

Simulation of MSE signals, as a diagnostic tool for MHD instabilities

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Abstract

Mode identification of MHD instabilities is often based on external signals, typically, by Mirnov loops placed at some distance from the plasma. With a sufficient number of coils it is possible to identify the dominant poloidal and toroidal harmonics at the plasma boundary. The internal structure can be determined from the electron cyclotron emission^a. Recently^b, the motional Stark effect, MSE, diagnostic has proved to be a complementary alternative, as it measures the local field perturbation within the discharge. The predicted mode structure for global instabilities is easily obtained from codes, such as PEST, from which, we can simulate the expected MSE signal. We examine a variety of plasma pressure and safety-factor profiles and geometry to determine the expected MSE signal. Comparisons with experiment will be presented.

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^aOkabayashi et al. Nuclear Fusion Vol. 38, (1998) 1149

^bJayakumar et al. BAPS Vol. 48(2003) QP1.045



Approach

- Equilibrium from EFIT/TRANSP
 - Best fit to experiment
- Stability analysis using ideal MHD model
 - PEST : Determine eigenvalue and eigenfunction

$$\omega^2, \xi(\mathbf{x})$$

- Determine perturbed field

$$Q = \nabla \times \mathbf{x} \times B$$

- Determine MSE signal



The fluctuation in the MSE signal is determined using the equilibrium field, B , and perturbed field, Q .

$$Q \equiv \nabla \times (\mathbf{x} \times B) = \tilde{B}$$

$$Q_q \equiv \frac{Q \cdot \nabla q}{|\nabla q|}$$

$$Q_y \equiv \frac{Q \cdot \nabla y}{|\nabla y|}$$

$$Q_f \equiv \frac{Q \cdot \nabla f}{|\nabla f|}$$

$$\text{Pitch} \equiv P = \frac{B_z}{B_f} = \frac{B_q \text{ (mid-plane)}}{B_f}$$

$$\text{Pitch perturbation} \quad dP = \frac{B_q + Q_q}{B_f + Q_f} - \frac{B_q}{B_f}$$



Normalizing the Signal-I

- PEST provides a linear solution with an arbitrary scale factor – to fix its value:
- Determine the magnetic scalar potential in the vacuum region
- Use the plasma surface value and apply Green's 2nd identity to compute the solution at the Mirnov probe location
- Match the field to the experimental measurement to normalize the amplitude of the perturbation



Normalizing the Signal-II

- A crude approximation
- Use the cylindrical solution

$$B(r) = B(r_s) \left(\frac{r_s}{r} \right)^{(m+1)}$$

$$r_s = R(\text{surface}) - R(\text{axis})$$

$$r = R(\text{loop}) - R(\text{axis})$$

m = poloidal mode number



Reference to work presented at APS DPP'03

Observation of MHD modes using the Motional Stark Effect Diagnostic

R. J. Jayakumar (1), A.M. Garofalo (3), H. Reimerdes (3), T. Rhodes (5), S. Allen (1),
M. Austin (6), M. Makowski (1), M. Okabayashi(4), T. Strait (2) and A.D. Turnbull (2)

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(4) Princeton Plasma Physics Lab, (5) University of California Los Angeles,
(6) University of Texas at Austin



APS/DPP 03



Summary

- An MHD instability has been observed using the Motional Stark Effect (MSE) diagnostic on the DIII-D tokamak.
 - Possibly for the first time, the local oscillating poloidal field component in a hot plasma device has been measured.
 - First measurement of oscillating poloidal field component using MSE.
- The oscillating poloidal field component associated with a Resistive Wall Mode (RWM) has been measured in a plasma where the mode was rotated using internal field coils.
- The radial profile of the amplitude of oscillation poloidal field is in general agreement with profile of the Electron Cyclotron Emission (ECE) oscillation amplitude.
- The peak flux surface displacement obtained from MSE observations is in agreement with expectation.

APS/DPP 03



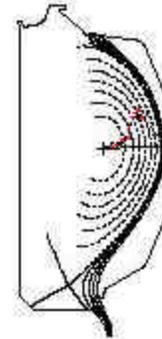
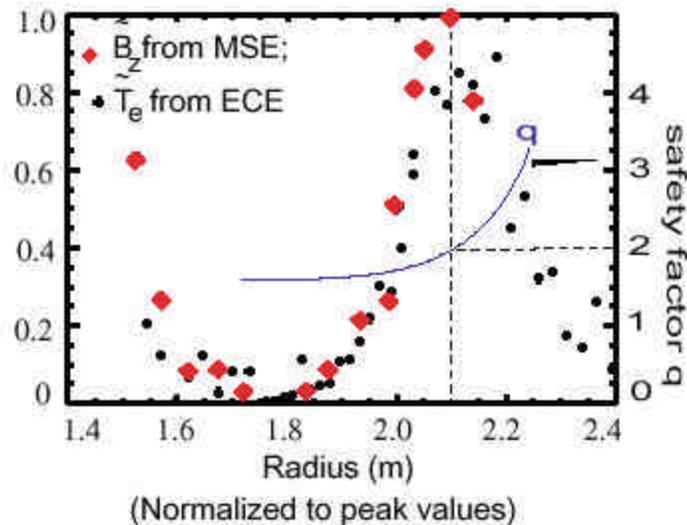
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Mode structure reconstructed from MSE

PROFILE OF THE OSCILLATION AMPLITUDE OF B_z (MSE) AND T_e (ECE) ARE SIMILAR



The oscillations are sum of $n=0$ and $n=1$ components.
(not separable)

$\tilde{T}_e \sim 85$ eV; $\tilde{B}_z \sim 55$ gauss (RMS)
Correspond to a peak flux surface
Displacement of ~ 0.55 cm

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Peak location is at the $q \sim 2$ radius.



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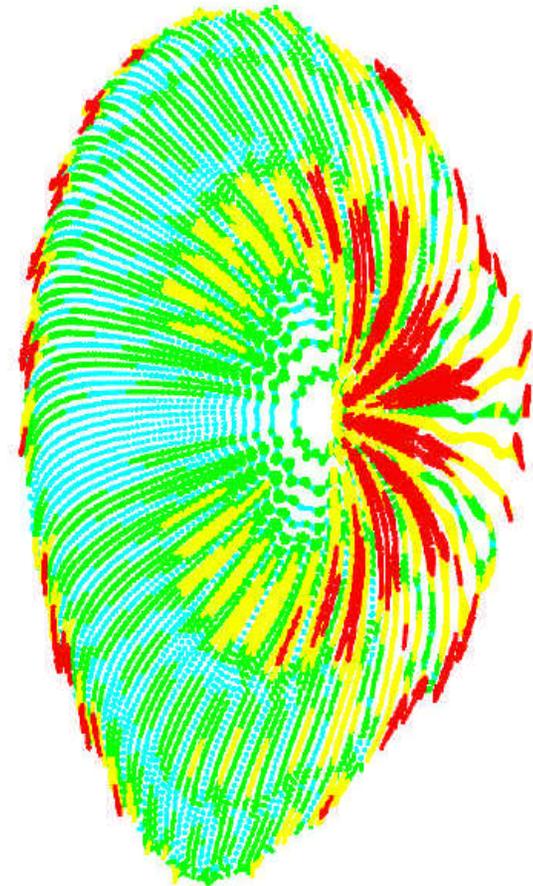
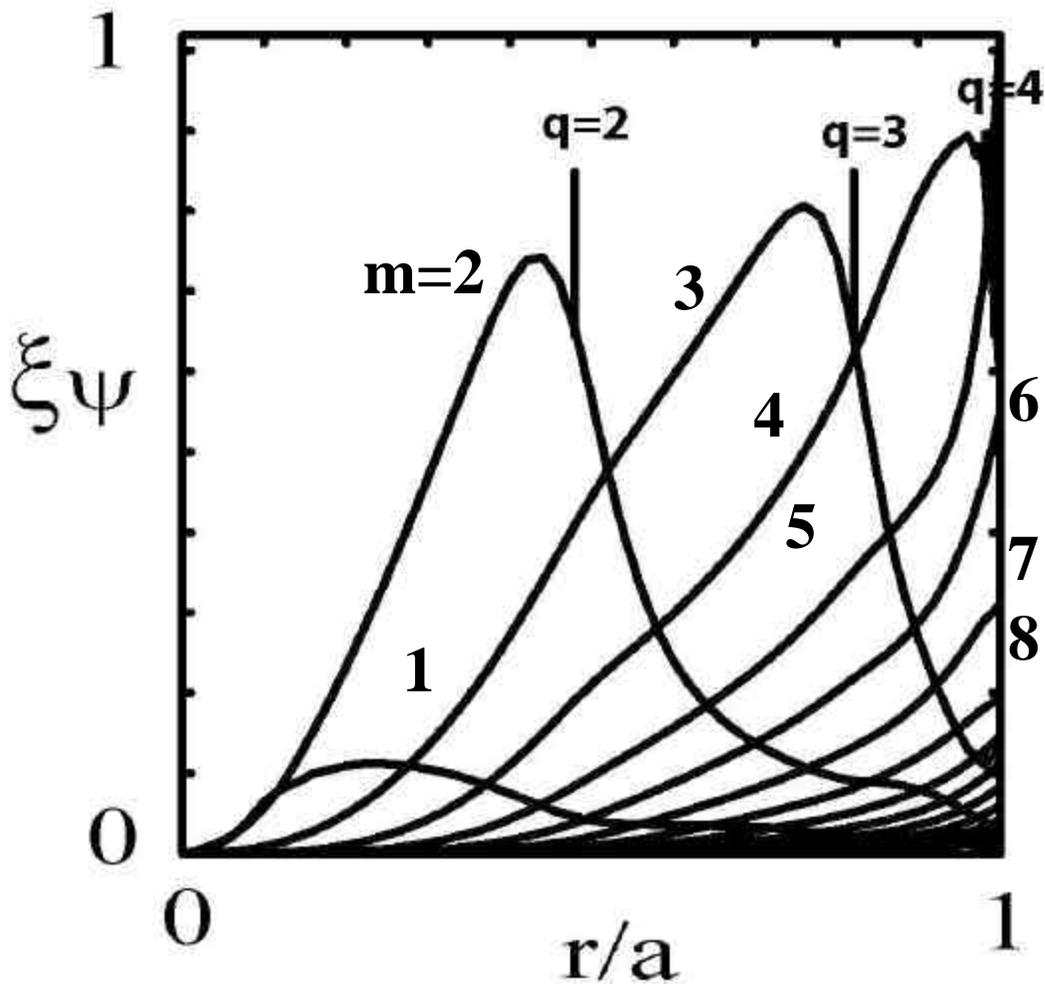


