

Measurements and Modeling of Plasma Flow Damping in the HSX Stellarator

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and the HSX Team



Outline

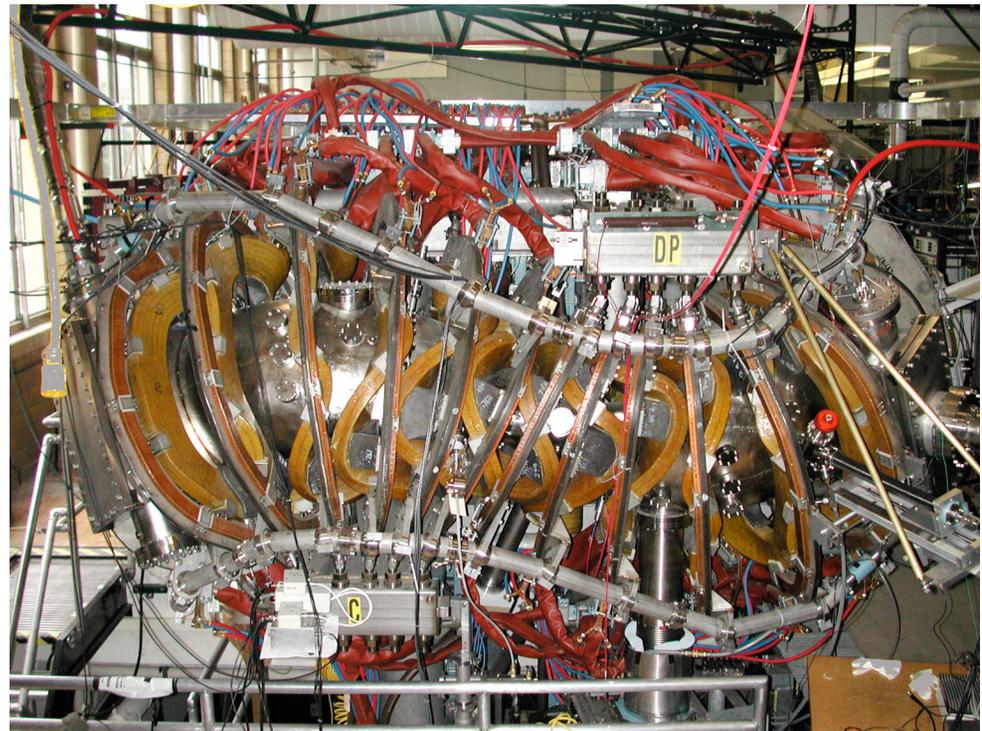
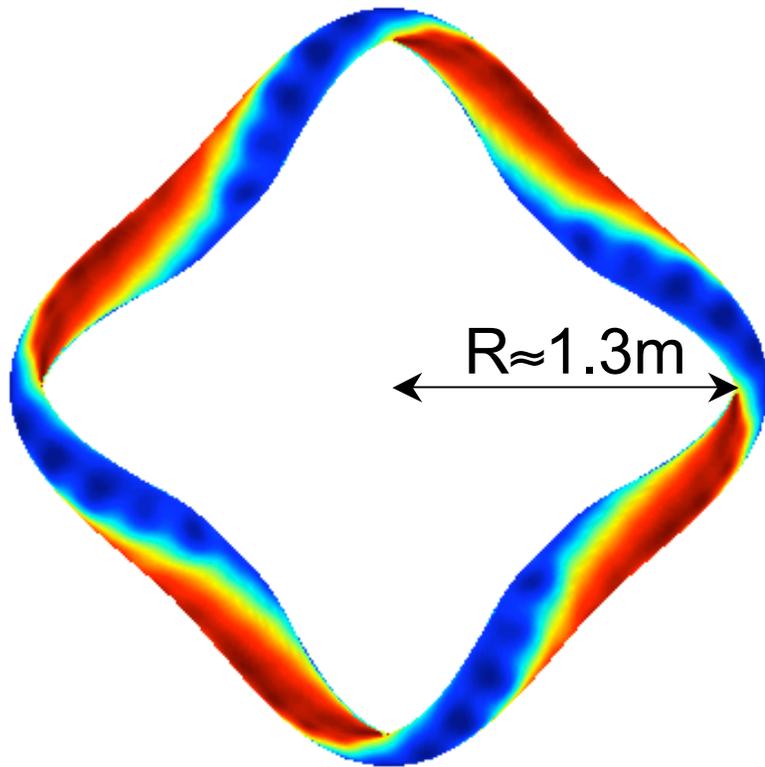
- Description of experiments and diagnostics
- Studies of flow and electric field evolution
 - Asymmetries between the spin-up and relaxation
 - Two time-scale flow evolution
 - Reduced damping with quasisymmetry
- Neoclassical modeling of flow damping
 - Original model for the spin-up
- Measurements/modeling comparison
 - Reduced flow damping in quasisymmetric configurations
 - Flow damping larger than the neoclassical prediction



Modular Coils Lead to Experimental Flexibility

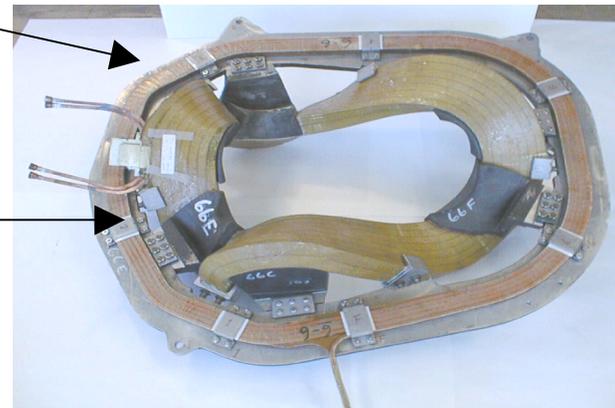
$B=0.5T$

$P_{ECH} < 200 \text{ kW @ } 28 \text{ GHz}$



Modular Coil

Auxiliary Coil

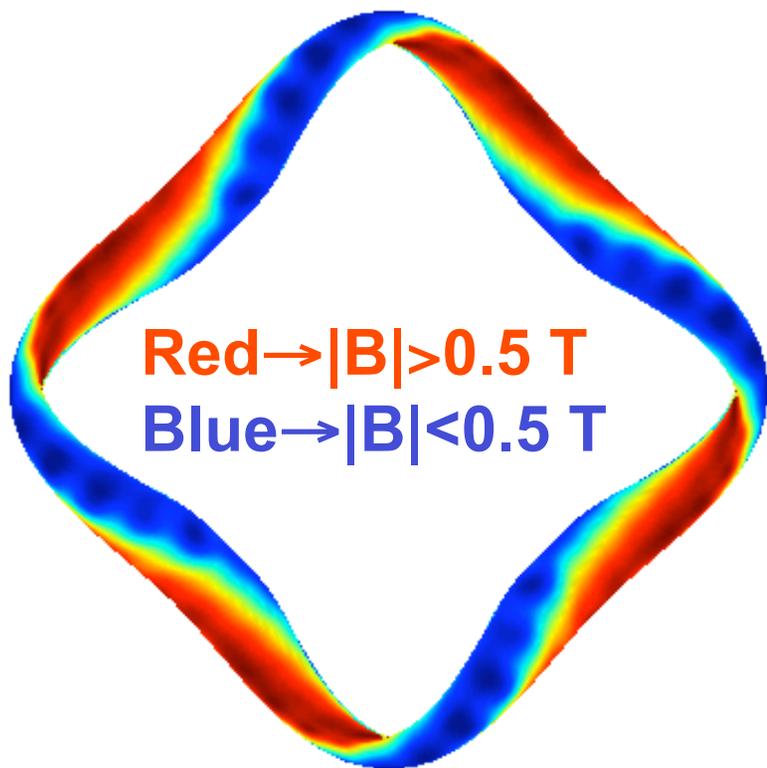


HSX is located at the University of Wisconsin-Madison



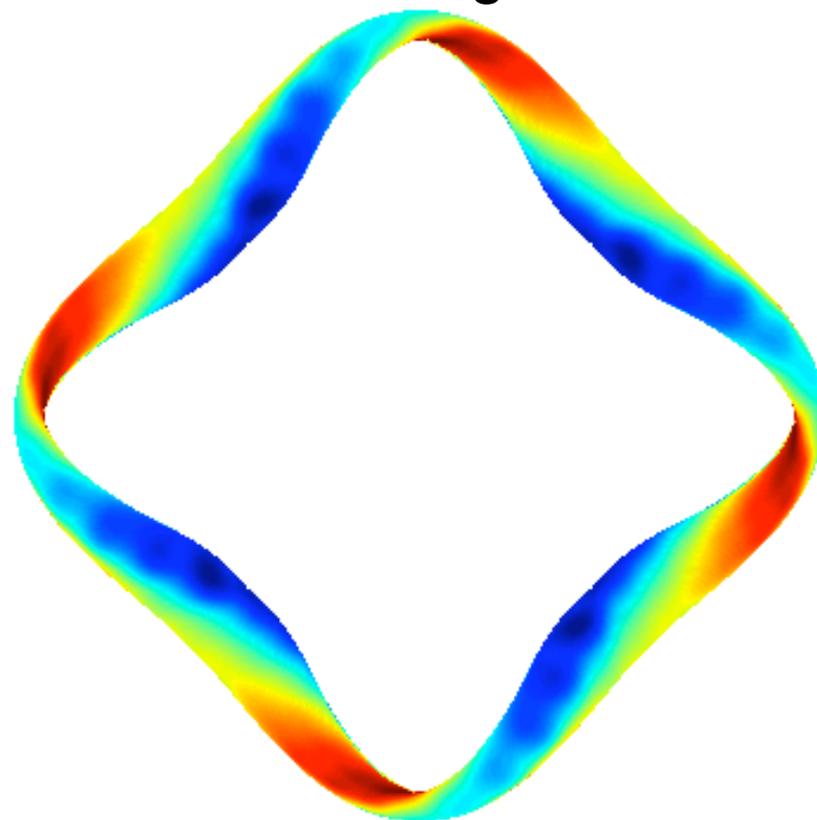
HSX Provides Access to Configurations With and Without Symmetry

