

Calculations of Linear Two-Fluid Plasma Response to Nonaxisymmetric Fields

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Understanding Time-Independent 3D Plasma Response Requires Non-Ideal Physics

- **When is response ideal? When is it vacuum?**
- **Does the response vary smoothly between these two limits?**
- **If so, what is the relevant dimensionless parameter? $\tau_R \Omega$?**
- **To what extent are other effects important?**
 - Viscosity, two-fluid, FLR, etc..
- **How can this inform our interpretation of experimental results?**

Progress in 3D Response Since Last Meeting

- **Implementation of time-independent method**
 - Response to Boozer's concerns
- **Better understanding of results, especially in two-fluid case**
- **Application to several DIII-D experiments**
 - Lanctot, Mordijck, Buttery

Time-Dependent Method

- $\mathbf{B}(t) = \mathbf{B}_0 + \mathbf{B}_1(t)$
 - \mathbf{B}_0 is the axisymmetric equilibrium field
 - $\mathbf{B}_1(0)$ is the “vacuum field” from non-axisymmetric coils (I-coils).
- **Conducting-wall boundary condition**
 - $\mathbf{n} \cdot \mathbf{B}$ is held constant in time on simulation domain boundary (approximately vacuum vessel).
- **Simulation is time-advanced until the steady-state is reached.**
- **Final \mathbf{B}_1 is applied field + plasma response.**
- **Caveats:**
 - Equilibrium must be stable (otherwise response is ∞)
 - More computationally intensive than time-independent calculation

Time-Independent Method

- **Time-dependent time-step:**

$$\frac{\mathbf{x}^{n+1} - \mathbf{x}^n}{\delta t} = A \cdot [\theta \mathbf{x}^{n+1} + (1 - \theta) \mathbf{x}^n]$$

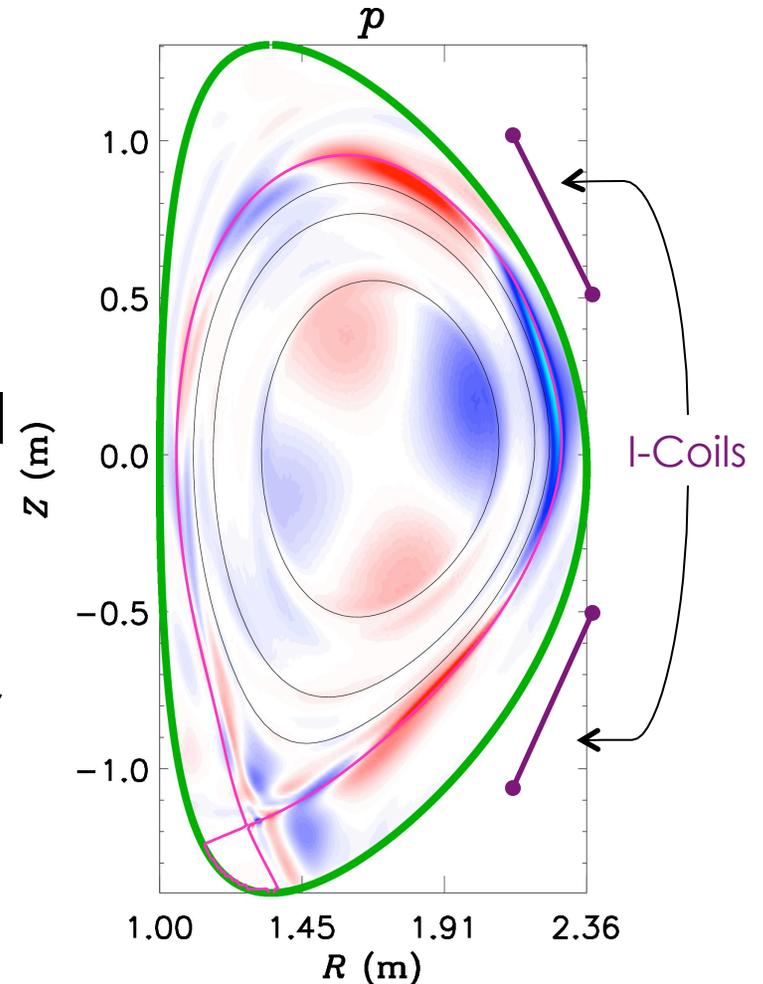
- **Time-independent solution ($\mathbf{x}^{n+1} = \mathbf{x}^n$) is**

$$A \cdot \mathbf{x}^{n+1} = 0 \quad (\text{BCs make some elements of RHS nonzero})$$