Outstanding issues for comparison of theory with experiment

- Theory
 - Normalization
 - Linear mode structure – nonlinear effects?
 - Accuracy of equilibrium reconstruction
 - Modeling of external boundary
- Experiment
 - Separation of $n=0$ and $n=1$ components
 - Noise level as high as $\sim 30\text{-}40\%$
 - Poloidal mode number identification
 - Radial resolution
 - Rotation



Unstable eigenfunction of displacement vector for DIII-D Shot:113550 at 1725 ms.

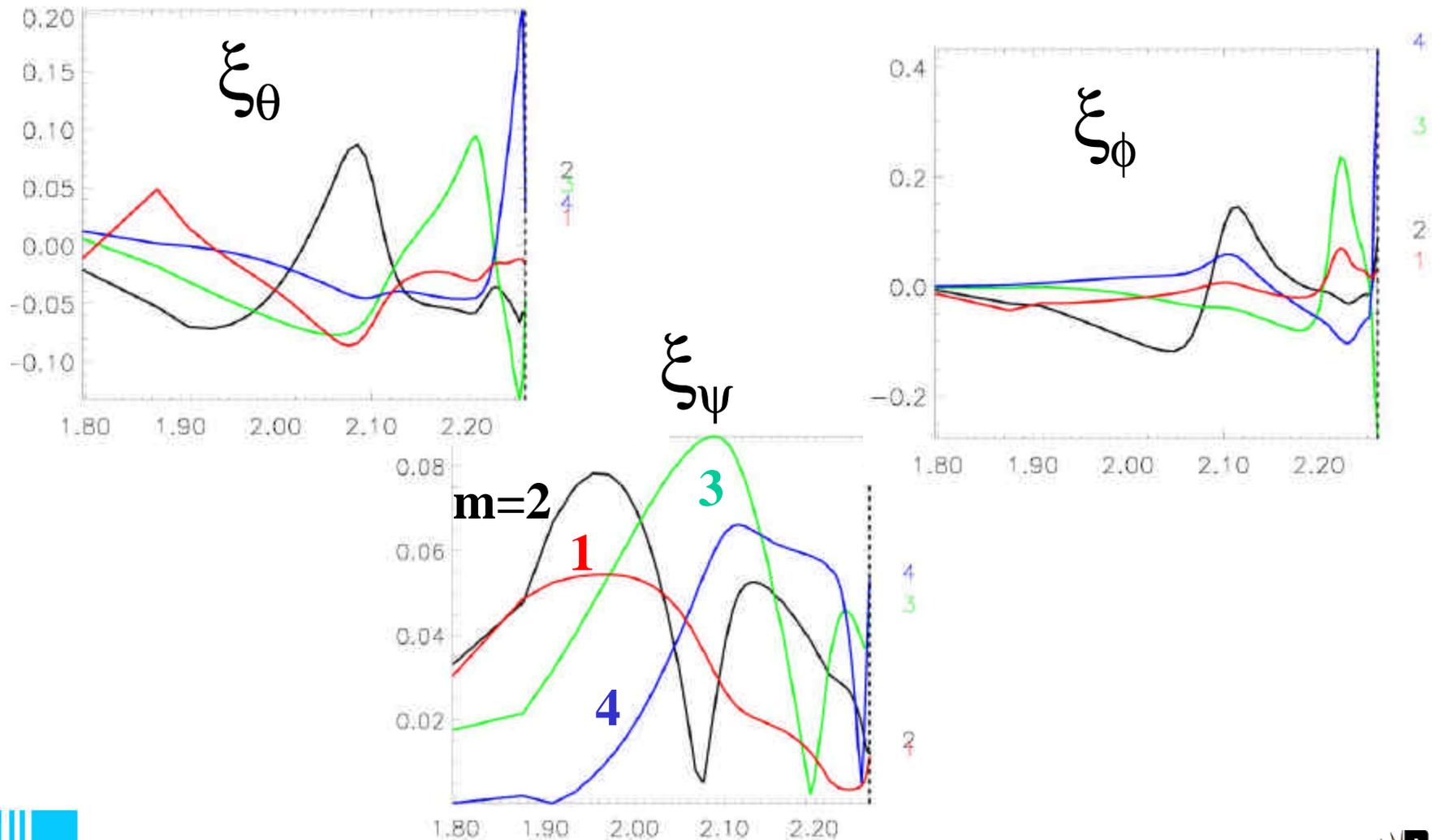


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Unstable eigenfunction of perturbed field

(higher harmonics not shown)