QHS Configuration



QHS: Helical Bands of Constant $|B|$

Mirror Configuration

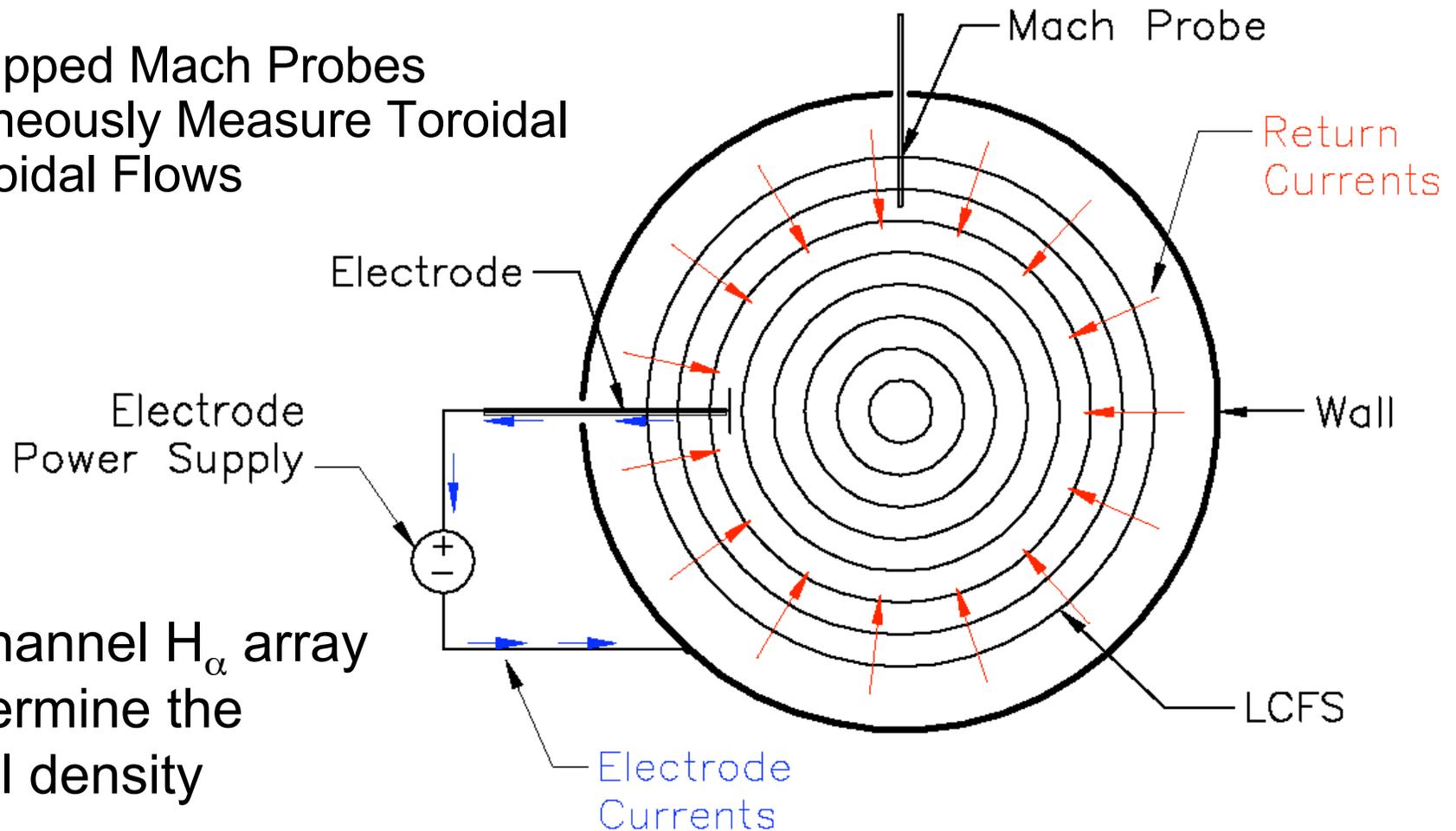


Mirror: Helical Bands are Broken



Probes and Electrodes Used to Study Flow Damping

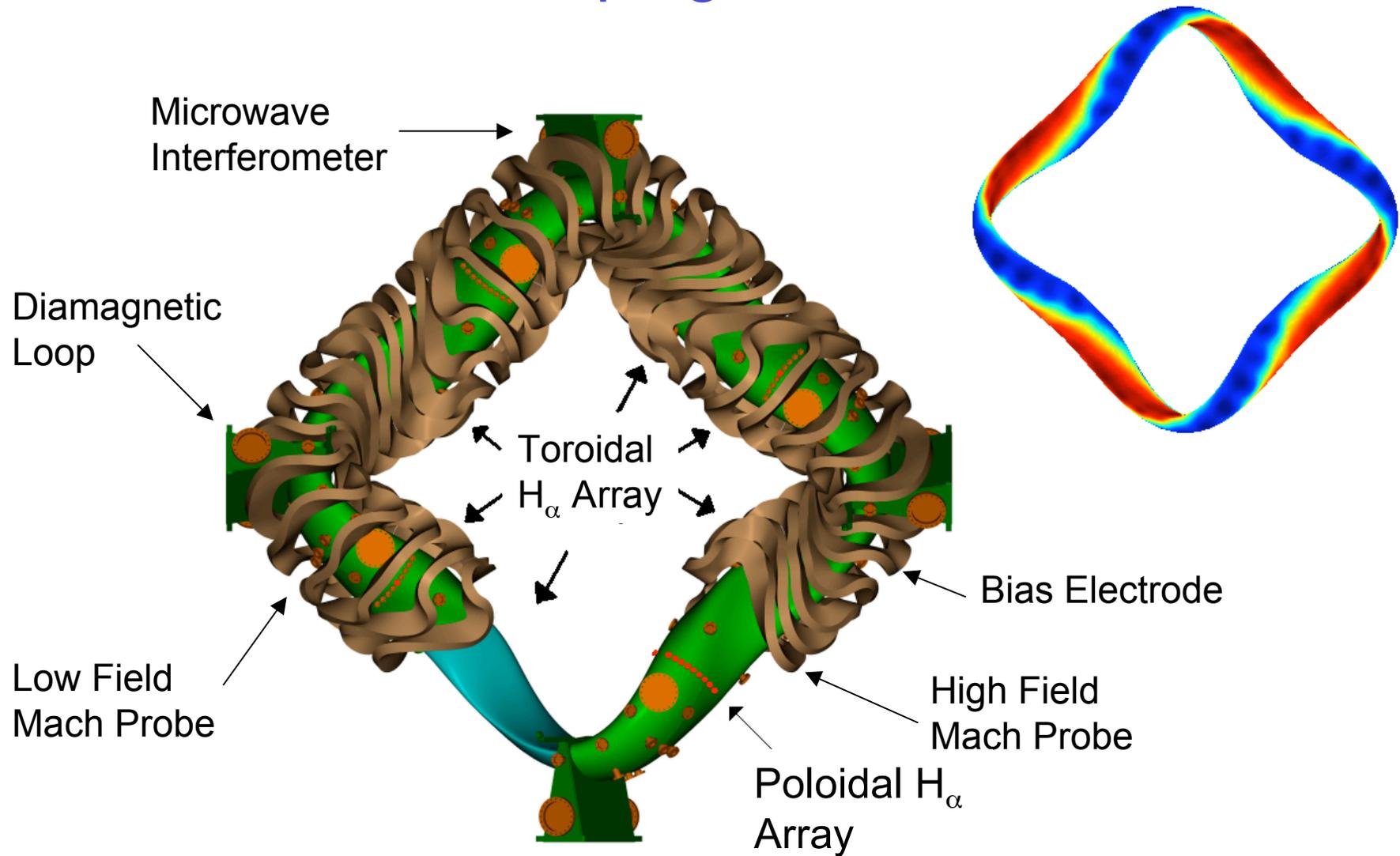
- Bias Electrode to Drive Flows
- Multi-Tipped Mach Probes Simultaneously Measure Toroidal and Poloidal Flows



- 16 channel H_{α} array to determine the neutral density



Comprehensive Diagnostics For Flow Damping Studies



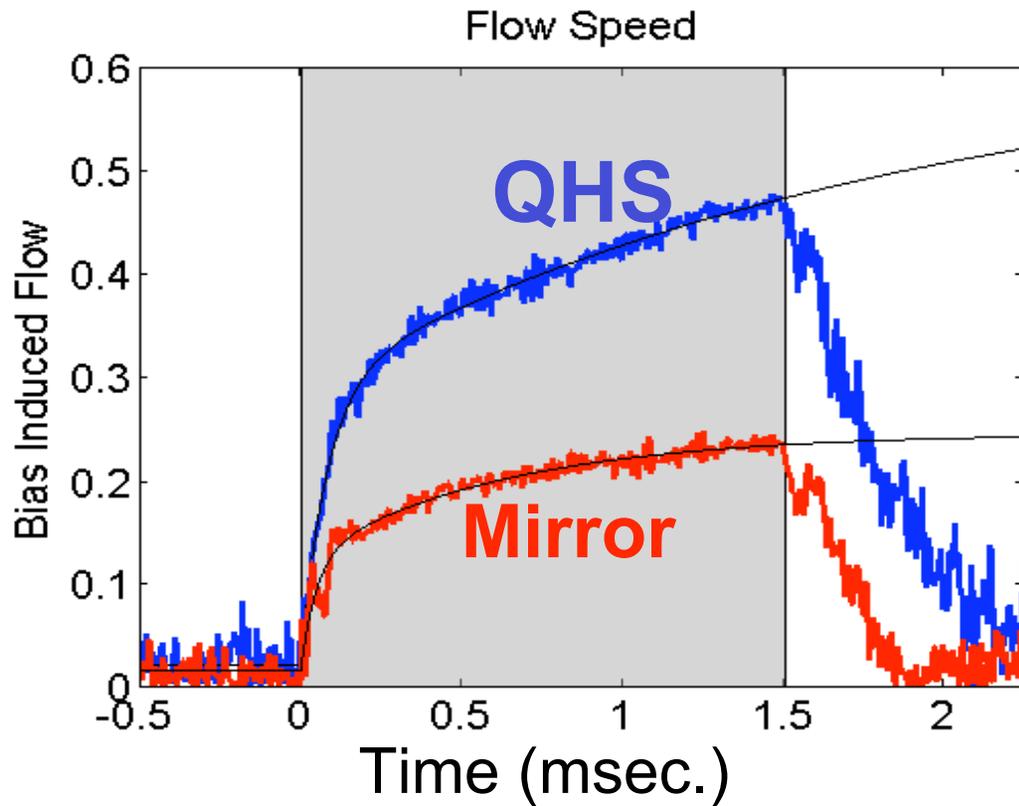
Biased Electrode Experiments

Demonstrate New Flow Phenomena:

- 1) Reduced Flow Damping with Quasisymmetry
- 2) Two Time-Scale Flow Evolution



Preview: QHS Flows Damp More Slowly, Goes Faster For Less Drive



**QHS: 8 A of
electrode current**

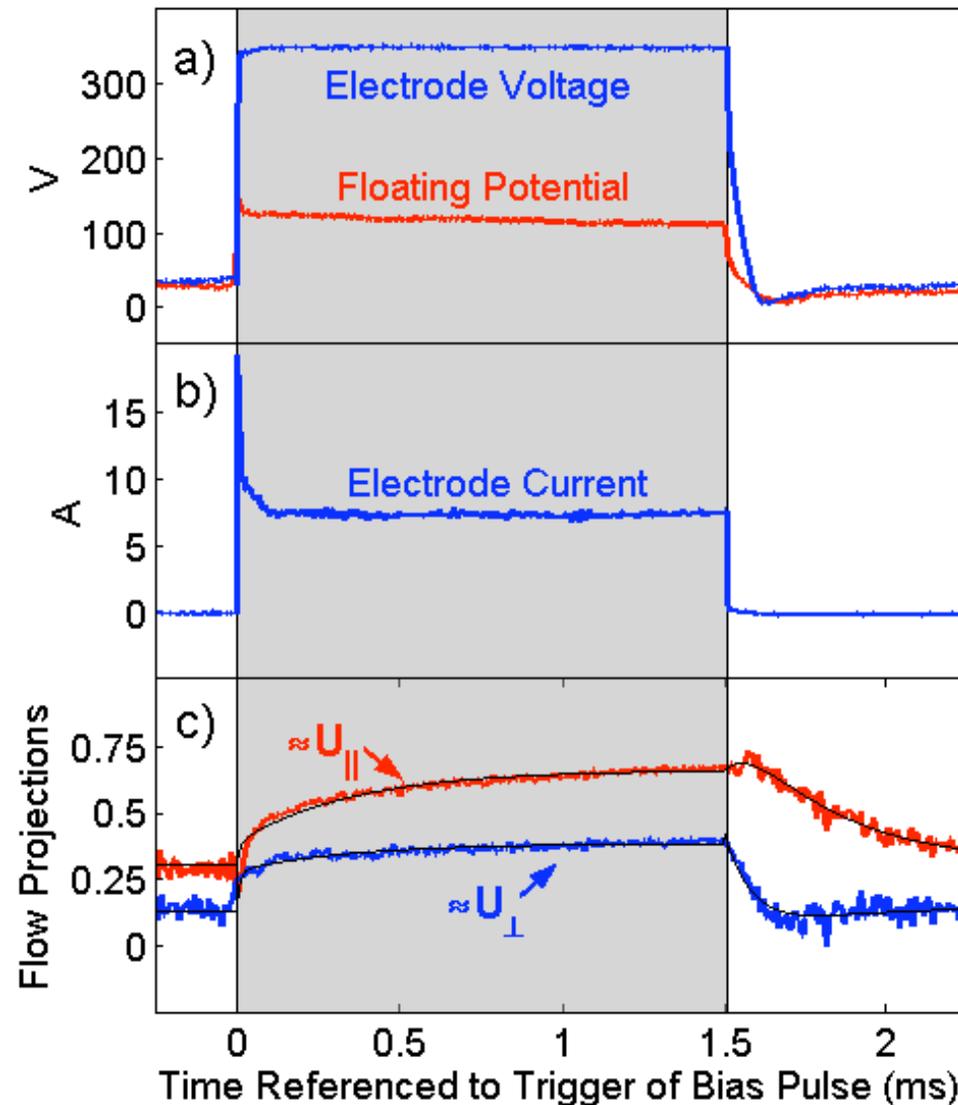
**Mirror: 10 A of
electrode current**

All other parameters ($n_e=1 \times 10^{12} \text{cm}^{-3}$, $n_n \approx 1 \times 10^{10} \text{cm}^{-3}$, $T_i \approx 25 \text{eV}$, $B=0.5 \text{T}$, $P_{\text{ECH}}=50 \text{ kW}$) held constant.



