

# Local Electron Thermal Transport in the MST Reversed-Field Pinch

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Recent diagnostic developments, including an upgrade of the MST Thomson scattering system, have facilitated the first local power balance analysis of electron thermal transport in the MST. Experiments have focused on acquiring profile data for  $T_e$  and  $n_e$  under standard and enhanced confinement (PPCD) operating modes. Analysis shows that the thermal conductivity is lower during PPCD, and that the profile remains flat through the plasma core, decreasing sharply in the edge where the gradient of  $T_e$  is largest. Power deposition is calculated from the MSTFIT reconstructed current density profile and estimates of the neoclassical (Hirshman-Sigmar) resistivity (trapped particle corrections to resistivity are found to be significant for 2D RFP equilibria.) A Monte Carlo uncertainty analysis is done for all results of the transport analysis.

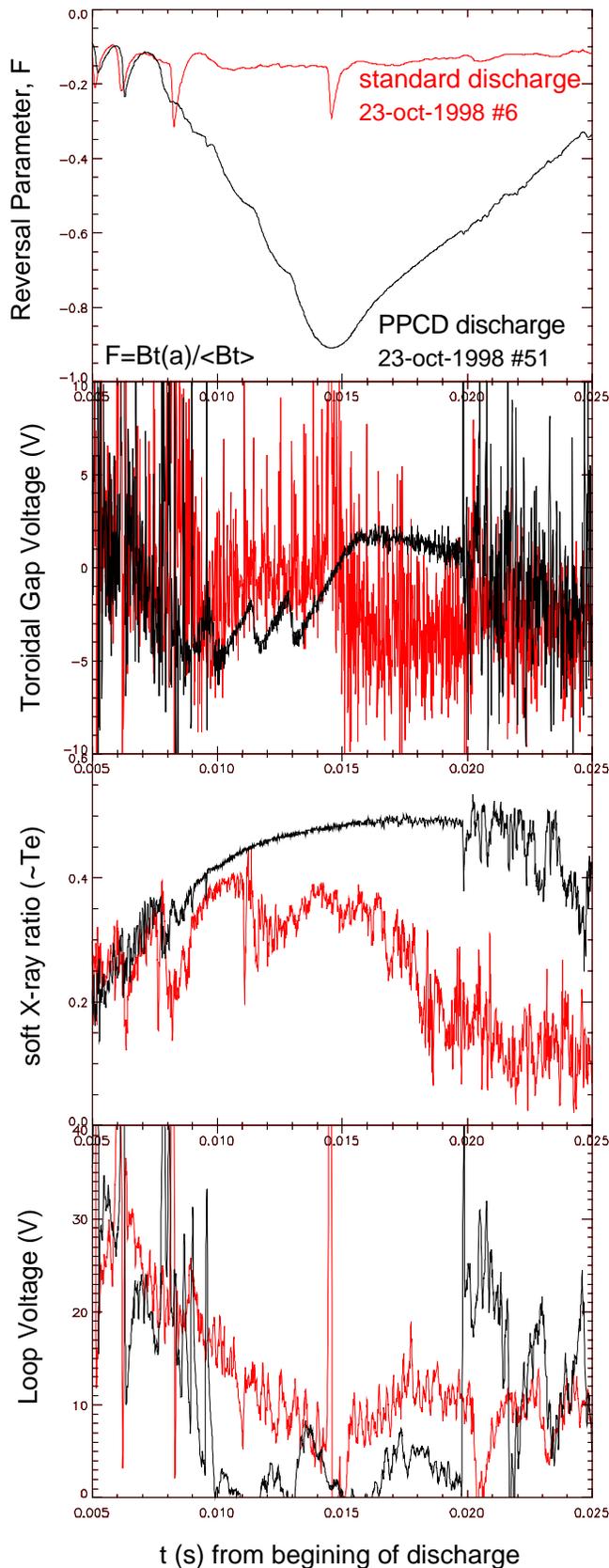
This work was supported by the U.S. D.O.E.

# Motivation

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- Ongoing improvements to the diagnostic arsenal of the Madison Symmetric Torus have initiated new investigations into the plasma physics of MST transport.
- The use of Pulsed Poloidal Current Drive (PPCD) in the past has demonstrated improved confinement in the MST, based on central electron temperature measurements and a hypothesized profile shape.
- Recent upgrades of the Thomson scattering system have facilitated measurements of the  $T_e$  profile in the MST out to  $r/a=0.88$ .
- Addition of an NIR Bremsstrahlung array has allowed an estimation of  $Z_{\text{eff}}$  in the MST. **We assume  $Z_{\text{eff}}=2$ .**
- These measurements, when coupled with density profiles from FIR Interferometry and MSTFIT reconstructed equilibria, have made it possible to calculate many transport quantities, including the electron thermal conductivity  $\kappa_e$  and the energy confinement time  $\tau_E$ .
- Monte Carlo analysis is used to ascribe error bands for transport quantities, rather than attempting to nonlinearly propagate the experimental uncertainty in the measured data.

# What is Pulsed Poloidal Current Drive?

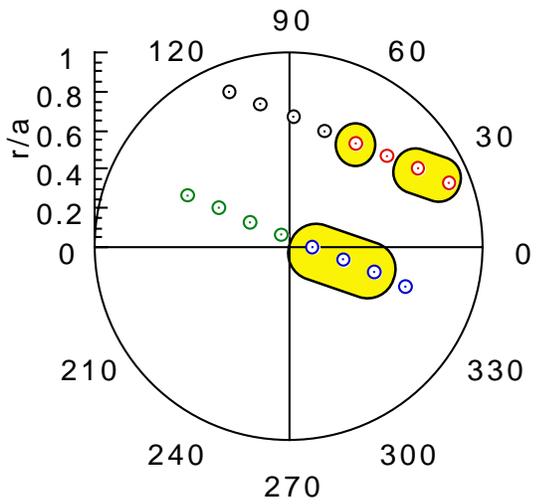
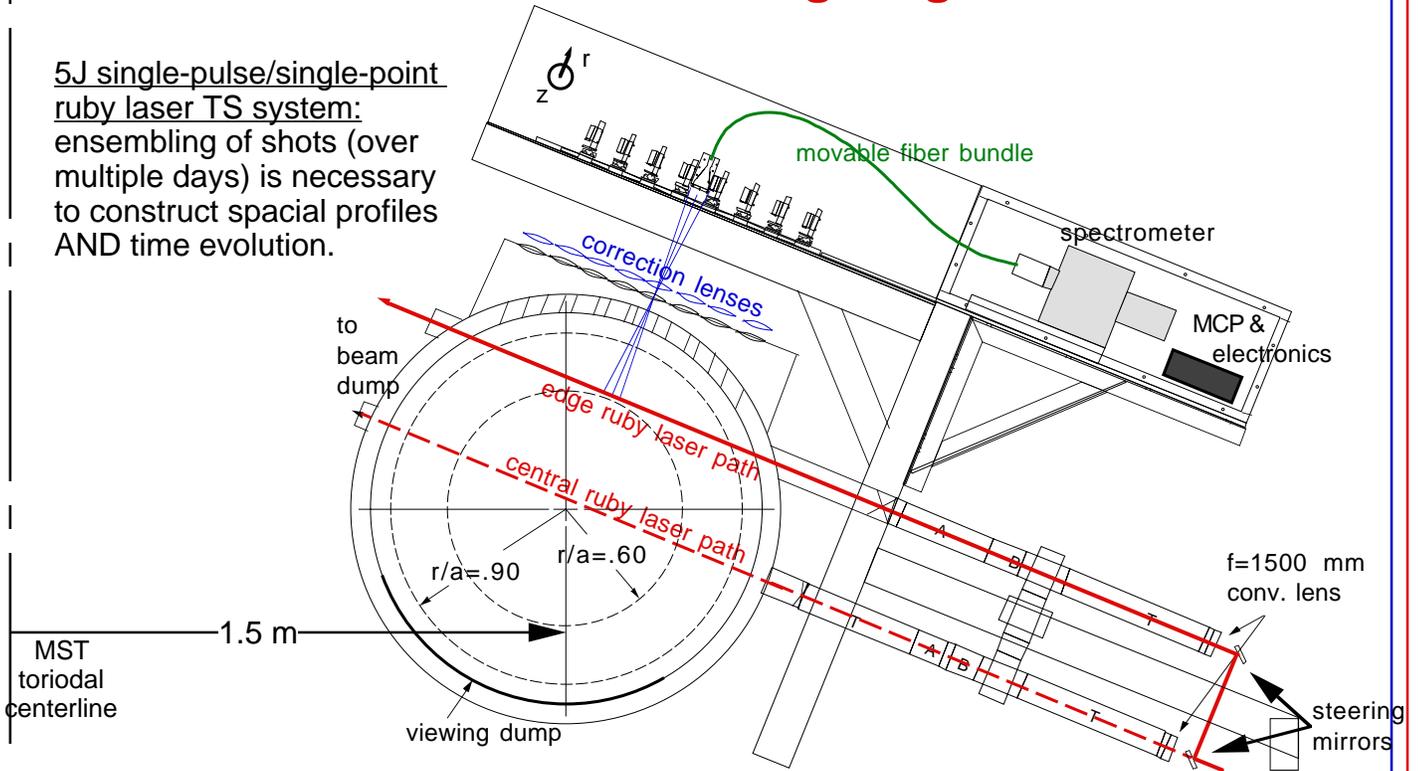