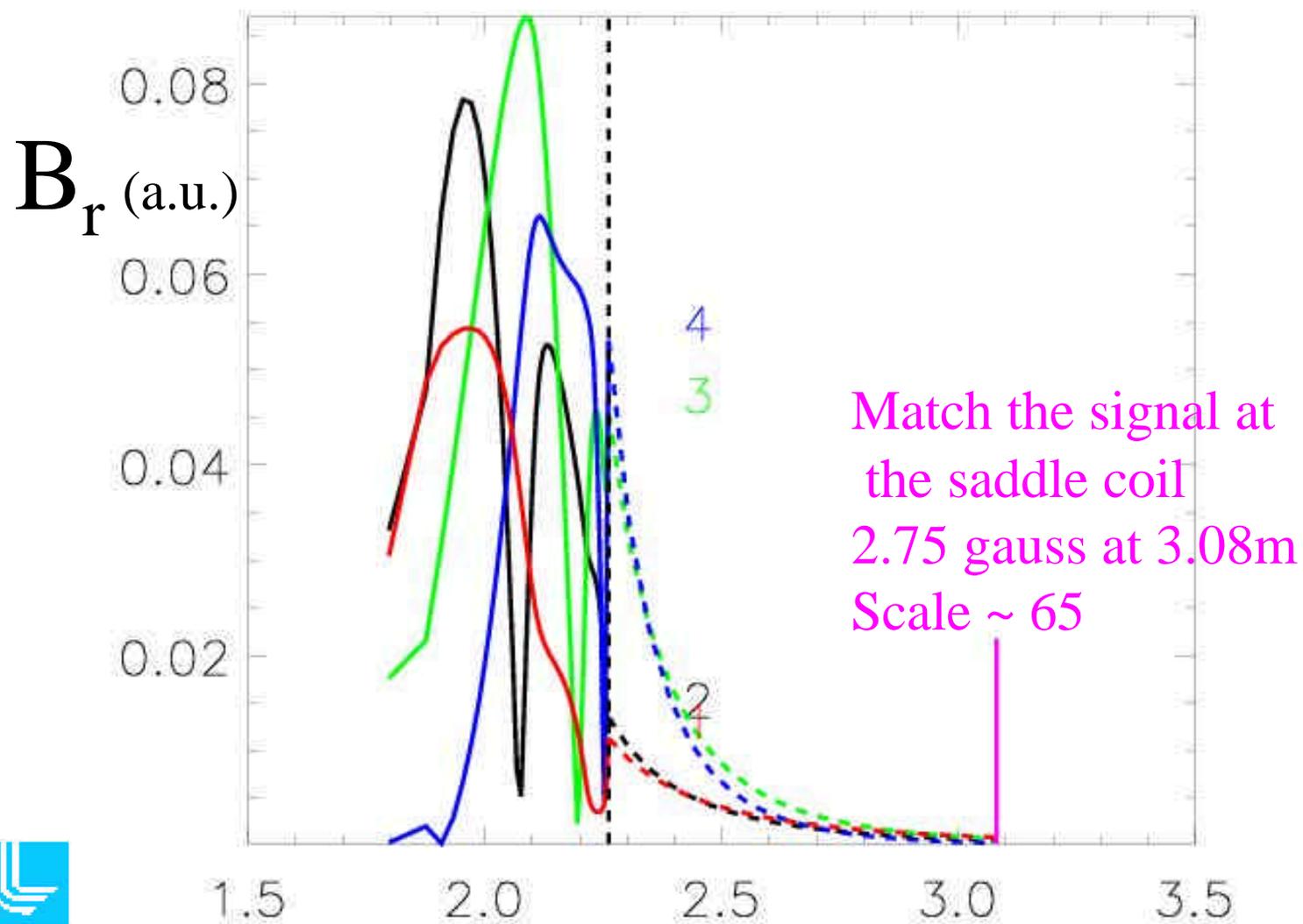


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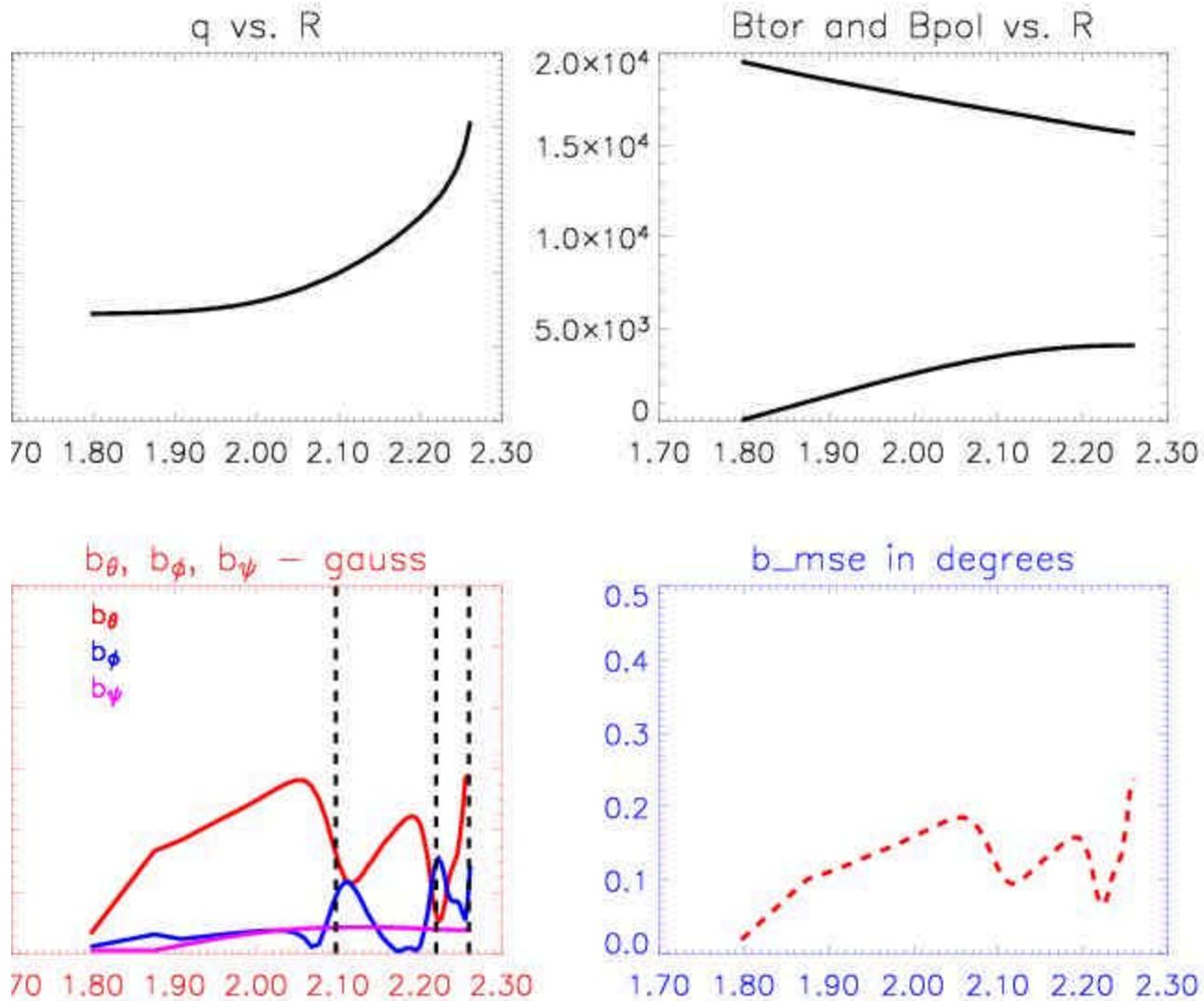
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Normalizing the Signal-III



