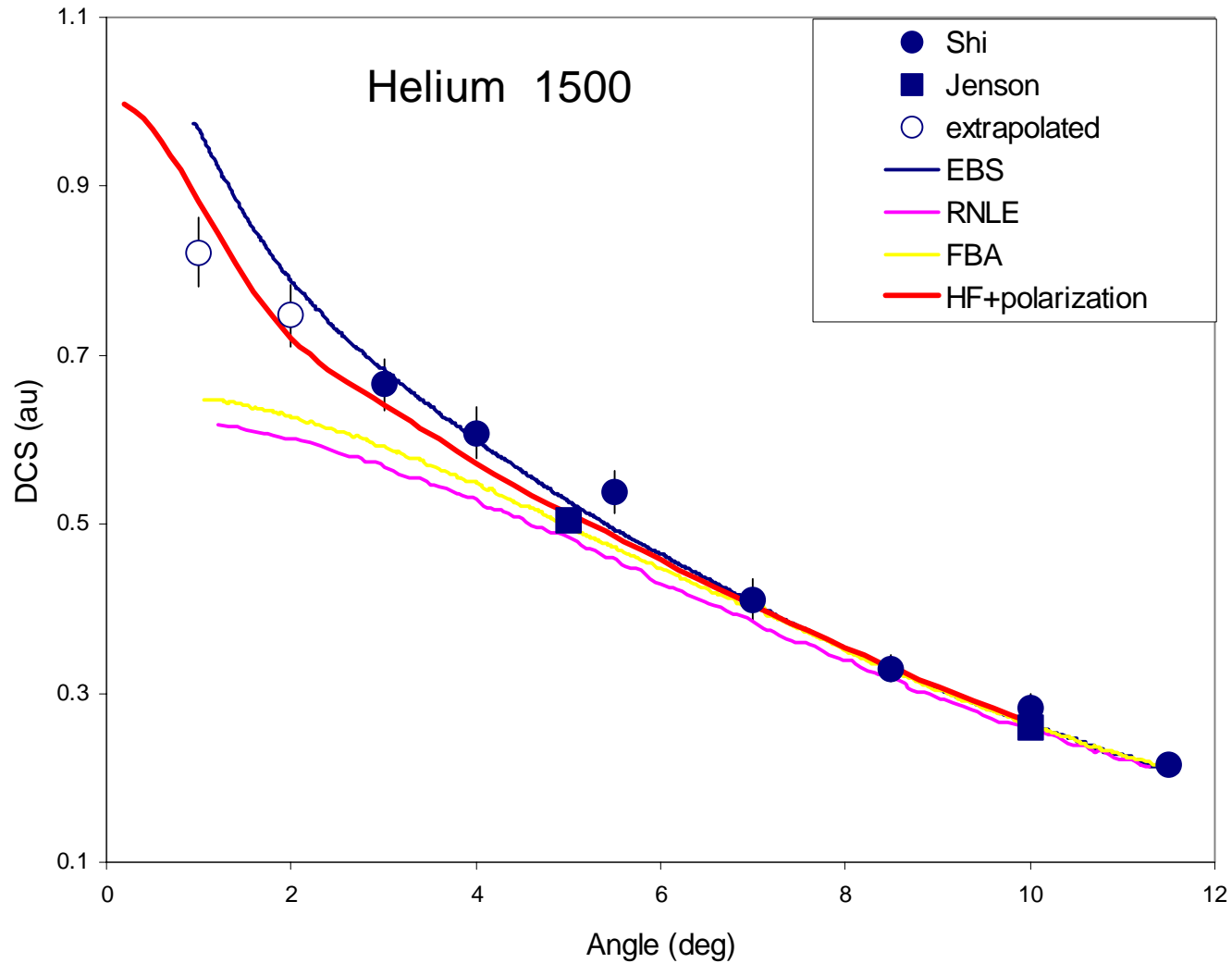
$$V_{\text{exch}}(r) = \frac{1}{2} \left[ |E - V_s(r)| - \left( |E - V_s(r)|^2 + 4\pi\rho(r) \right)^{1/2} \right]$$

$$V_s(r) = -\frac{2}{r} [Z - Y(r)]; \quad Y(r) = \sum_{i=1}^N \left\{ \int_0^r u_i^2(r') dr' + r \int_r^\infty [u_i^2(r') / r'] dr' \right\}.$$

## Core polarization effects

$$V_{\text{pol}}(r) = -\frac{\alpha_d}{2(r^2 + d^2)^2} \quad d = \langle r \rangle_{np} + \frac{1}{3} \ln(E / \text{Ry})$$

# Some examples/results



# Griffin and Pindzola, Phys. Rev A **53**, 1915 (1996)

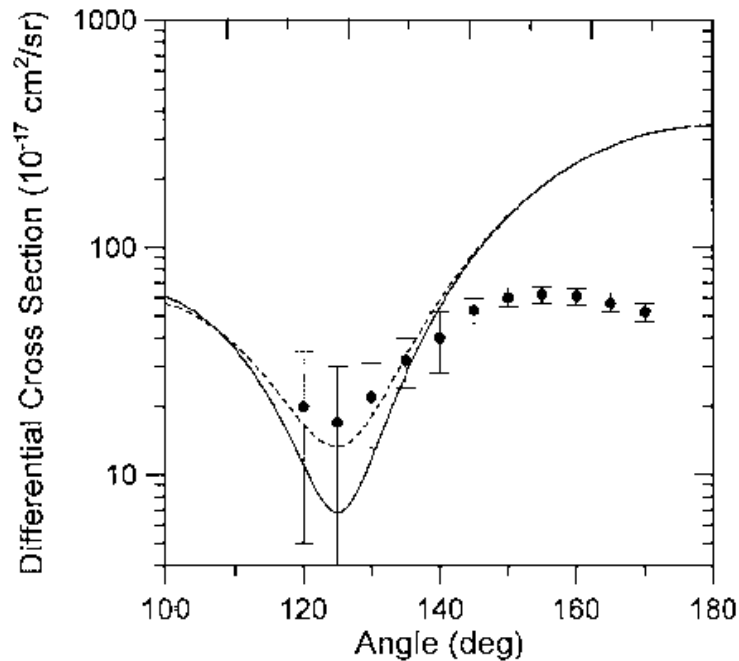


FIG. 1. Angular differential cross section for elastic scattering in  $\text{Ar}^+$  at 3.3 eV. Solid curve: from a 17-state  $R$ -matrix calculation including polarized pseudostates; dashed curve: from a 17-state calculation convoluted over energy and angle using the energy and angular widths given in Ref. [1]; experimental points are from Ref. [1].

