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We Need a Predictive Model for Plasma Response to 3D Fields in Tokamaks

• Effects of 3D fields need to be evaluated for ITER:

- RMP ELM suppression
- Changes in particle/heat flux to wall
 - Edge displacements
 - Divertor footprints
- Fast ion transport
- Mode locking
- We now have the tools and the measurements to evaluate our capabilities to model some of these things
 - Modeling of edge displacements generally finds good agreement with experiment
- Transport in 3D fields is next step towards predictive models

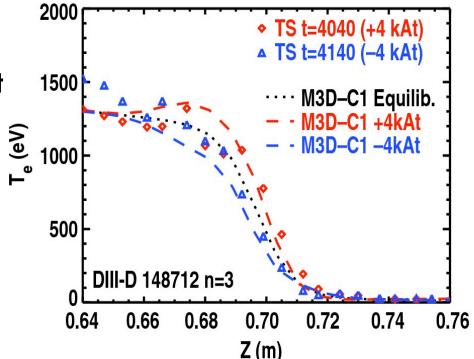


Edge Displacement



Measurements of Edge Response to 3D Fields Are Generally in Good Agreement With Two-Fluid Modeling

- T_e and n_e profiles are "displaced" by the application of 3D fields
- Edge displacements are a robust feature of 3D plasma response
 - Provide a measurement for validating codes
 - Provide an indication of internal plasma response
 - May cause problems in ITER
 - Focus of ITPA WG (Chapman)

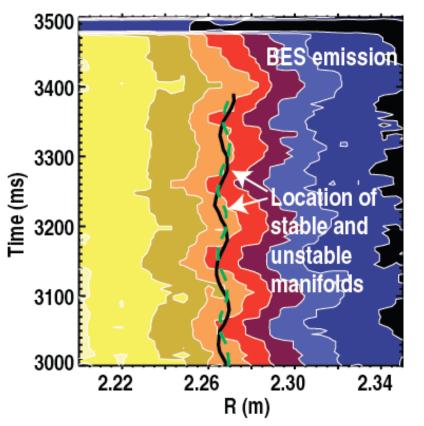


 We find generally good agreement between two-fluid modeling (M3D-C1) and measurements of edge response



Rotating *n*=1,2 Fields Sweeps Structures Past Diagnostics

- On DIII-D, the toroidal phase of n=1 and n=2 fields can be smoothly rotated
- Displacement is phase dependent
- Two possibilities
 - Displacement is 3D
 - Displacement is 2D, but phase dependent
 (i.e. there are significant error fields)

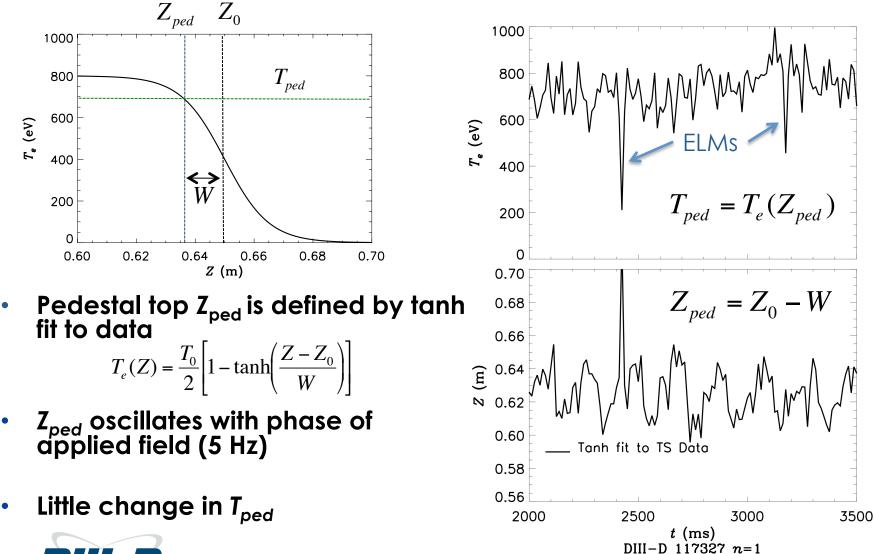


 Measured displacement is generally larger than calculated displacement of separatrix manifolds from vacuum fields



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Displacement Can Be Quantified By The Change In The Location Of The Pedestal Top