- **This is equivalent to our time-dependent calculation, if we remove d/dt terms and set $\theta = 1$**
- **Caveats:**
 - Doesn't work with semi-implicit methods
 - Only makes sense for linear calculations
 - System may be very poorly conditioned
 - Says nothing about dynamics

Response Calculated For DIII-D Discharges

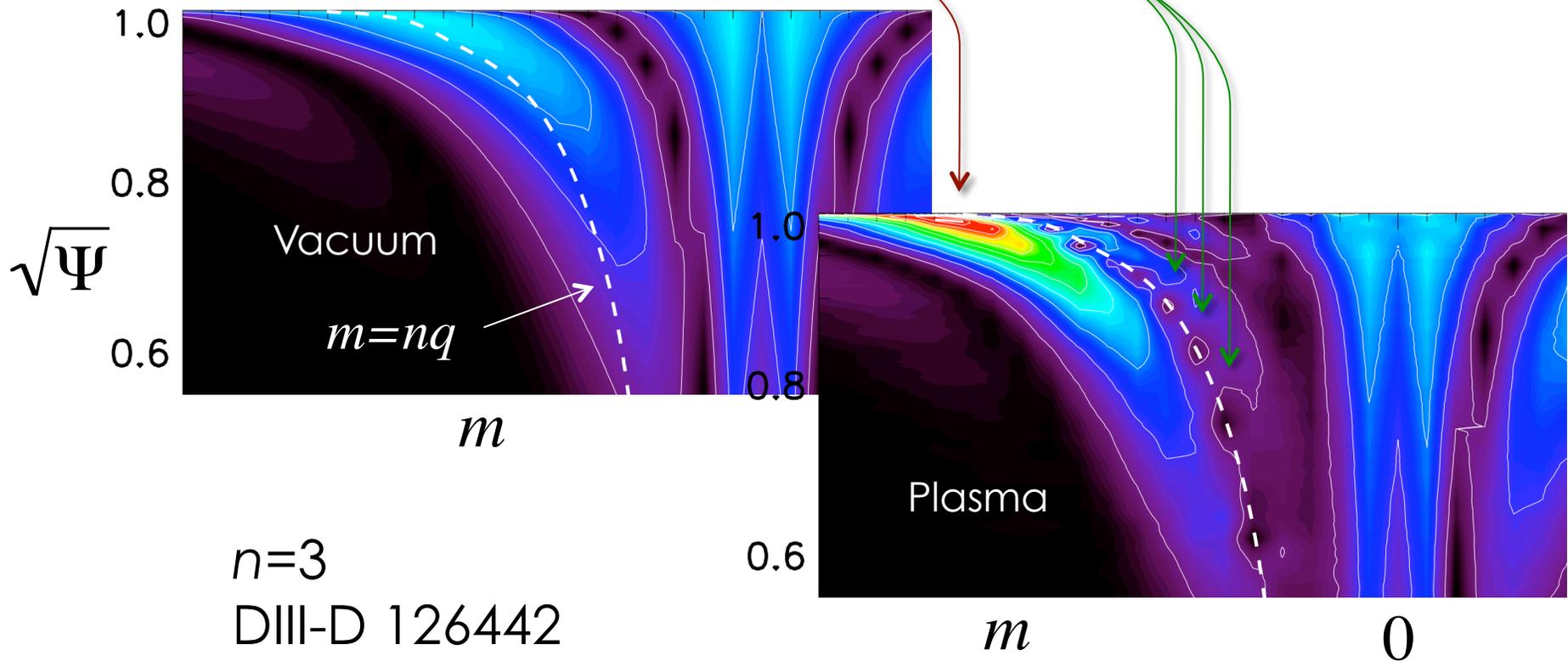
- Use EFIT reconstructions of DIII-D shots 135758--73
- Equilibrium is recalculated self-consistently using $\rho \Omega^2 = \alpha p_0$
- Numerical considerations require I-Coils to be outside wall (unlike experiment)
- Wall excludes plasma response, invalidating magnetics diagnostics