Predicted MSE signal

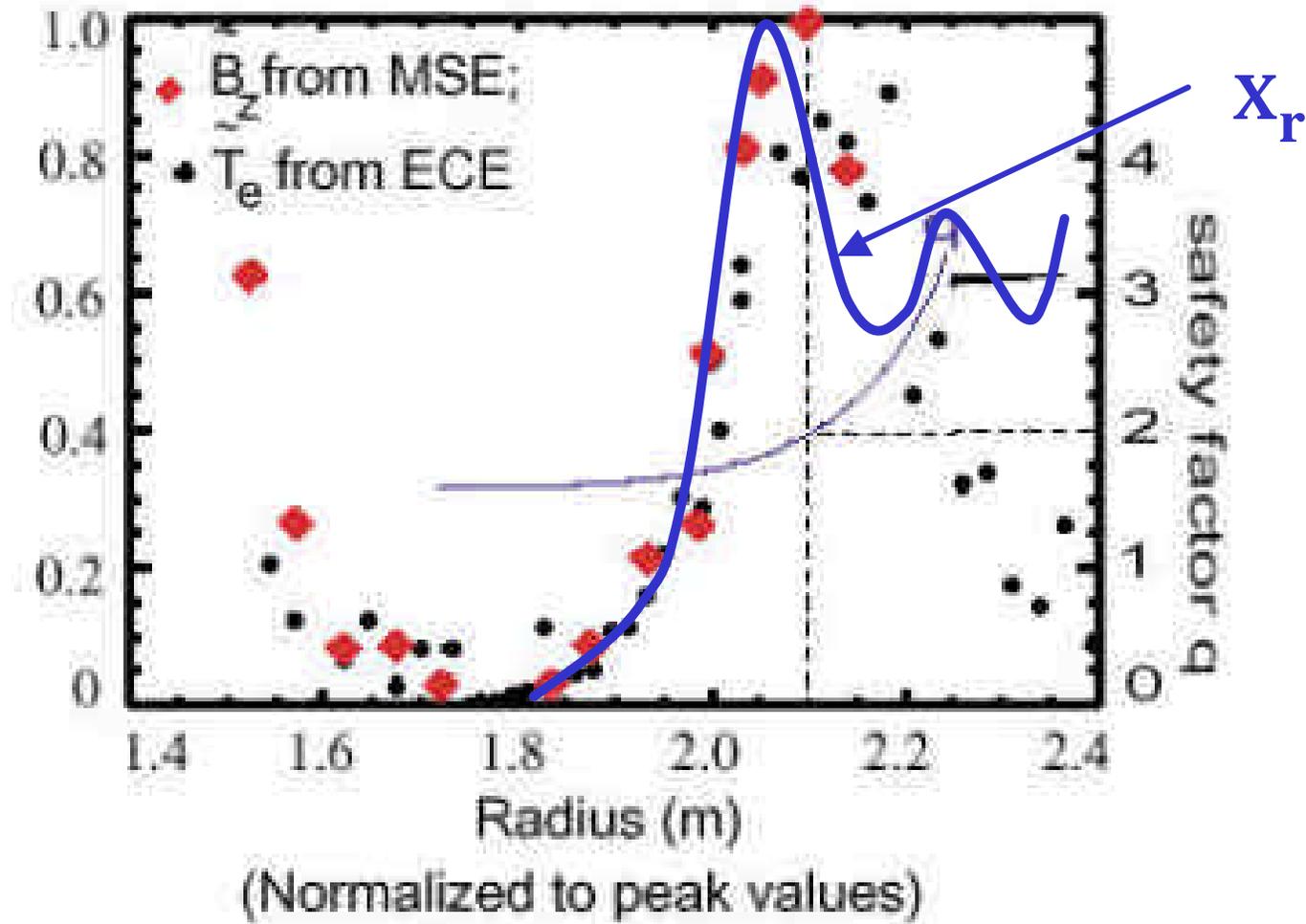


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Predicted Displacement

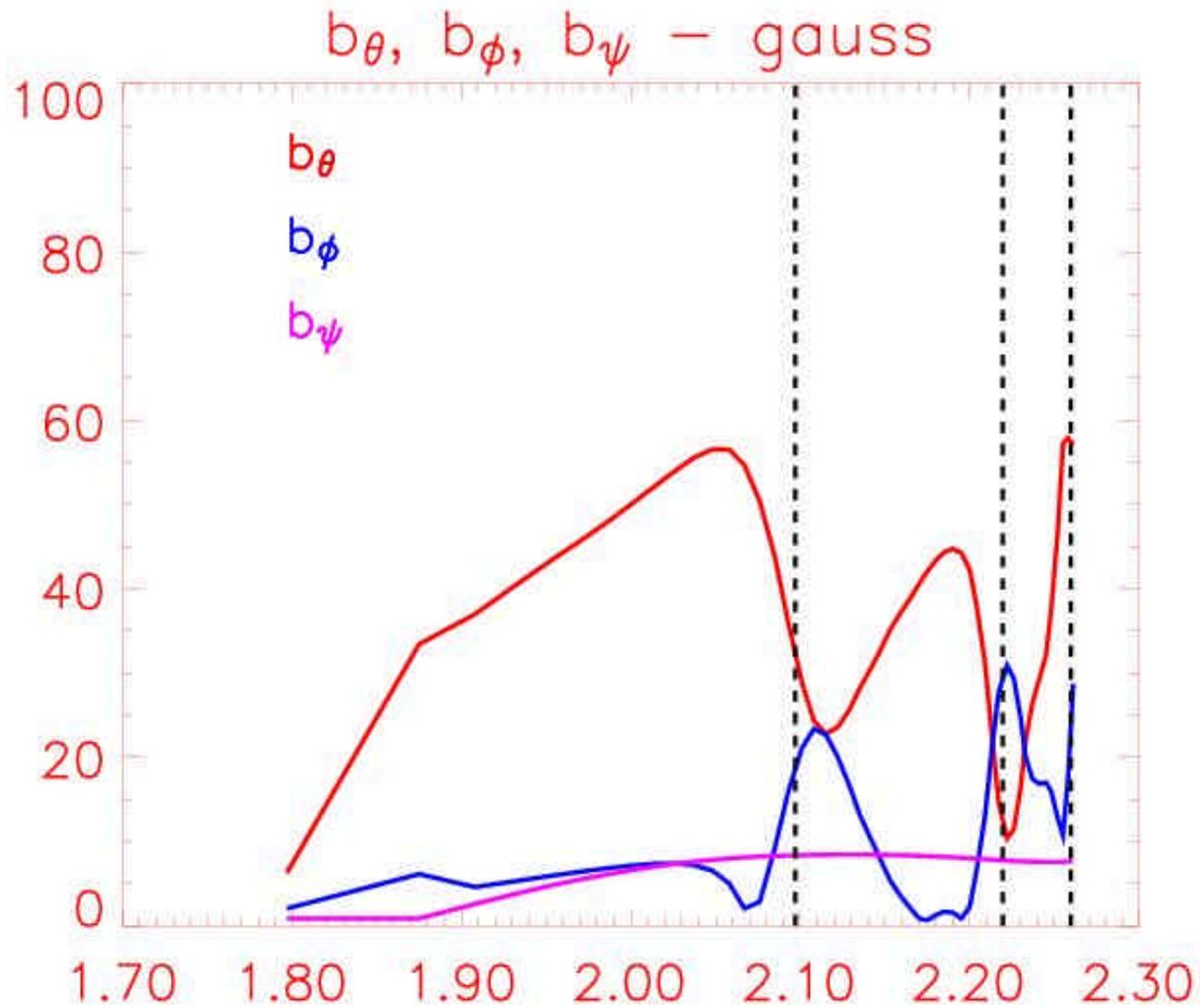


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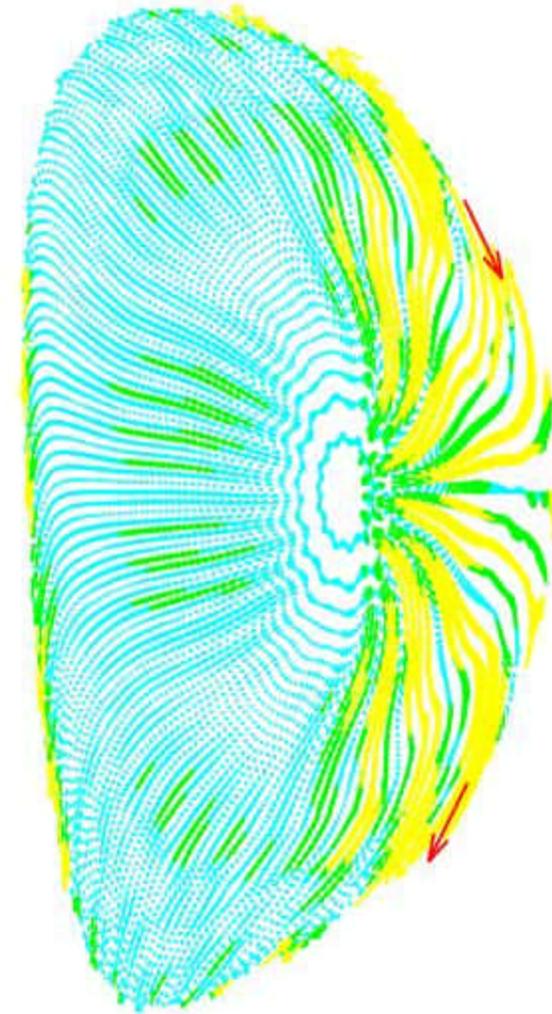
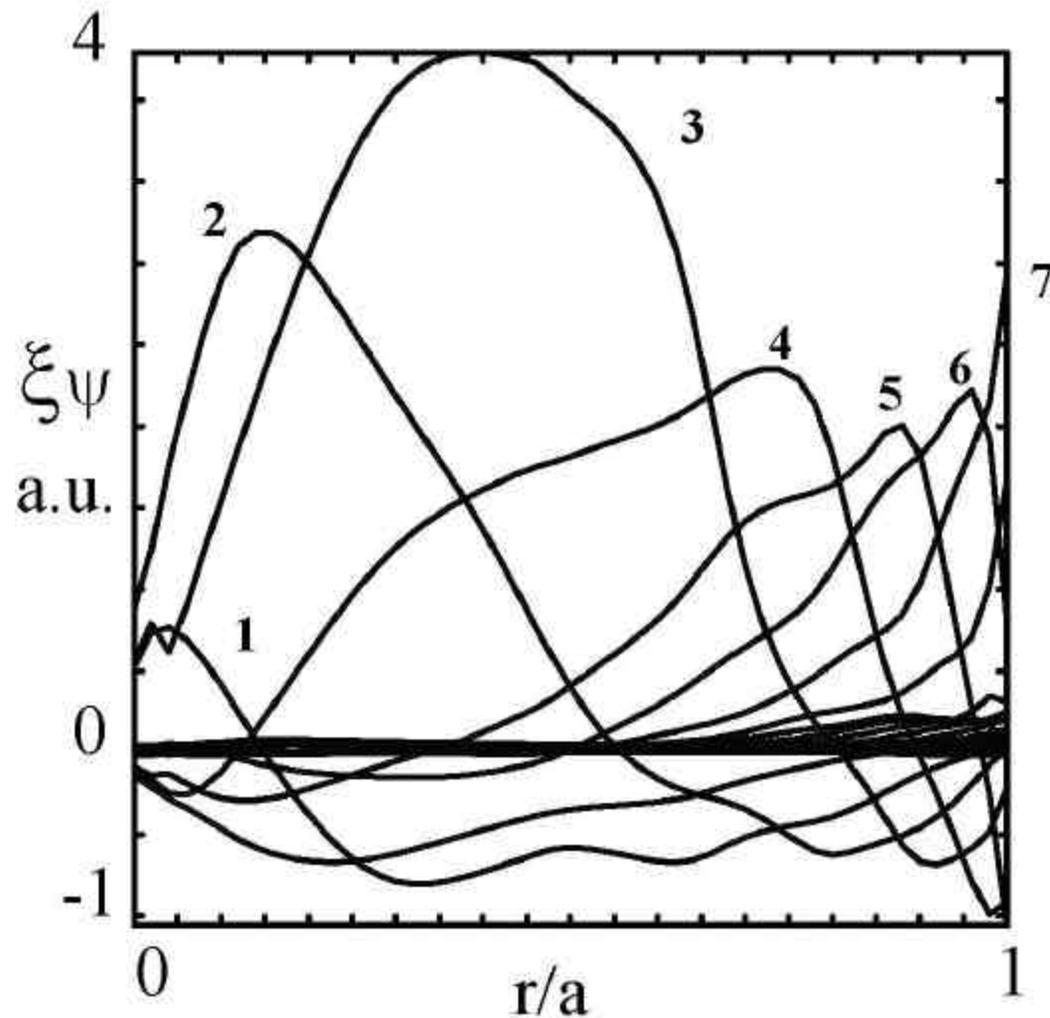
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Radial distribution of perturbation



