Progress on NIMROD disruption mitigation modeling

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Recent efforts focus on many aspects of the runaway electron problem

Three topics:

- D₂ Dilution cooling as an "optimal" runaway suppression mechanism (vast majority)
- Effects of plasma elongation on MHD and runaway confinement (very brief)
- Direct calculations of suprathermal electron acceleration and confinement (also brief)





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Overview of the runaway electron problem

- The crux of the problem:
 - Disruptions \rightarrow large E fields \rightarrow high energy runaway electrons
 - Avalanche amplification of runaways:

 $A = exp(\gamma t) \approx exp(2.5I_p)$ (Huge in ITER)

- Solutions?
 - Collisional suppression

$$E_{crit} = 0.12n_{e,20}$$

 Confinement time shorter than acceleration time





D₂ dilution should be optimal for collisional suppression

$$E/E_{crit} = \eta(n,T)j/0.12n_{e,20} \propto T^{-3/2}n^{-1}$$

- The least cooling for the largest density increase minimizes E/E_{crit}
- Imagine cooling purely by D₂ dilution, neglecting all radiation and atomic physics. Then T~n⁻¹. In this case, $E/E_{crit} \sim n^{1/2}$
- When the temperature drops more strongly due to radiative cooling, then E/E_{crit} rises more sharply with density. Since the thermal quench precedes the current quench, E/E_{crit} always gets worse before it gets better
- In normal DIII-D operation, assume T=4keV, n=8x10¹⁹/m³, j=2x10⁶A/m². Then we have **E/E_{crit}=0.09**
- For ITER nominal parameters of T=8.9keV, n=10²⁰/m3, j=1.4x10⁶A/m², this gives us E/E_{crit}=0.01





DIII-D Simulation assumes 100x density increase, in-situ carbon



A uniform carbon density of 1% of the pre-dilution core electron density (8.6x10¹⁷/m³) is assumed. At this initial T_e (~40eV), the physical value of Spitzer resistivity can be used for the simulation without numerical difficulty





Both resistivity and thermal conduction impact occurrence/amplitude of MHD



Two perpendicular heat transport models are considered:

Pfirsch-Schluter-

$$\chi_{\perp}(m^{2}/s) = \frac{1.411 \times 10^{-21} n \ln \Lambda_{i}}{B^{2} T^{1/2}} (1+1.6q^{2})$$

Braginskii-
$$\chi_{\perp}(m^{2}/s) = \frac{2kT}{m_{i} \omega_{ci}^{2} \tau_{i}} \approx \frac{2.4 \times 10^{-20} n}{B^{2} T^{1/2}}$$

Difference is ~3-4x

Two values of resistivity, Spitzer and a factor of 2.3 higher





Significant current peaking can occur



Spitzer resistivity, Braginskii transport





Flux surfaces tend to re-heal in the current quench

The case with the most significant MHD (P-S transport, 2.3 times Spitzer resistivity) sees nearly all flux surfaces destroyed after 1/1 crash, but the flux surfaces mostly reform by the end of the current quench







Rosenbluth ratio in DIII-D is bad news







ITER simulation has 150x density multiplication

- An ITER equilibrium generated by L. Lao is used
- D_2 dilution cooling by a factor of 150 is assumed for the initial condition, where the post dilution density is assumed to be a uniform value of 1.5×10^{22} /m³



•A uniform beryllium density of 1% of the predilution electron density (10¹⁸/m³) is assumed. The beryliium radiation is comparable to the bremsstrahlung in some regions, but does not dominate the overall radiated power

•Simulation is run at actual Spitzer resistivity, P-S transport





Thermal quench is MHD-free







More uniform cooling, little current peaking







Rosenbluth ratio is much better in ITER







Dilution cooling conclusions

- You can't really beat the Rosenbluth criteria in DIII-D, but demonstration of massive particle injection in the core maybe sufficient for ITER
- ITER shows less propensity for MHD in this mitigation scenario. Also, flux surfaces heal during the current quench in DIII-D
- Particle loss in the NIMROD simulations is the biggest issue for ITER
 but I have no idea if this is real





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Does elongation effect runaway electron confinement during disruptions?

Some tokamaks tend to observe RE's during some current quenches:

FTU, Tore-Supra, TEXTOR all run circular, limited plasmas

JET ran only limited plasmas for a number of years before the installation of its first divertor. Disruption runaways were much more prevalent back then, compared to now.

JT-60U is diverted, with low elongation

Some tokamaks don't see RE's during the current quench (except perhaps during killer pellet experiments):

DIII-D, ASDEX-U, and C-Mod run diverted, elongated plasmas (vertically unstable)

This suggests that elongation and/or vertical stability might have something to do with generation of runaways during a disruption.





C-Mod low elongation simulation in progress



Previous high elongation C-Mod simulation with Ne gas jet. Higher n modes grow first, fast growth of n=1 tends to trigger thermal quench



Low elongation simulation with other wise similar plasma parameters and Ne jet parameters in progress– no conclusions yet





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Postprocessing to determine runaway confinement time in NIMROD results

Experimental Motivation:

 C-Mod experiment seeds plasma with suprathermal electrons to study runaway conversion and confinement during the disruption.

Procedure:

- Assume (initially) that suprathermal/runaway electrons follow the field lines perfectly
- Initialize suprathermal electrons with given positions and velocities
- Run nimfl to track electron trajectories, but advance electron velocity and time using F=eE-mvµee





Rapid loss of electrons during the thermal quench



Summary

- D₂ Dilution cooling as an "optimal" runaway suppression mechanism
 - \rightarrow Mostly wrapped up, some interesting results, probably a short paper
- Effects of plasma elongation on MHD and runaway confinement
 - \rightarrow Very preliminary, eventual experimental comparison w/ C-Mod
- Direct calculations of suprathermal electron acceleration and confinement
 - $\rightarrow\,$ Some results, but more physics to include, possible inclusion in nimrod during run time