* Lanctot et al., *Phys. Plasmas* **17**:030701 (2010)

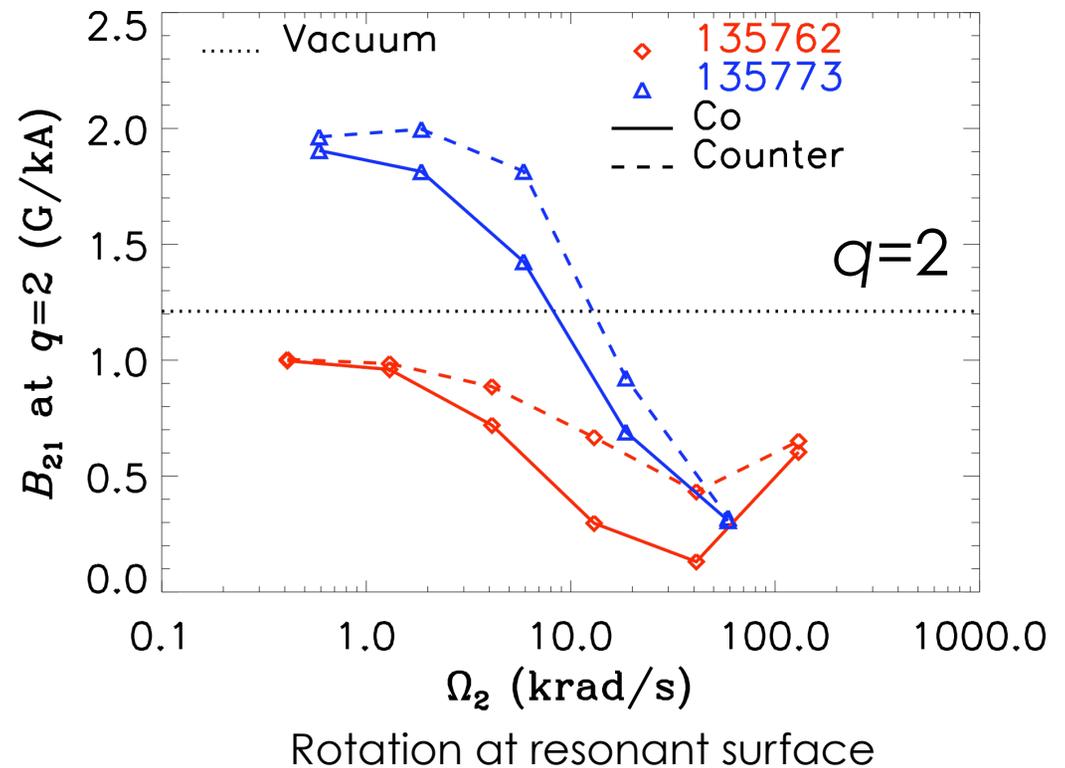
Plasma Screens Resonant Fields, Enhances Non-Resonant Fields