Asymmetries and Multiple Time-Scales Observed in Flow Evolution

- Potentials:
Fast Rise and Slow Decay
- Electrode Current:
Large Spike and Fast Termination
- Plasma Flows:
Fast and Slow Time-Scales at Rise and Decay



Neoclassical Modeling

Goal: Assess the flow damping caused by

- 1) Symmetry breaking ripples
- 2) Ion-neutral friction



Solve the Momentum Equations on a Flux Surface

- Two time-scales/directions come from the coupled momentum equations on a surface

$$m_i N_i \frac{\partial}{\partial t} \langle \mathbf{B}_P \cdot \mathbf{U} \rangle = - \frac{\sqrt{g} B^\zeta B^\alpha}{c} \langle \mathbf{J}_{\text{plasma}} \cdot \nabla \psi \rangle - \langle \mathbf{B}_P \cdot \nabla \cdot \Pi \rangle - m_i N_i \langle v_{\text{in}} \mathbf{B}_P \cdot \mathbf{U} \rangle$$

$$m_i N_i \frac{\partial}{\partial t} \langle \mathbf{B} \cdot \mathbf{U} \rangle = - \langle \mathbf{B} \cdot \nabla \cdot \Pi \rangle - m_i N_i \langle v_{\text{in}} \mathbf{B} \cdot \mathbf{U} \rangle$$

- Use Hamada coordinates, linear neoclassical viscosities, neglect heat fluxes
- Steady state solution yields radial conductivity

$$\langle \mathbf{J}_{\text{plasma}} \cdot \nabla \psi \rangle = \sigma_\perp \left(\langle \mathbf{E}_r \cdot \nabla \psi \rangle - \frac{\langle \nabla p_i \cdot \nabla \psi \rangle}{e N_i} \right)$$



Spin-Up and Spin-Down are Treated Differently in Modeling

- At bias turn-on, switches put voltage on the electrode ($\sim 1 \mu\text{sec.}$).
- Measurements show electric field is established on the electrode voltage-rise time-scale.
- **Spin-Up** Model: Flows and radial current respond to the electrode **potential** rise.
- At bias turn-off, switches break the electrode current ($\sim 1 \mu\text{sec.}$).
- **Relaxation** Model: Flows and electric field respond to the electrode **current** termination.



Flow Rise: Electric Field is Turned on Quickly

- Assume that the electric field, $d\Phi/d\psi$, is turned on quickly

$$\frac{\partial\Phi}{\partial\psi} = \begin{cases} E_{r0} & t < 0 \\ E_{r0} + \kappa_E \left(1 - e^{-t/\tau}\right) & t > 0 \end{cases}$$

- **ExB** flows and compensating Pfirsch-Schlueter flow grow on the electric field time-scale
- Parallel flow grows at a “Hybrid rate” v_F determined by viscosity and ion-neutral friction

$$v_F = \nu v_\alpha + v_\zeta + v_{in}$$

Toroidal Damping

Poloidal Damping

- Two time-scales/two direction flow evolution

$$\mathbf{U}(t) \approx U_E^\alpha \left(1 - e^{-t/\tau}\right) \mathbf{e}_\alpha + \mathbf{U}_\parallel \left(1 - e^{-v_F t}\right)$$



Flow Decay: External Radial Current is Quickly Turned Off

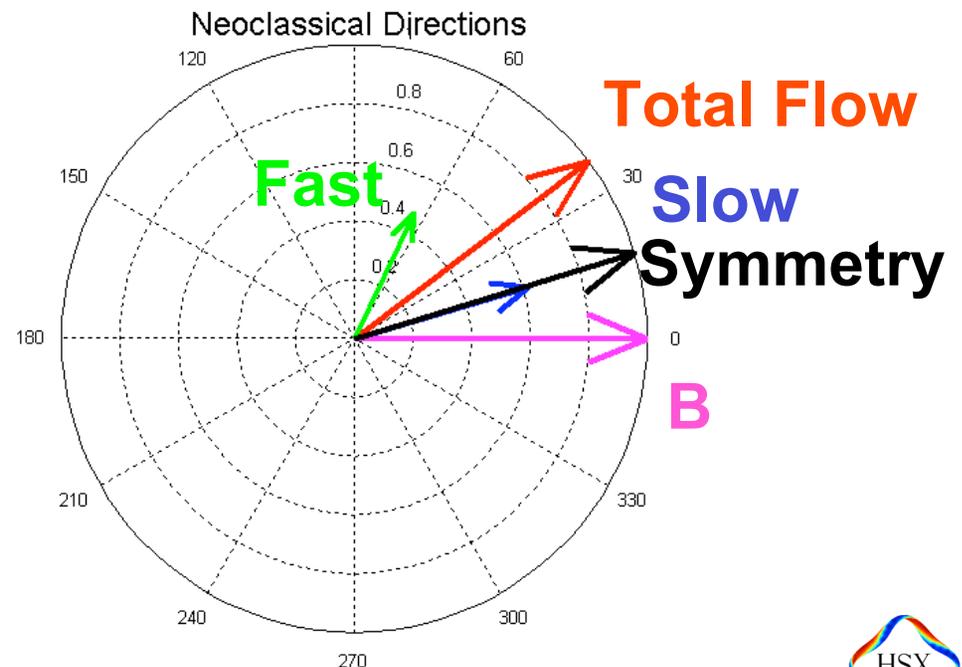
- $\gamma_f(\psi)$ (fast), and $\gamma_s(\psi)$ (slow rate) are flux surface quantities related to the geometry and ion-neutral collision frequency.
- Break the flow into parts damped on each time-scale:

$$\mathbf{U} = \underbrace{e^{-\gamma_f(t-t_0)} \mathbf{f}}_{\text{Fast}} + \underbrace{e^{-\gamma_s(t-t_0)} \mathbf{s}}_{\text{Slow}}$$

➤ Large neutral density ($n_n = 1 \times 10^{12} \text{ cm}^{-3}$) in this calculation.

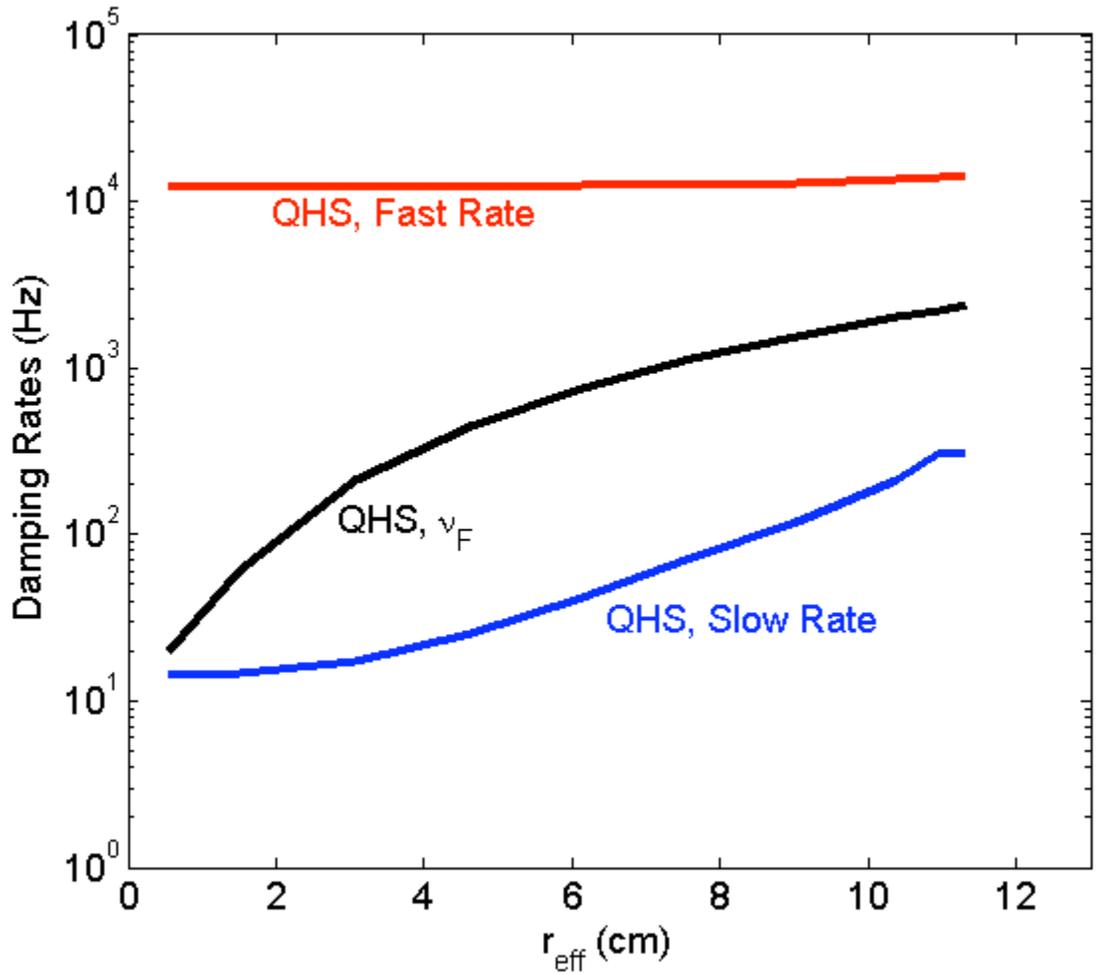
➤ **Slow rate** corresponds to flows in the direction of symmetry.

➤ Numerically calculated Hamada basis vectors used in this figure.



The Hybrid Rate is Intermediate to the Fast and Slow Rate

Fast Rate
is faster than
Hybrid Rate, ν_F
is faster than
Slow Rate



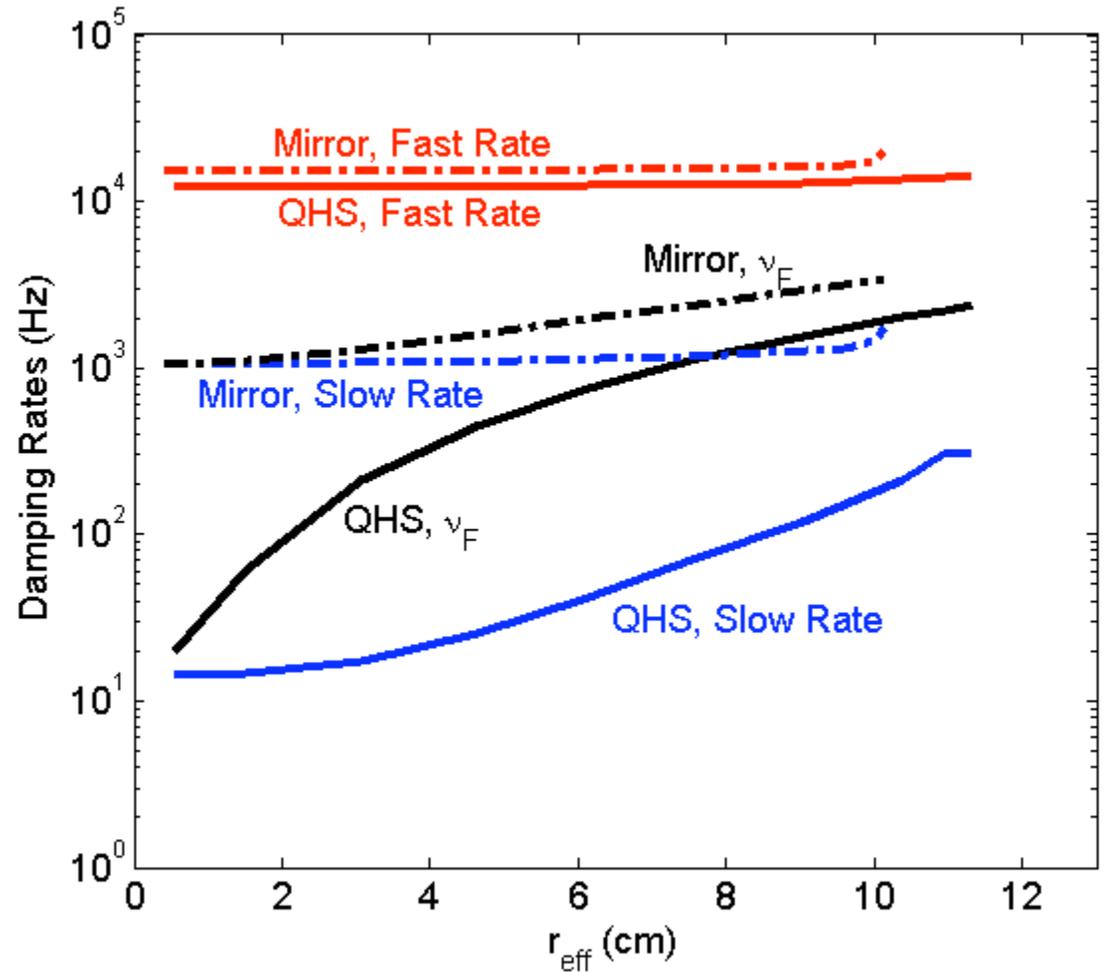
Mirror Shows Increased Neoclassical Damping Compared to QHS