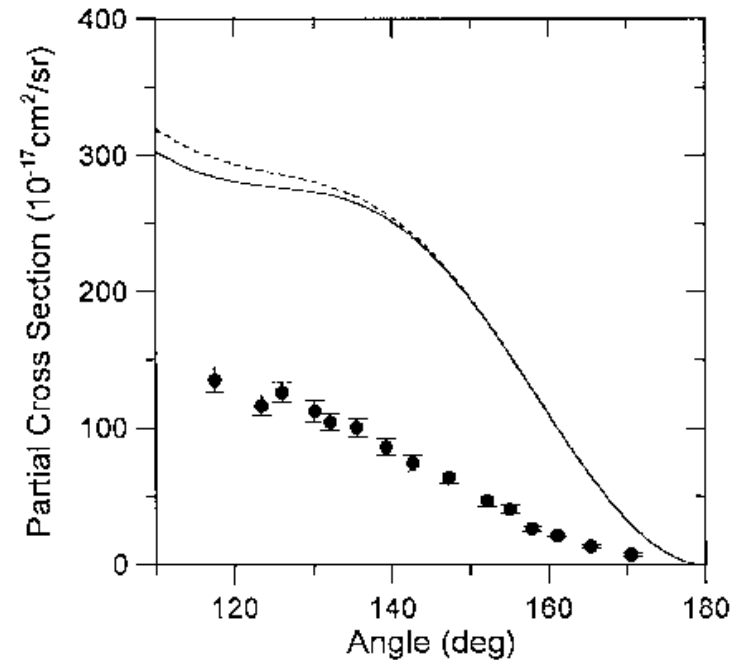
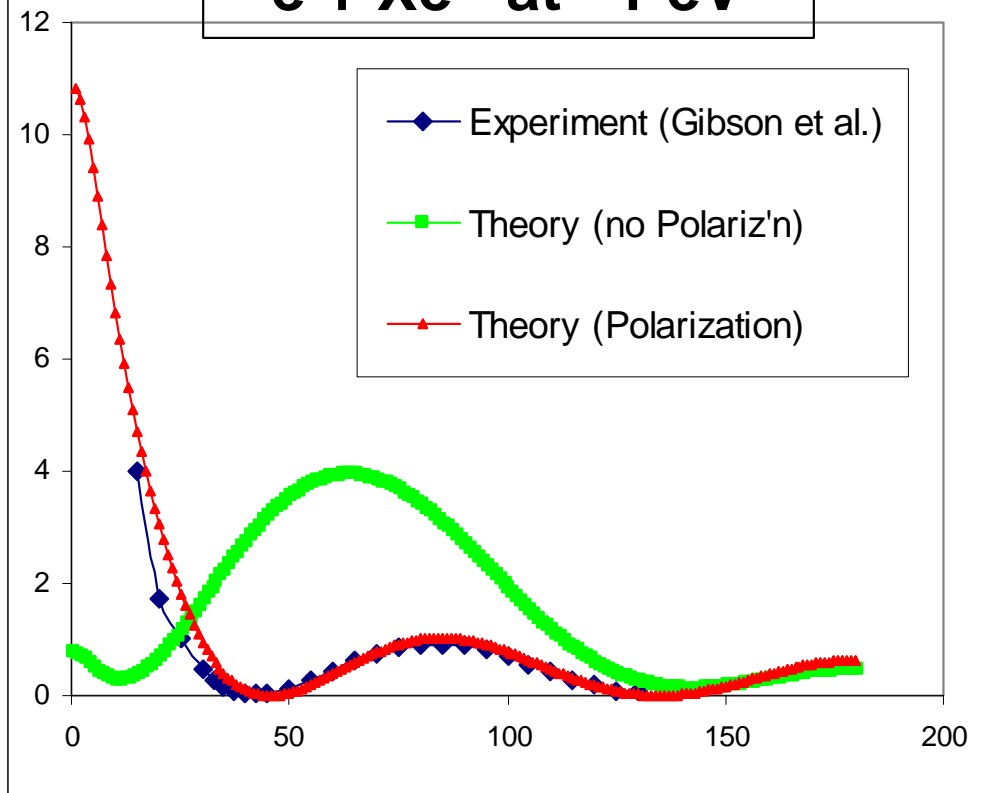
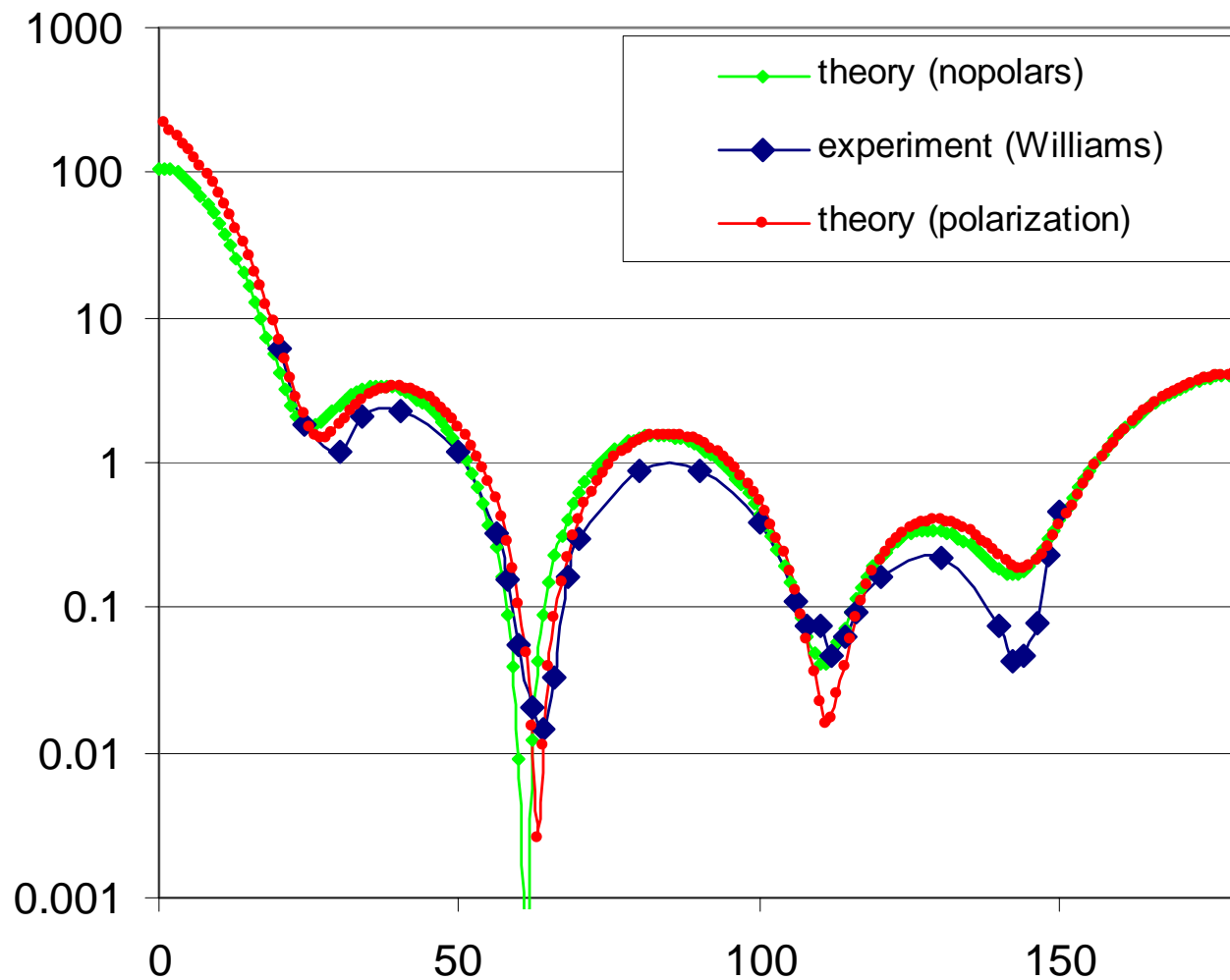


FIG. 2. Partial differential cross section for elastic scattering in  $\text{Ar}^+$  at 3.3 eV from an initial angle through  $180^\circ$ , as a function of the initial angle. Solid curve: from 17-state  $R$ -matrix calculation; dashed curve: from a 17-state calculation convoluted over energy and angle before integration over the angle; experimental points: from Ref. [1].

# e + Xe at 1 eV



e + Xe at 125 eV



# Atoms studied completely

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He Ne Ar Kr Xe

Sr Ba

Be Mg Ca

Cd Zn Hg

# Parametric fitting of results

- Experimental data were fitted by varying the cutoff distance  $d$ .
- After the fit, the trend of the values of  $d$  was examined.
- We found:  
for  $Ne, Ar, Kr, Xe, Ba, Sr, Be, Mg, Ca, Zn, Cd, Hg$

$$d = \langle r \rangle_{np} + \frac{1}{3} \ln(E / Ry)$$

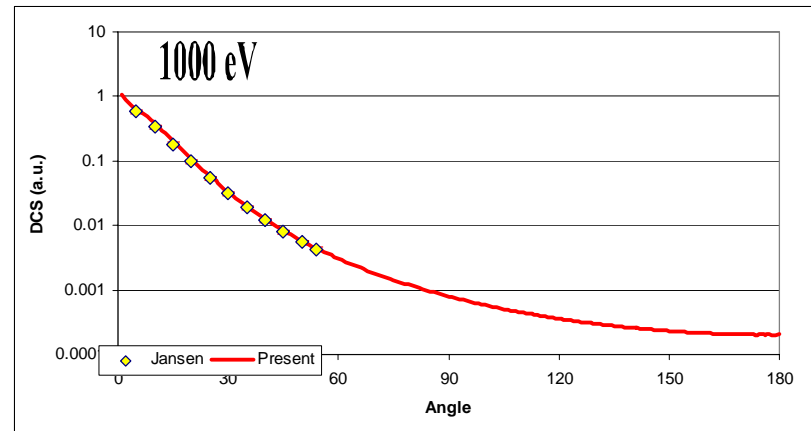
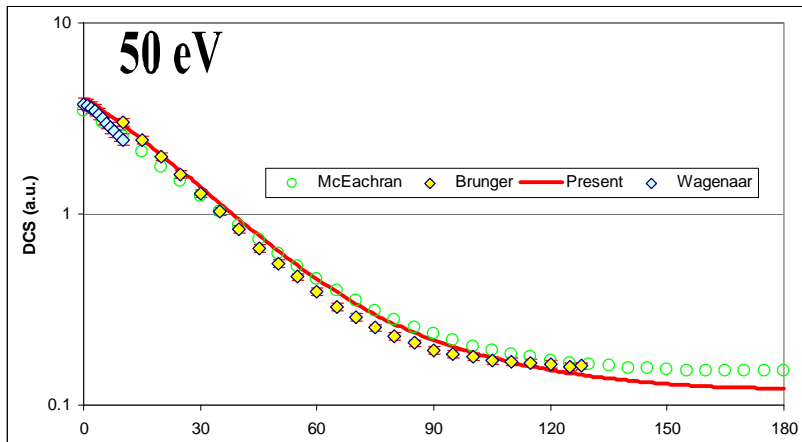
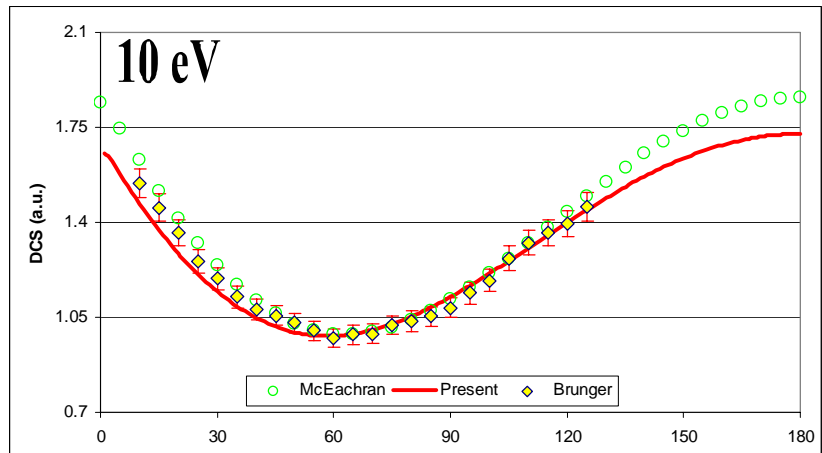
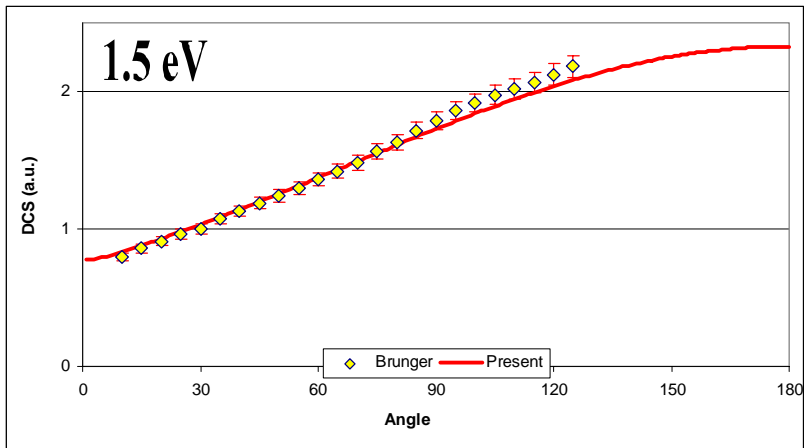
and a linear relationship for  $He$ ,

