



STUDIES OF NONLINEAR RESISTIVE AND EXTENDED MHD IN ADVANCED TOKAMAKS USING THE NIMROD CODE

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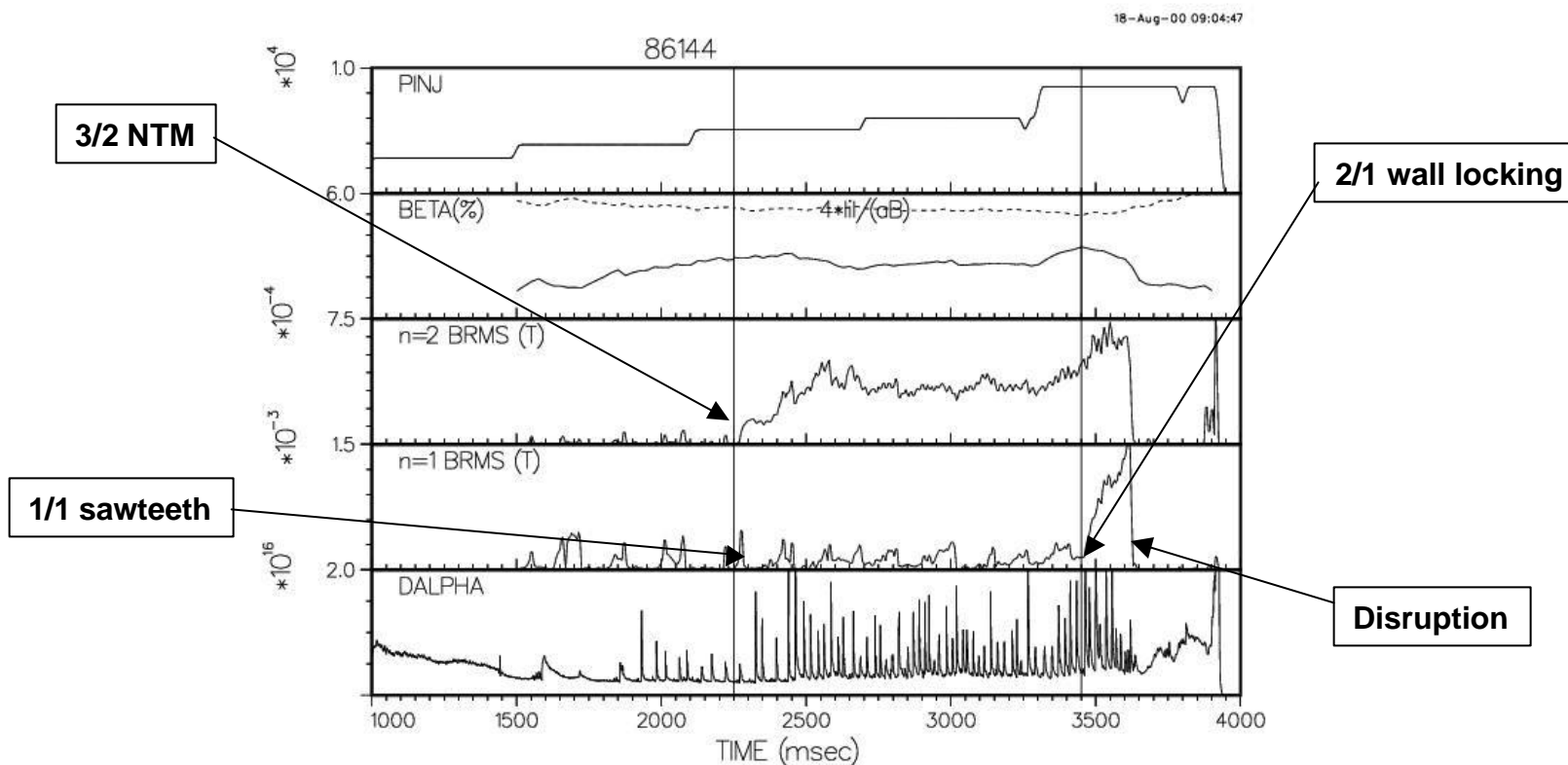
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MODERN TOKAMAKS ARE RICH IN MHD ACTIVITY

Example: DIII-D shot 86144



- Sawtooth discharge
- 3/2 NTM triggered at 2250 msec
- 2/1 locks to the wall

IMPORTANCE OF REALISTIC MODELING

- Cost of next generation of fusion experiments estimated to be at least several billion \$\$
- Cost proportional to volume: $\$ \sim V$
- Power density proportional to square of max. pressure:
 $P/V \sim p_{\max}^2$
 - $\Rightarrow \$ \sim 1/p_{\max}^2$ for fixed P and B (engineering constraints)
- Physics uncertainties limit max. pressure to $\sim 2/3$ theoretical p_{\max}

Uncertainties in nonlinear physics account for $\sim 1/2$ the cost of advanced fusion experiment!
- Predictive fluid modeling with realistic parameters has high leverage to remove this uncertainty

MODELING REQUIREMENTS

- **Slow evolution**
Nonlinear fluid model required
- **Plasma shaping**
Realistic geometry required
- **High temperature**
Realistic S required
- **Low collisionality**
Extensions to resistive MHD required
- **Strong magnetic field**
Highly anisotropic transport required
- **Resistive wall**
Non-ideal boundary conditions required