QHS/Mirror Comparison

Fast rates are comparable

Mirror ν_F is larger by a factor of 2-3

Mirror slow rate is larger by 1-2 orders of magnitude

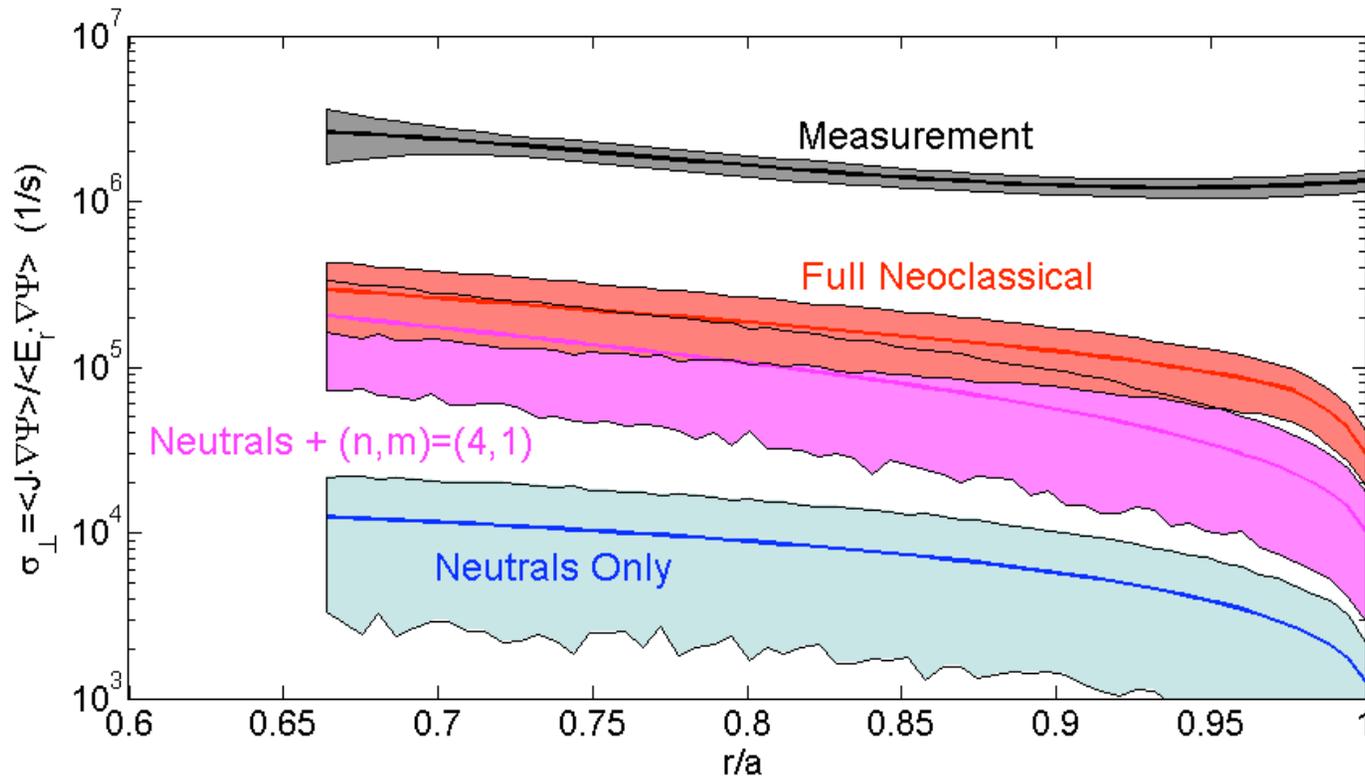


Comparison of Neoclassical Theory with Measurements

- 1) Reduced Flow Damping with Quasisymmetry
- 2) Evidence of Anomalous Flow Damping



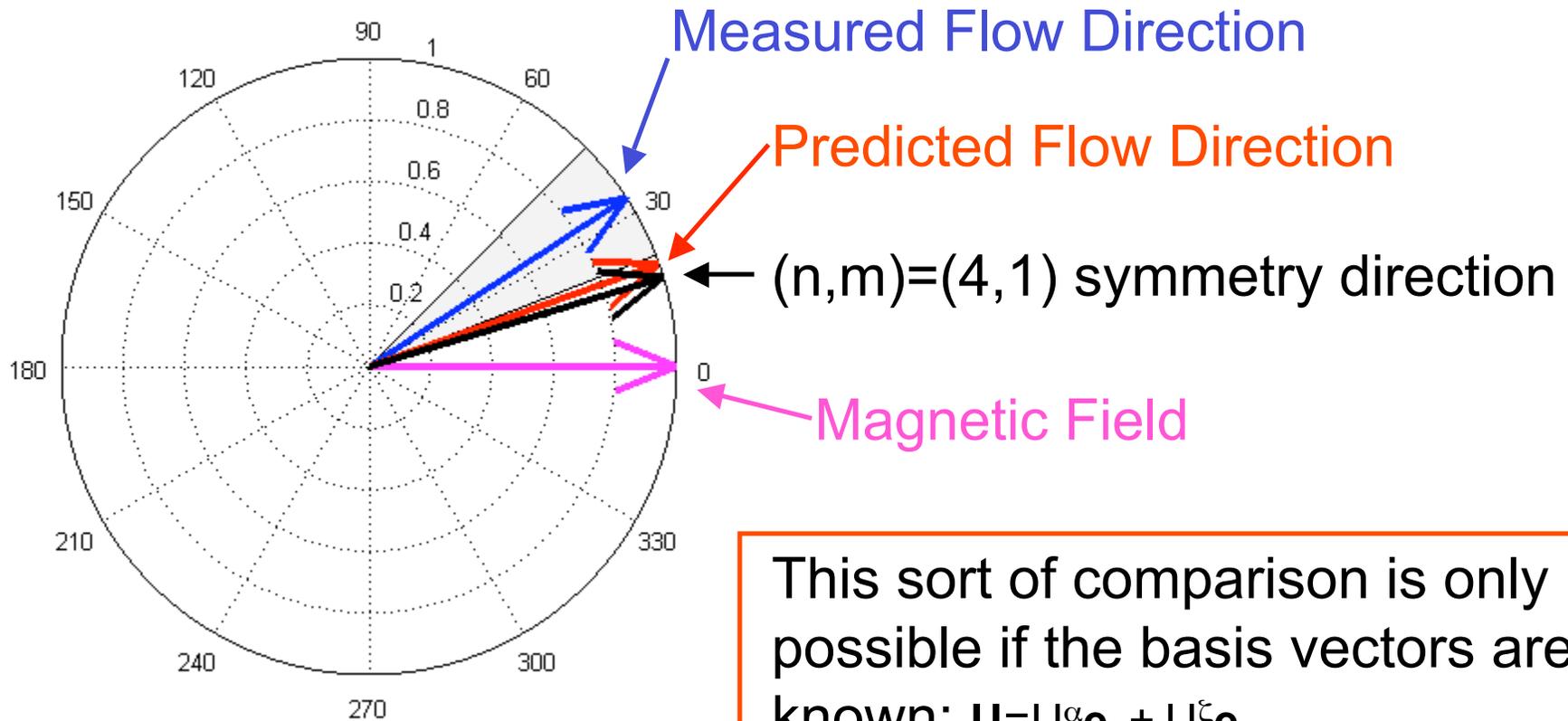
QHS Radial Conductivity is Larger than the Neoclassical Prediction



$$\langle \bar{\mathbf{J}}_{plasma} \cdot \bar{\nabla} \psi \rangle = \sigma_{\perp} \left(\langle \bar{\mathbf{E}}_r \cdot \bar{\nabla} \psi \rangle - \frac{\langle \bar{\nabla} p_i \cdot \bar{\nabla} \psi \rangle}{eN_i} \right)$$



Steady State Flow Direction Differs Somewhat from Neoclassical Prediction



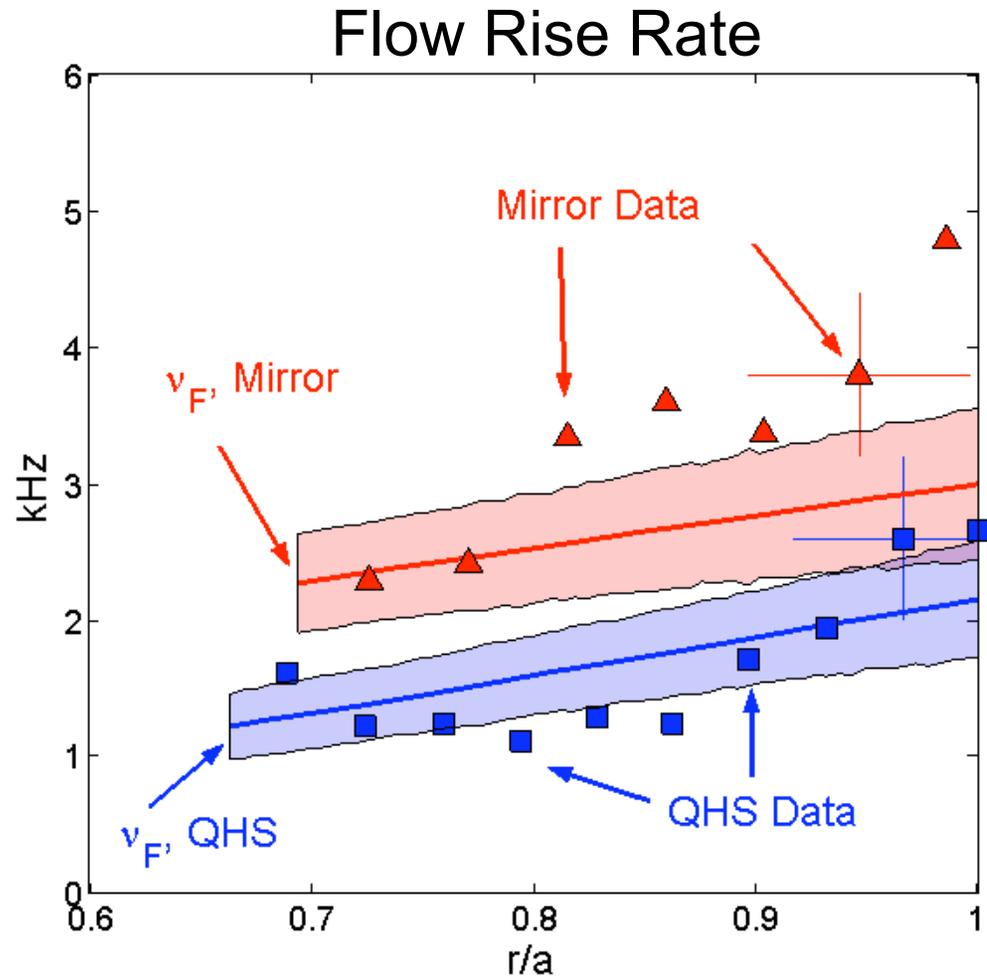
This sort of comparison is only possible if the basis vectors are known: $\mathbf{U} = U^\alpha \mathbf{e}_\alpha + U^\zeta \mathbf{e}_\zeta$



Modeling Predicts the Difference in the QHS and Mirror Slow Rise Rates

➤ Mirror flows rise more quickly than QHS.