$$d = 0.73 + 0.04(E / Ry)$$



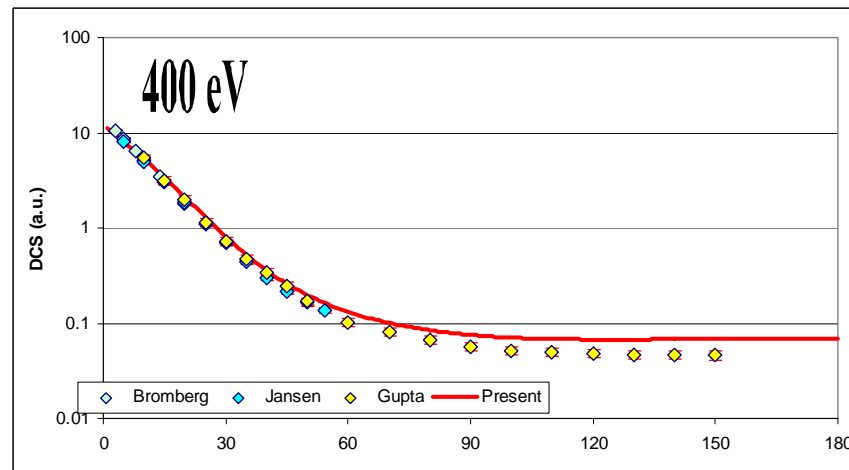
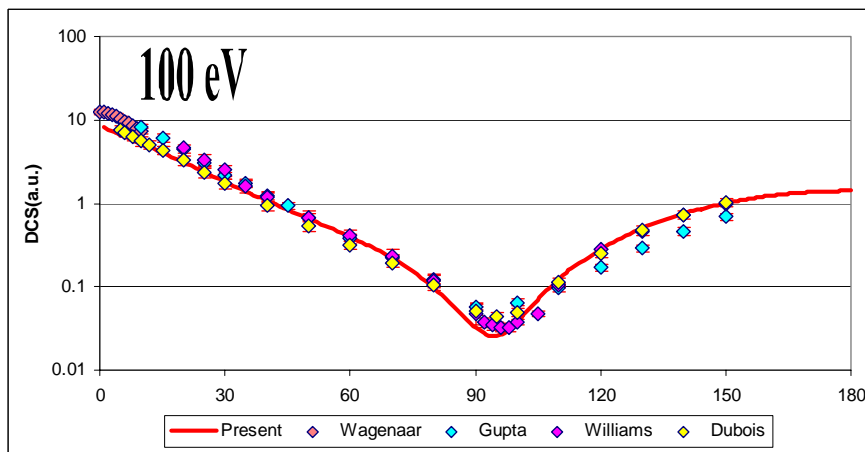
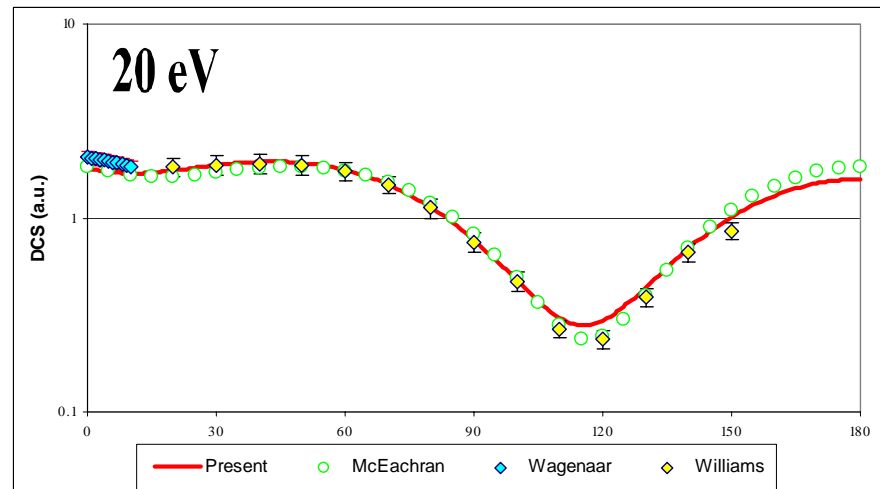
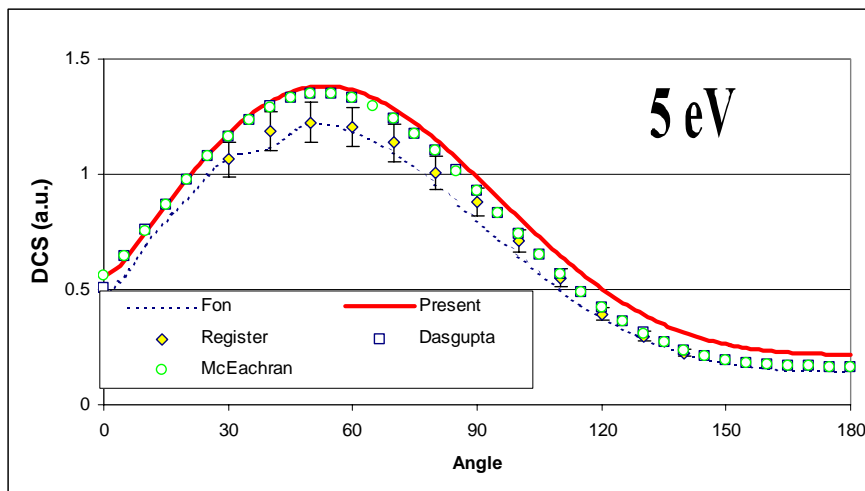
He

# Helium



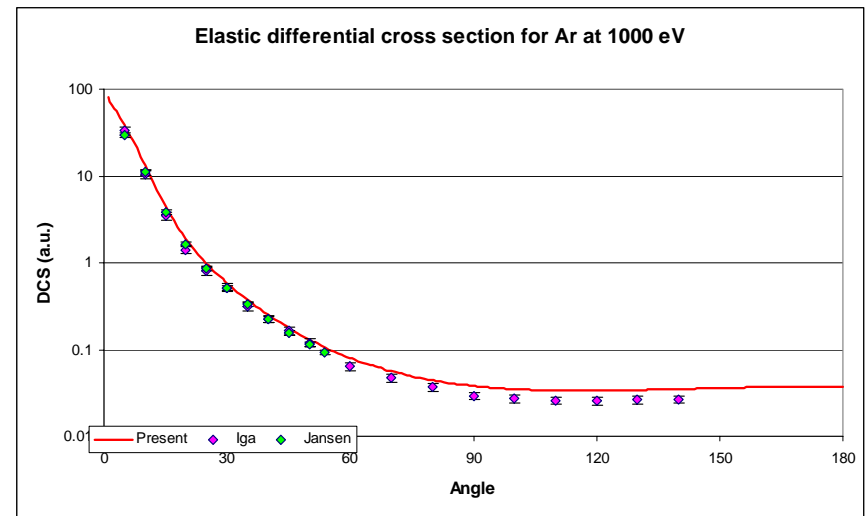
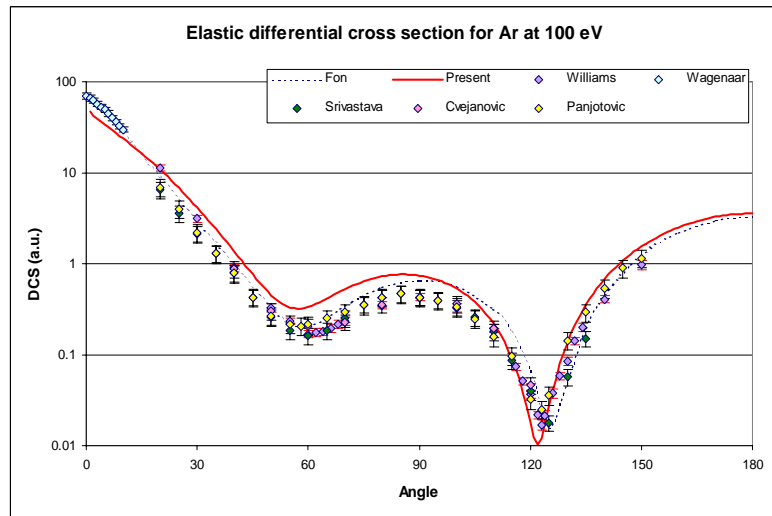
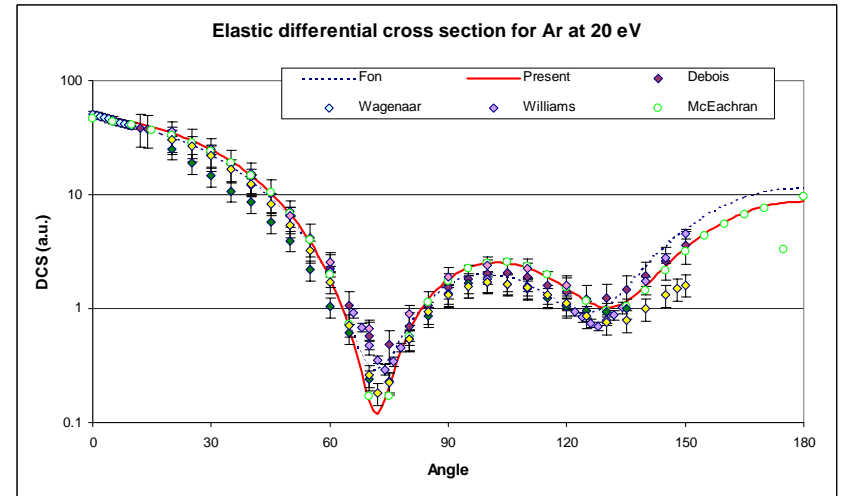
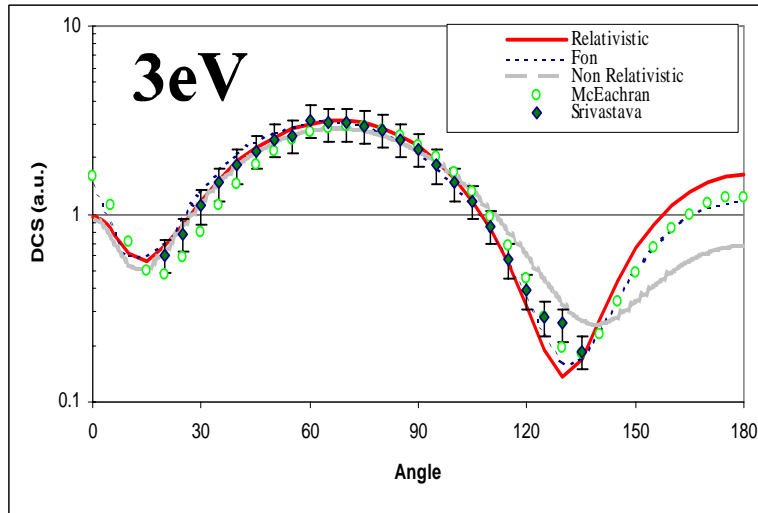
Ne