Unstable eigenfunction of displacement vector for ITER – $b_N = 3.4$

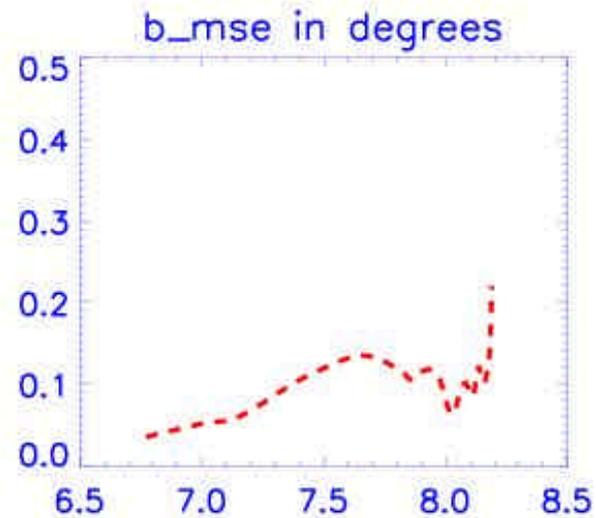
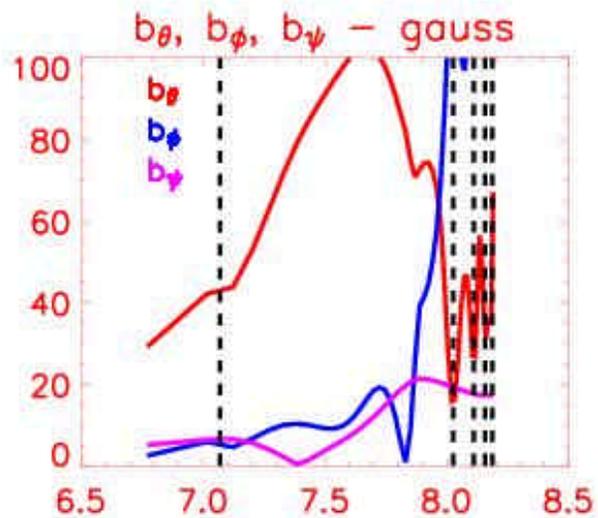
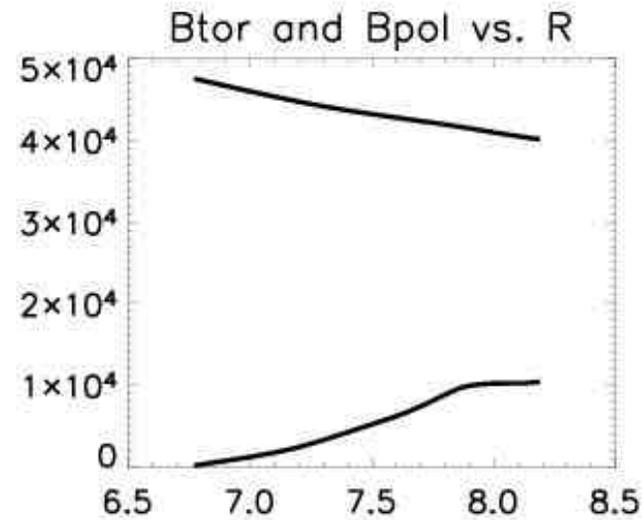
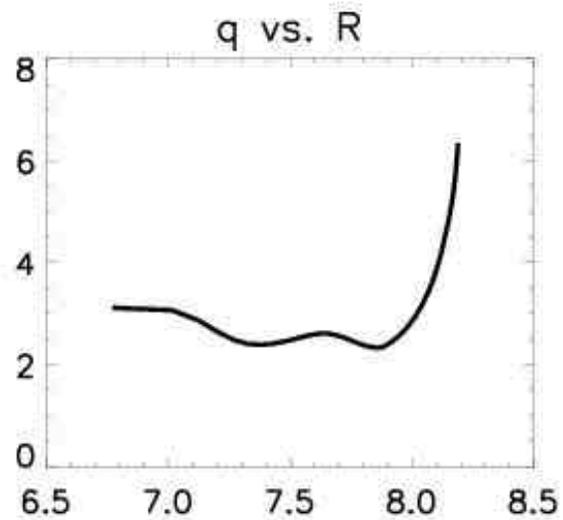


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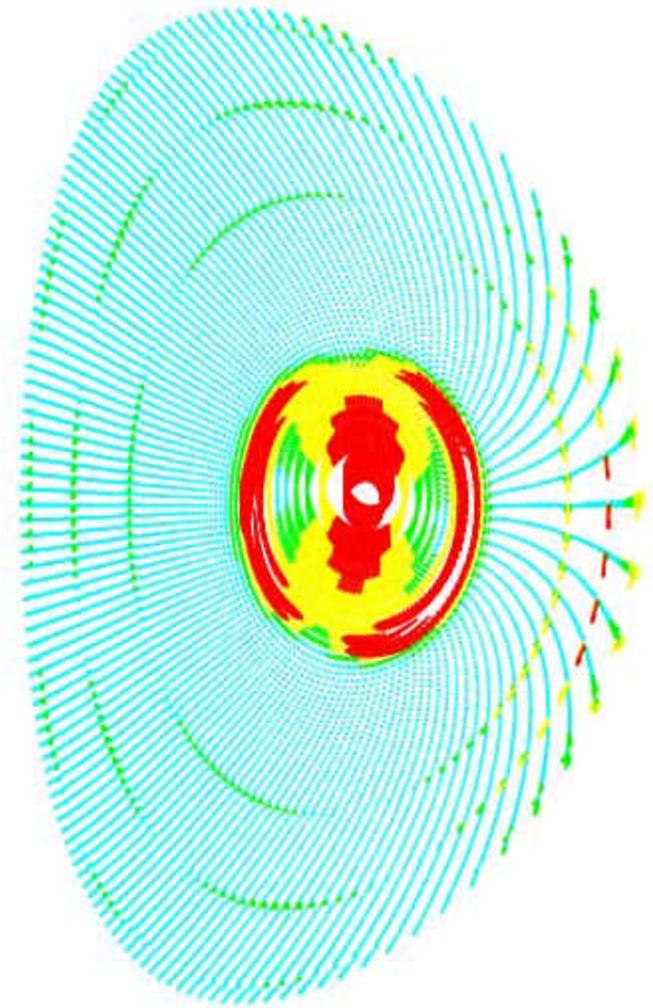
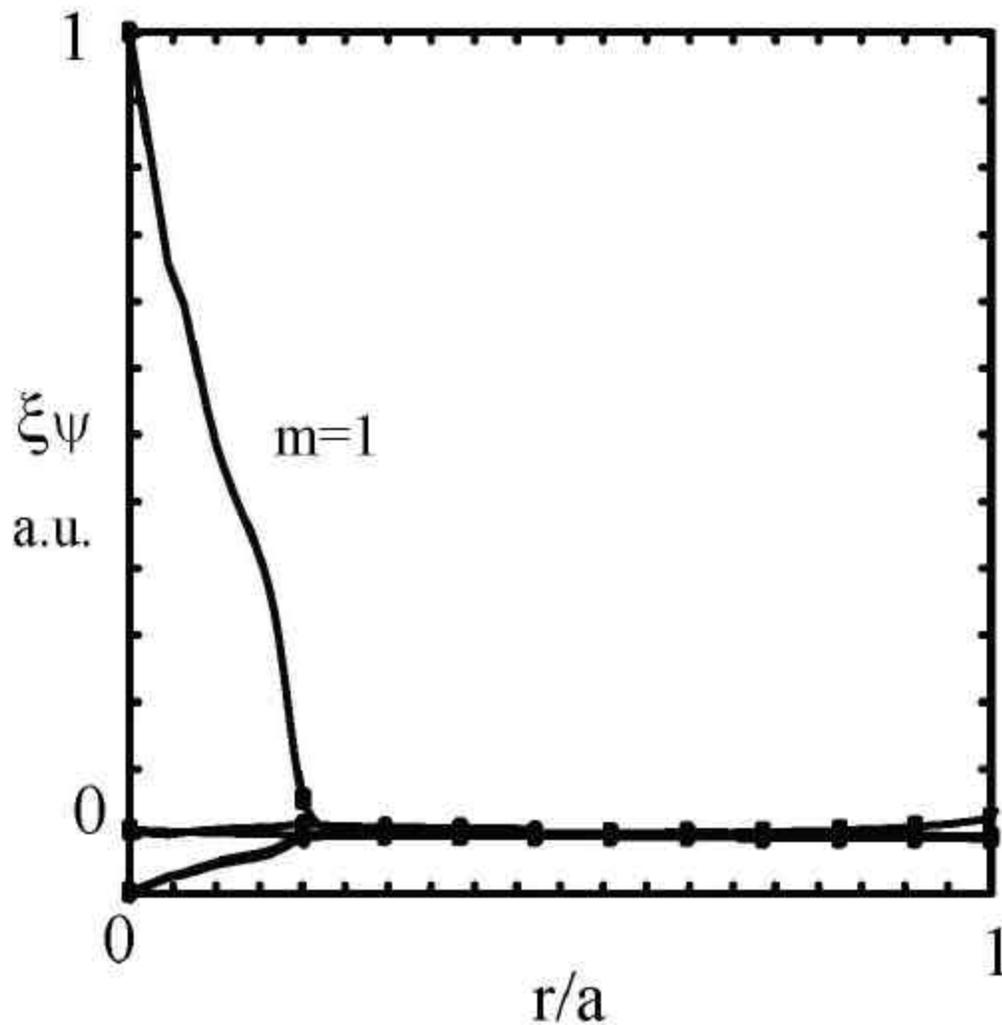
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Predicted MSE signal - ITER