➤ Neoclassical hybrid time ν_F shows good agreement with the measurements.

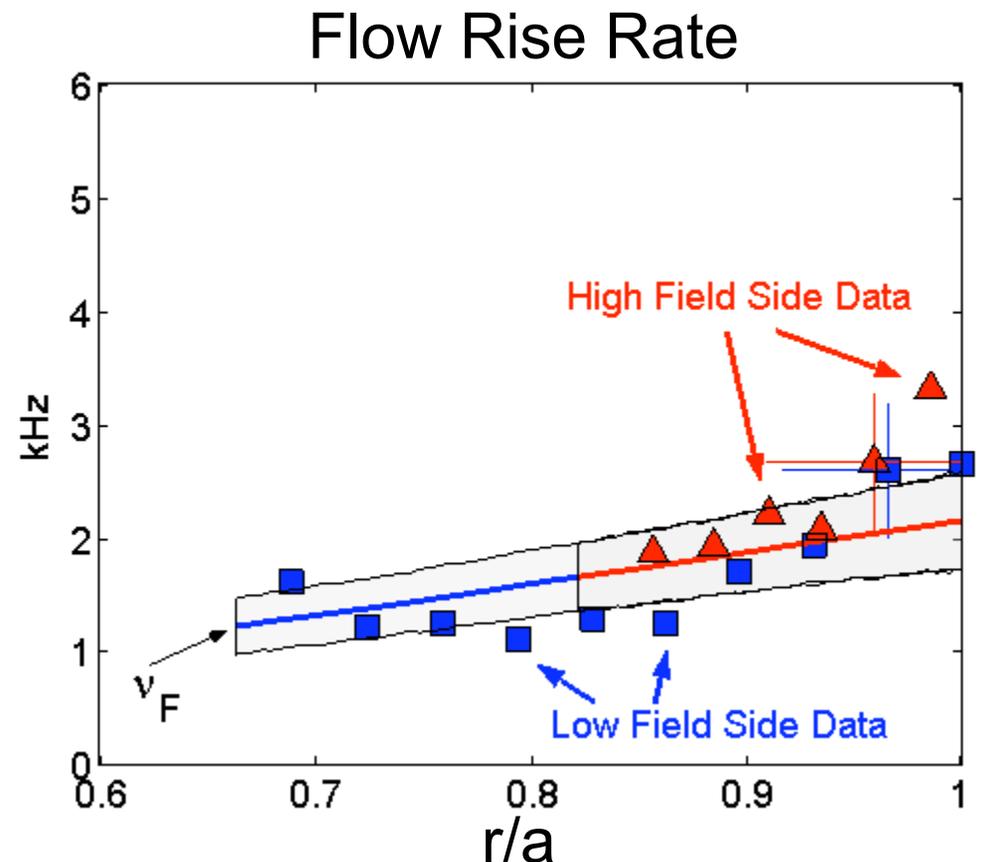


Similar Flow Rise Rates Simultaneously Measured at High and Low Field Locations

All relevant time-scales are similar on high and low field sides

- Slow Flow Rise Time
- Floating Potential Decay Time
- Fast Flow Decay Time
- Slow Flow Decay Time

Floating Potential and J_{sat} profiles similar as well.

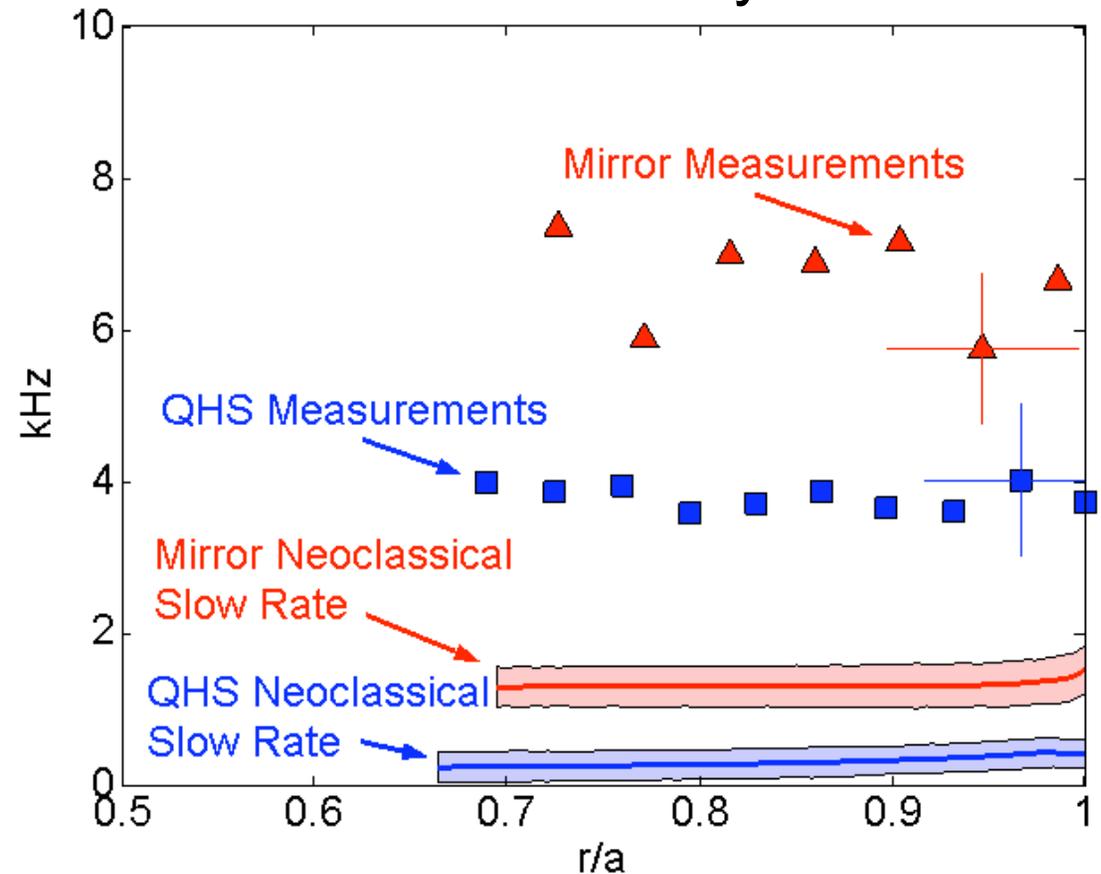


Flow Decay Rates Show Reduced Damping with Quasisymmetry

Slow Flow Decay Rate

➤ Neoclassical model predicts a much slower decay than the measurements (Factor of 10 in QHS, factor of 3-5 in Mirror).

➤ *Difference* between measurements is comparable to the *difference* between the models.



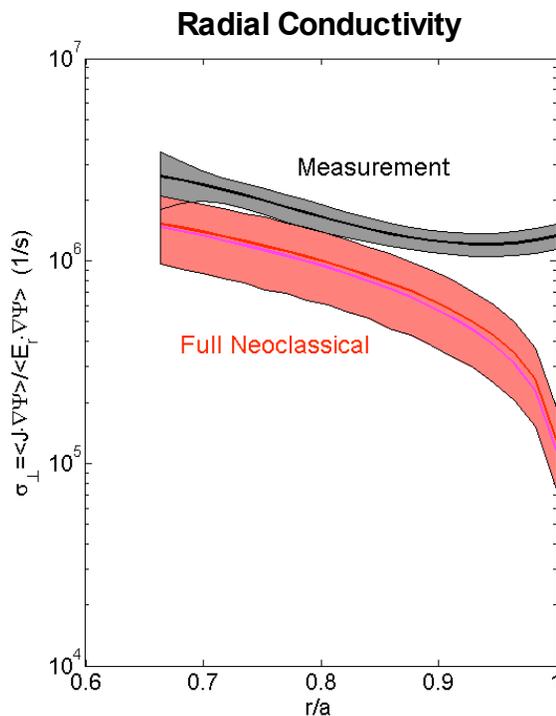
Conclusion

Quasisymmetry reduces flow damping, even in the presence of some anomalous damping.



Artificially Increased Damping Improves Theory/Experiment Comparison

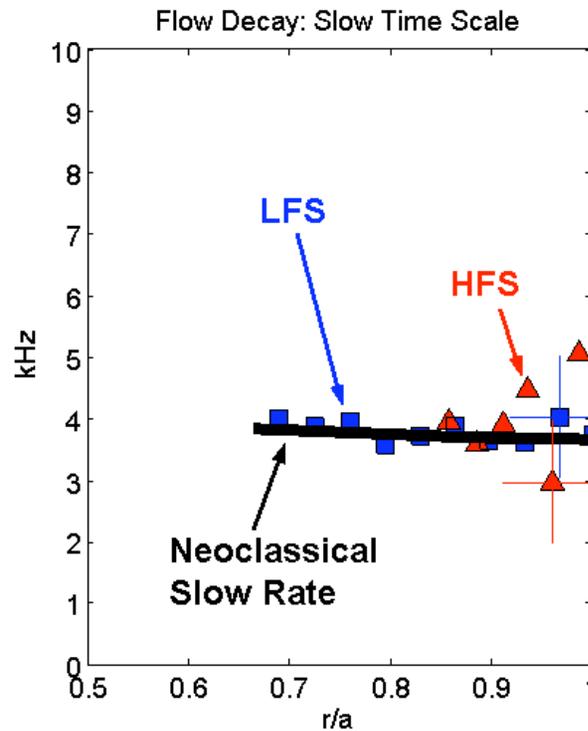
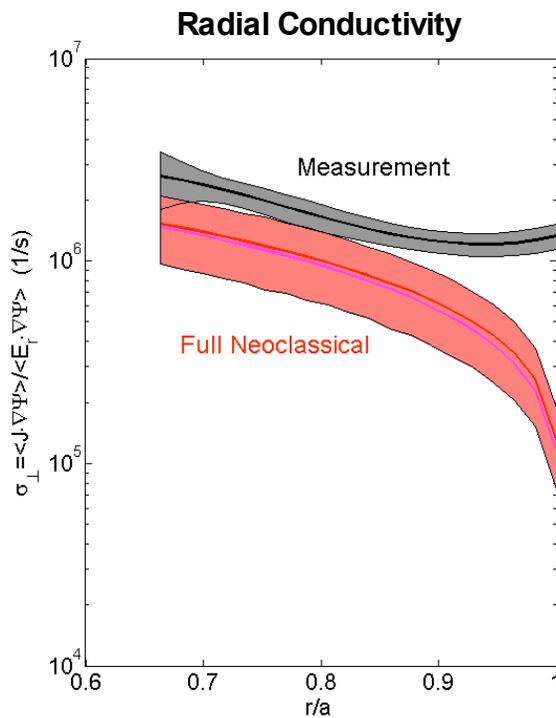
Increase the neutral density to *simulate* extra damping. $\nu_{in} \rightarrow \nu_{eff} \approx 3.6\text{kHz}$

$$\frac{a^2}{4\tau} \approx \frac{3600(0.11^2)}{4} \approx 10 \frac{\text{m}^2}{\text{s}}$$