During a "normal" plasma discharge a series of 5 voltage pulses are applied to the MST shell to produce a magnetic field which is in the opposite sense to the established toroidal magnetic field, forcing flux out of the MST. This opposing  $B$  is thought to cause a poloidally driven edge current.

PPCD is observed to improve machine performance. It is believed that magnetic modes in the MST derive their free energy from the gradient in the current profile, which is steepest at the edge. By inductively driving current in the edge, PPCD could reduce the gradient and remove free energy from the magnetic modes, resulting in decreased magnetic fluctuations and hence improved particle and energy confinement.

# MST Thomson Scattering Diagnostic

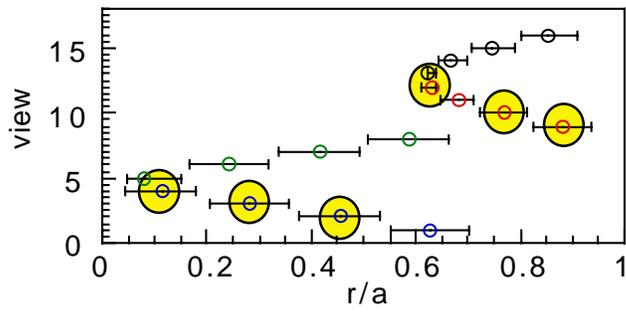
5J single-pulse/single-point ruby laser TS system: ensembling of shots (over multiple days) is necessary to construct spacial profiles AND time evolution.



## Thomson Scattering Views

view	r/a	$\phi$
1	0.627	-18.46
2	0.455	-16.93
3	0.283	-13.53
4	0.115	0
5	0.080	123.91
6	0.244	147.06
7	0.415	151.39
8	0.587	153.19
9	0.882	22.28
10	0.769	31.41
11	0.681	43.26
12	0.630	57.76
13	0.625	73.61
14	0.666	88.59
15	0.746	101.10
16	0.854	110.82

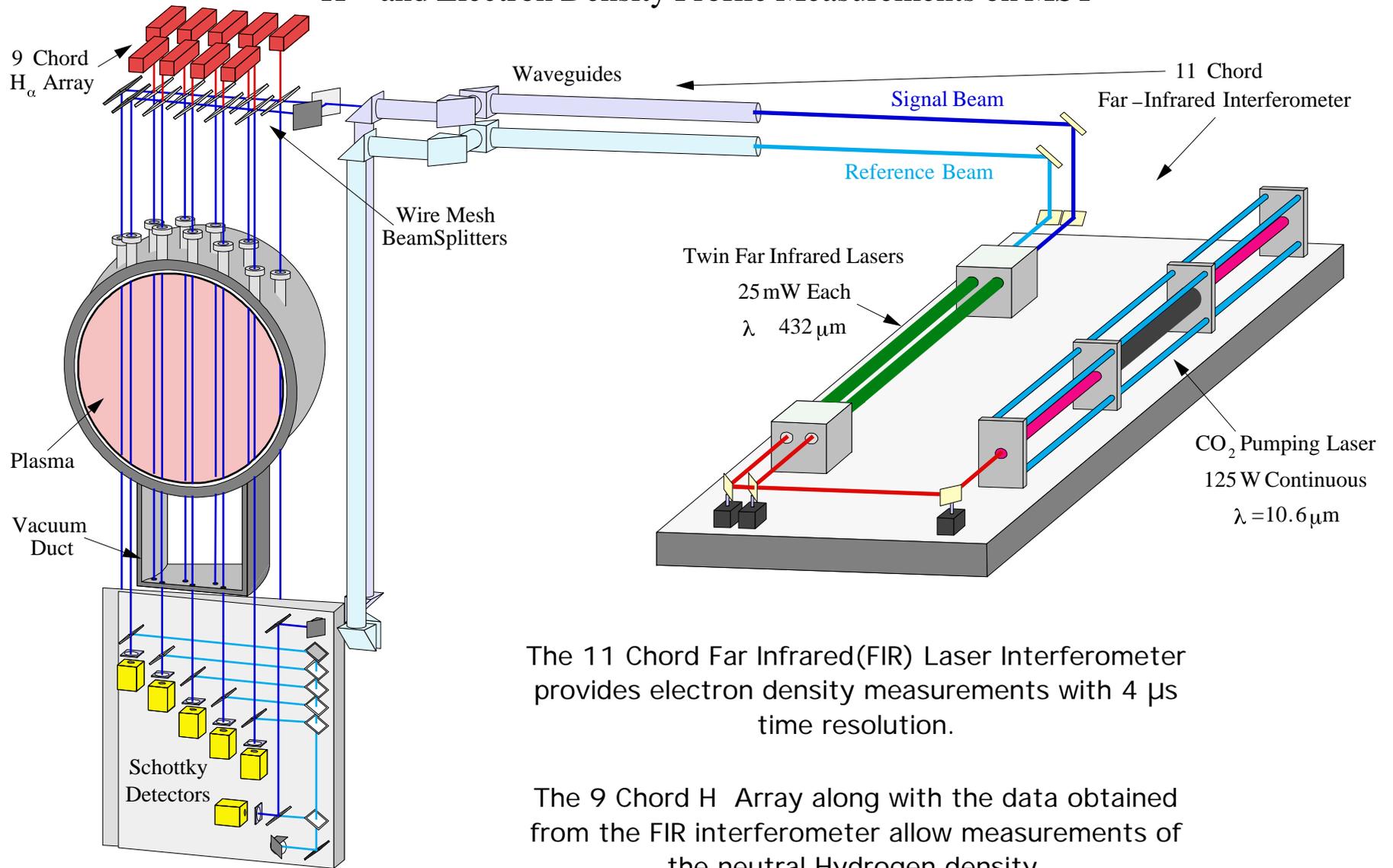
chords used



error bars indicate radial range of scattering volume on each view

# Far Infrared Interferometer Diagnostic

## H and Electron Density Profile Measurements on MST



The 11 Chord Far Infrared(FIR) Laser Interferometer provides electron density measurements with 4  $\mu\text{s}$  time resolution.