- Generally, two types of responses are seen
 - Screening
 - Kink excitation

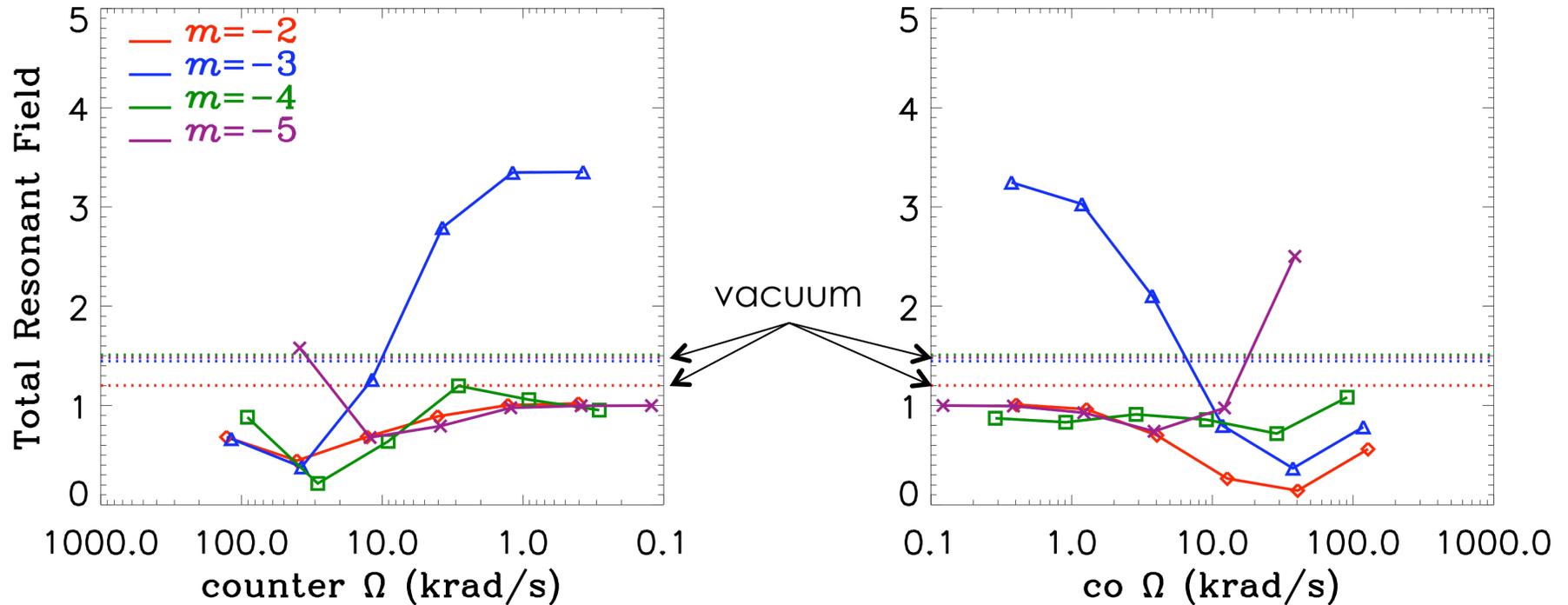


Rotation (Usually) Improves Screening; Response Depends on Beta

- Co-current rotation is found to screen better than counter-current (even in resistive 1F model)
- In low- β case, plasma screens even without rotation
- In high- β case without rotation plasma amplifies resonant fields