# Neon



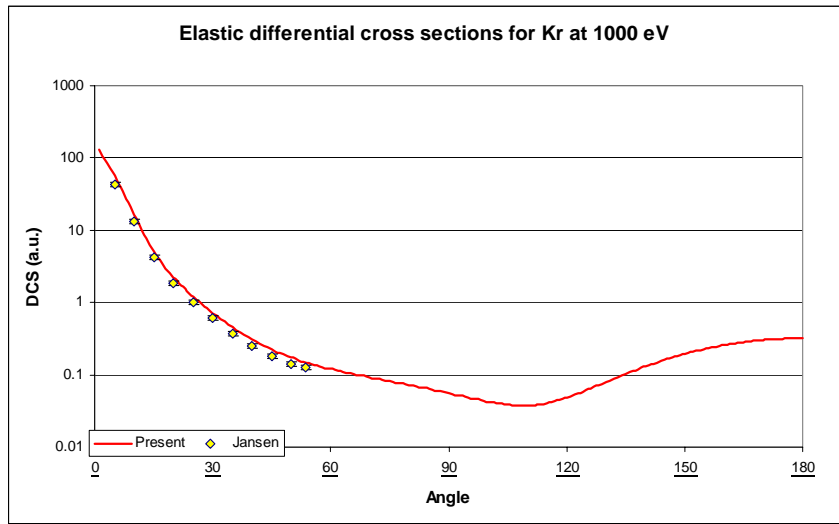
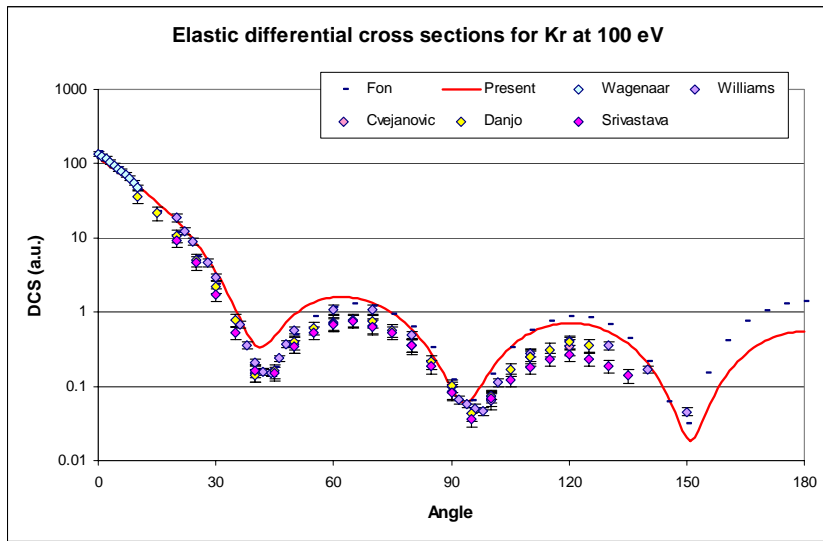
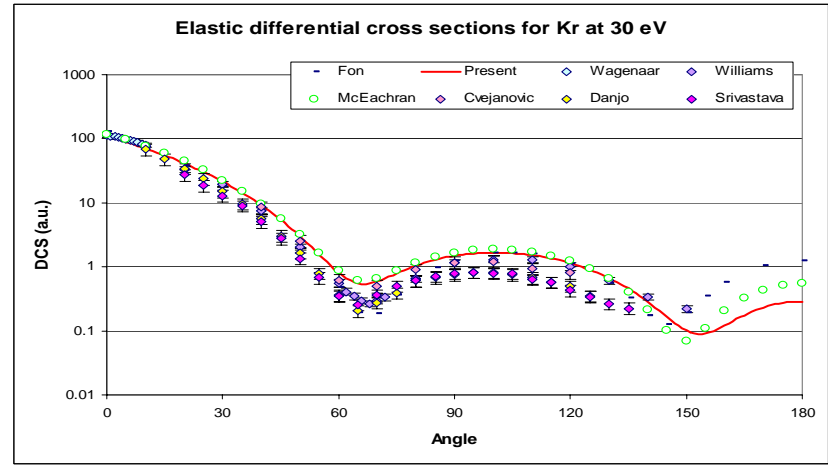
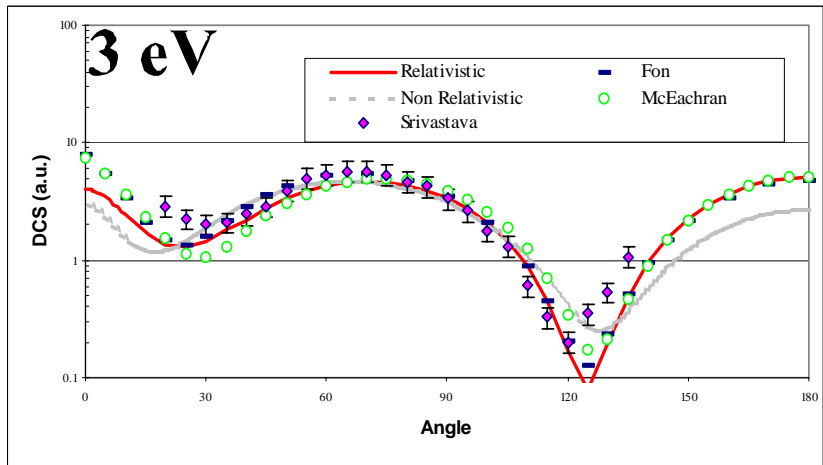
Ar