The 9 Chord H Array along with the data obtained from the FIR interferometer allow measurements of the neutral Hydrogen density.

# MSTFIT and Monte Carlo Error Analysis

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- MSTFIT is a fully toroidal (2D) equilibrium reconstruction code that utilizes measured data to constrain the reconstruction.
- Plasma Parameters:  $F$ ,  $\beta$ , and  $I_p$  are entered along with edge magnetic probe signals to manage the reconstruction of equilibrium quantities.
- Once an equilibrium is found, it becomes easier to invert chord averaged quantities, such as FIR measured electron density. The inversions are done in flux coordinate geometry, which can differ significantly from machine geometry.
- Diagnostic data such as  $T_e$ ,  $T_i$ ,  $n_e$ ,  $dT_e/dt$ ,  $dn_e/dt$ , and  $Z_{\text{eff}}$  are then combined with reconstructed quantities to ultimately calculate profiles of transport coefficients.
- By randomly varying the measured quantities within their error bars and collecting the calculated transport quantities, a Monte Carlo type "error band" is derived.

# MSTFIT is a fully toroidal, 2D equilibrium reconstruction code

Green's function technique is used to relate current and flux

Green's function relates flux to current distribution

$$\psi_{i,j} = \sum_{\text{plasma } m,n} G_{i,j,m,n} I_{m,n} + \sum_{\text{vessel } k} G_{i,j,k} I_k$$

$$G(\mathbf{x}, \mathbf{x}') = \frac{\mu_0}{2\pi} \left( \left( 1 - \frac{1}{2} k^2 \right) K(k) - E(k) \right); k^2 = \frac{4 R R'}{(R+R')^2 + Z^2}$$

plasma current is modelled by two flux functions

$$I_p(\mathbf{x}) = \frac{F(\hat{\psi}(\mathbf{x}))F'(\hat{\psi}(\mathbf{x}))}{\mu_0 R(\mathbf{x})} + R(\mathbf{x})P'(\hat{\psi}(\mathbf{x}))$$

data are also linearly related to current distribution for example

$$\text{flux: } \psi_v = \sum_{\text{plasma } m,n} G_{v,m,n} I_{m,n} + \sum_{\text{vessel } k} G_{v,k} I_k$$

$$\text{probes: } B_{\theta_p} = \sum_{\text{plasma } m,n} \left. \frac{\partial G}{\partial n} \right|_{p,m,n} I_{m,n} + \sum_{\text{vessel } k} \left. \frac{\partial G}{\partial n} \right|_{p,k} I_k$$

Current profile is specified by Grad-Shafranov equation

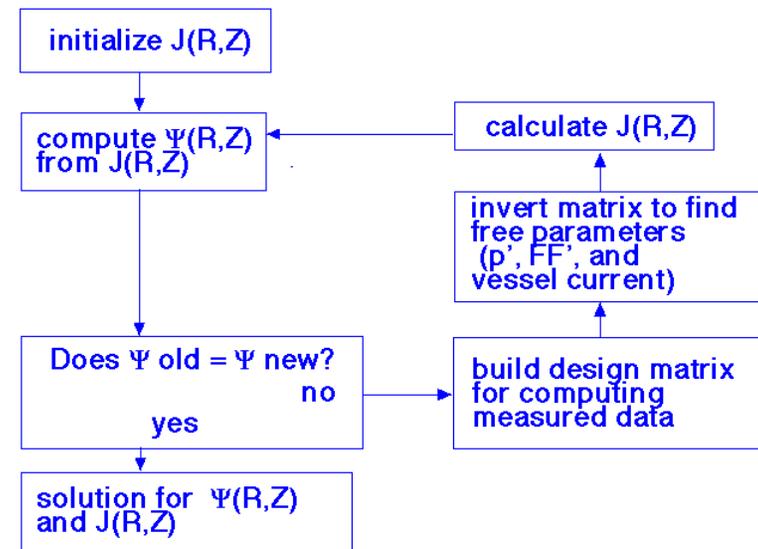
$$J_\varphi = \frac{FF'}{\mu_0 R} + R P'$$

- FF' and p' are the free functions, profiles need to be specified to determine toroidal current density
- MSTFit uses free parameters to determine each profile using basis functions.

$$F'(\hat{\psi}) = F'_0(\hat{\psi}) + \sum_n a_n F'_n(\hat{\psi}) \quad P(\hat{\psi}) = \sum_n b_n P_n(\hat{\psi})$$

MSTFIT also uses free parameters to specify the current distribution in the vessel

Solution is obtained by successive approximations for flux on plasma grid

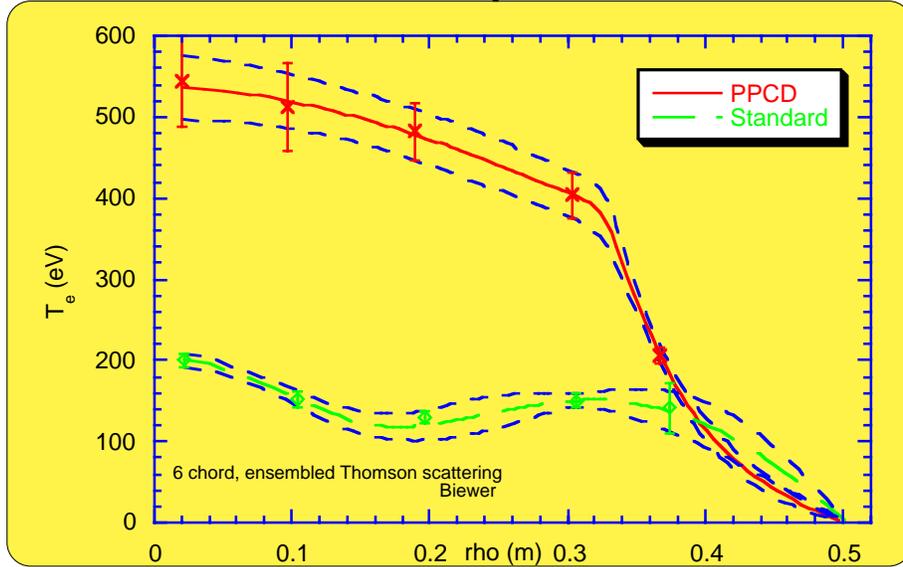


Splines are used as the basis functions to construct profiles

- The i-th basis function is given by the spline of a unit height at the i-th knot location with all other knot values set to 0.
- Knot locations are pre-determined to well represent typical RFP profiles.
- The linear combination of the basis functions identically reproduces the spline interpolation of all values.