Unstable eigenfunction of displacement vector for internal kink $R/a=2$



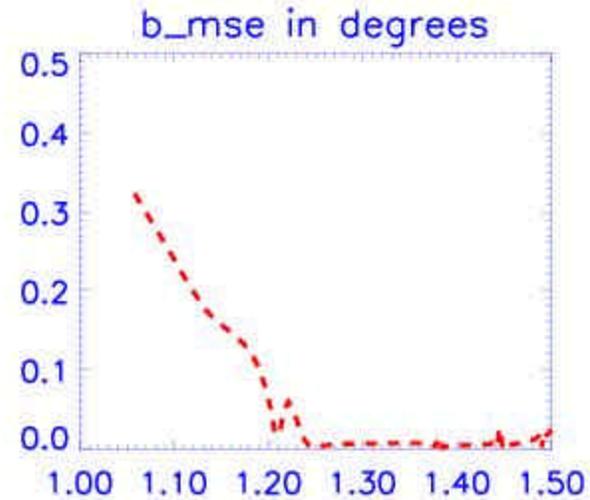
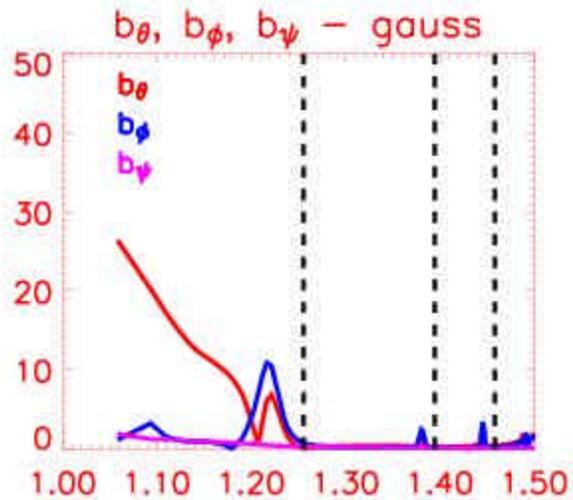
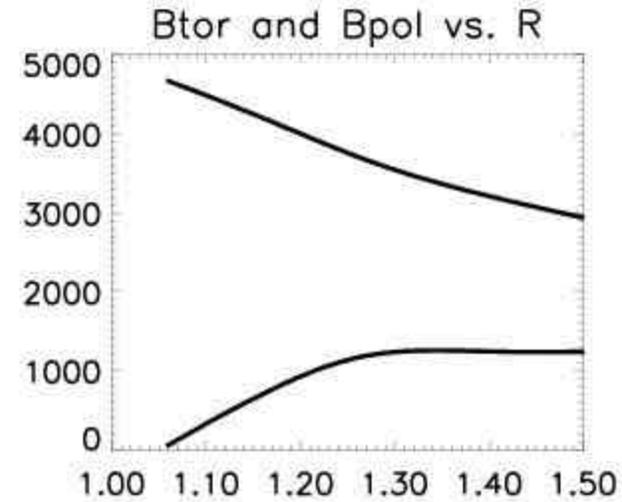
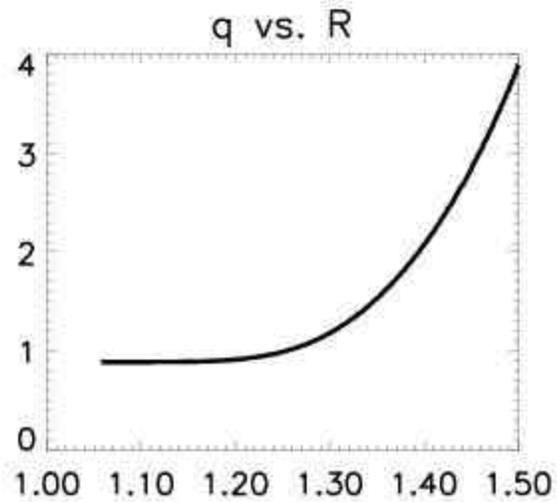
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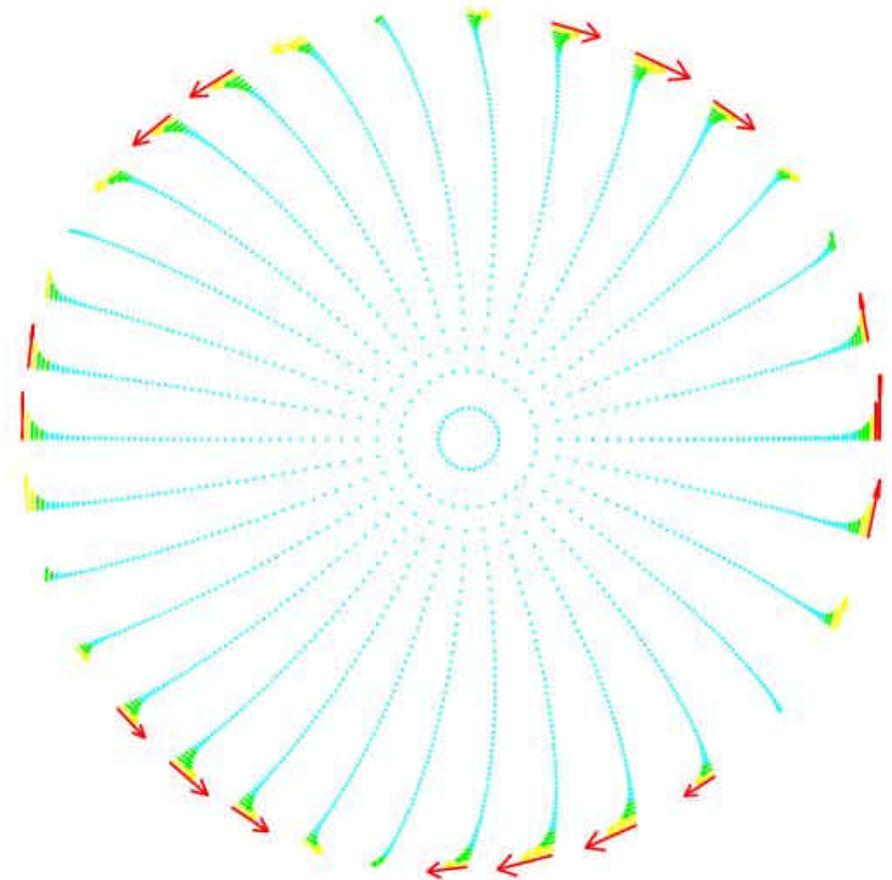
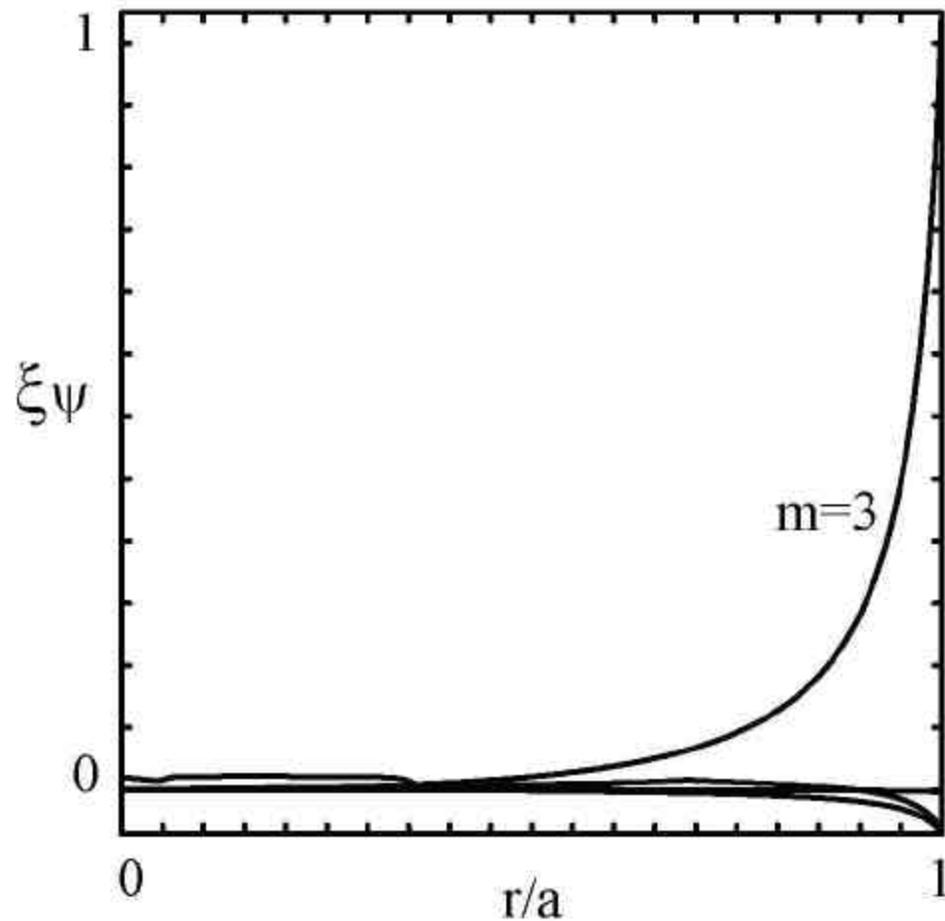