# Argon



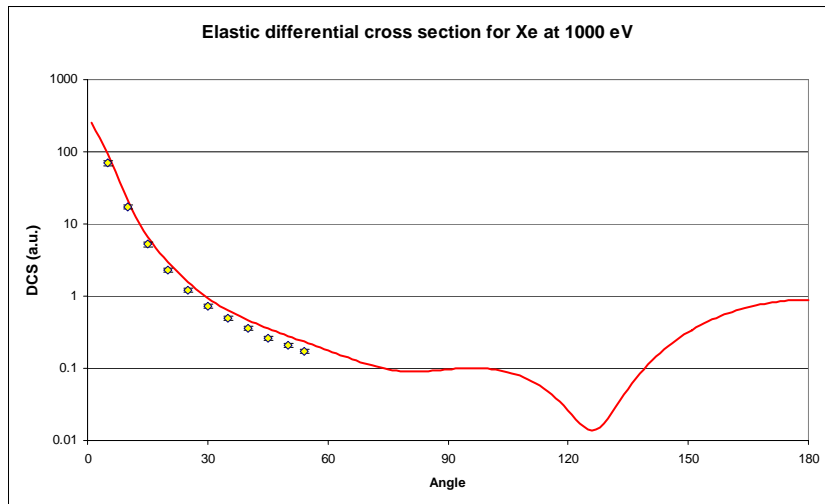
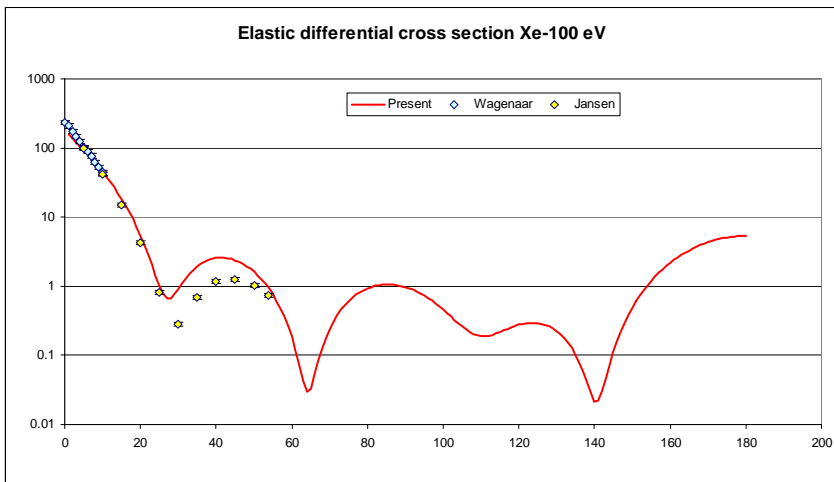
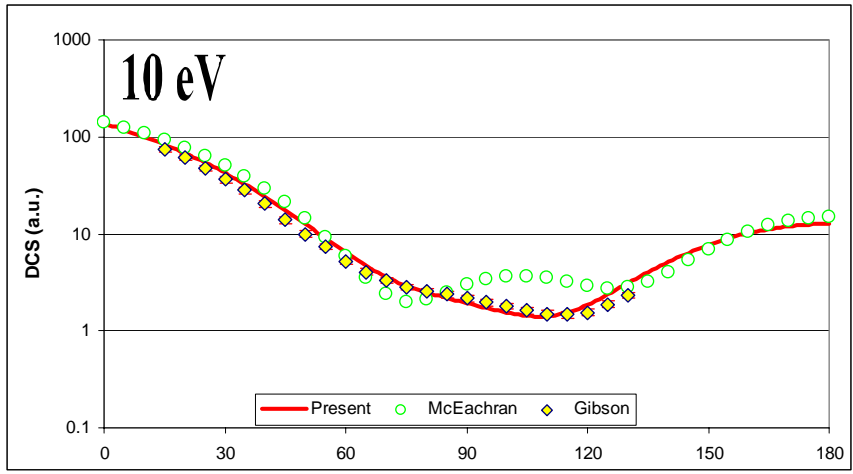
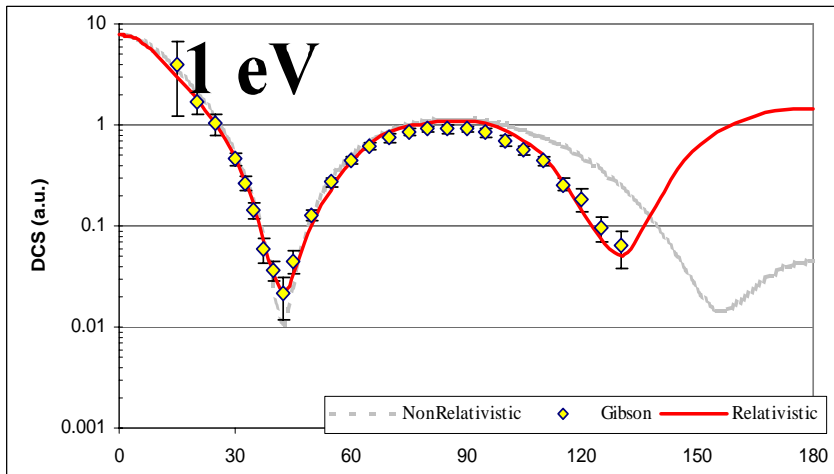
Kr

# Krypton



Xe

# Xenon



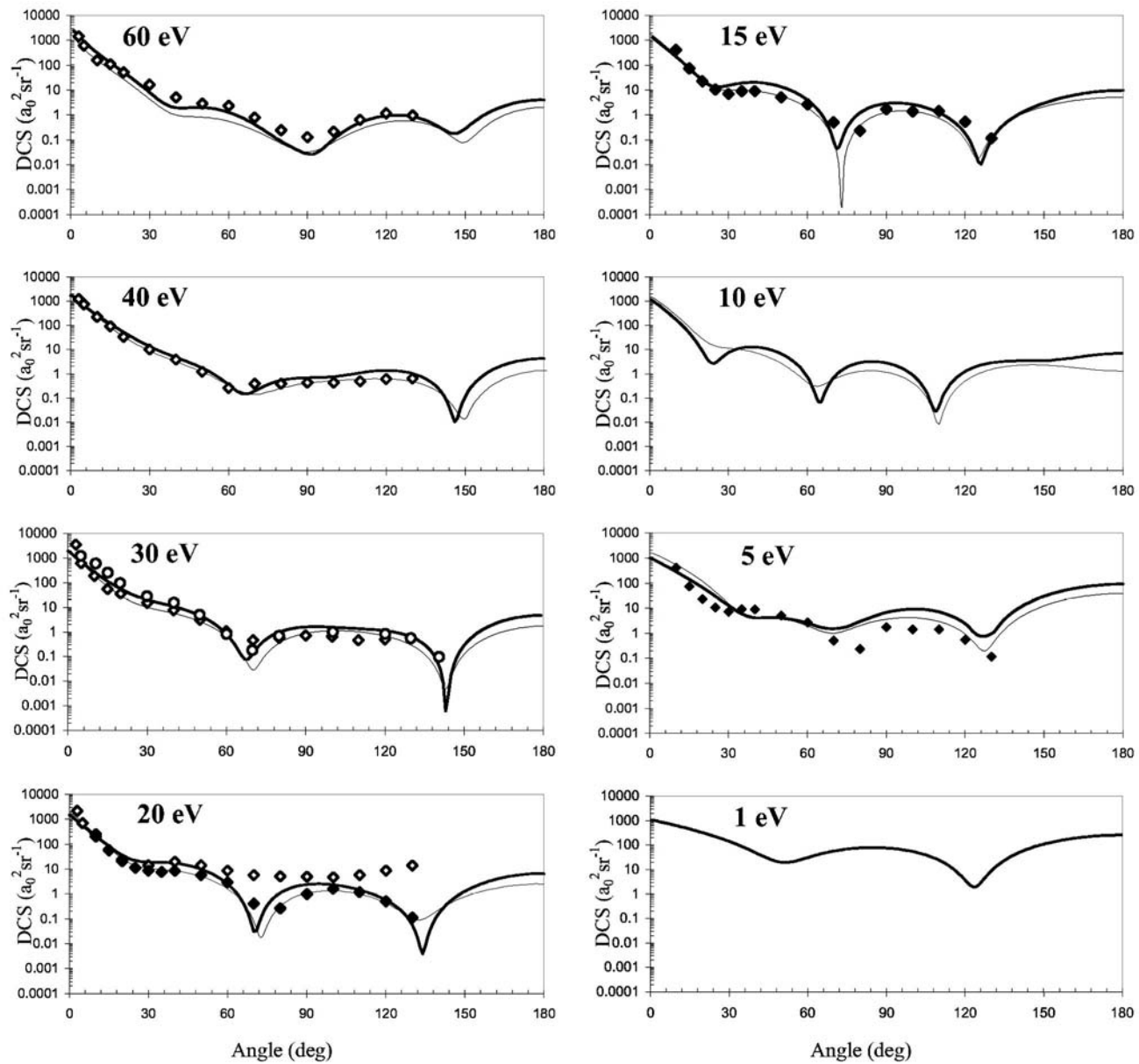


FIG. 2. Differential cross sections for elastic electron scattering by Ba from 1 to 60 eV: thick solid line, present work; thin solid line, CCC115 calculations of Fursa and Bray [22]. Experimental data:  $\diamond$ , Jensen *et al.* [13];  $\circ$ , Trajmar [17];  $\blacklozenge$ , Wang *et al.* [14].

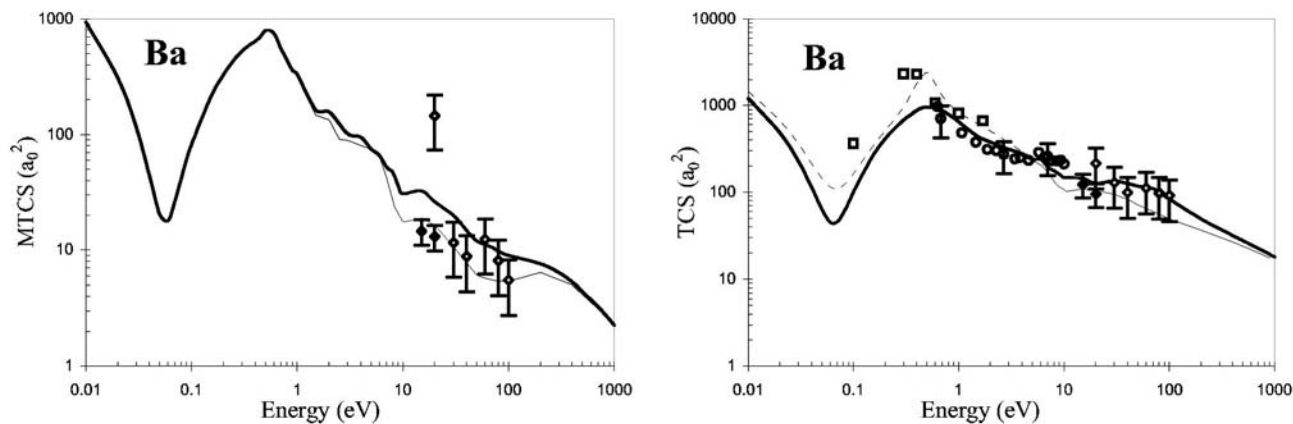


FIG. 7. Momentum transfer and total cross sections for elastic electron scattering from barium: thick solid line, present work. Experimental data:  $\circ$ , Romanyuk *et al.* [5];  $\diamond$ , Jensen *et al.* [13];  $\blacklozenge$ , Wang *et al.* [14]; Other theoretical: thin solid line, Fursa *et al.* [23]; dotted line, Yuan and Zhang [10];  $\square$ , Kelemen *et al.* [20].