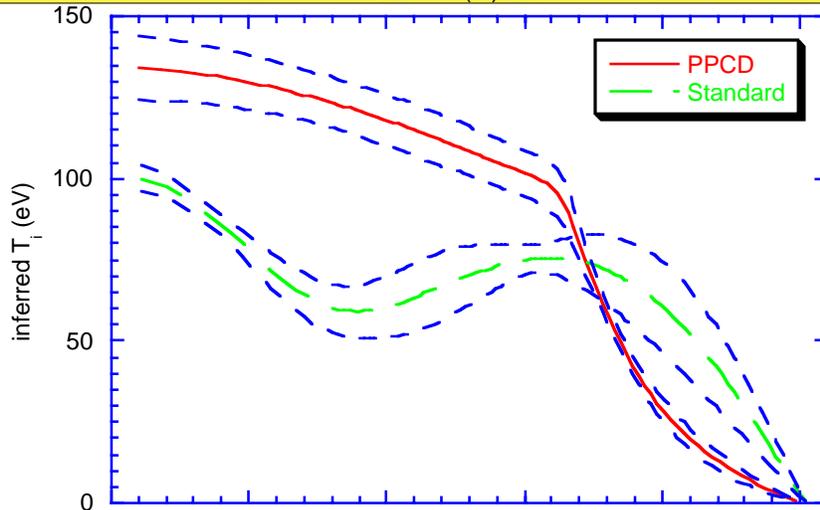
# Measured Profiles

$I_p \sim 200$  kA

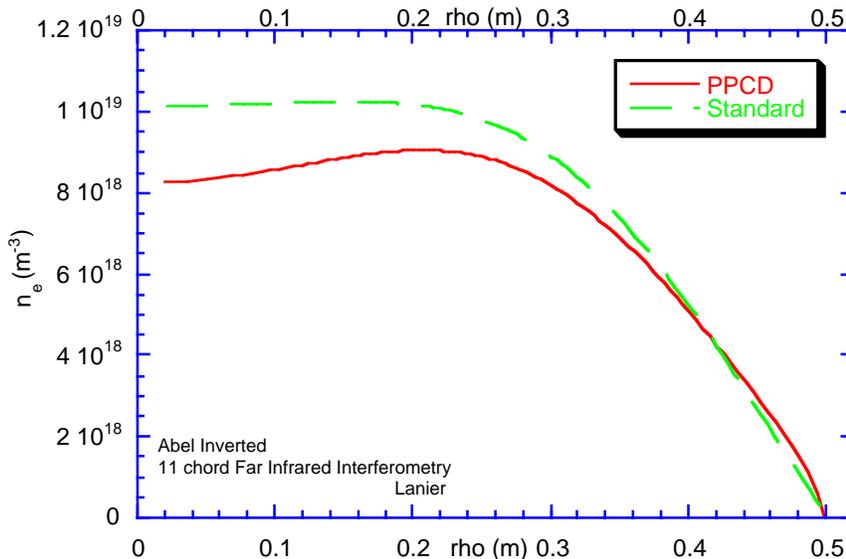


First 6 point electron temperature profiles measured on the MST.

Thomson Scattering measurements are assembled over multiple days and are susceptible to changes in machine conditions.



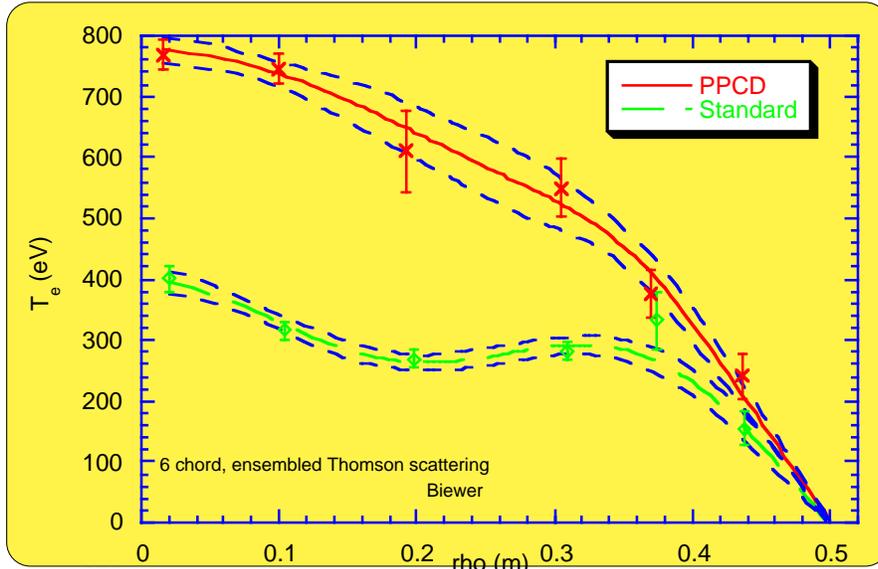
The assumptions that  $T_i = 1/4 T_e$  for PPCD and  $T_i = 1/2 T_e$  for Standard discharges are supported by measurements from majority and minority impurity diagnostics.



FIR Interferometry measured density profiles are Abel inverted using MSTFIT, a fully toroidal equilibrium reconstruction code.

