

# M3D Sawtooth Update

Josh Breslau

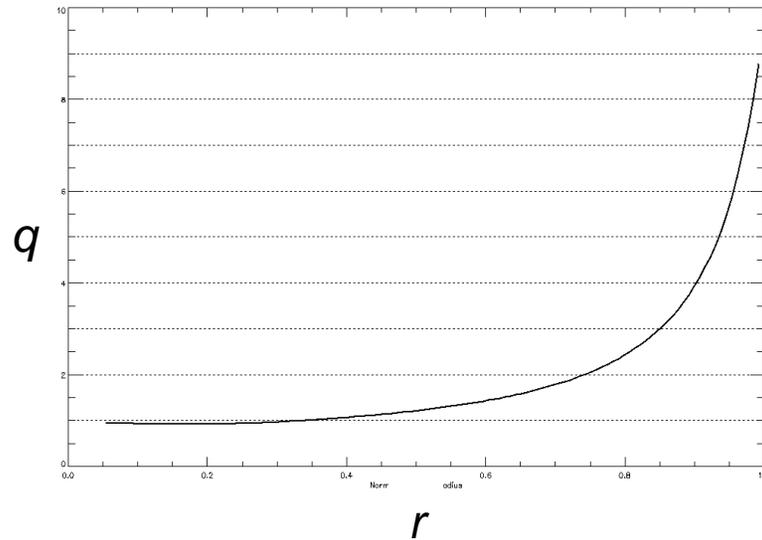
CEMM Meeting

Annapolis, MD

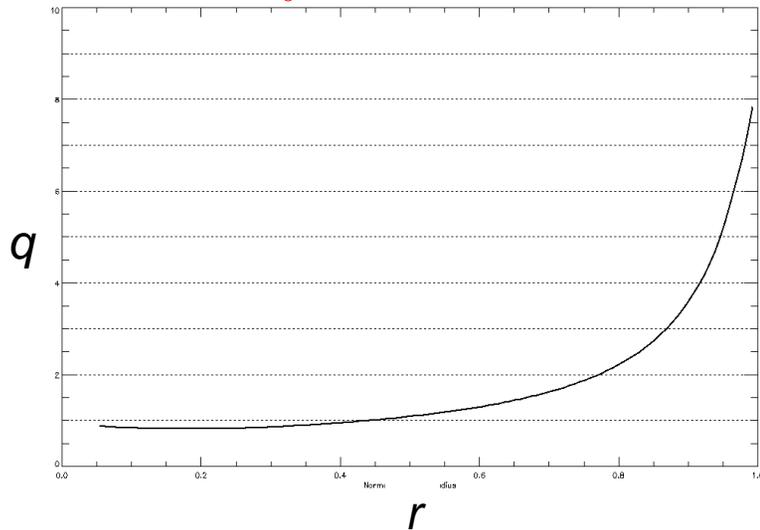
April 22, 2007

# CDX Equilibria

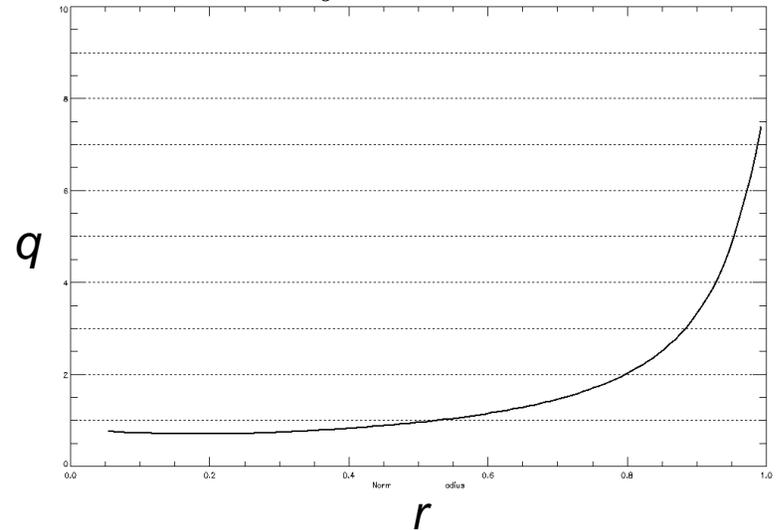
Original: time 11:  $q_0 = 0.92$ ;  $q=1$  at  $r=0.33$



time 19:  $q_0 = 0.82$ ;  $q=1$  at  $r=0.44$



time 29:  $q_0 = 0.71$ ;  $q=1$  at  $r=0.53$



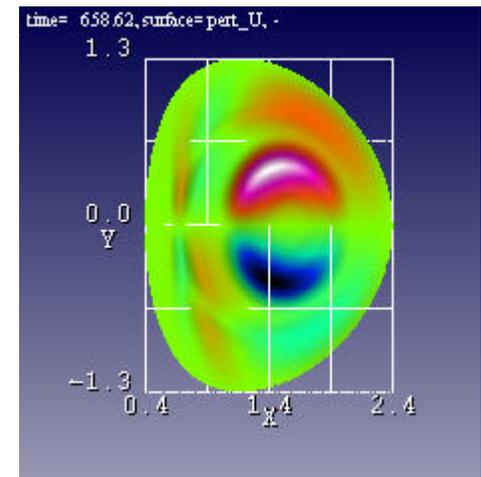
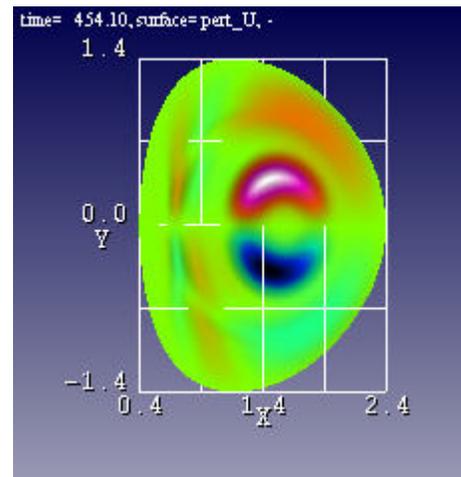
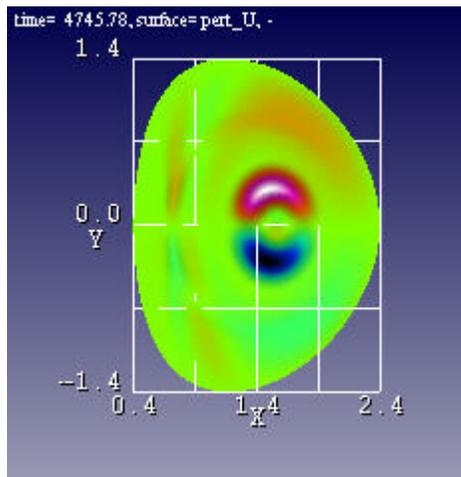
# n=1 Eigenmodes