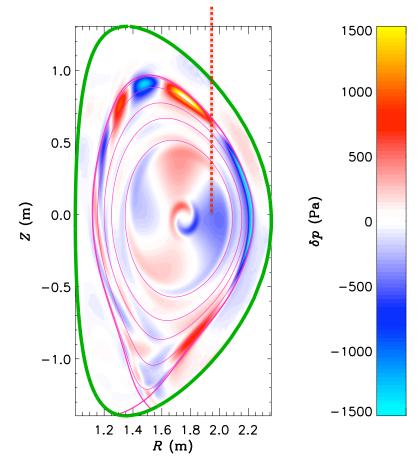
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Time-Independent Response is Calculated as Boundary-Value Problem

Boundary Conditions:

- Normal component of magnetic field is fixed equal to applied field
- No-slip, no pressure perturbation
- Linear time-independent solution is solved directly (not by initial value calculation)

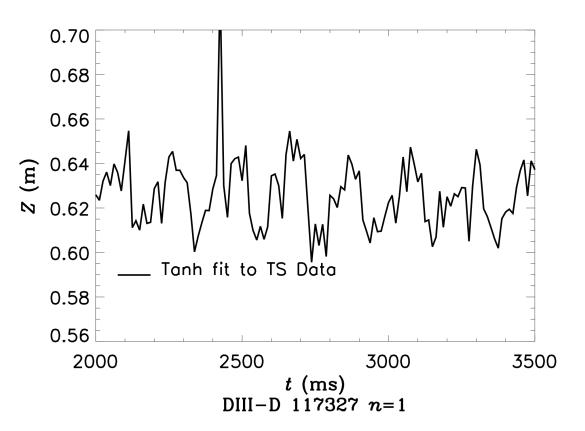
Core Thomson chord





Two-Fluid Modeling Reproduces Phase and Magnitude of Displacement

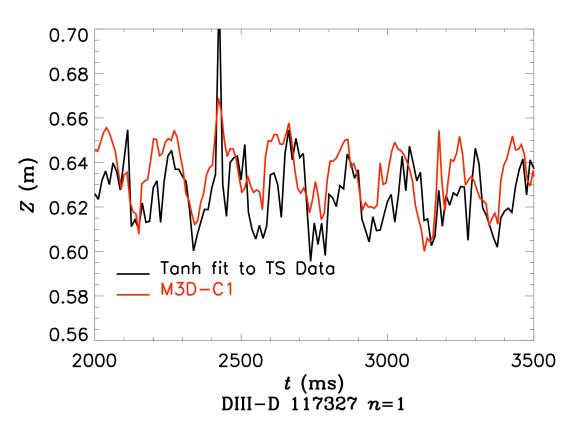
- In the experiment, the peak-to-peak displacement is ~4 cm
- Vacuum modeling finds few mm





Two-Fluid Modeling Reproduces Phase and Magnitude of Displacement

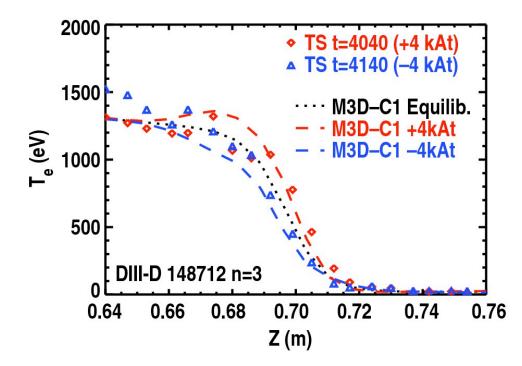
- In the experiment, the peak-to-peak displacement is ~4 cm
- Vacuum modeling finds few mm
- M3D-C1 Modeling finds good agreement in phase and magnitude of displacement





n=3 Fields Yield Smaller Displacements Than n<3

- *n*=3 fields cannot be rotated on DIII-D, but can be flipped
- Flipping n=3 fields yields displacement of $\sim 1-2$ cm

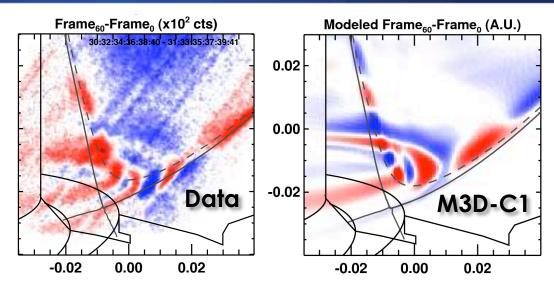


M3D-C1 finds agreement through much of pedestal



X-Ray Data Reveals Field-Aligned 3D Structure

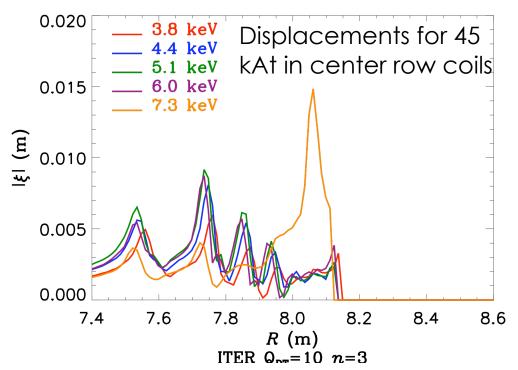
 Data is obtained by flipping I-coil fields and taking difference between signals