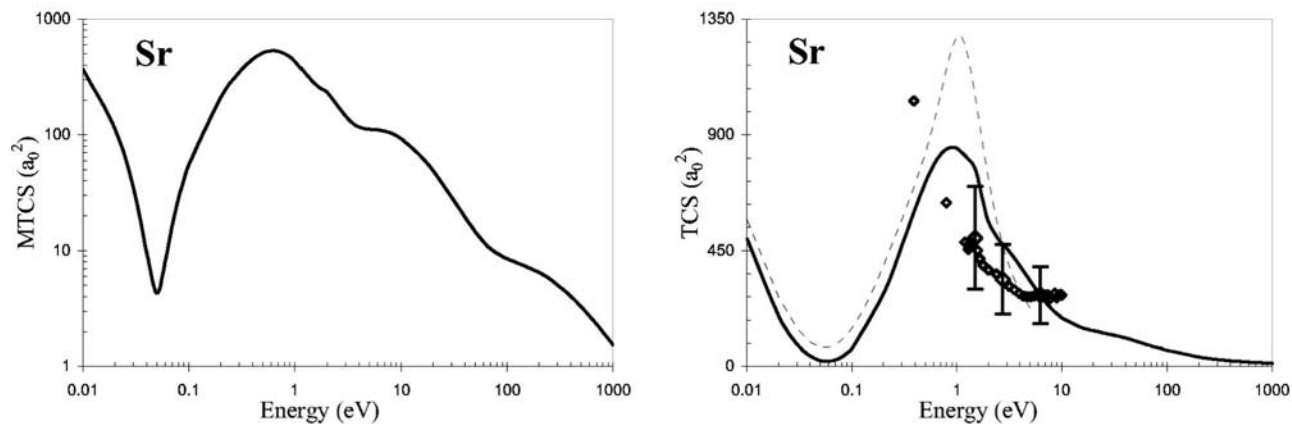


FIG. 8. Momentum transfer and total cross sections for elastic electron scattering from strontium: thick solid line, present work;  $\diamond$ , experimental data of Romanyuk *et al.* [5]; Other theoretical: dotted line, Yuan and Zhang [10].

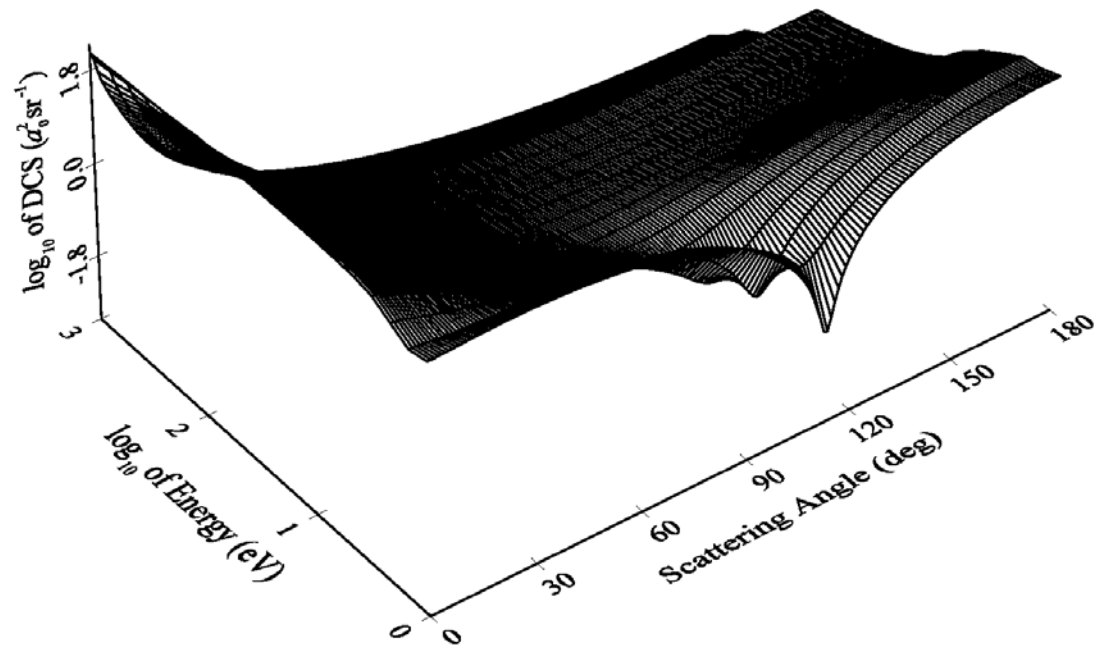


FIG. 3: a three-dimensional view of differential cross section for elastic electron scattering from beryllium.

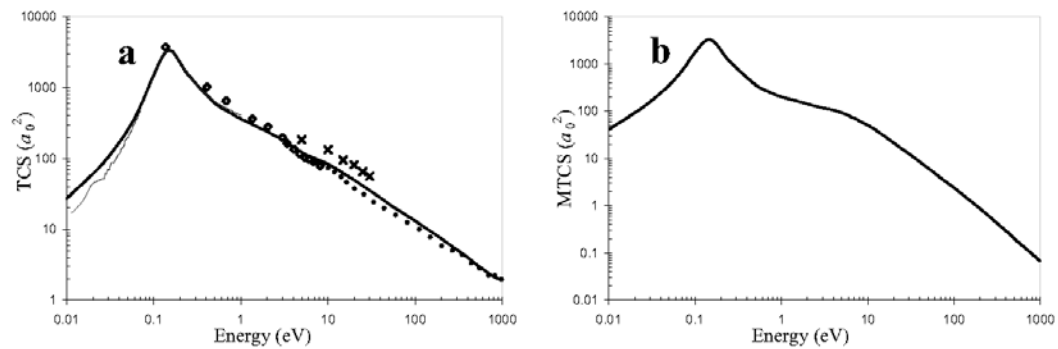


FIG. 4: Total (a) and momentum transfer (b) cross sections for elastic electron scattering from beryllium: thick solid line, present work; Other theoretical: thin solid line, Yuan and Zhang [11];  $\times$ , Kaushik *et al.* [10];  $\diamond$ , Fabrikant [9];  $\bullet$ , Fursa and Bray[4, 5].



Mg

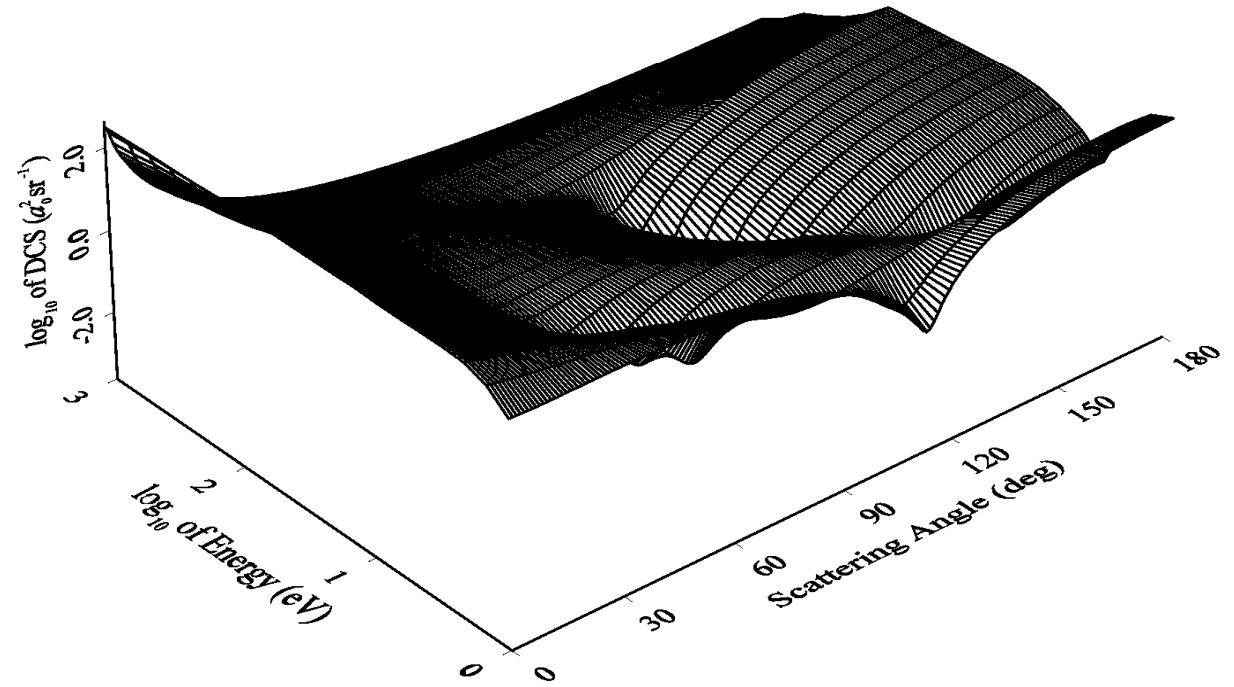


FIG. 7: a three-dimensional view of differential cross section for elastic electron scattering from magnesium.