time 11  
 $q_0 = 0.92$

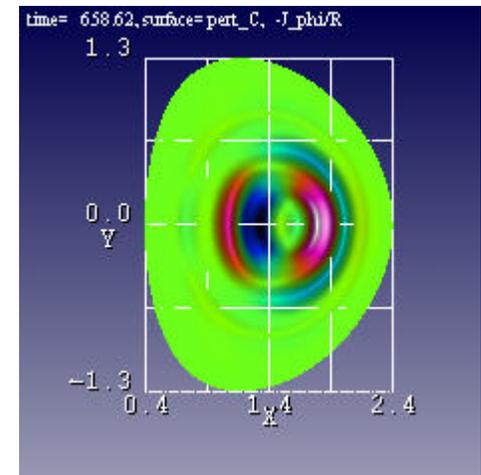
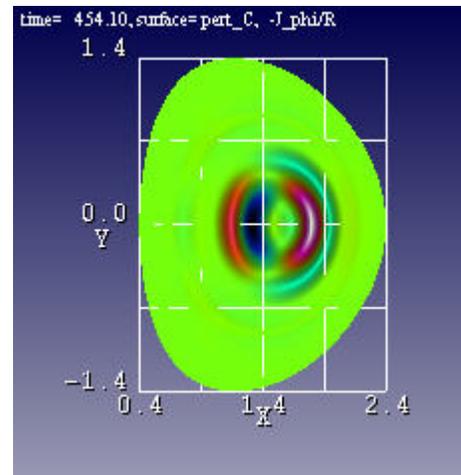
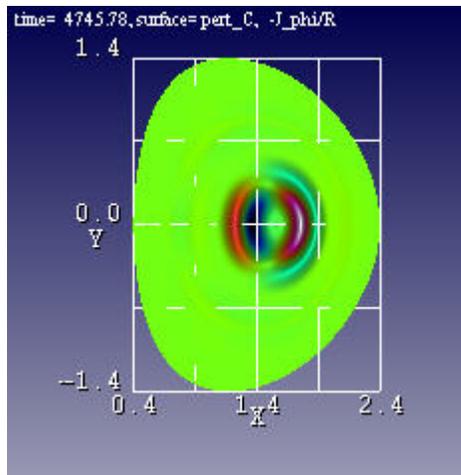
time 19  
 $q_0 = 0.82$