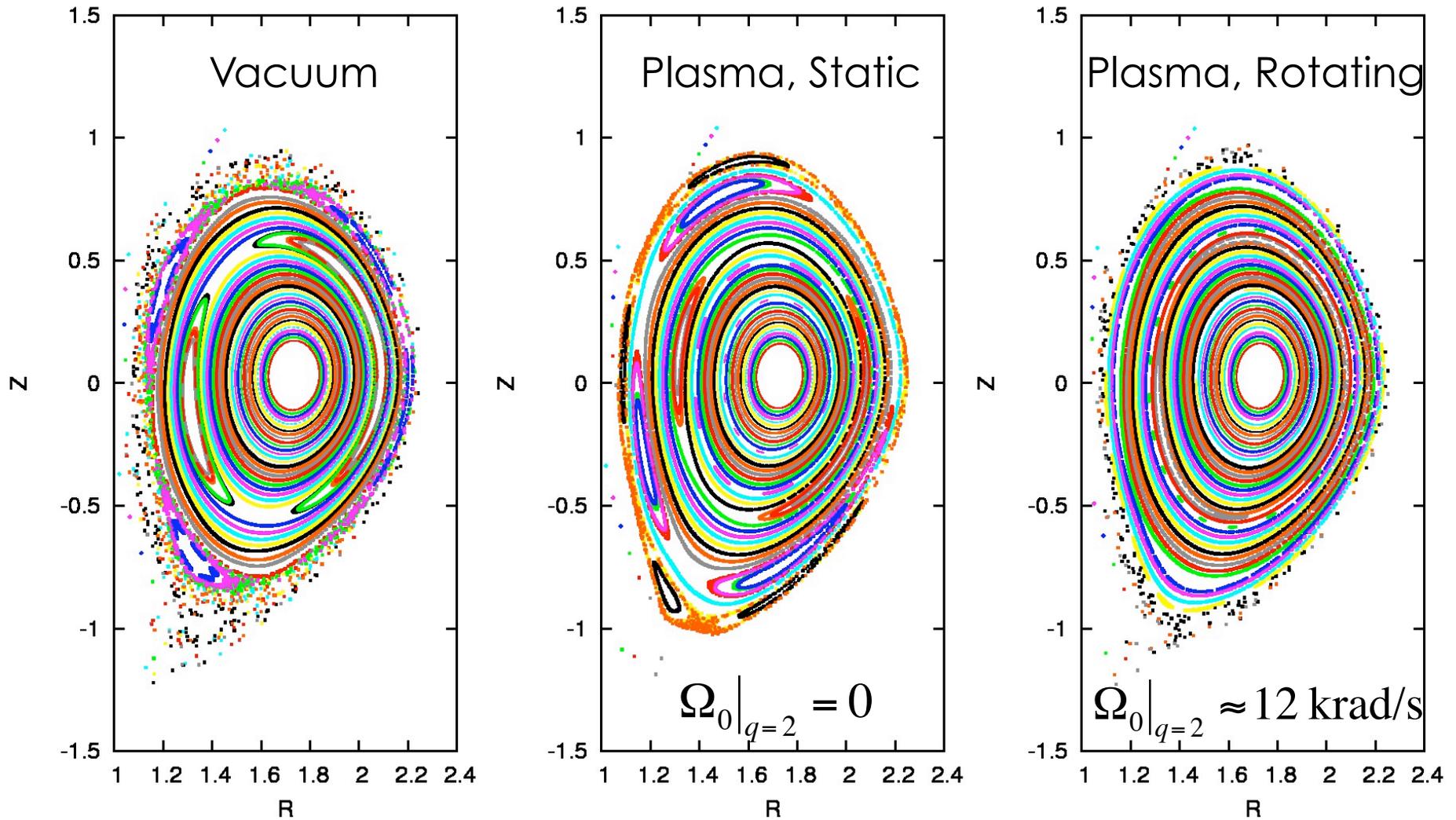


Shot 135762 Responds Most Strongly at $q=3$ Surface



- **Rotation generally suppresses resonant response**
- **Strong rotation enhances resonant response near edge (especially $q=5$)**