Artificially Increased Damping Improves Theory/Experiment Comparison

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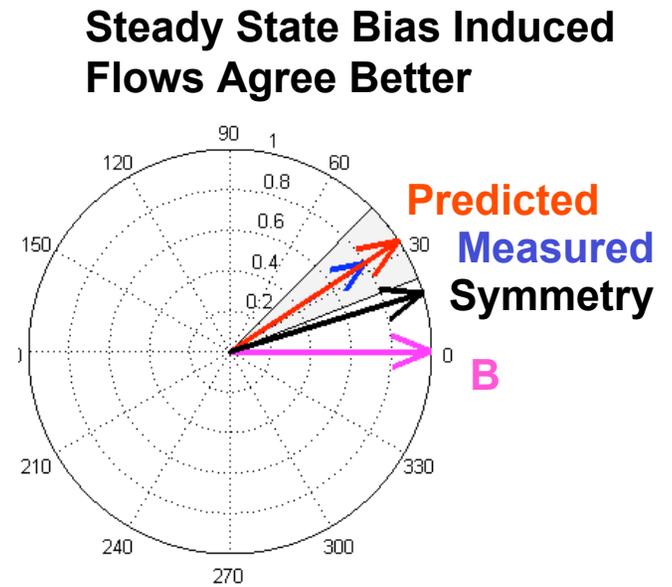
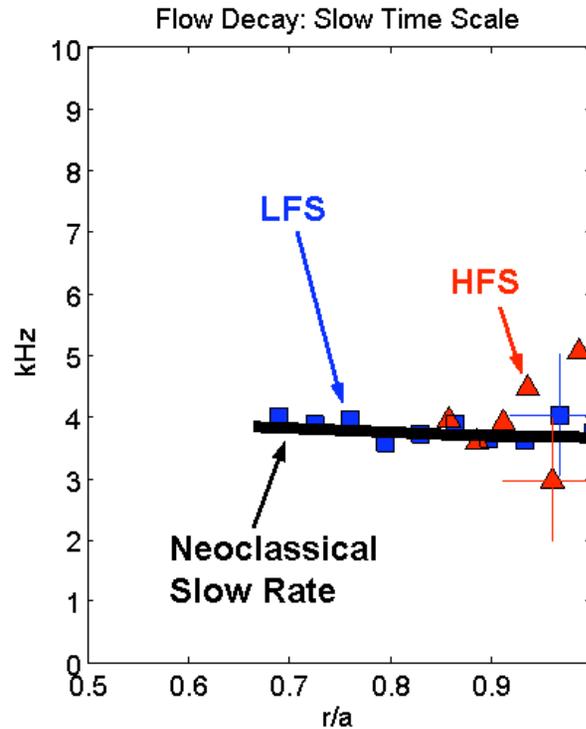
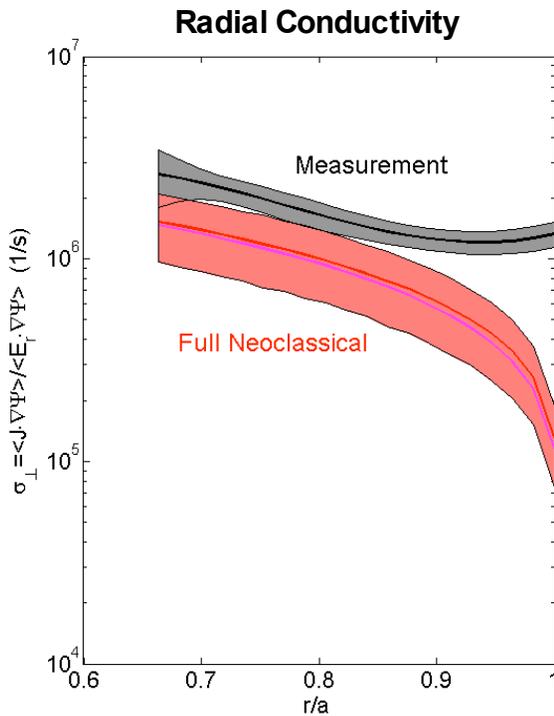
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Summary

- We have observed 2 time-scale flow evolution in HSX.
- An original model for the spin-up reproduces many of the features in the measurement.
- The damping in the symmetry direction appears to be larger than the neoclassical prediction with neutrals.
- The QHS configuration exhibits reduced damping compared to a configuration with the symmetry broken.



The End

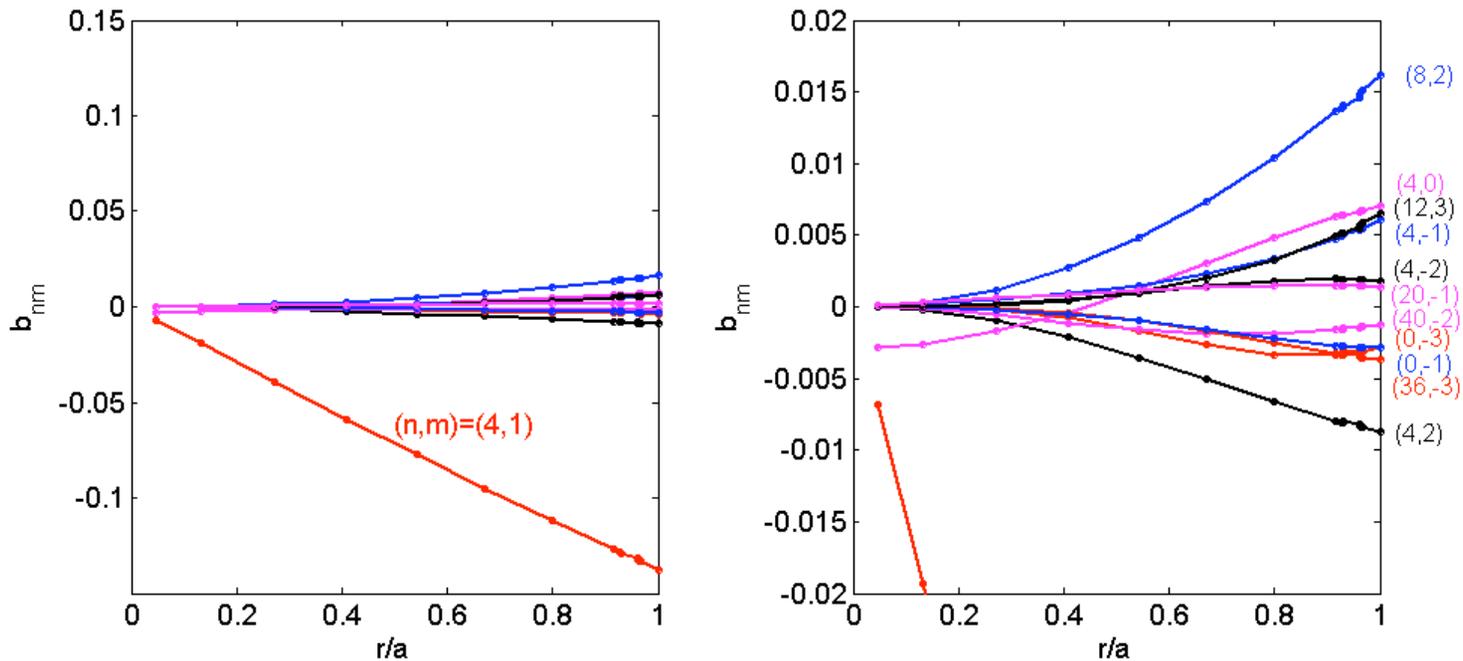


APS-DPP 2004



Neoclassical Theory, Including Neutrals, is a Candidate to Explain Flow Damping in HSX

- Near the edge, there are a number of growing symmetry breaking terms in the Hamada spectrum.

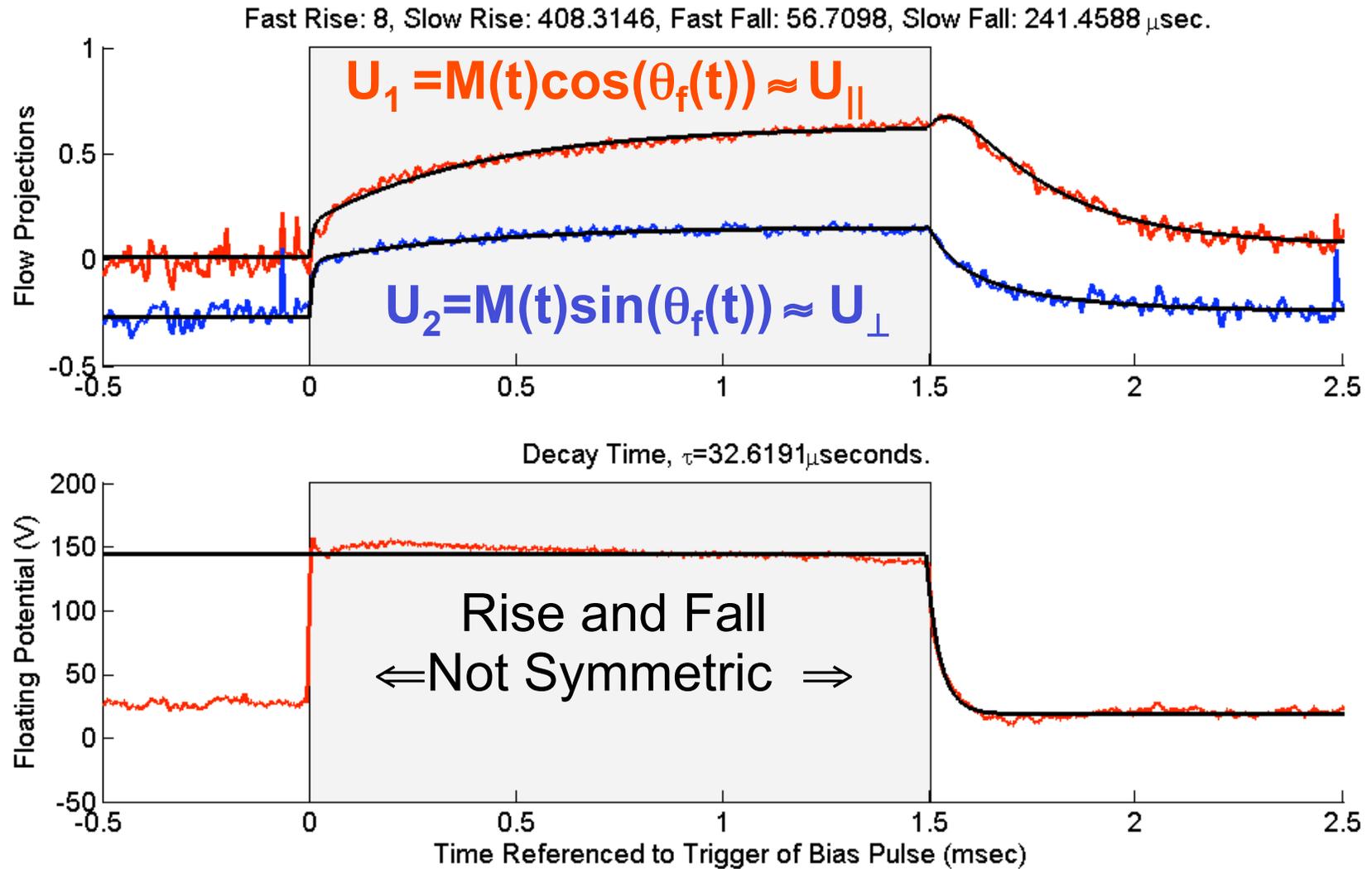


- Low density plasma allows significant neutral penetration.

$$\lambda_{\text{mfp,H}} = \frac{\sqrt{\frac{2E_H}{m}}}{n_e \langle \sigma v \rangle_{\text{H+e} \rightarrow \text{p}+2\text{e}}} = \frac{\sqrt{\frac{2 \cdot 3 \cdot 1.6 \times 10^{-19}}{1.67 \times 10^{-27}} \left(\frac{\text{m}}{\text{s}} \right)}}{10^{12} (\text{cm}^{-3}) \cdot 2.5 \times 10^{-8} (\text{cm}^3 \text{s}^{-1})} \approx 1\text{m}$$