time 29  
 $q_0 = 0.71$

Poloidal  
velocity  
stream  
function

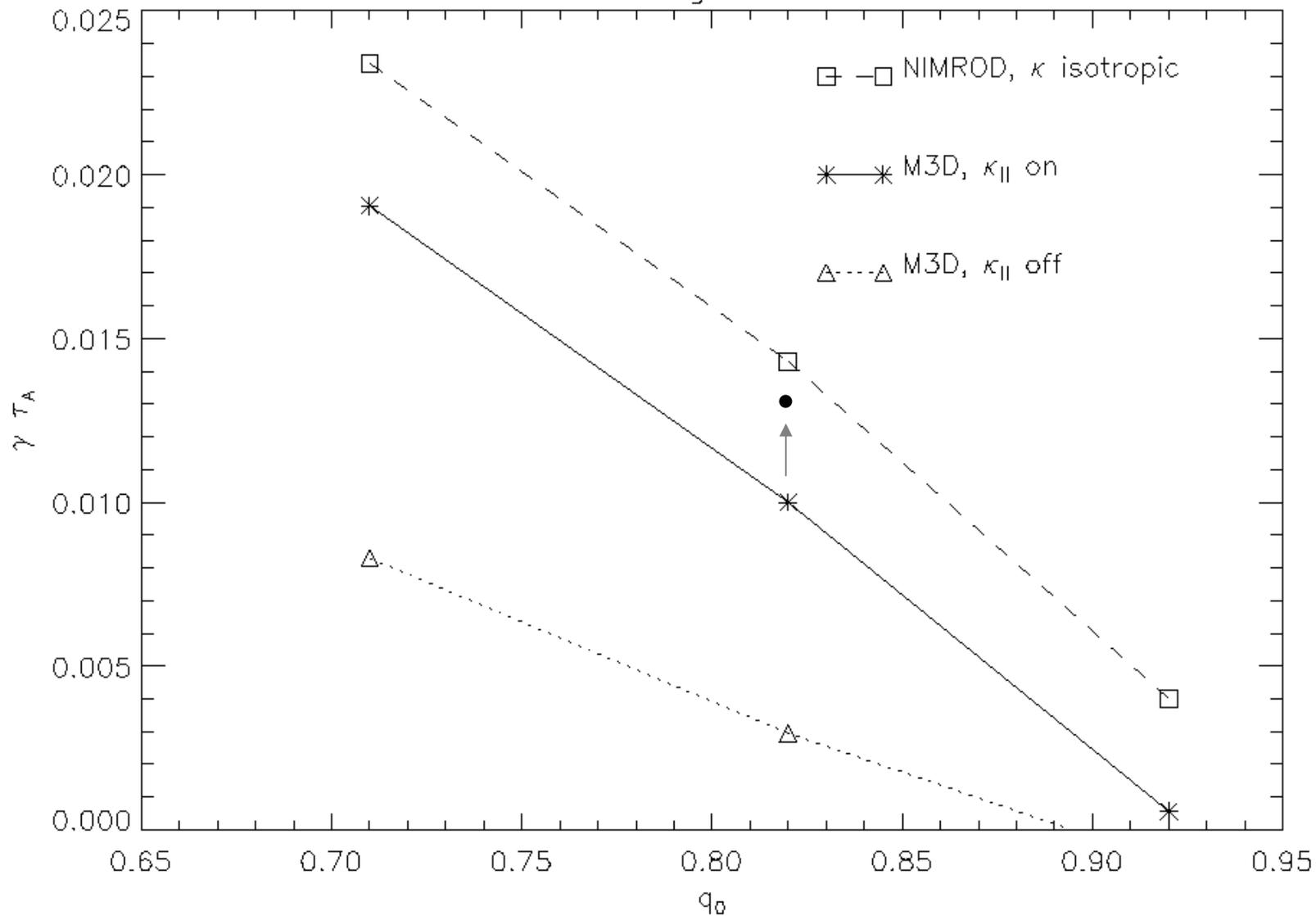


Toroidal  
current  
density

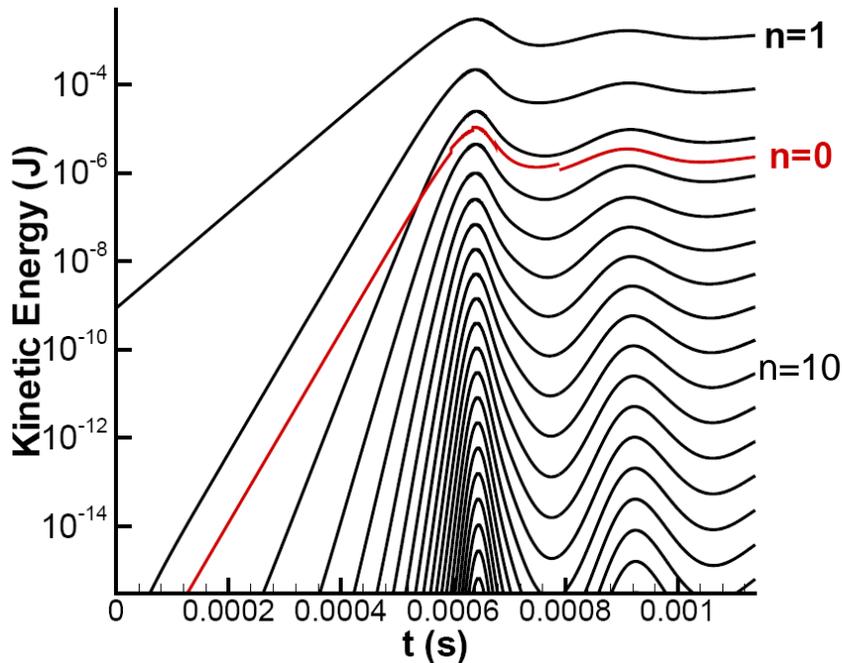


# Linear $n=1$ Growth Rates

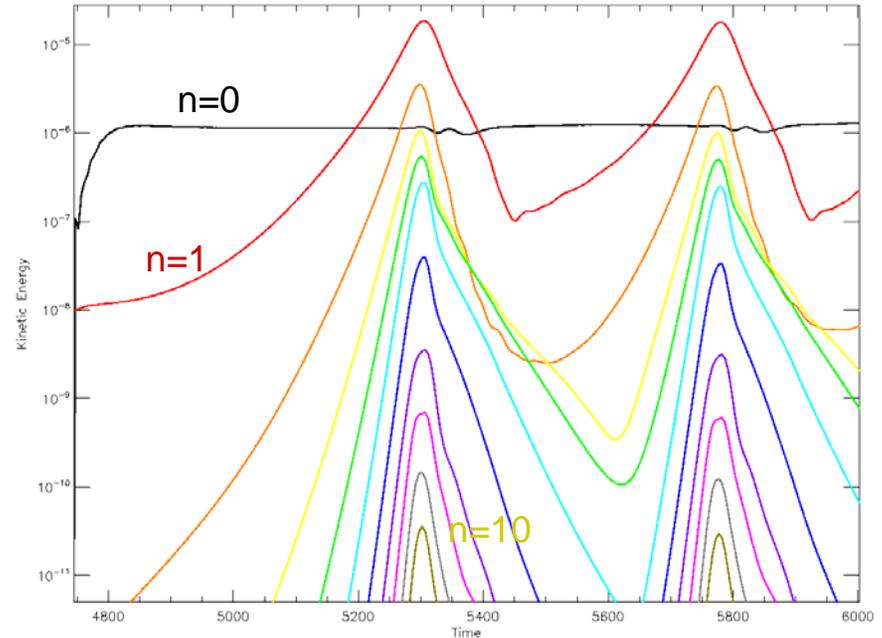
Converged rates



# Differences Between Initial NIMROD and M3D CDX-U Nonlinear Results (time 11)



NIMROD



M3D

- M3D  $n=1$  growth rate quickly exceeds linear due to rapid drop in  $q_0$ .
- M3D has much higher  $n=0$  energy compared to other modes.
- Periods between crashes differ:  $\sim 800 \tau_A$  for NIMROD vs.  $480 \tau_A$  for M3D.
- 2<sup>nd</sup> crash energy is diminished more in NIMROD than in M3D.

# Understanding the M3D Result

- Linear growth rate of  $n=1$  increases in nonlinear run because  $q_0$  is steadily dropping.
- Rate of change of  $q_0$  is directly proportional to average velocity of  $n=0$  flow, which shows up as large  $n=0$  kinetic energy. It varies inversely with  $dt$ .
- $n=0$  flow arises from discretization error; converges to zero as poloidal resolution increases.
- Culprit: interaction of flow with toroidal current density ( $C$ ) source term exposes poor conservation properties of existing formulation of  $C$  equation.

# Original C Equation

Poloidal flux (aeqn, numerically unstable, boundary problems):

$$\frac{\partial \psi}{\partial t} = \frac{R}{R_0} [U, \psi] + \frac{R}{R_0} (U, F) - (\chi, \psi) + [\chi, F] + \eta C + \frac{\partial \Phi}{\partial \phi}. \quad (1)$$

Toroidal current density (ceqn, poor conservation properties):

$$C_a \equiv \Delta^* \psi :$$

$$\begin{aligned} \frac{\partial C_a}{\partial t} = & \frac{R}{R_0} \left\{ [U, C_a] + [\Delta^\dagger U, \psi] + 2 \left[ \frac{\partial U}{\partial R}, \frac{\partial \psi}{\partial R} \right] + 2 \left[ \frac{\partial U}{\partial z}, \frac{\partial \psi}{\partial z} \right] \right\} + \frac{2}{R_0} \left[ U, \frac{\partial \psi}{\partial R} \right] + \frac{2}{R_0 R} \frac{\partial U}{\partial z} \frac{\partial \psi}{\partial R} \\ & + \frac{R}{R_0} \left\{ (U, \nabla_\perp^2 F) + (\Delta^\dagger U, F) + 2 \left( \frac{\partial U}{\partial R}, \frac{\partial F}{\partial R} \right) + 2 \left( \frac{\partial U}{\partial z}, \frac{\partial F}{\partial z} \right) \right\} + \frac{1}{R_0} \left( \frac{\partial F}{\partial R}, U \right) - \frac{1}{R_0 R} \frac{\partial F}{\partial z} \frac{\partial U}{\partial z} \\ & - \left\{ (\psi, \nabla_\perp^2 \chi) + (C_a, \chi) + 2 \left( \frac{\partial \psi}{\partial R}, \frac{\partial \chi}{\partial R} \right) + 2 \left( \frac{\partial \psi}{\partial z}, \frac{\partial \chi}{\partial z} \right) \right\} + \frac{1}{R} \left( \frac{\partial \chi}{\partial R}, \psi \right) + \frac{1}{R^2} \frac{\partial \psi}{\partial R} \frac{\partial \chi}{\partial R} \\ & + \left\{ [\nabla_\perp^2 \chi, F] + [\chi, \nabla_\perp^2 F] + 2 \left[ \frac{\partial \chi}{\partial R}, \frac{\partial F}{\partial R} \right] + 2 \left[ \frac{\partial \chi}{\partial z}, \frac{\partial F}{\partial z} \right] \right\} - \frac{1}{R} \left\{ \left[ \frac{\partial \chi}{\partial R}, F \right] + \left[ \chi, \frac{\partial F}{\partial R} \right] \right\} \\ & + \frac{\partial}{\partial \phi} (\nabla_\perp^2 \Phi) - \frac{1}{R} \frac{\partial^2 \Phi}{\partial \phi \partial R} \end{aligned} \quad (1')$$

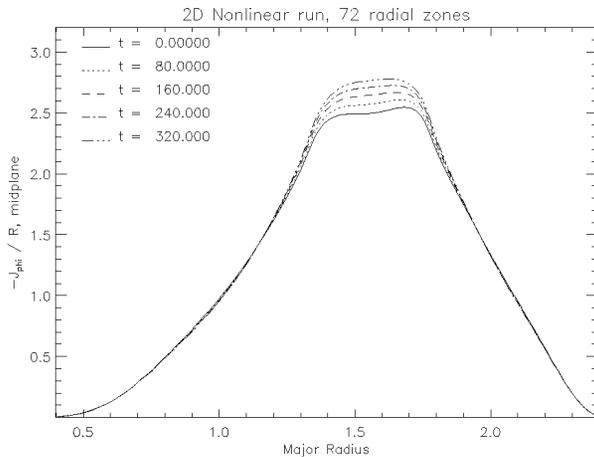
# Conservative C Equation

If  $\dot{\psi} = b$ , then  $\lambda_i \nabla_{\perp} \cdot \left( \frac{1}{R^2} \nabla \dot{\psi} \right) = \frac{\lambda_i}{R^2} \dot{C}_a = \lambda_i \nabla_{\perp} \cdot \left( \frac{1}{R^2} \nabla b \right)$

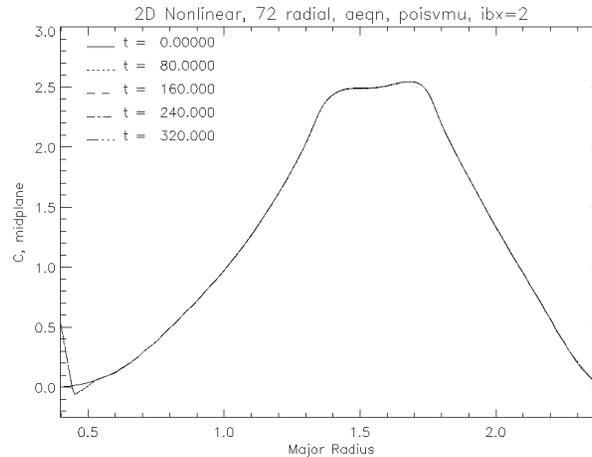
Integrating over the domain results in  $\vec{M} \cdot \left( \frac{\dot{C}_a}{R^2} \right) = \vec{S} \cdot \left( \frac{b}{R^2} \right)$ ,

or approximately  $\vec{M} \cdot \dot{C}_a = \vec{S} \cdot b$ .