Predicted MSE signal for internal kink

$R/a=2$



Unstable eigenfunction of displacement vector for external kink $R/a=20$



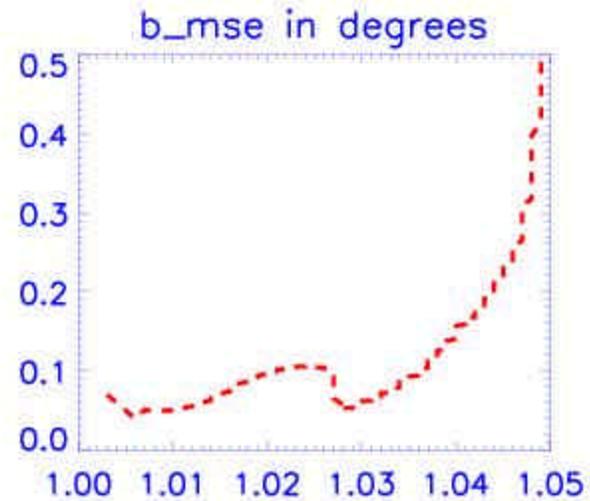
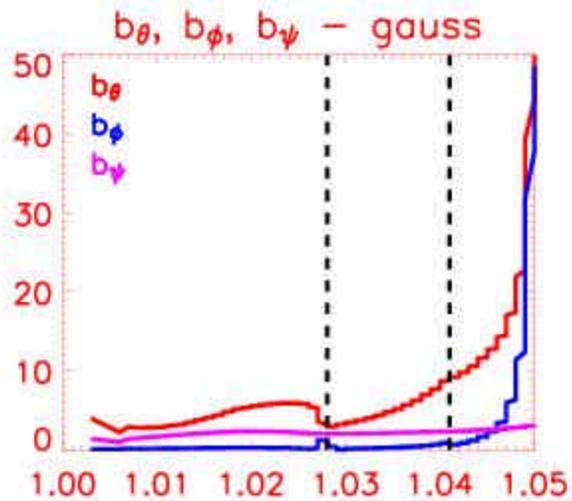
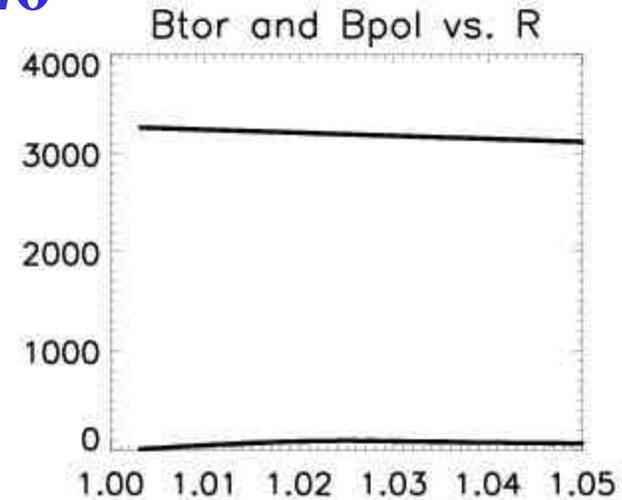
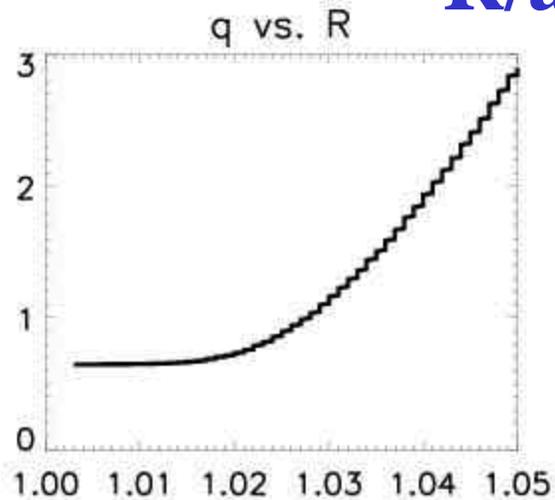
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Predicted MSE signal for external kink