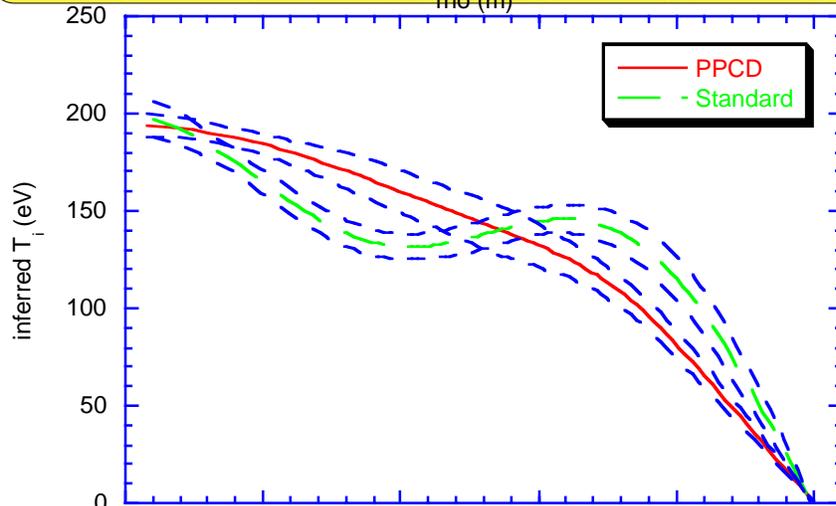
# Measured Profiles

## $I_p \sim 400$ kA

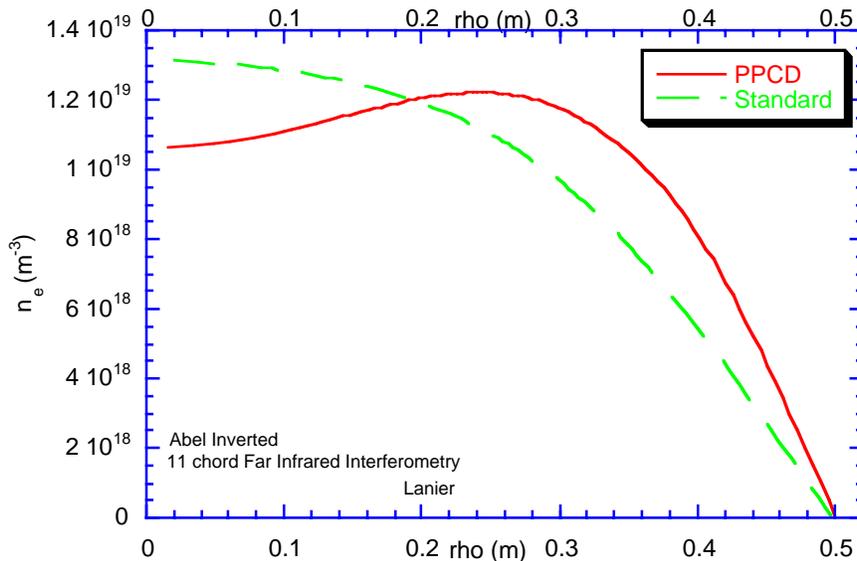


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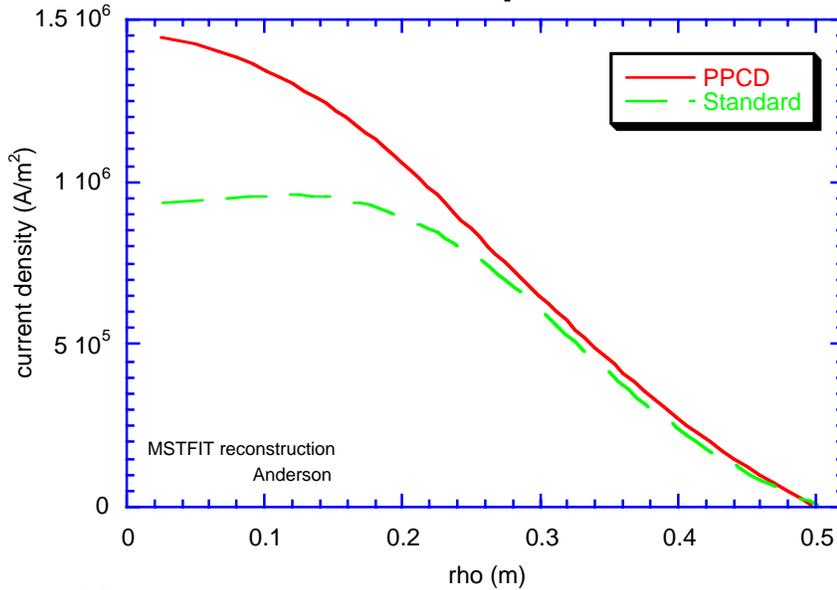
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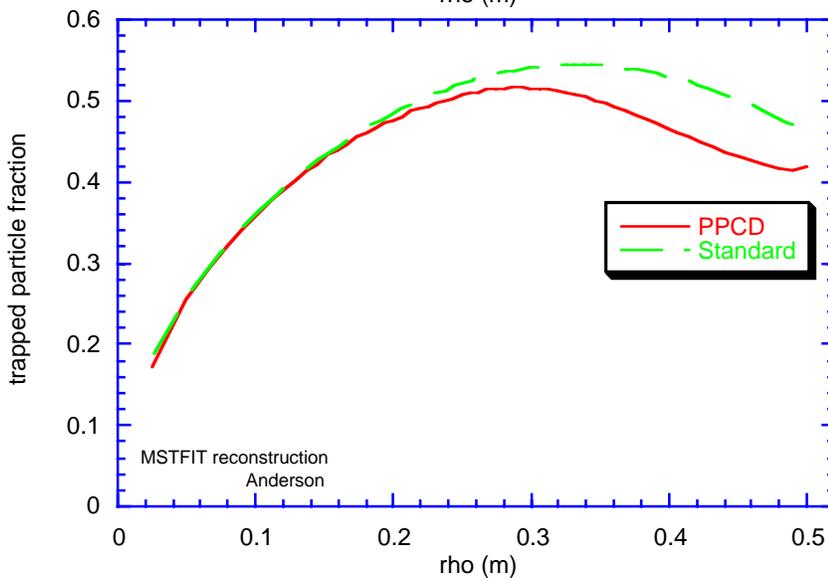
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# Reconstructed Profiles

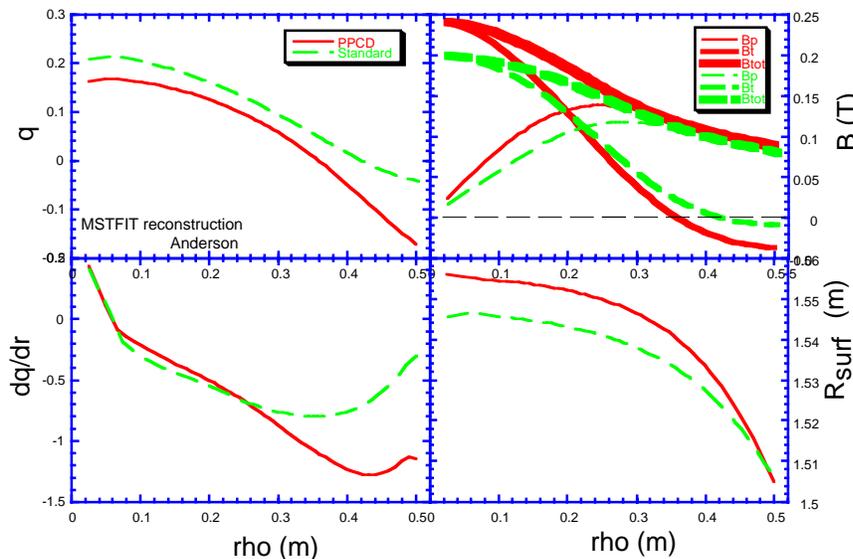
## $I_p \sim 200$ kA



The current density on axis clearly increases, whereas in the edge there is little apparent change.



MSTFIT predicts that a substantial fraction of the particles are magnetically trapped in the MST during both PPCD and Standard discharges.

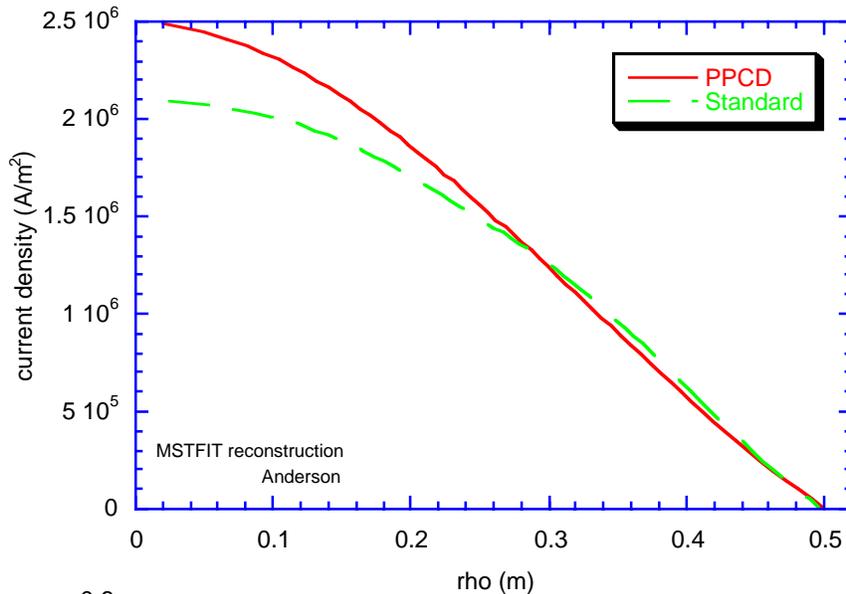


The application of PPCD clearly shifts the flux reversal surface inward, as shown by observing the zero crossing of the toroidal magnetic field.