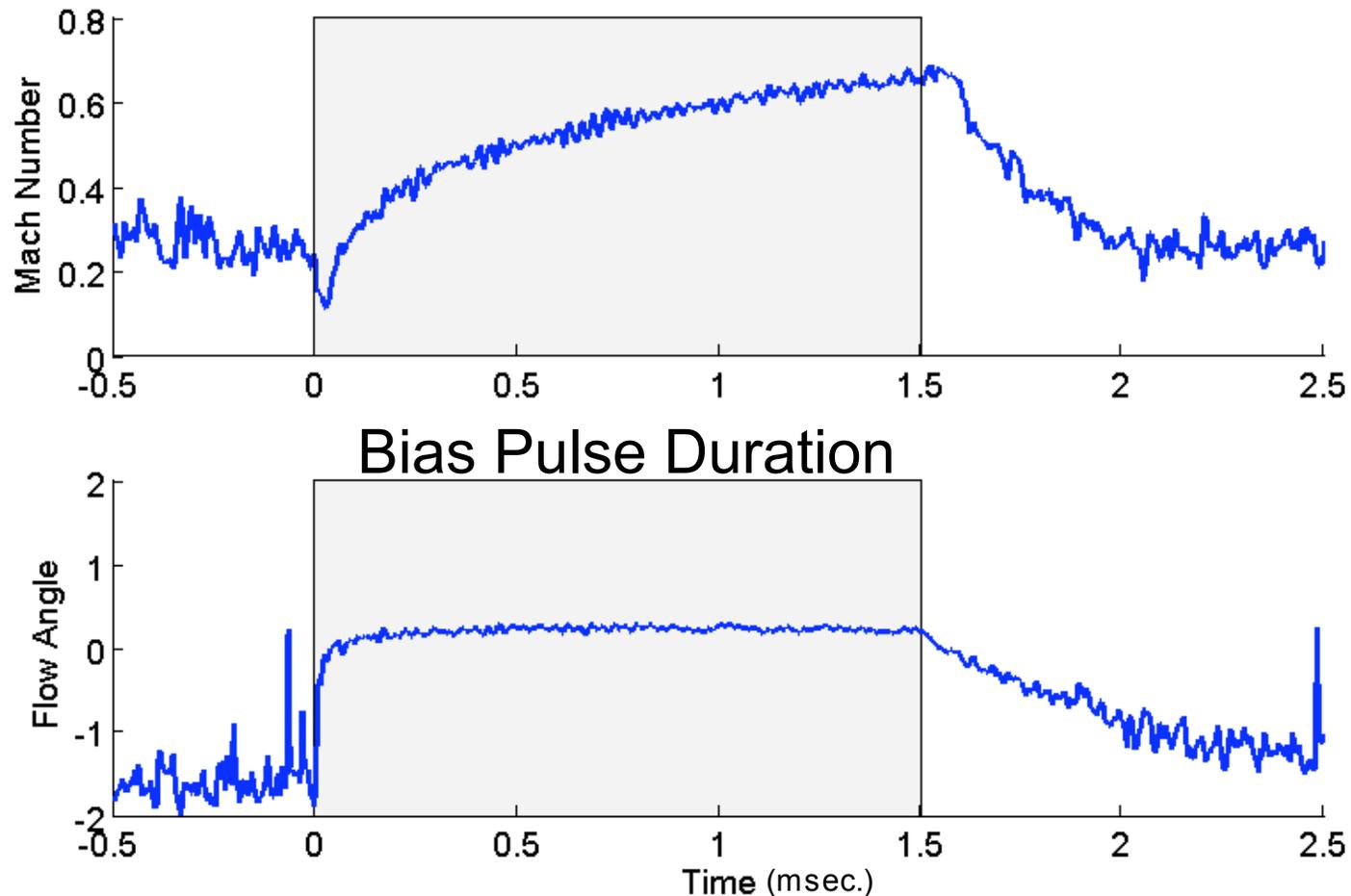
Two Time-Scale Model Fits Flow Evolution



Similar time-scales measured by LFS and HFS probes.



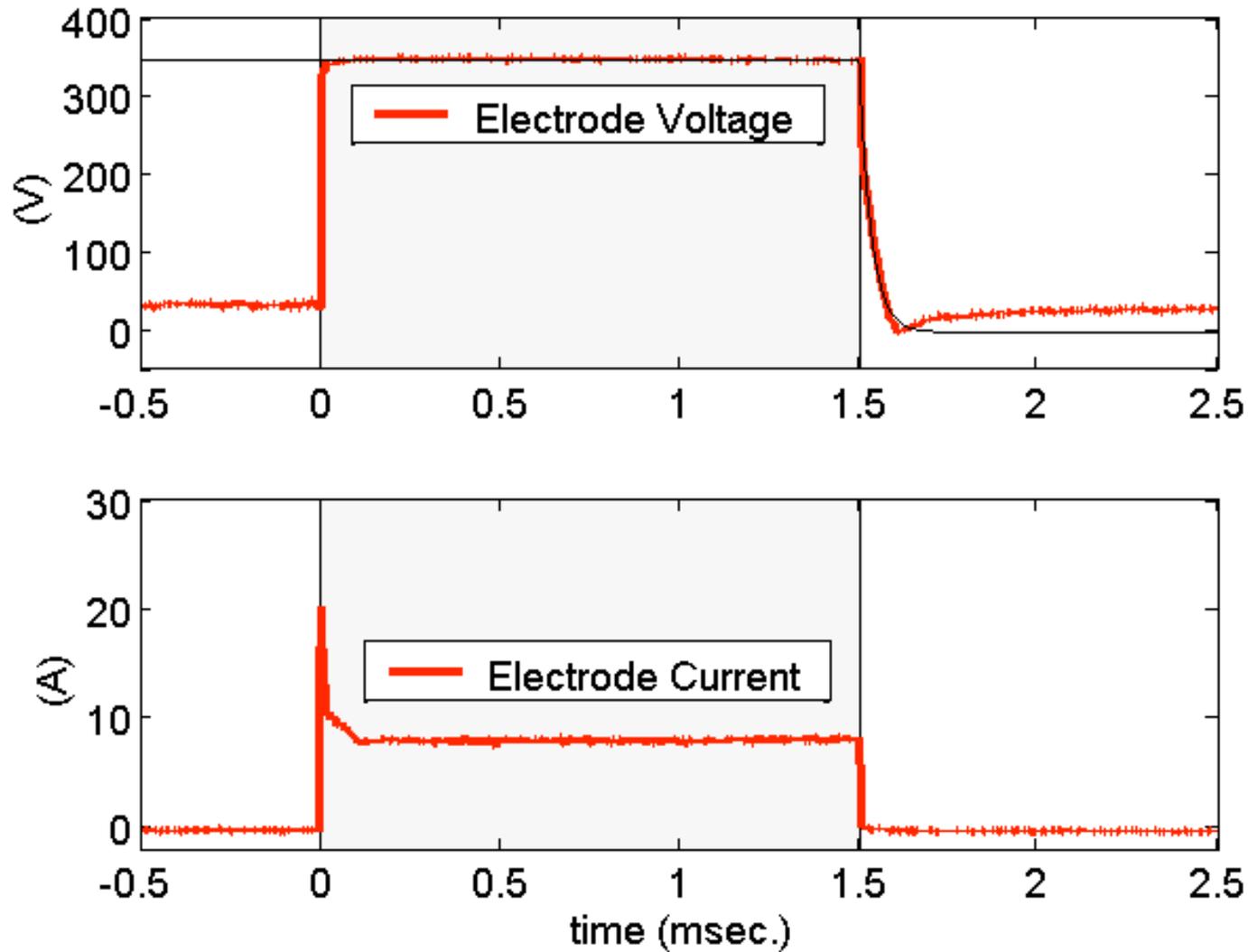
Both Flow Speed and Direction Evolve over the Electrode Pulse



Need to extract the time-scales and directions.

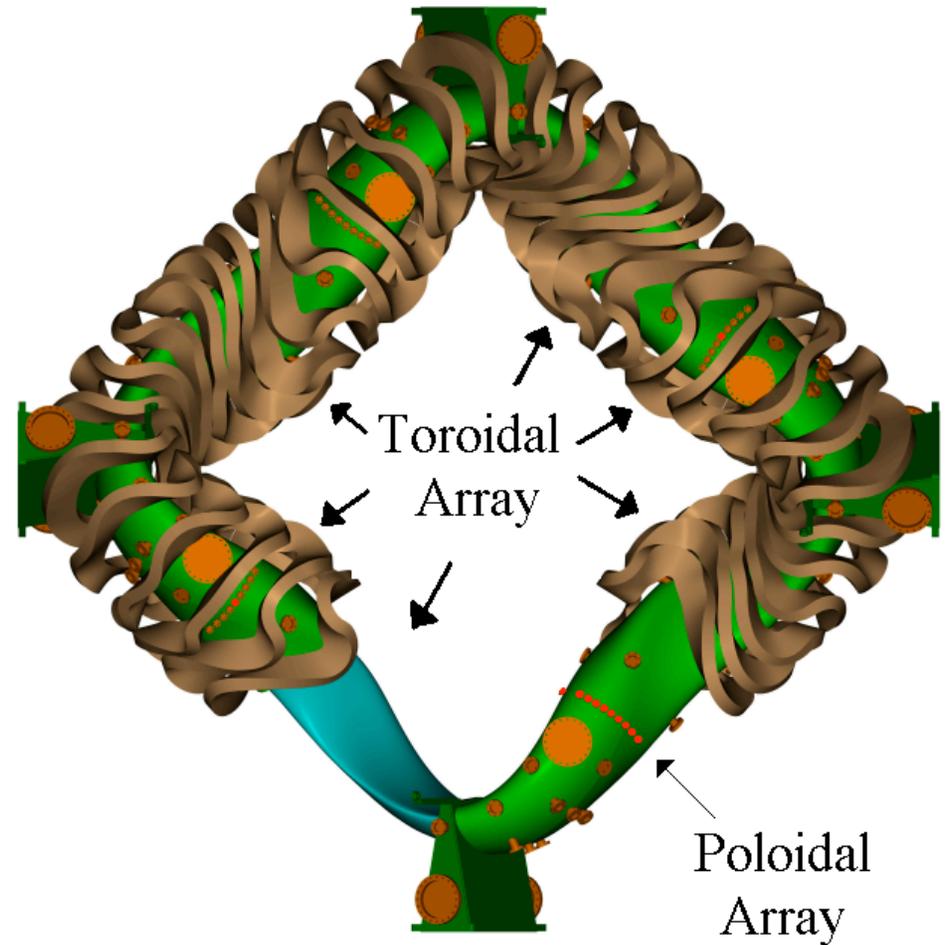
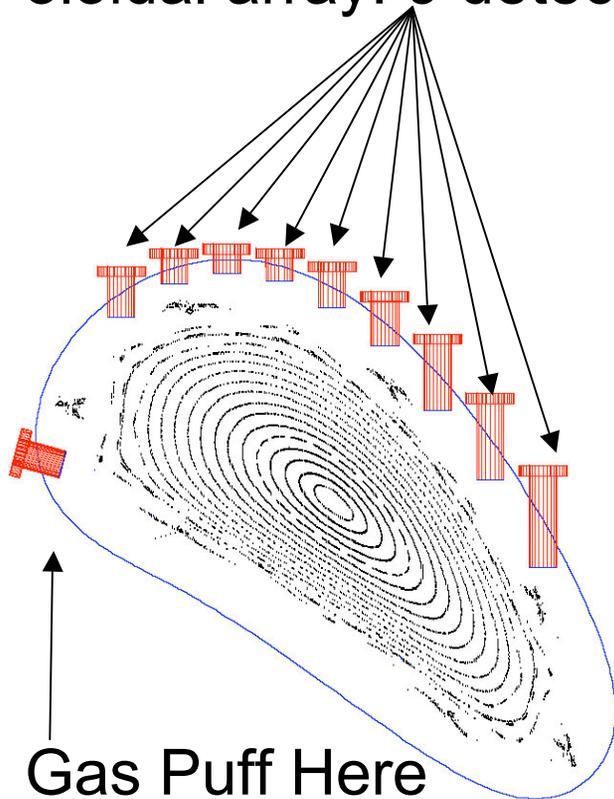


Voltage Application Initiates the Rise, Current Termination Initiates the Decay



Developed a Comprehensive Set of H_{α} Detectors for Neutral Density Measurements

- Toroidal array: 7 detectors on magnetically equivalent ports
- Poloidal array: 9 detectors



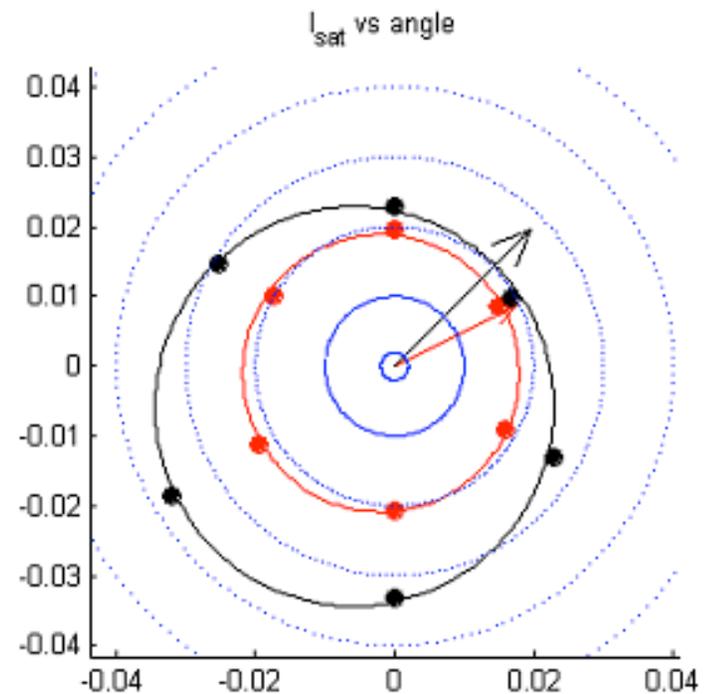
- All detectors absolutely calibrated
- Analysis done by J. Canik using DEGAS code



Mach Probes Used to Measure Time-Dependent Plasma Flows

- 6 tip mach probes measure plasma flow speed and direction on a magnetic surface.
- 2 similar probes are used to simultaneously measure the flow at high and low field locations, both on the outboard side of the torus.
- Data is analyzed using the unmagnetized model by Hutchinson.
- Time response of $\sim 10\text{-}20\mu\text{s}$

Looking \perp To The Magnetic Surface



$$I_{\text{sat}}(\theta) = A \exp\left(\left(\frac{M}{2}\right) \left[.64(1 - \cos(\theta - \theta_F)) + .7(1 + \cos(\theta - \theta_F))\right]\right)$$

- Probe measures V_f with a proud pin.