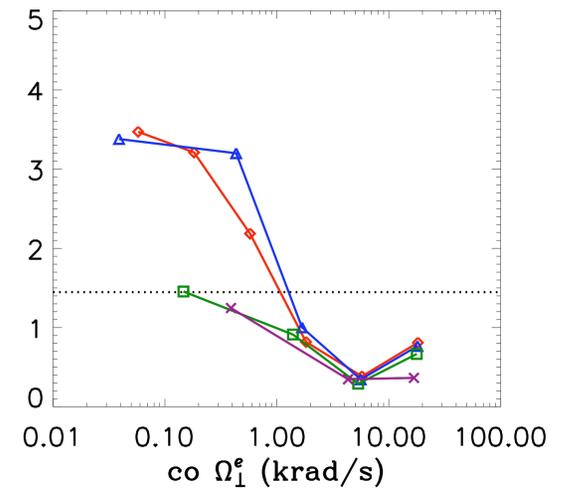
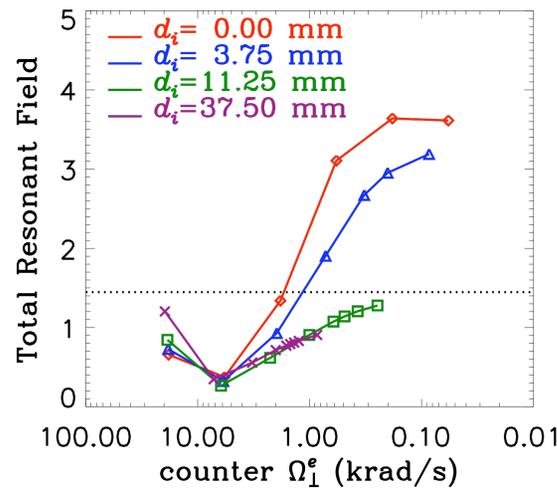
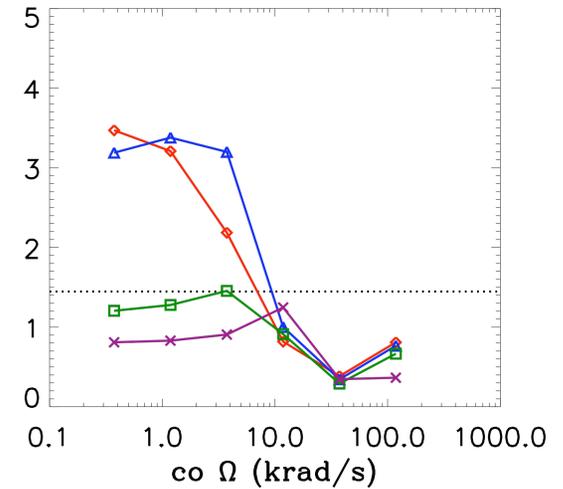
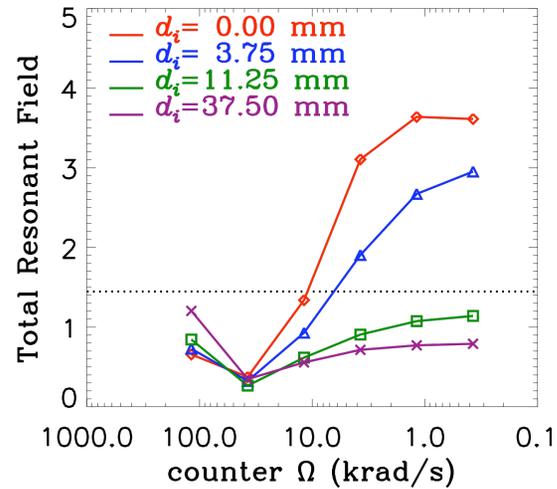
Rotation Improves Core Screening; But Stochasticizes Edge



Core Penetration Depends on Perpendicular Electron Velocity

- 2F terms introduced by raising d_i
- Core penetration correlates better with **perp. electron rotation (bottom)** than with **ion rotation (top)**.

$q=3$



What is “Perpendicular” Electron Velocity?

- **For incompressible flow, $\mathbf{v} = R^2 \omega(\psi) \nabla \varphi + \lambda(\psi) \mathbf{B}$**
 - Note that $\omega(\psi) \neq \Omega = \mathbf{v} \cdot \hat{\varphi} / R$ unless $\lambda(\psi) = 0$
- **The perpendicular angular velocity is defined as**

$$\Omega_{\perp} = \frac{\mathbf{v} \cdot \mathbf{B} \times \nabla \psi}{R |\mathbf{B} \times \nabla \psi|}$$