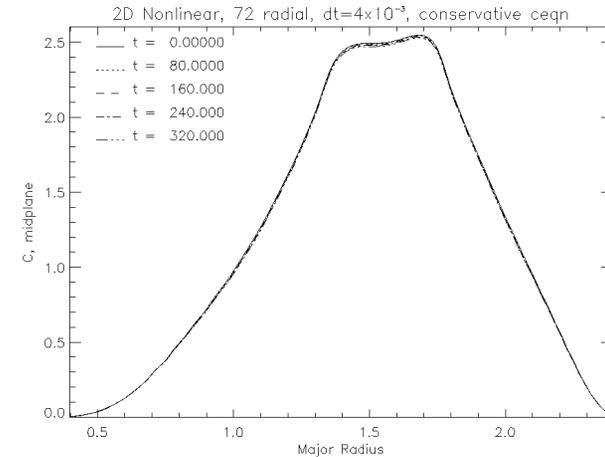
Original ceqn



Original aeqn

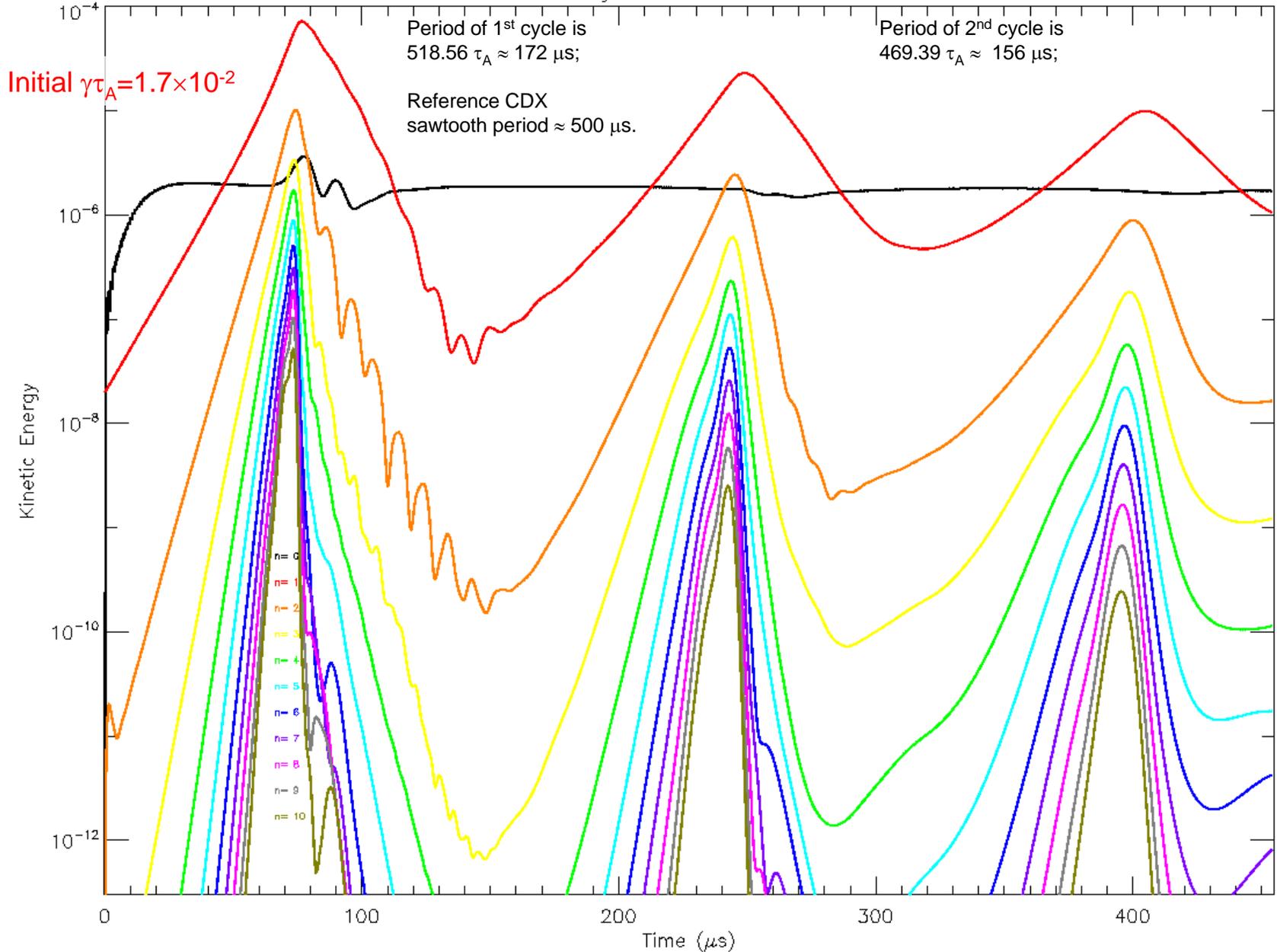


Conservative ceqn

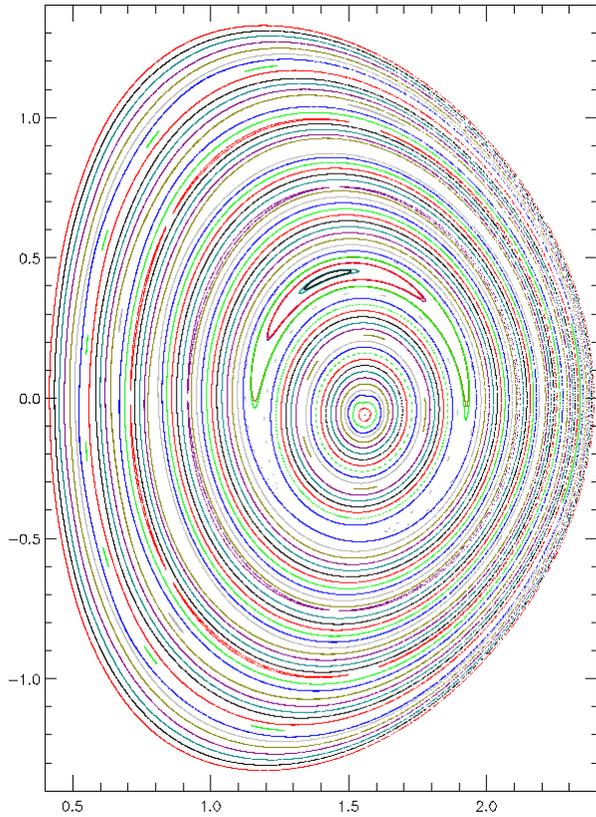


# Mode History

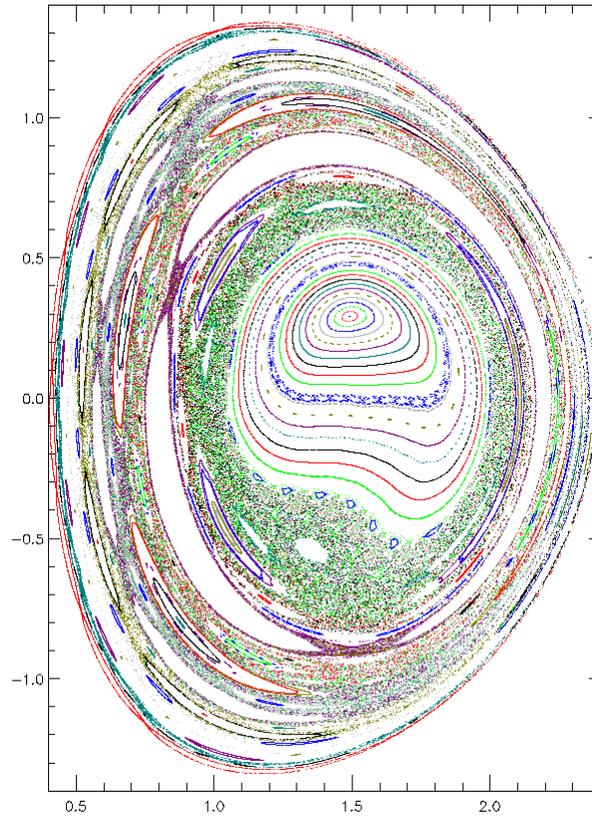
By Mode Number



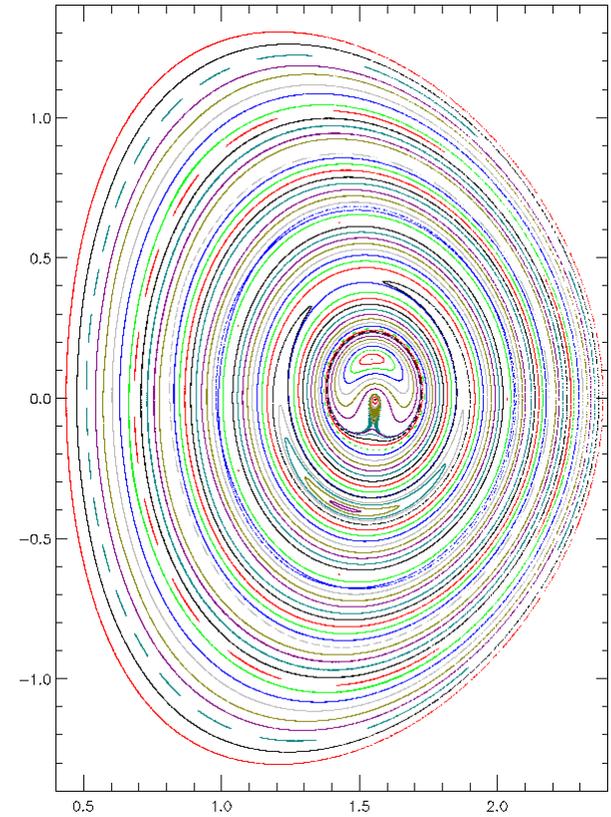
# Poincaré Plots



$t = 38 \mu\text{s}; \varphi = \pi/2$   