We Have Developed a Method to Calculate the Hamada Basis Vectors

- Method involves calculating the lab frame components of the contravariant basis vectors along a field line, similar to that by V.V. Nemov.

$$B^\Psi = \vec{B} \cdot \vec{\nabla} \Psi = 0 \quad \longleftarrow \text{Radial Basis Vector}$$

$$B^\zeta = \vec{B} \cdot \vec{\nabla} \zeta = \frac{1}{2\pi\sqrt{g}} \quad \longleftarrow \text{Toroidal Basis Vector}$$

$$B^\alpha = \vec{B} \cdot \vec{\nabla} \alpha = \frac{t}{2\pi\sqrt{g}} \quad \longleftarrow \text{Poloidal Basis Vector}$$

- Need initial condition on the basis vectors to complete this integration.
- Knowing (\sqrt{g}, t, B_α) at outboard symmetry plane is sufficient for calculating the initial conditions.
- Use two methods of computing the Pfirsch-Schlueter current to derive initial condition...

$$\mathbf{J}_{\parallel} = h \frac{\partial p}{\partial \psi} \mathbf{B}$$

Method by Nemov¹, h is numerically calculated

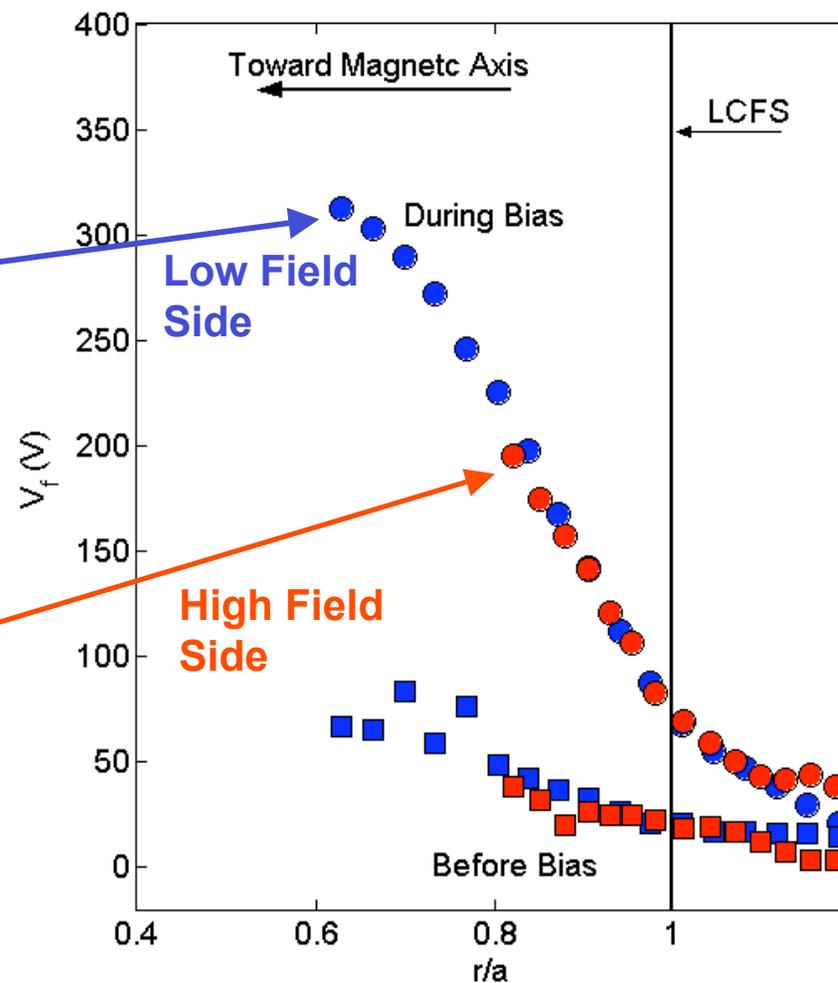
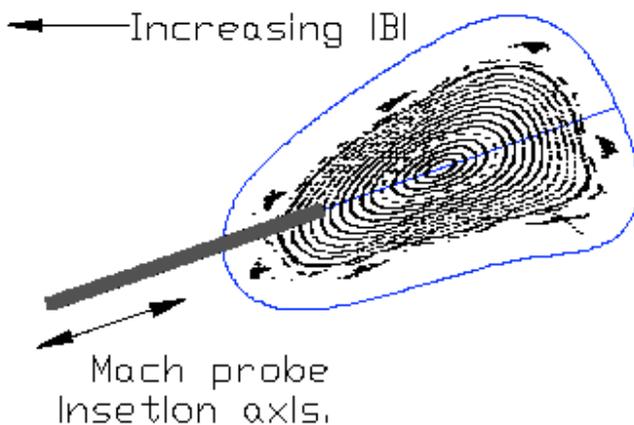
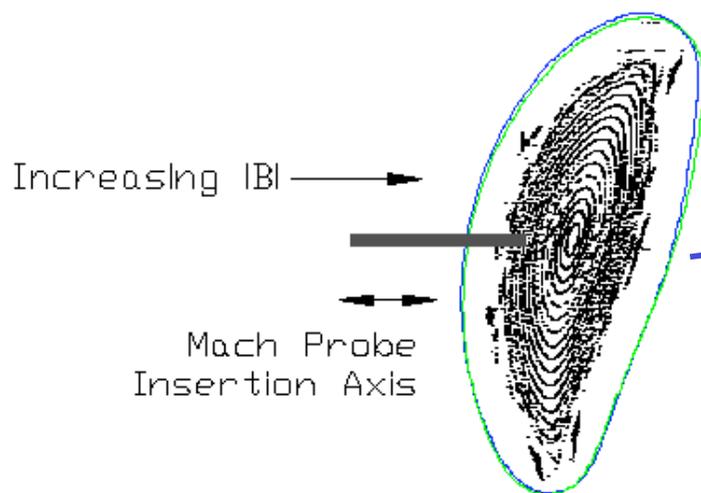
$$\mathbf{J}_{\parallel} = -\frac{B_\alpha}{B^2 B^\zeta \sqrt{g}} \frac{\partial p}{\partial \psi} \mathbf{B}$$

Method by Coronado and Wobig², B_α is the desired quantity

1) V.V. Nemov, Nuclear Fusion **30**, 927 (1990), 2) M. Coronado and H. Wobig Phys Fluids B **4**, 1294 (1992)

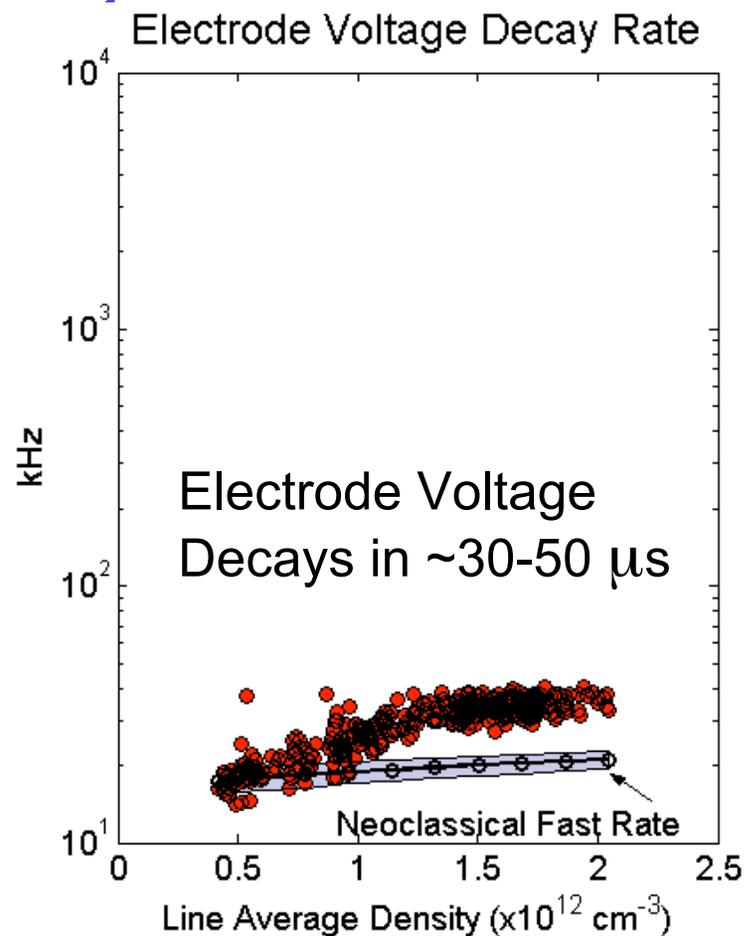
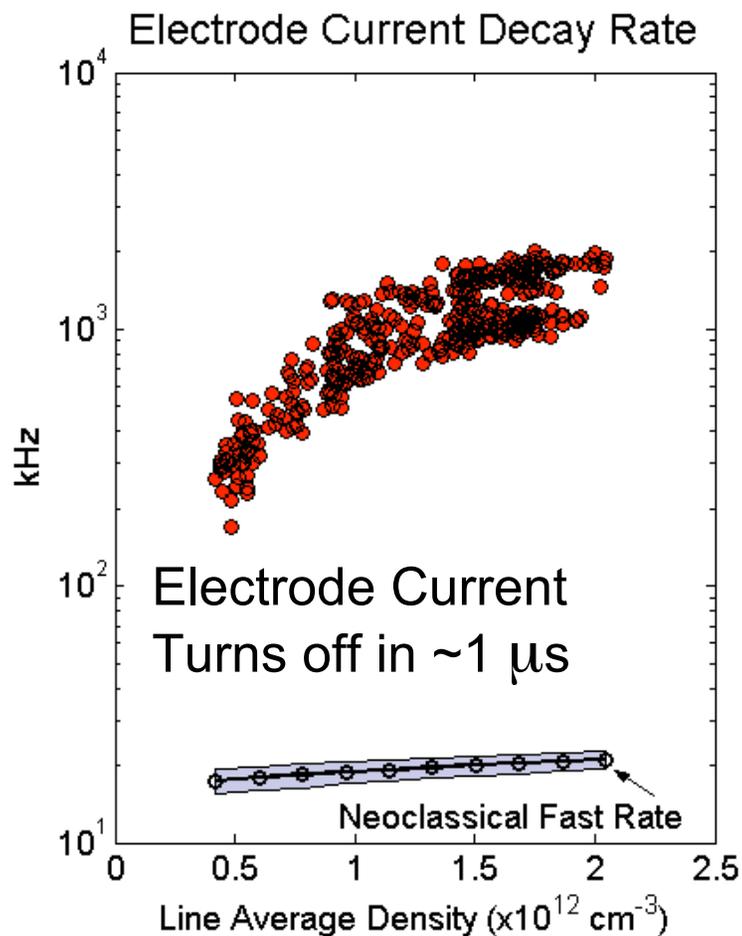


Floating Potential is a Flux Surface Quantity



Electrode Characteristics at Turn Off

Fit the Decay Model



Floating potential and fast component of flow decay on same time-scale as electrode voltage, in agreement with neoclassical fast rate.