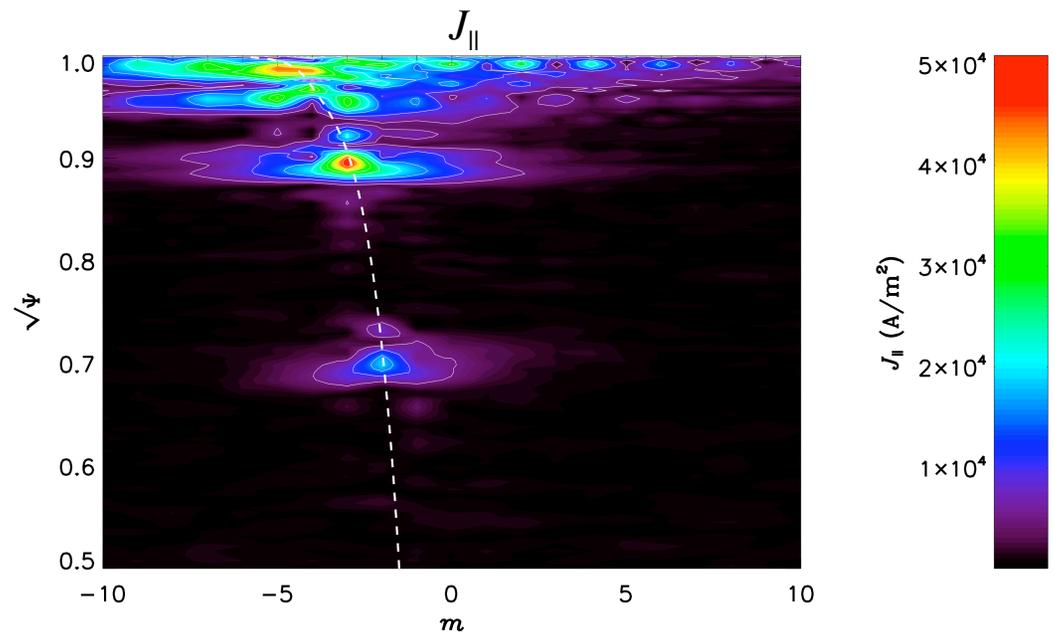
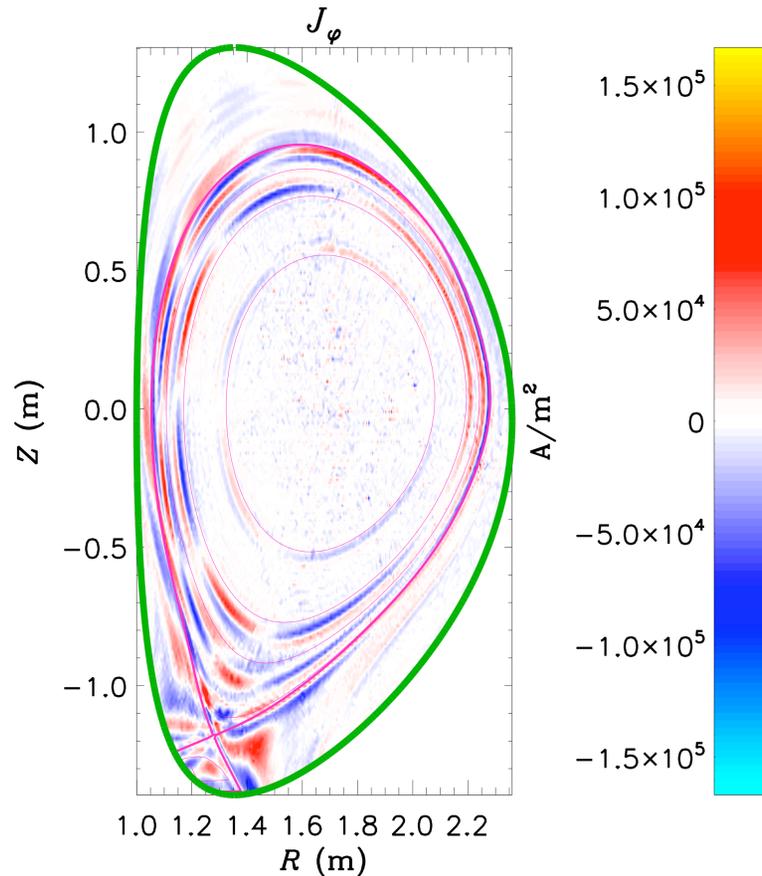
$$= \frac{|\nabla \psi \times \nabla \varphi|}{|B|} \omega(\psi)$$

- **From radial force balance: $\omega(\psi) = \phi'(\psi) + \frac{p'(\psi)}{nq}$**

$$\omega^e = \omega^i + \omega_*^e - \omega_*^i$$

- **Ω_{\perp} vanishes wherever ω vanishes, but also at nulls**
 - If Ω_{\perp} is the relevant quantity, toroidal rotation is less effective at screening near x-point

Parallel Currents Persist in Steady State



- These are from time-independent calculations, without rotation!
- Currents do not affect energy balance at first order (currents are periodic in ϕ)