(compare to NIMROD  $t = 129 \mu\text{s}$ )

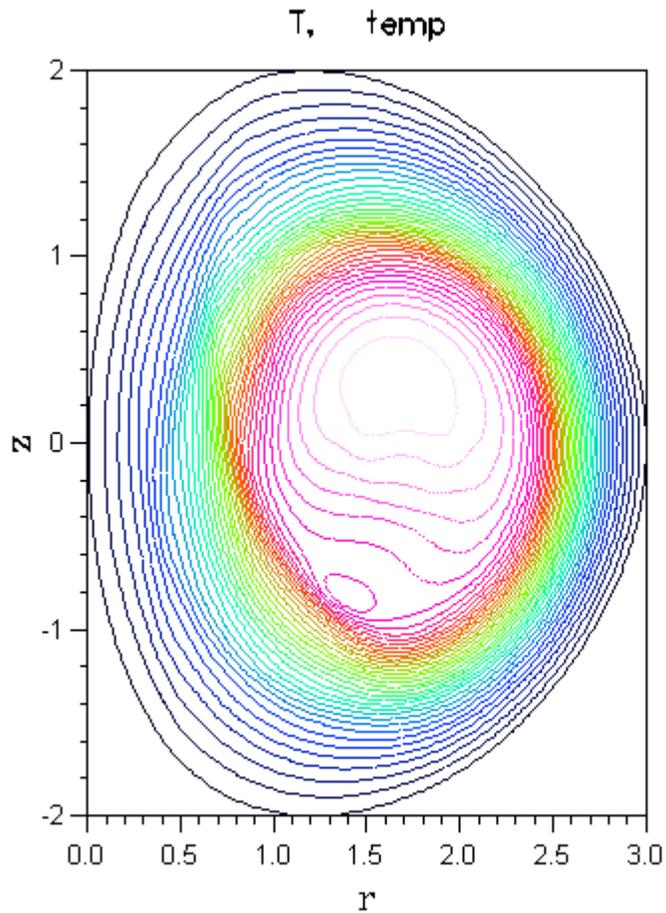


$t = 74.5 \mu\text{s}; \varphi = \pi/2$   
(compare to NIMROD  $t = 163 \mu\text{s}$ )

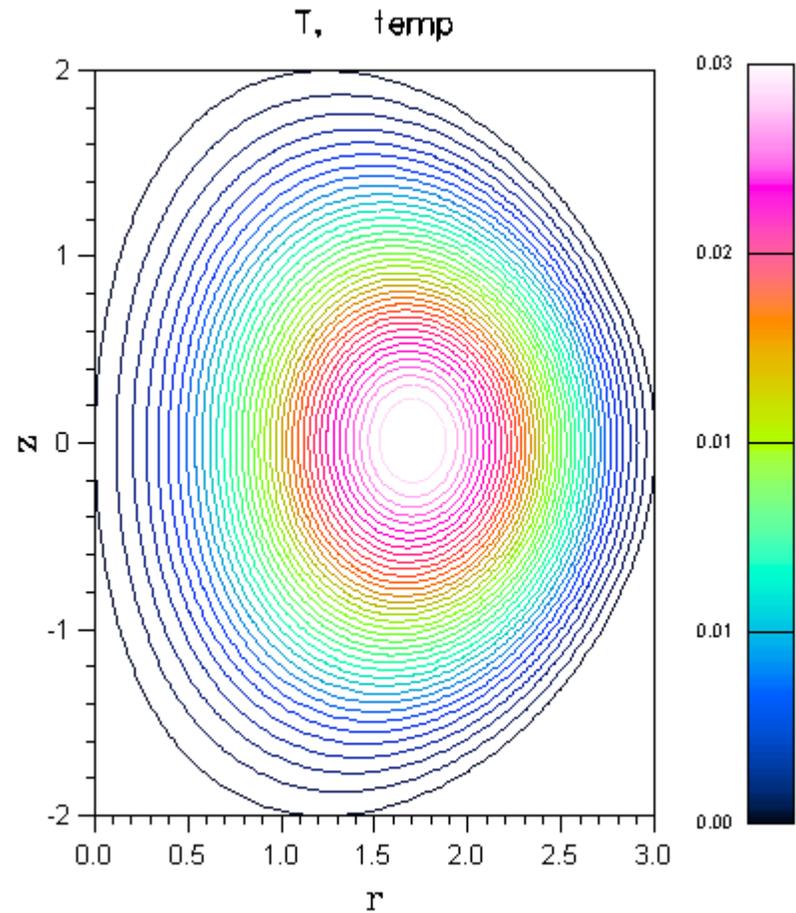


$t = 115.3 \mu\text{s}; \varphi = \pi/2$   
(compare to NIMROD  $t = 200 \mu\text{s}$ )