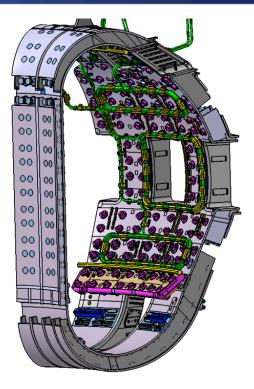


- The poloidal structure is strongly indicative of a <u>field-aligned</u> <u>helical response</u>
- Modeling agrees qualitatively with poloidal structure of response
- Radial localization indicates driven peeling-ballooning response



Preliminary Results Show Moderate Displacements for ITER





- Midplane edge displacements are found to be ~1/2 cm in Q_{DT}=10 scenarios with 45 kAt in the center row
 - Only center row considered (found to have strongest coupling)
 - ITER Q_{DT}=10 scenarios have ~10 cm outer gap



Linear Results Appear to be Valid In These Cases

 "Displacement" may be defined by movement of isotherms:

$$T_0(r+\xi) + \delta T(r+\xi) = T_0(r)$$

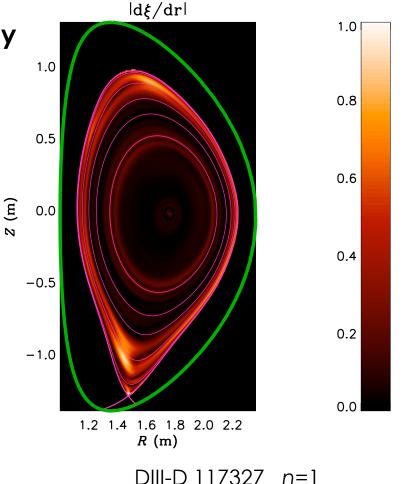
$$\begin{bmatrix} T_0(r) + \frac{dT_0}{dr}\xi \end{bmatrix} + \delta T(r) = T_0(r)$$

$$\xi = -\frac{\delta T}{dT_0/dr}$$

 Overlap of adjacent surfaces is possible, especially near moderational surfaces, edge, & x-point

Overlap criterion:

$$\left|\frac{d\xi}{dr}\right| > 1$$





Transport in 3D Fields



Predictive Modeling of Requires Calculation of Transport in 3D Fields

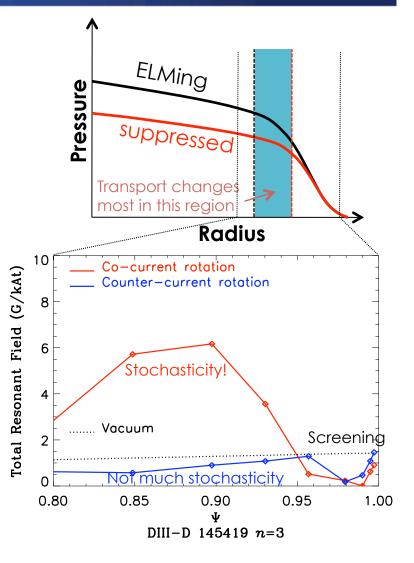
Transport in 3D fields is necessary for predictive modeling

- Fast ion loss
- Torque from RMP
- RMP ELM suppression
- New tools and interfaces are being developed for this purpose using fields from M3D-C1
 - Single-particle orbit calculations (ORBIT-RF, SPIRAL)
 - Choi GP8.00098
 - NTV torque calculations (Cole, Callen)
 - McCubbin JP8.00016
 - Flutter transport calculations (Callen)
 - Raum JP8.00017
 - Ballooning mode stability (Bird)



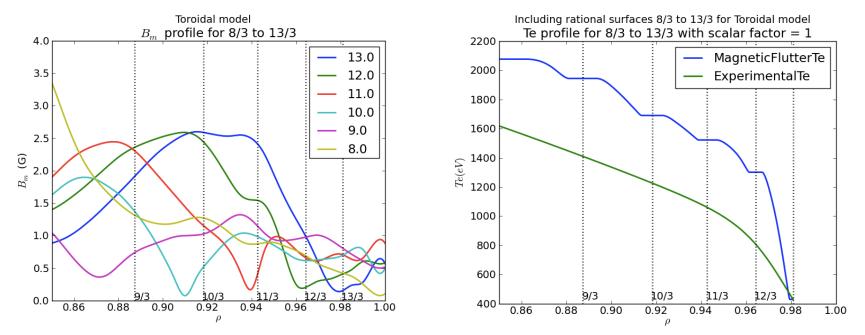
Island Formation at Pedestal Top is a Promising Hypothesis for RMP ELM Suppression