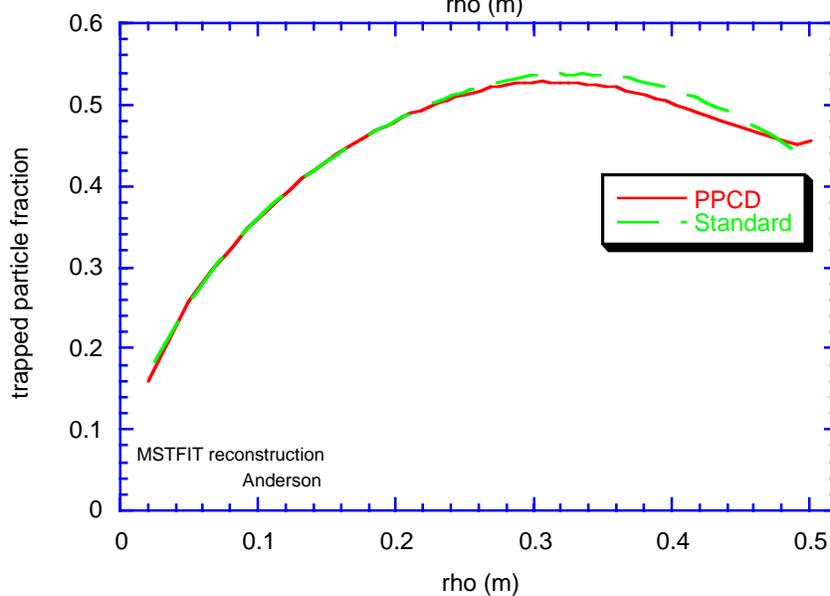
Also, PPCD leads to an increase in the amount of magnetic shear at the edge of the plasma.

# Reconstructed Profiles

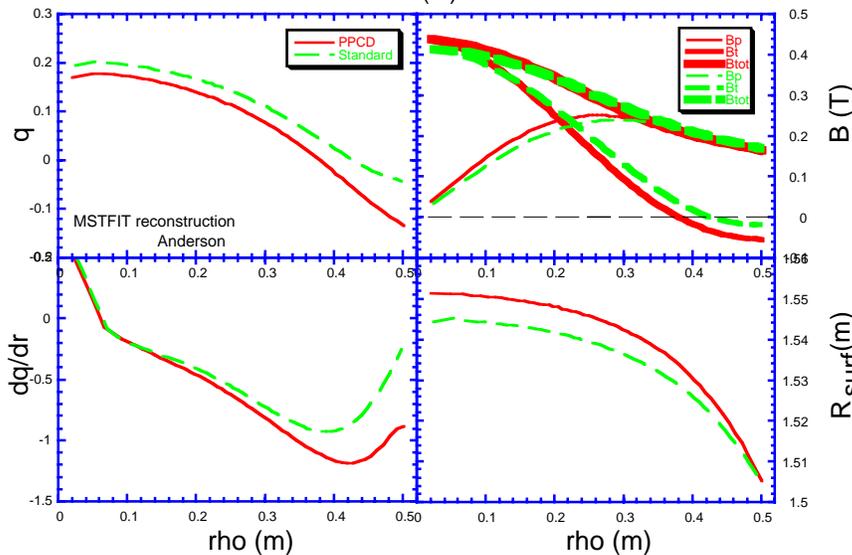
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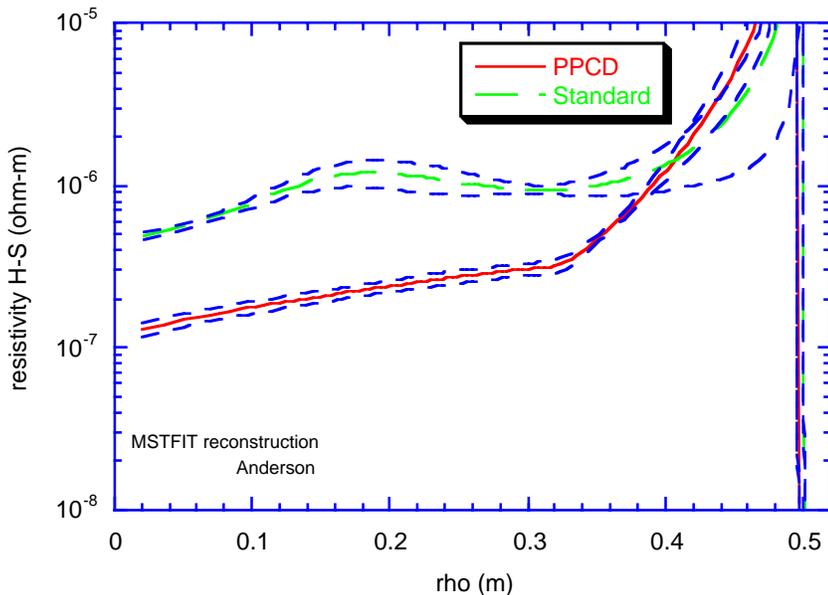


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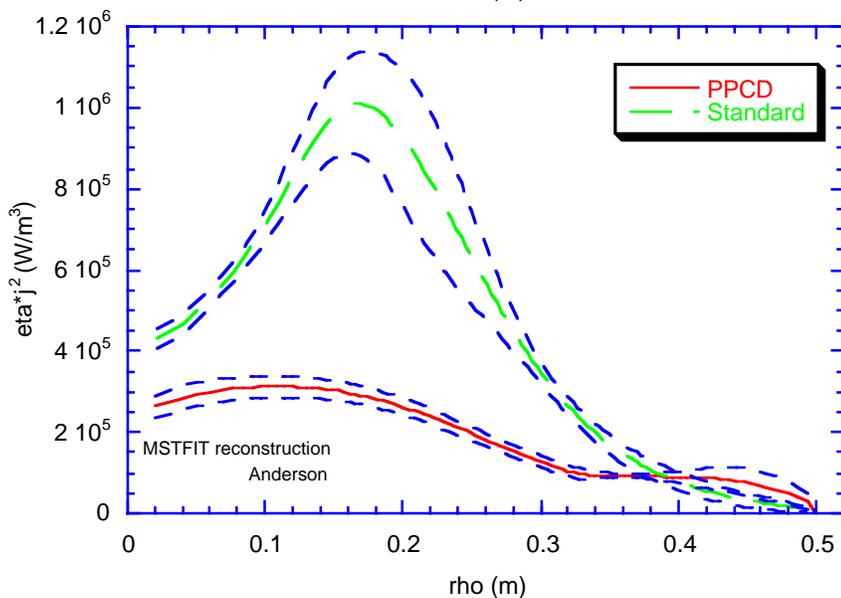
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# Calculated Profiles