$R/a=20$

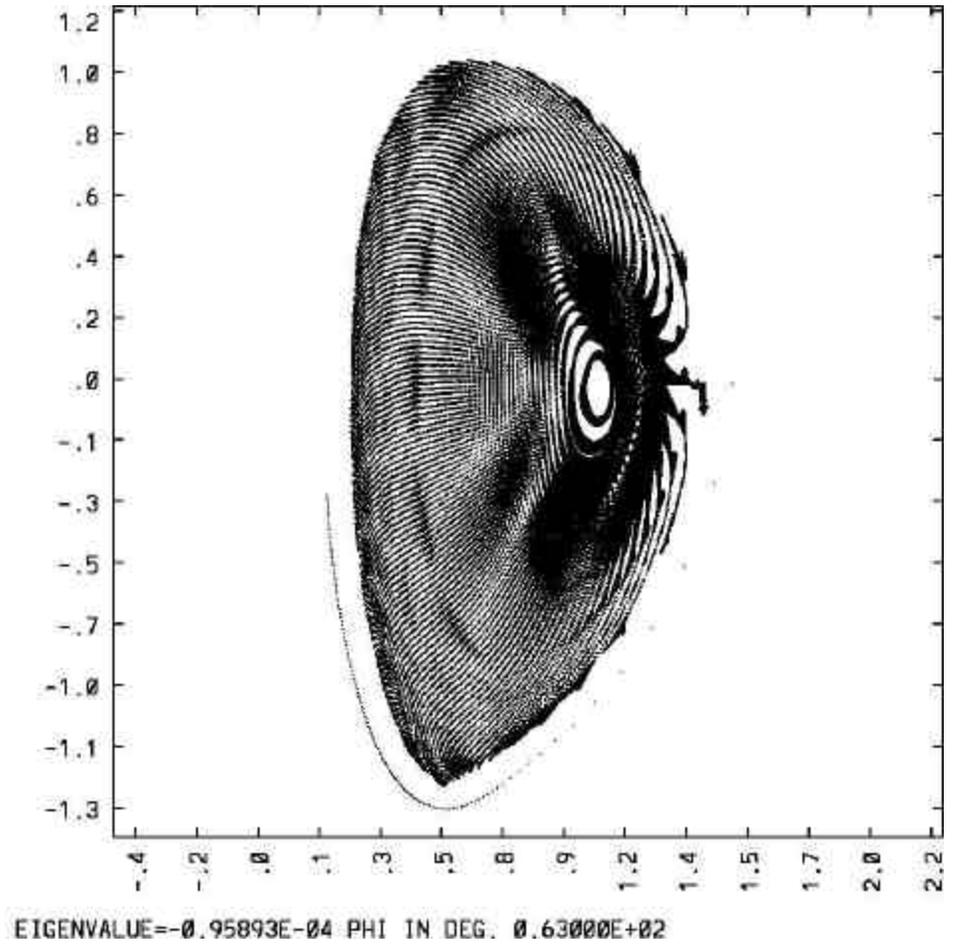
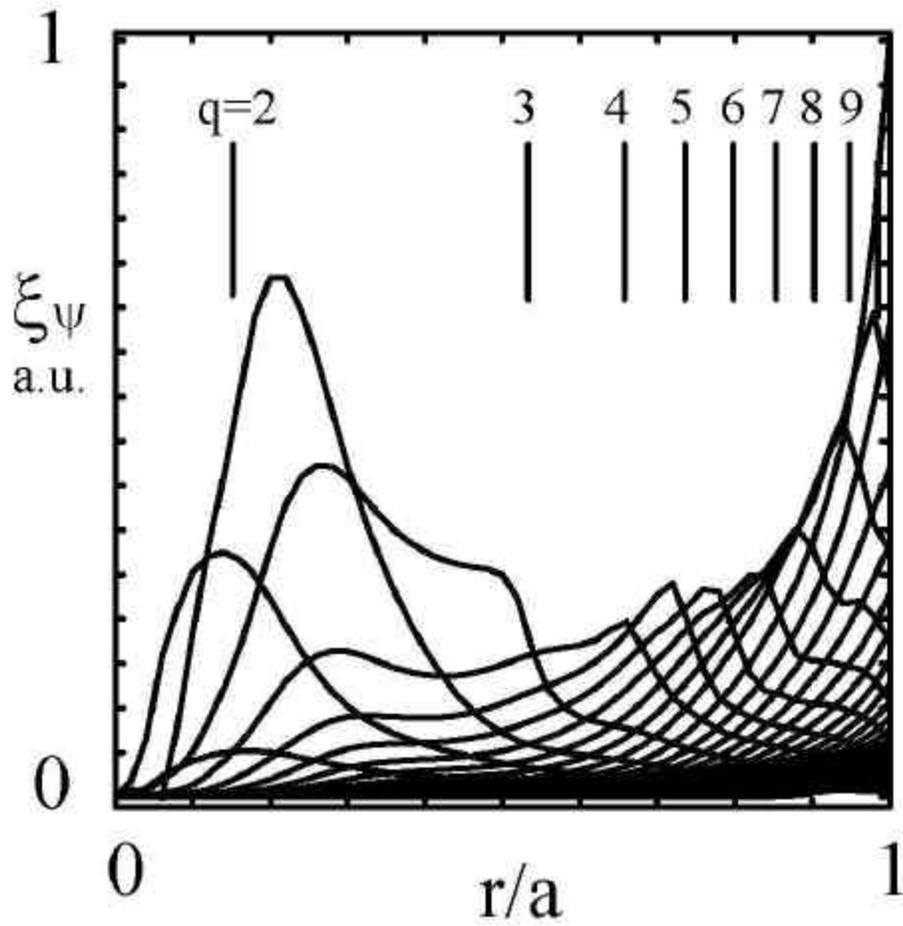


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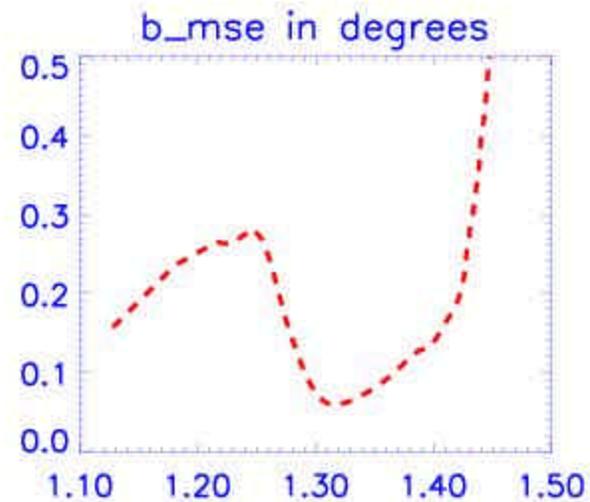
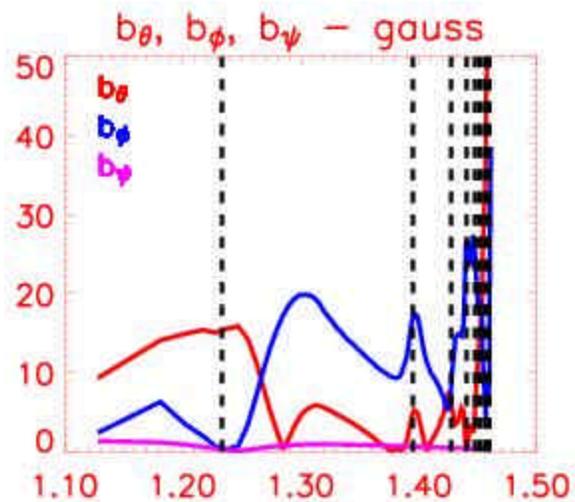
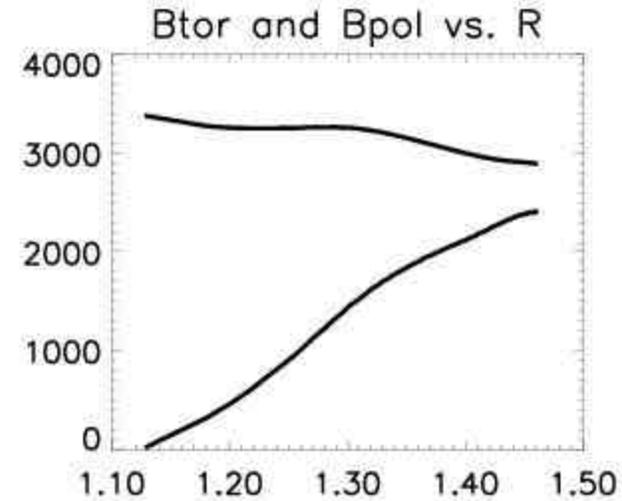
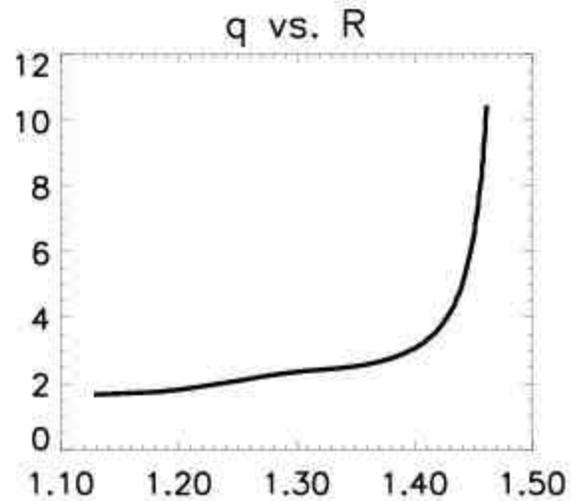
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Unstable eigenfunction of displacement vector for NSTX $-b_N \sim 6$



Predicted MSE signal for NSTX – $b_N \sim 6$



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Summary

- Investigated the use of the MSE signal as a diagnostic tool for internal MHD activity
- Developed a post-processor for the PEST code to predict radial profile of MSE fluctuations
- Started benchmarking with experimental data
- Need to refine the external model for predicting MHD signal at Mirnov loop location
- Applied procedure to a variety of geometries and instabilities