Summary and Conclusions

- **M3D-C1 can calculate linear two-fluid response with realistic values of resistivity, rotation, and perpendicular transport**
 - Time-independent and time-dependent methods agree
- **Time-independent parallel currents do exist in resonant layers**
 - *This is true even without rotation*
- **Screening is more complicated than ratio of rotational to resistive time-scales**
 - Depends crucially on proximity to marginal stability
 - Viscosity and thermal diffusion inhibit resonant response (both screening and resonant field amplification)
- **Perpendicular electron velocity appears to be the relevant rotation quantity for screening**

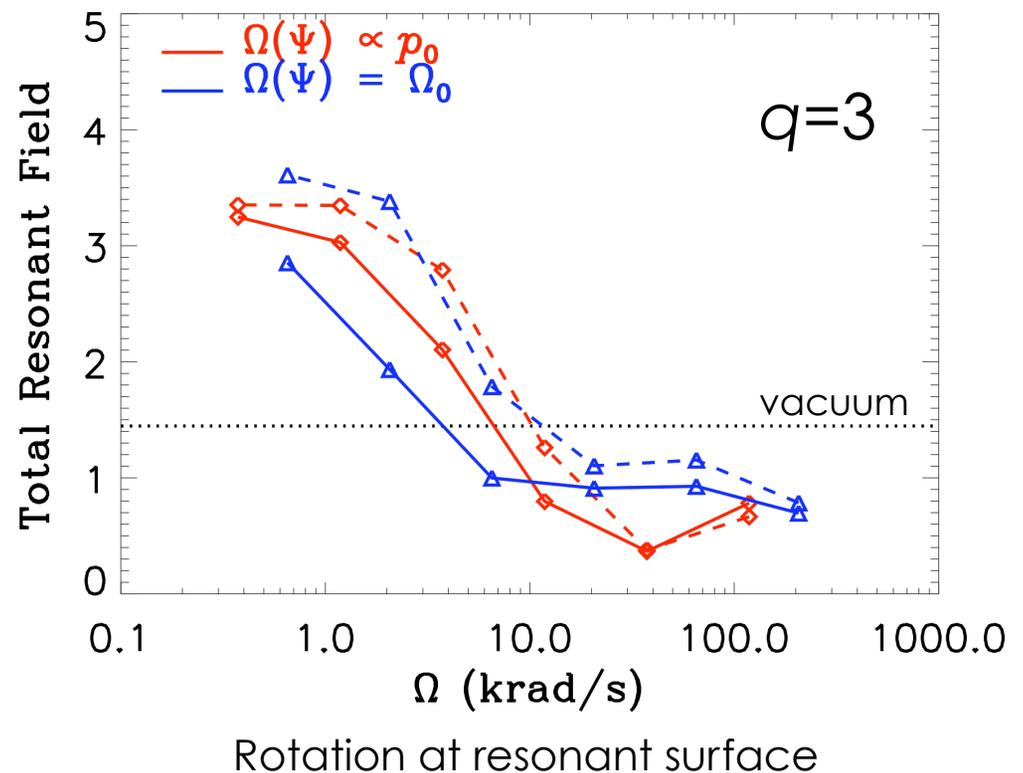
Future Work

- **Torque calculation**
 - General NTV models can be used to calculate torque generated by linear perturbation *a posteriori*
 - Torque can be calculated self-consistently using Braginskii-like NTV models
- **Nonlinear calculations**
 - M3D-C1 is now capable of 3D nonlinear calculations (see poster 1P24)
 - Unclear how to extend time-independent method to nonlinear model
- **How do we use this to optimize coils? Optimize for what?**

Extra Slides

Rotation Magnitude, Not Shear, Is Responsible For Screening

- **Red**: rotation profile proportional to pressure
- **Blue**: uniform rotation throughout plasma (*i.e.* no shear)
- **Rotation screens, even when there is no shear.**



Rotation-Driven Edge Response Depends on Ion Rotation (Not Electron Rotation)

- Edge mode correlates better with **ion rotation (top)** than with **perp. electron rotation (bottom)**

- Not “penetration,” but proximity to marginal stability.

$q=5$