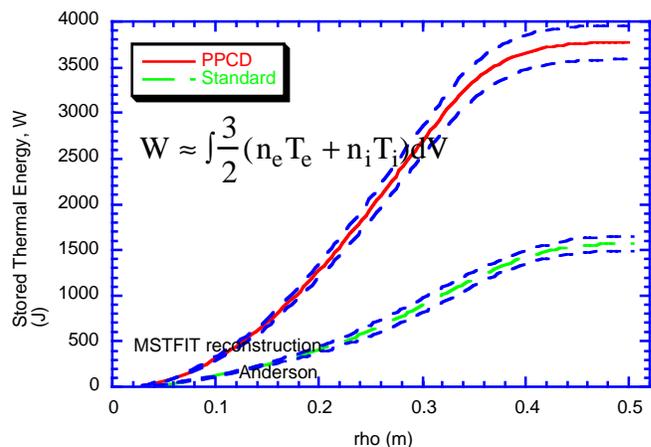
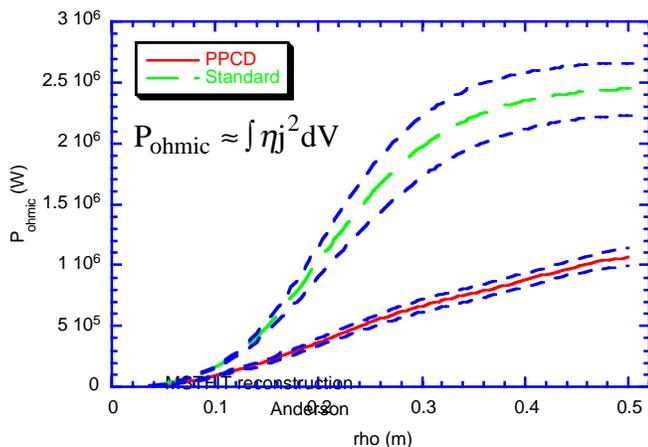
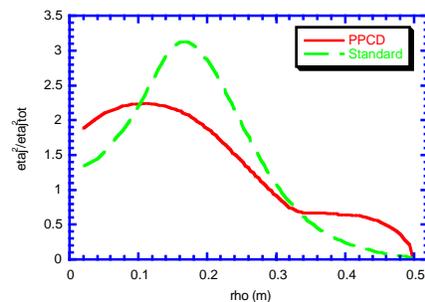
$I_p \sim 200$  kA



Neoclassical effects become very important when modeling in 2D. Using Hirshman-Sigmar resistivity (rather than Spitzer) in "Standard" reconstructions matches the measured and calculated ohmic input power with a conservative estimate of  $Z_{eff}=2$ . I.e. No need of "anomalous" resistivity to account for MHD dynamo action.

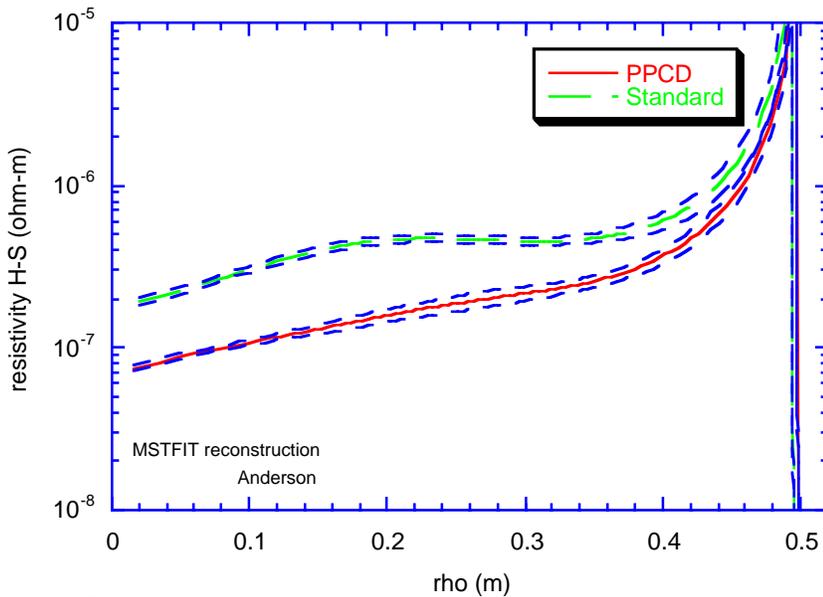


Normalizing to total input power shows that power deposition is **enhanced in the edge during PPCD** as compared to Standard discharges at this current.