# Temperature Contours

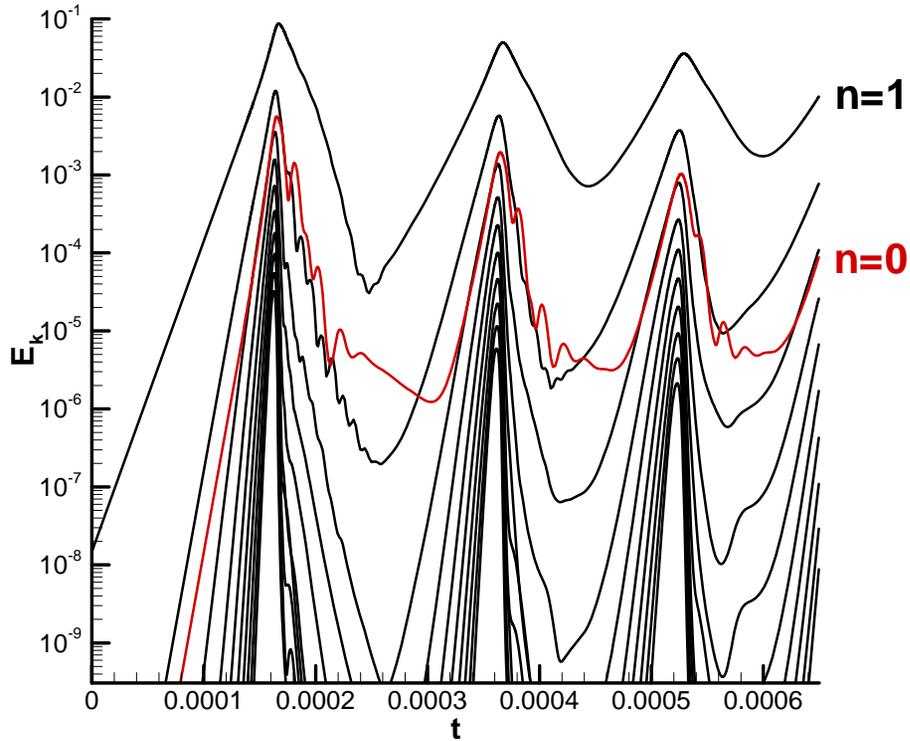


$t = 74.5 \mu\text{s}; \varphi = \pi/2$   
(compare to NIMROD  $t = 163 \mu\text{s}$ )

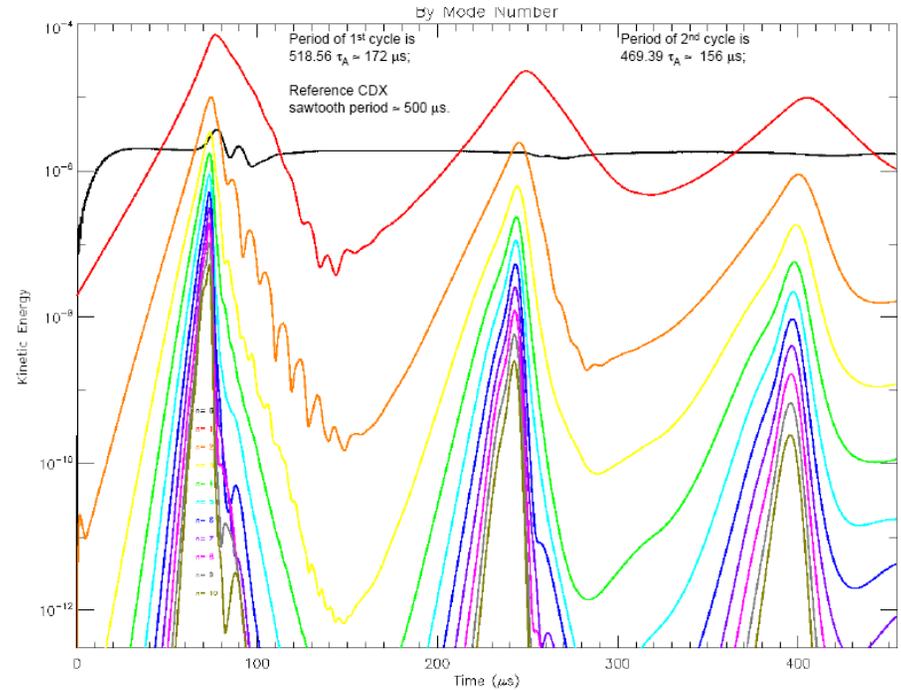


$t = 115.3 \mu\text{s}; \varphi = \pi/2$   
(compare to NIMROD  $t = 200 \mu\text{s}$ )