- Hypothesis: RMP ELM suppression is achieved by limiting pedestal width
 - Confinement is degraded by 3D fields at top of pedestal
- Does response at top of pedestal cause ELM suppression?
 - With co-current rotation, a large response is expected at pedestal top (where ω_e crosses zero)
 - RMP ELM suppression not definitively observed with counter-current rotation (since ω_e never crosses zero)
- What is the source of the additional transport?





Flutter Transport Calculations Show Enhanced Thermal Transport Near Pedestal Top



- Magnetic response from M3D-C1 can serve as basis for transport calculations
 - TRIP3D \rightarrow parallel thermal transport
 - ORBIT-RF / SPIRAL \rightarrow fast ion transport
 - New post-analysis tools for NTV & flutter transport
 - 3D KBM stability

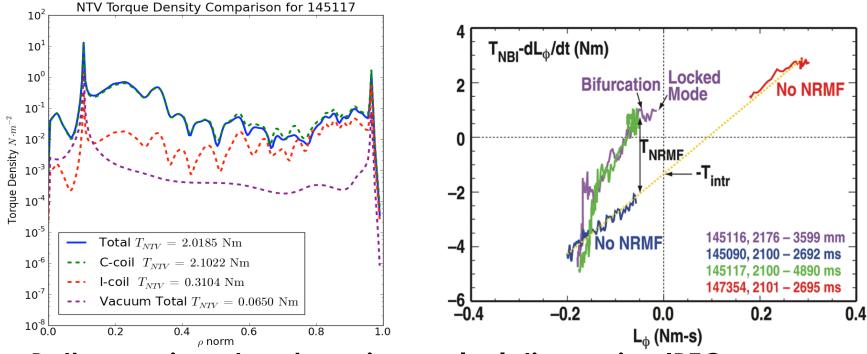
Raum JP8.00017 Callen BP8.00160



NTV Calculation Using M3D-C1 Fields Finds Agreement with Experiment

• Evaluating Cole's formula for NTV finds:

- 0.065 Nm using vacuum fields
- 2.0185 Nm using plasma response



 Both experiment and previous calculations using IPEC response find 2—3 Nm
 McCubbin JP8.00016

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Summary

- Plasma response calculations yield good agreement with experimental measurements of edge displacement
- Edge displacements are largely helical, not (just) axisymmetric
 - M3D-C1 response is purely helical, and agrees with experiment
 - X-ray data shows clear helical response
- Displacements may be strongly enhanced by plasma response (*i*.e. stable mode driven to finite amplitude)
- Transport calculations in 3D fields are being integrated with M3D-C1 calculations





Extra Slides



Two-Fluid Model Implemented in M3D-C1

$$\frac{\partial n}{\partial t} + \nabla \cdot (n\mathbf{u}) = 0 \qquad \mathbf{E} = -\mathbf{u} \times \mathbf{B} + \eta \mathbf{J} + \left[\frac{d_i}{n} (\mathbf{J} \times \mathbf{B} - \nabla p_e)\right] \\ n\left(\frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u}\right) = \mathbf{J} \times \mathbf{B} - \nabla p - \nabla \cdot \Pi \qquad \Pi = -\mu \left[\nabla \mathbf{u} + (\nabla \mathbf{u})^T\right] \\ \frac{\partial p}{\partial t} + \mathbf{u} \cdot \nabla p = -\Gamma p \nabla \cdot \mathbf{u} - \left[\frac{d_i}{n} \mathbf{J} \cdot \left(\Gamma p_e \frac{\nabla n}{n} - \nabla p_e\right)\right] \qquad \mathbf{q} = -\kappa \nabla p - \kappa_{\parallel} \mathbf{b} \mathbf{b} \cdot \nabla \left(\frac{p_e}{n}\right) \\ -(\Gamma - 1) \nabla \cdot \mathbf{q} \qquad \mathbf{J} = \nabla \times \mathbf{B} \\ \frac{\partial \mathbf{B}}{\partial t} = -\nabla \times \mathbf{E} \qquad p_e = p/2 \end{cases}$$

- **Two-fluid** terms scale with ion skin depth (d_i)
- Time-independent equations may be solved directly for linear response
- Boundary conditions: normal B from external coils is held constant at boundary