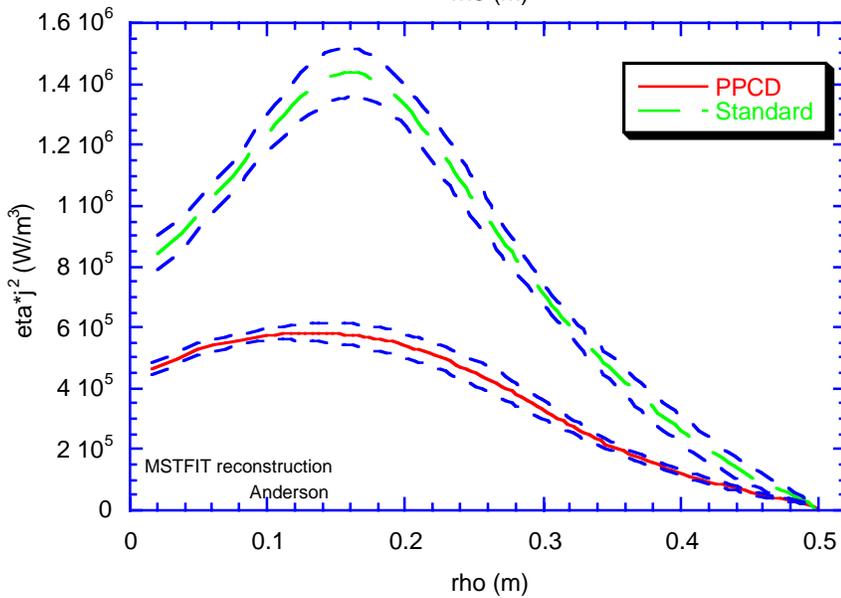


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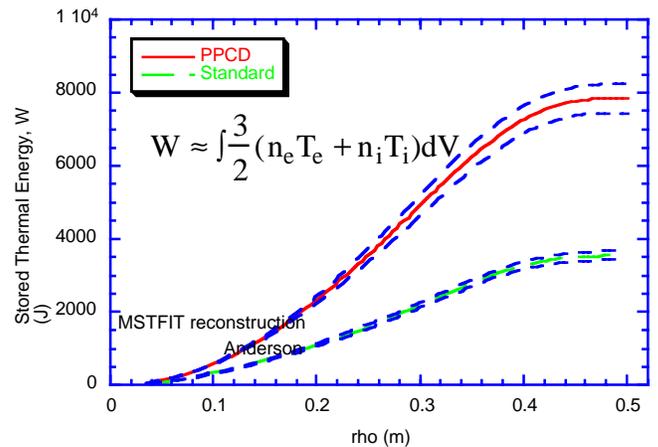
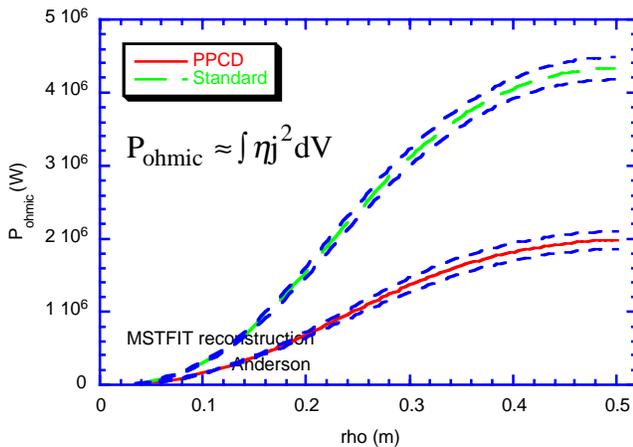
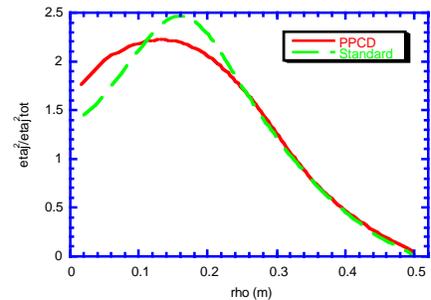
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Normalizing to total input power shows that the power deposition profile is nearly identical during PPCD and Standard discharges at this current.



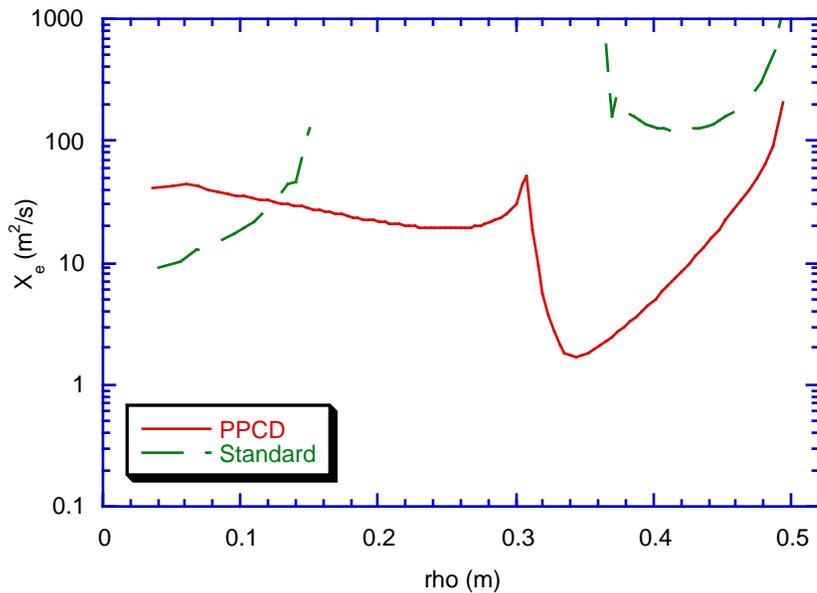
# Main Results

200 kA

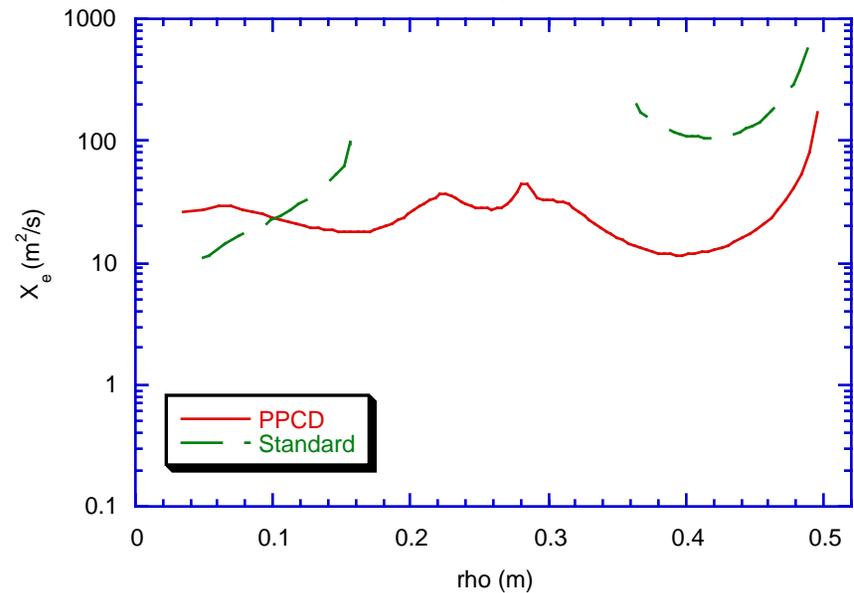
400 kA

	Standard		PPCD			Standard		PPCD		
Pohmic	2.45 +/- 0.22	1.07 +/- 0.08	MW	Pohmic	4.34 +/- 0.16	1.99 +/- 0.12	MW			
W	2.0 +/- 0.1	4.2 +/- 0.2	kW	W	4.4 +/- 0.2	8.8 +/- 0.4	kW			
dW/dt		383.1 +/- 28.6	kW/dt	dW/dt		133.1 +/- 38.1	kW/dt			
E	0.81 +/- 0.10	6.45 +/- 1.28	ms	E	1.02 +/- 0.07	4.83 +/- 0.65	ms			
pol	6.36 +/- 0.33	13.28 +/- 0.66	%	pol	3.29 +/- 0.11	7.52 +/- 0.38	%			

$$\tau_E = \frac{W}{P_{ohmic} - \frac{W}{t}}$$



$$\beta_p = \frac{\frac{1}{V} (n_e T_e + n_i T_i) dV}{\frac{1}{2\mu_0} B_p^2}$$



$$e \frac{1}{n_e T_e} \frac{-\eta_{HS} j^2 dV + \frac{\partial}{\partial t} \frac{3}{2} n_e T_e dV}{4\pi^2 R_0 r} - \frac{5}{2} D_{meas.} T_e n_e$$

$$D_{Class.} = \frac{\eta_{HS} (n_e T_e + n_i T_i)}{B_{tot}^2} \quad 0.005 \text{ m}^2/\text{s}$$

for comparison:  $D_{Bohm} = \frac{T_e}{16eB_{tot}} \quad 100 \text{ m}^2/\text{s} \quad D_{meas.}$

# Conclusions

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- For the first time, 6 point **profiles of  $T_e$  have been measured** during PPCD in the MST.
- **Low  $T_e$  in the edge is coincident with enhanced magnetic shear in the edge.** During PPCD the flux surface where toroidal magnetic field reverses direction is forced inward (i.e. F is deepened), leading to a steepening of the q profile, particularly in the edge.
- It is **unnecessary to invoke any anomalous resistivity during "standard" MST discharges** if the fraction of trapped particles is taken into account, along with a neoclassical (as per Hirshman and Sigmar) assessment of resistivity. With a very reasonable estimation of  $Z_{\text{eff}}=2$  the reconstructed input power matches the measured input power.
- **The energy confinement time greatly improves during PPCD,** presumably due to the reduction of magnetic fluctuations.
  - »  $\tau_E \sim 6 \text{ ms}$  (if  $Z_{\text{eff}} \sim 2$ ) from 1 ms
- **$\beta_p$  reaches 13% during low current PPCD.**