# Comparison with NIMROD Results

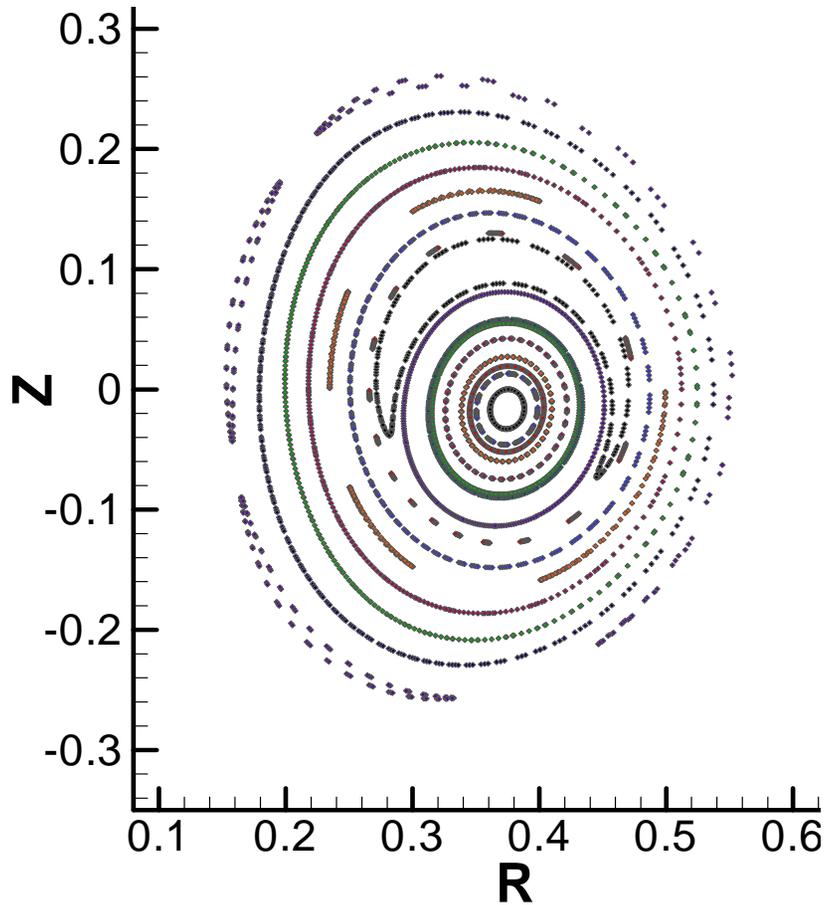


NIMROD Kinetic energy in first 10 modes

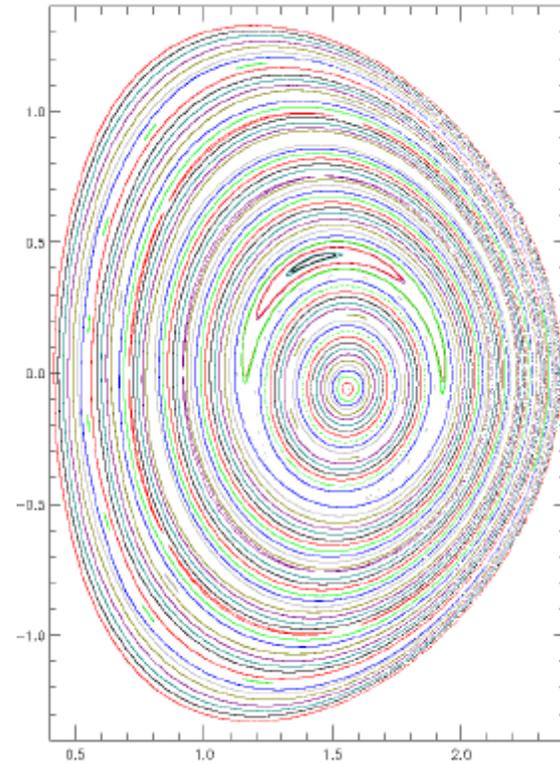


M3D Kinetic energy in first 10 modes

# Poincaré Plots

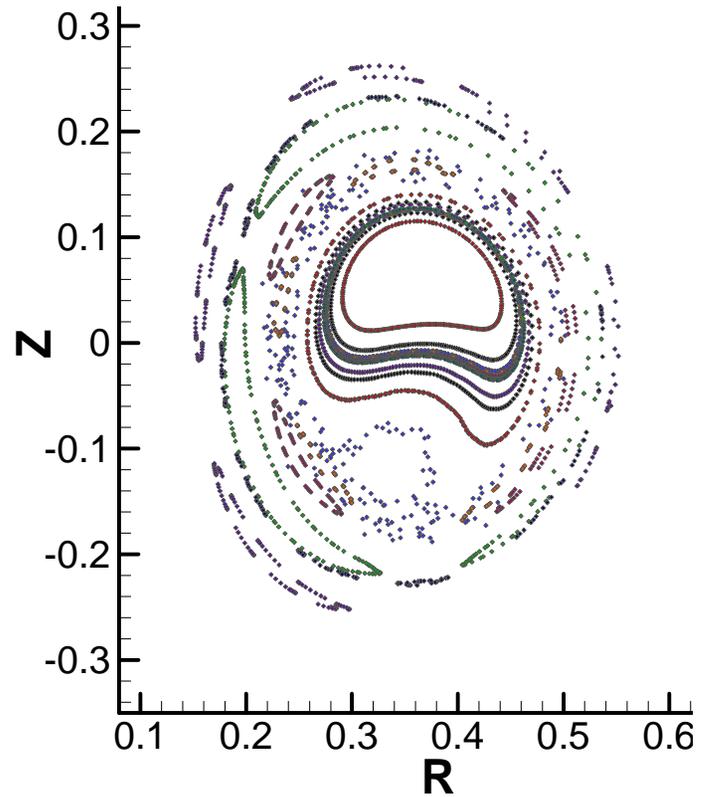


NIMROD 129  $\mu\text{s}$

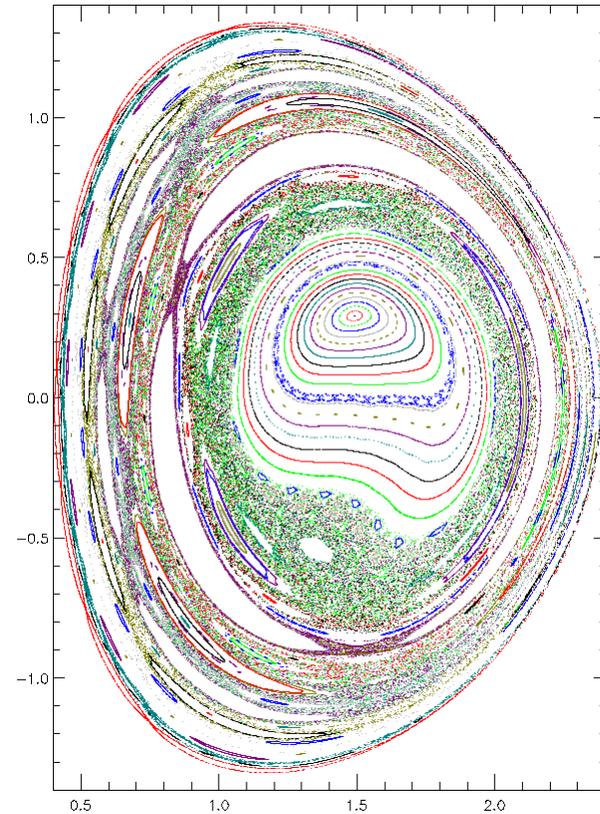


M3D 38  $\mu\text{s}$

# Poincaré Plots

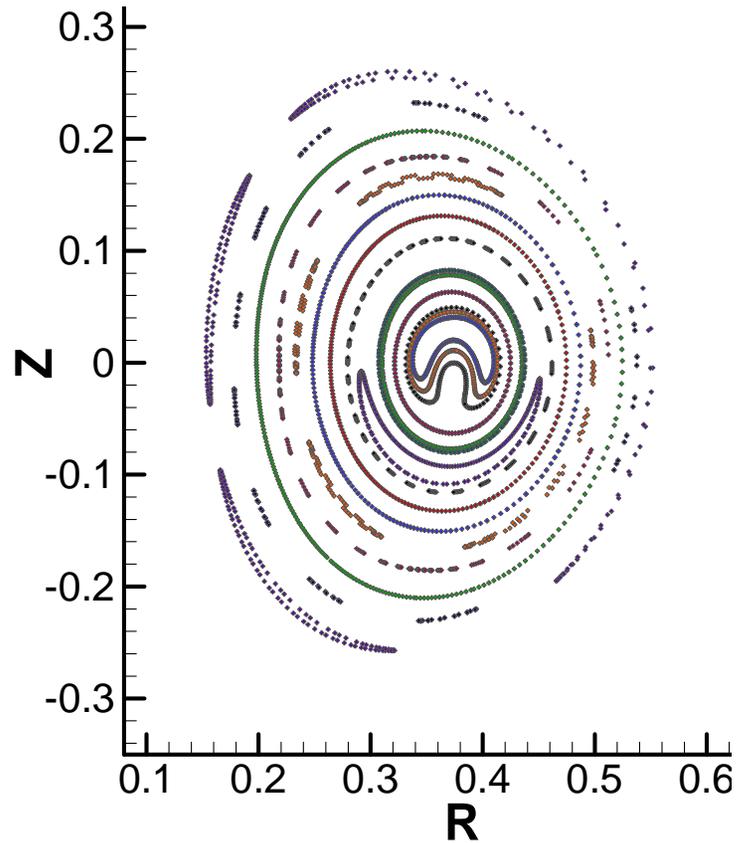


NIMROD 163  $\mu\text{s}$

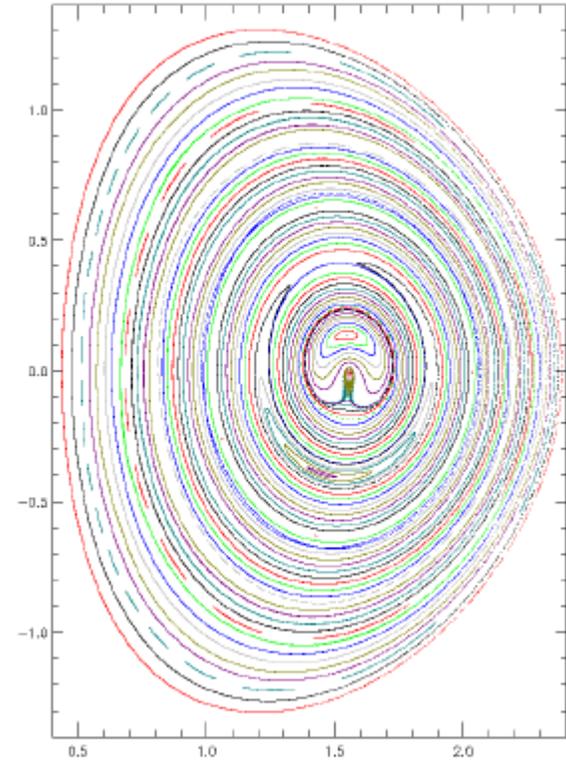


M3D 74.5  $\mu\text{s}$

# Poincaré Plots

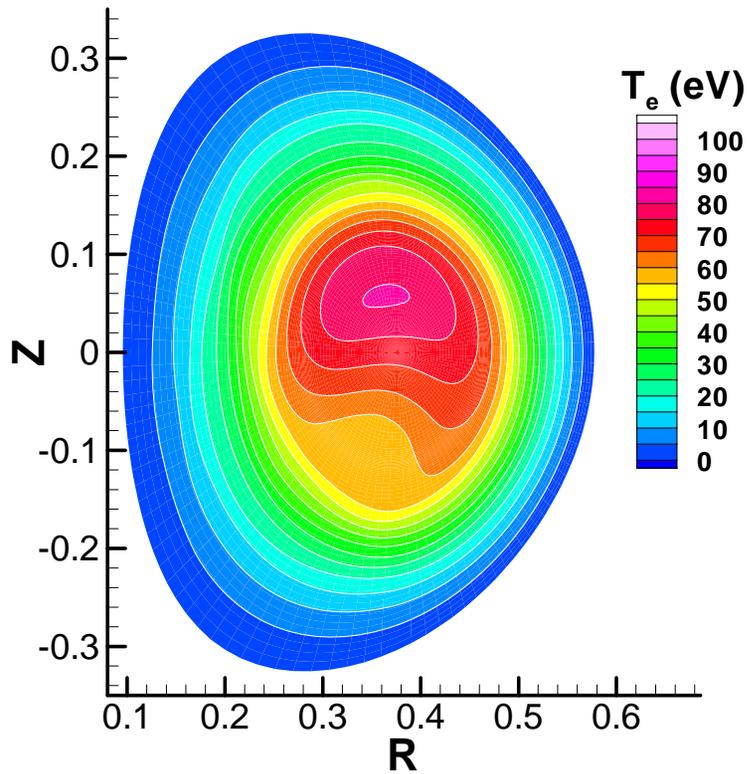


NIMROD  $200 \mu\text{s}$