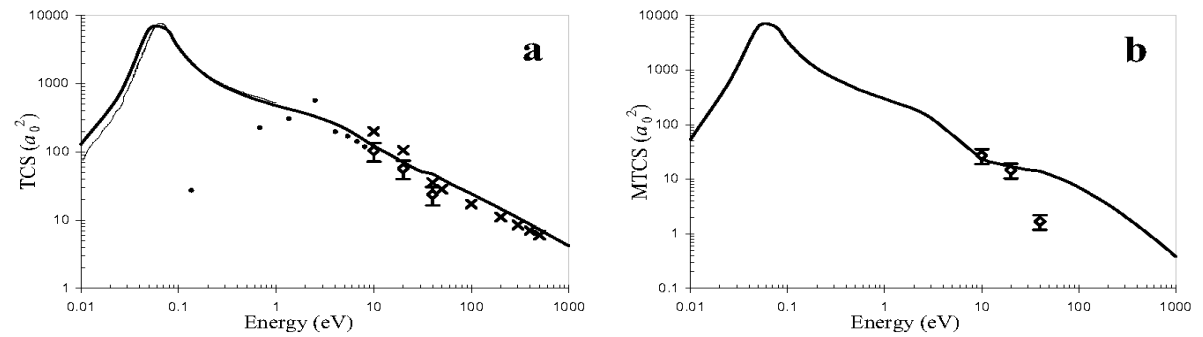


FIG. 8: Total (a) and momentum transfer (b) cross sections for elastic electron scattering from magnesium: thick solid line, present work; Experimental data:  $\diamond$ , Williams and Trajmar [15]; Other theoretical: thin solid line, Yuan and Zhang [11];  $\times$ , Khare *et al.* [20];  $\bullet$ , Fabrikant [9, 17].

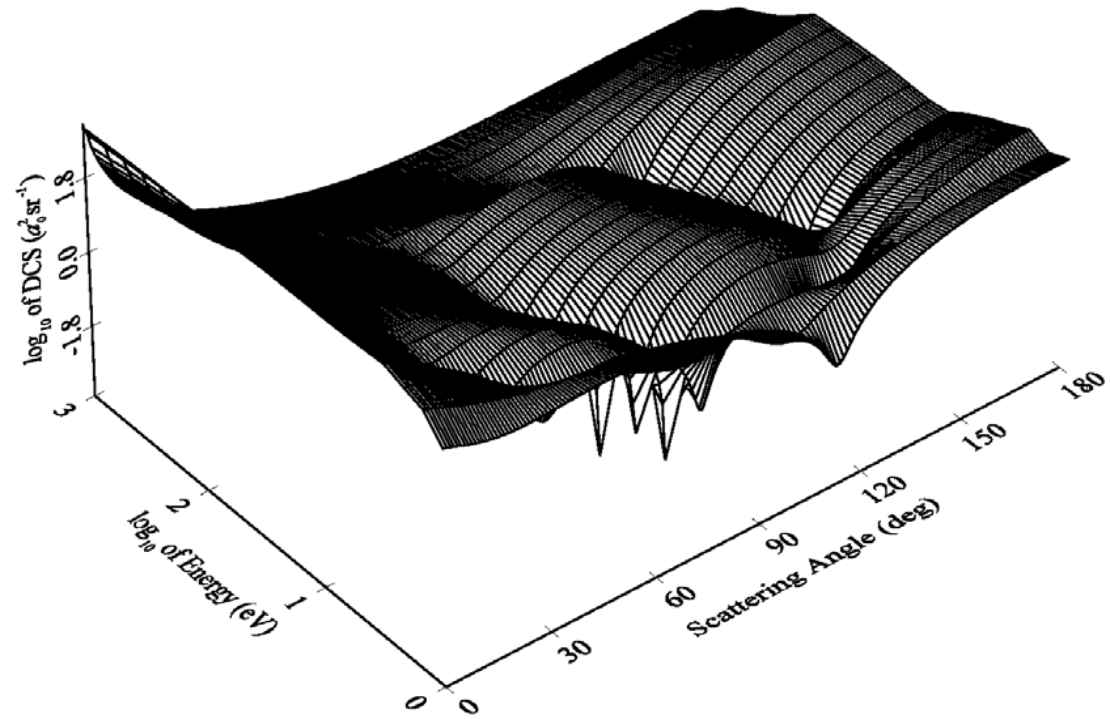


FIG. 11: a three-dimensional view of differential cross section for elastic electron scattering from calcium.

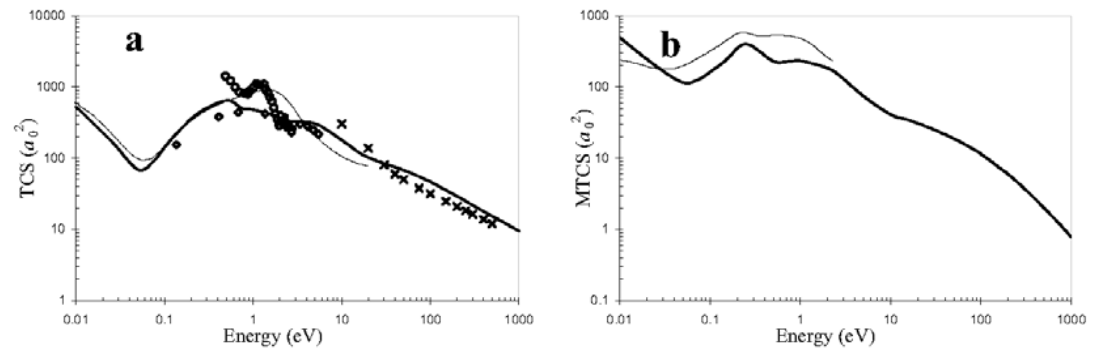
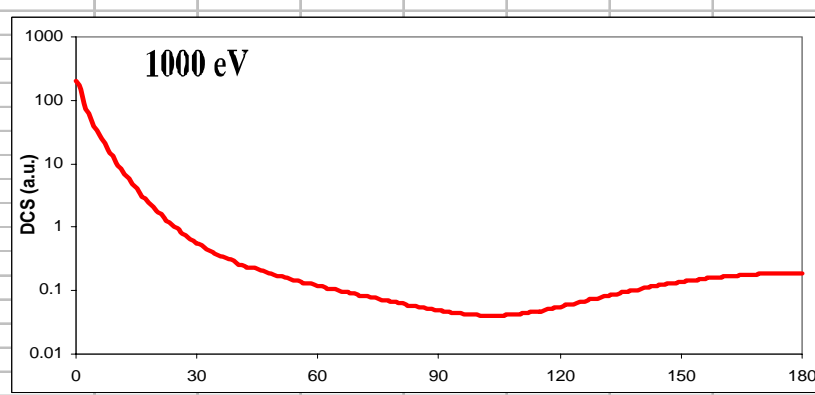
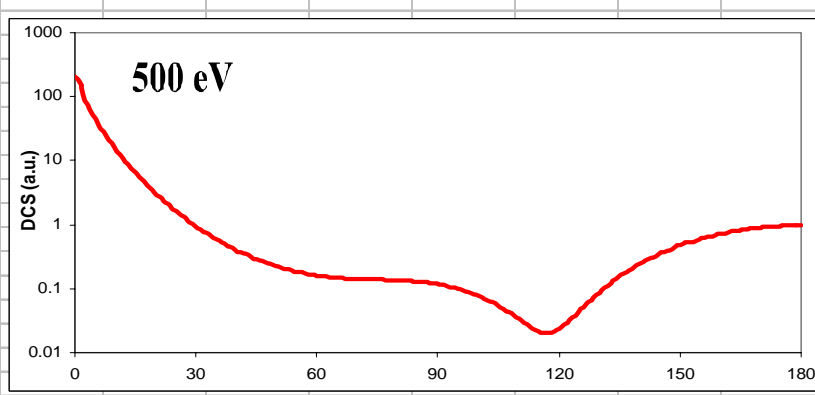
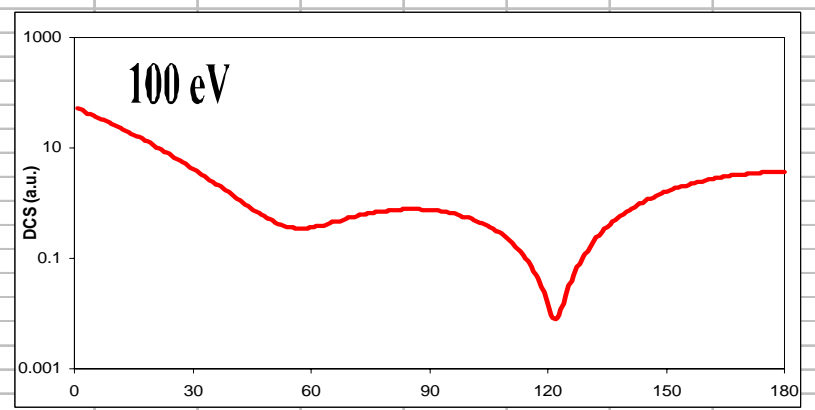
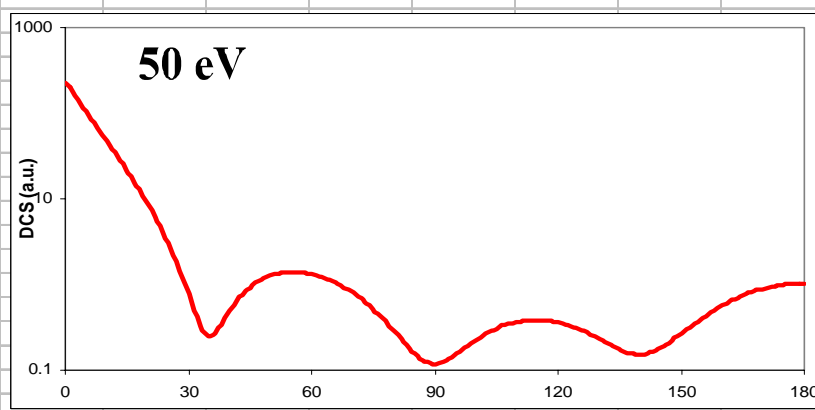
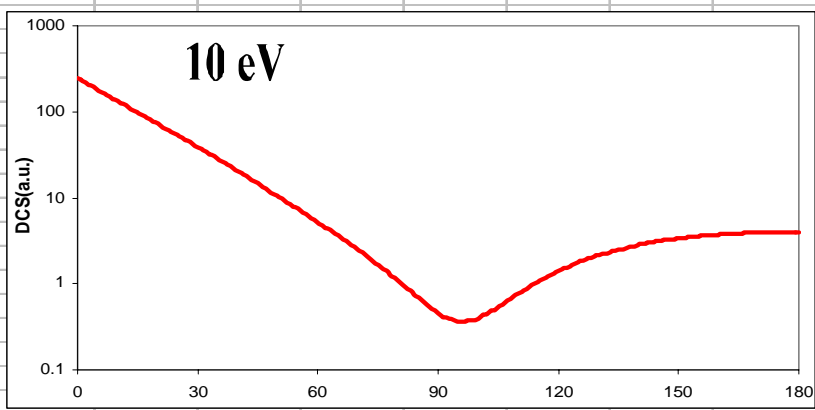
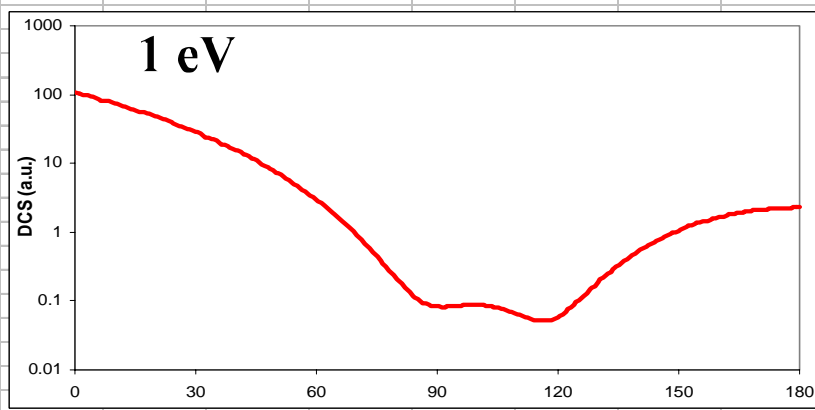