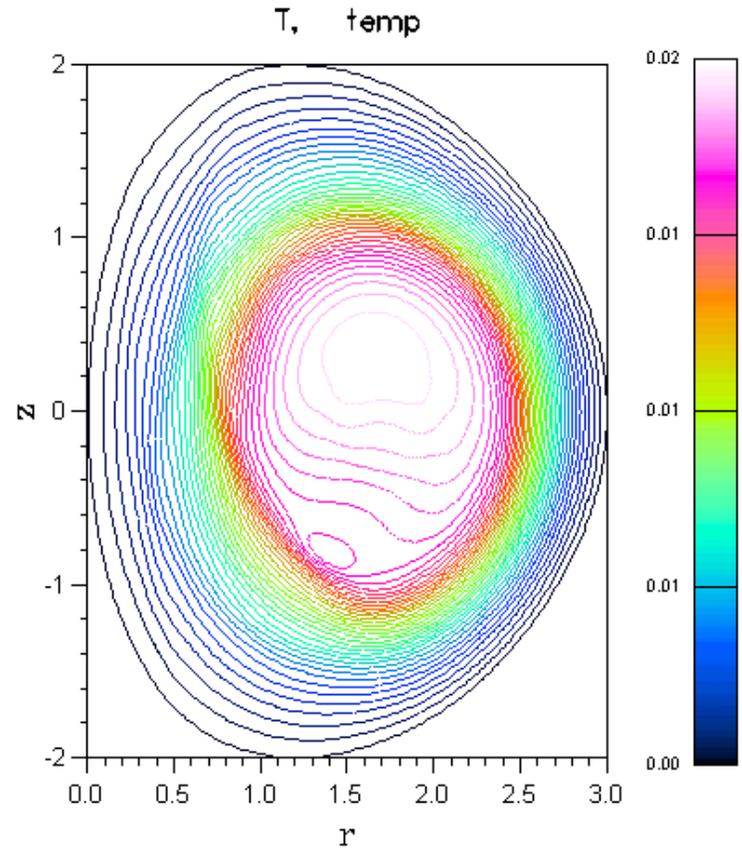


M3D  $115.3 \mu\text{s}$

# Temperature Contours

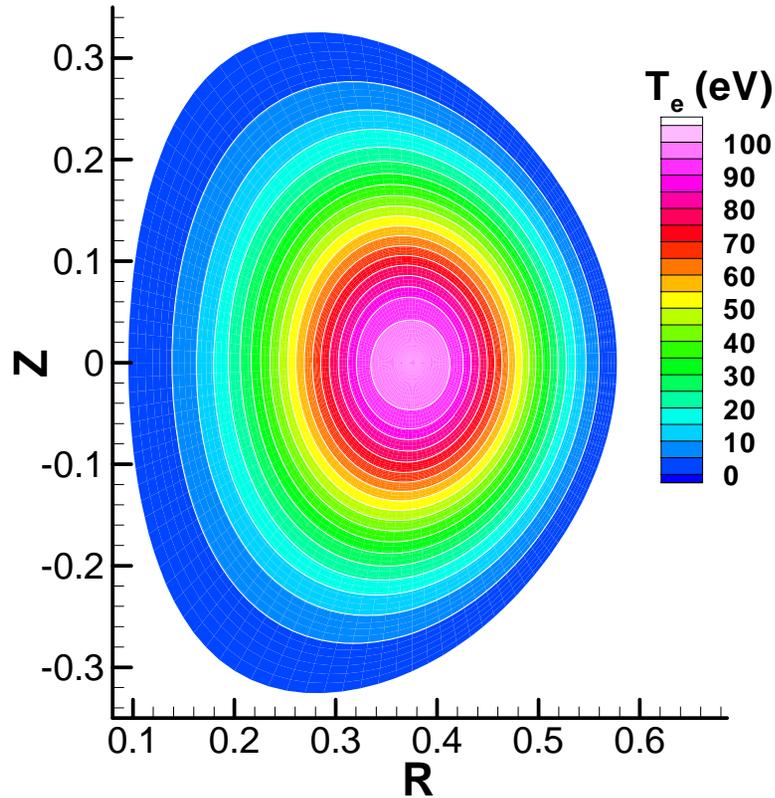


NIMROD 163  $\mu\text{s}$

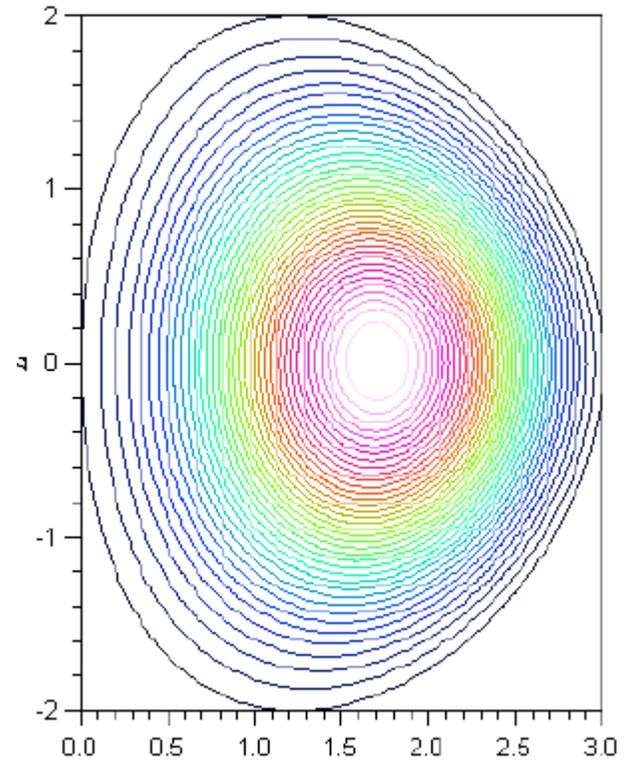


M3D 74.5  $\mu\text{s}$

# Temperature Contours



NIMROD  $200 \mu\text{s}$



M3D  $115.3 \mu\text{s}$

# Conclusions

- M3D and NIMROD are now in substantial agreement on the nonlinear CDX sawtooth benchmark.
- The next important step is validation: running a more physically accurate test case to try to achieve better agreement of both codes with experimental results.