FIG. 12: Total (a) and momentum transfer (b) cross sections for elastic electron scattering from calcium: **Legend for TCS graph:** thick solid line, present work; Experimental data:  $\circ$ , Romanyuk *et al.* [23]; Other theoretical: thin solid line, Yuan [26];  $\diamond$ , Fabrikant [9];  $\times$ , Khare *et al.* [24]; **Legend for MTCS graph:** thick solid line, present work; Other theoretical: thin solid line, Cribakin *et al.* [25].

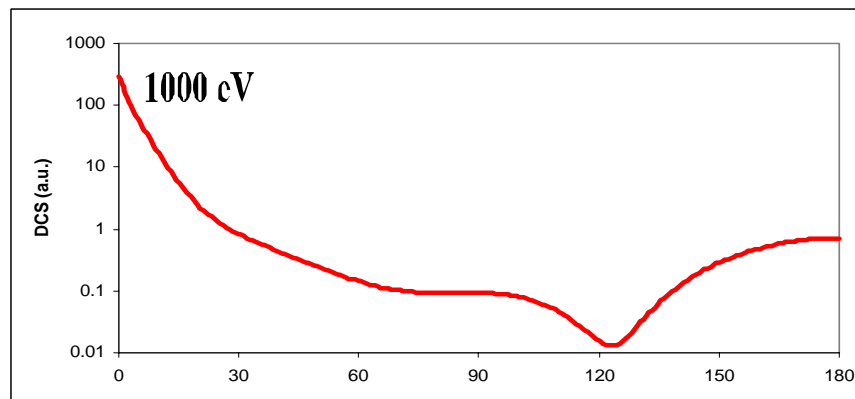
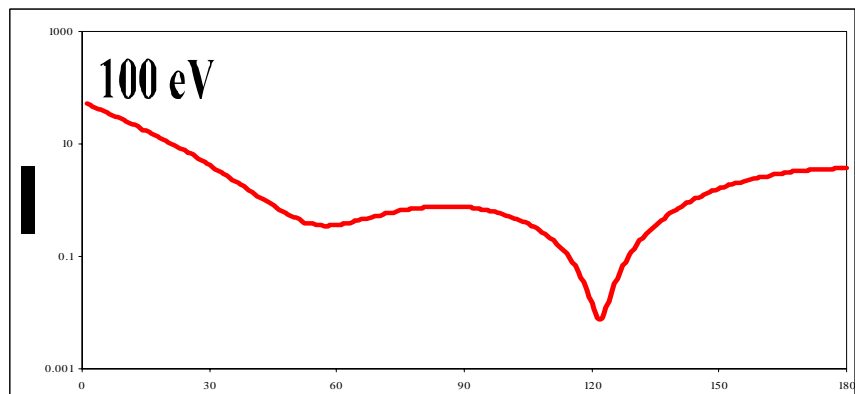
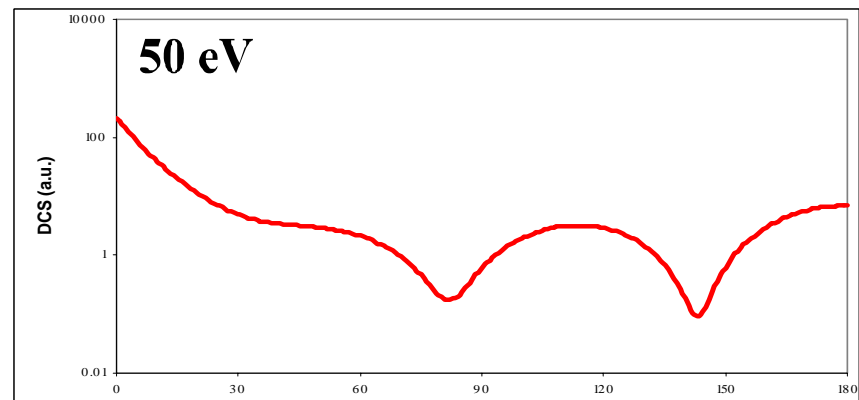
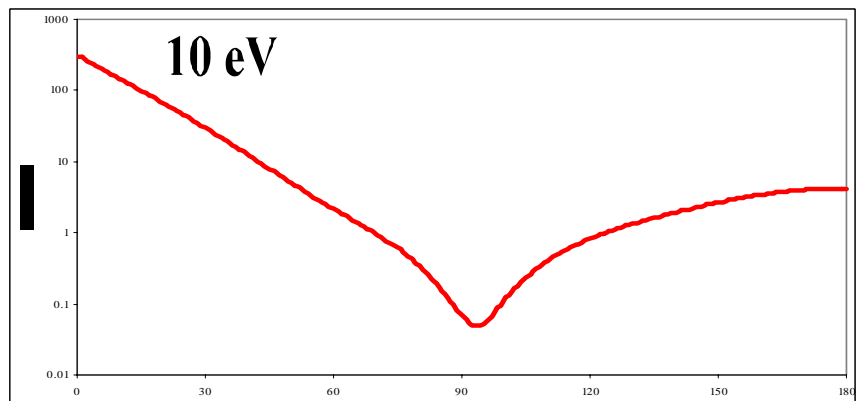
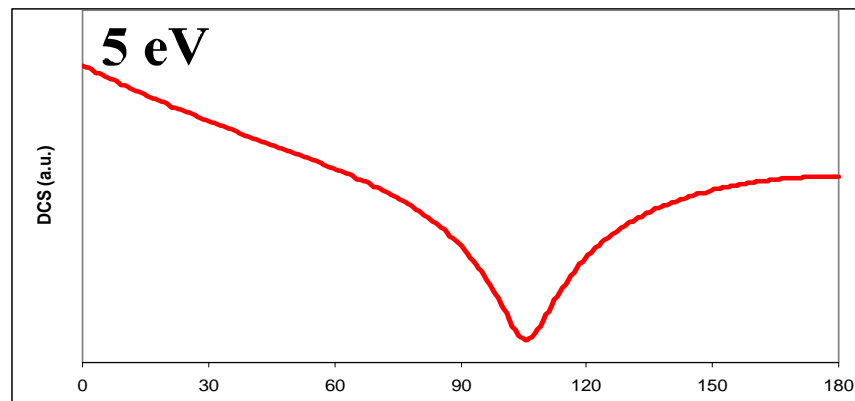
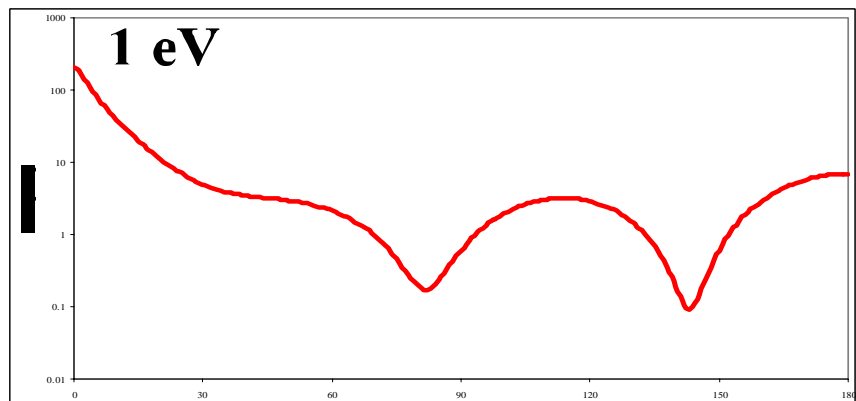
Zn

# Zinc



Cd

# Cadmium



Hg

# Mercury