2-FLUID MODEL

- Maxwell (no displacement current):

$$\frac{\mathcal{I}\mathbf{B}}{\mathcal{I}t} = -\nabla \times \mathbf{E} \quad , \quad \nabla \times \mathbf{B} = m_0 \mathbf{J} \quad ,$$

- Momentum, energy, and continuity for each species ($a = e, i$):

$$m_a n_a \left(\frac{\mathcal{I}\mathbf{v}_a}{\mathcal{I}t} + \mathbf{v}_a \cdot \nabla \mathbf{v}_a \right) = -\nabla \cdot \mathbf{P}_a + q_a n_a (\mathbf{E} + \mathbf{v}_a \times \mathbf{B}) + \sum_b \mathbf{R}_{ab} + \mathbf{S}_a^m$$

$$\frac{\mathcal{I}\rho_a}{\mathcal{I}t} + \mathbf{v}_a \cdot \nabla \rho_a = -\frac{3}{2} \rho_a \nabla \cdot \mathbf{v}_a - \mathbf{P}_a : \nabla \mathbf{v}_a - \nabla \cdot \mathbf{q}_a + \mathbf{Q}_a$$

$$\frac{\mathcal{I}n_a}{\mathcal{I}t} = -\nabla \cdot (n_a \mathbf{v}_a) + S_a^n$$

- Current and quasi-neutrality:

$$\mathbf{J}_a = n_a q_a \mathbf{v}_a, \quad n = n_e = Z n_i$$

SINGLE FLUID FORM

- Add electron and ion momentum equations:

$$r \left(\frac{\rho \mathbf{v}}{\rho t} + \mathbf{v} \cdot \nabla \mathbf{v} \right) = -\nabla \cdot \mathbf{P}' + \mathbf{J} \times \mathbf{B}$$

- Subtract electron and ion momentum equations (Ohm's law):

$$\mathbf{E} = - \underbrace{\mathbf{v} \times \mathbf{B}}_{\text{Ideal MHD}} + \underbrace{\frac{h\mathbf{J}}{ne}}_{\text{Resistive MHD}} + \underbrace{\frac{1}{ne} \frac{1-n}{1+n} \mathbf{J} \times \mathbf{B}}_{\text{Hall Effect}} - \underbrace{\frac{1}{ne(1+n)} \nabla \cdot (\mathbf{P}'_e - n\mathbf{P}'_i)}_{\text{Pressure Effects and Closures}} + \underbrace{\frac{1}{e_0 w_{pe}^2 (1+n)} \left[\frac{\rho \mathbf{J}}{\rho t} + \nabla \cdot (\mathbf{v}\mathbf{J} + \mathbf{J}\mathbf{v}) \right]}_{\text{Electron Inertia}}$$

MODELING GOALS

- ***Analysis and interpretation of experimental data***
 - NIMROD can interface with common design and analysis codes
 - EFIT
 - TOQ
 - CHEASE
 - DIII-D has extensive data base
 - Study nonlinear phenomena
 - Tests both resistive and extended MHD models
- ***Analysis and evaluation of designs for advanced experiments***
 - Example: VDE in ITER

DIII-D EXPERIMENTAL DISCHARGES

- **Shot 87009**

- Highly shaped plasma
- Disruption when heated through b limit
 - Why is growth faster than simple exponential?
 - What causes disruption?
- Nonlinear resistive MHD

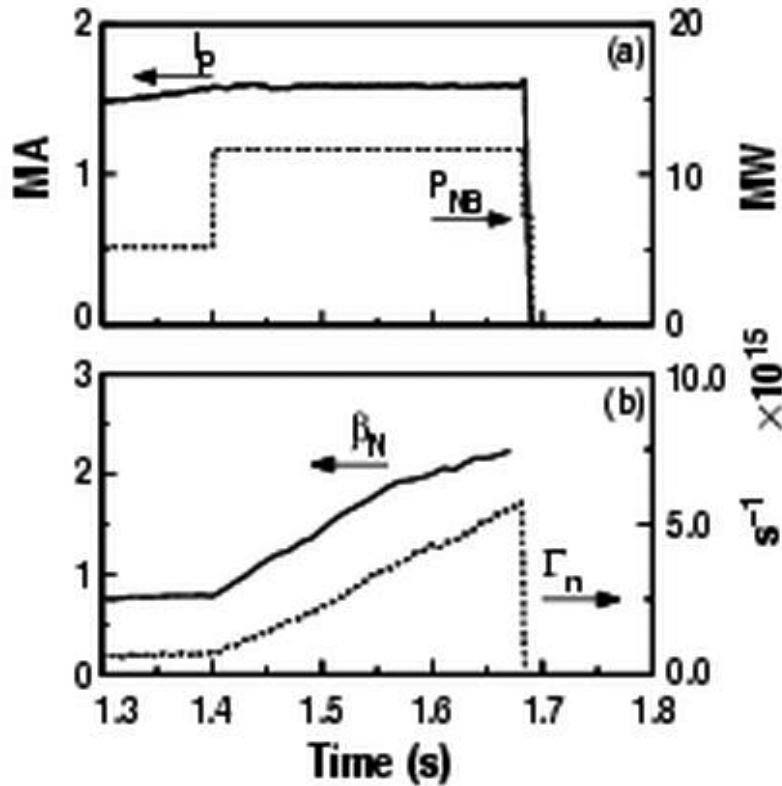
- **Shot 86144**

- ITER-like discharge
- Sawteeth
 - Nonlinear generation of secondary islands
 - Destabilization of NTM?
- Tests both resistive MHD and closure models for Extended MHD

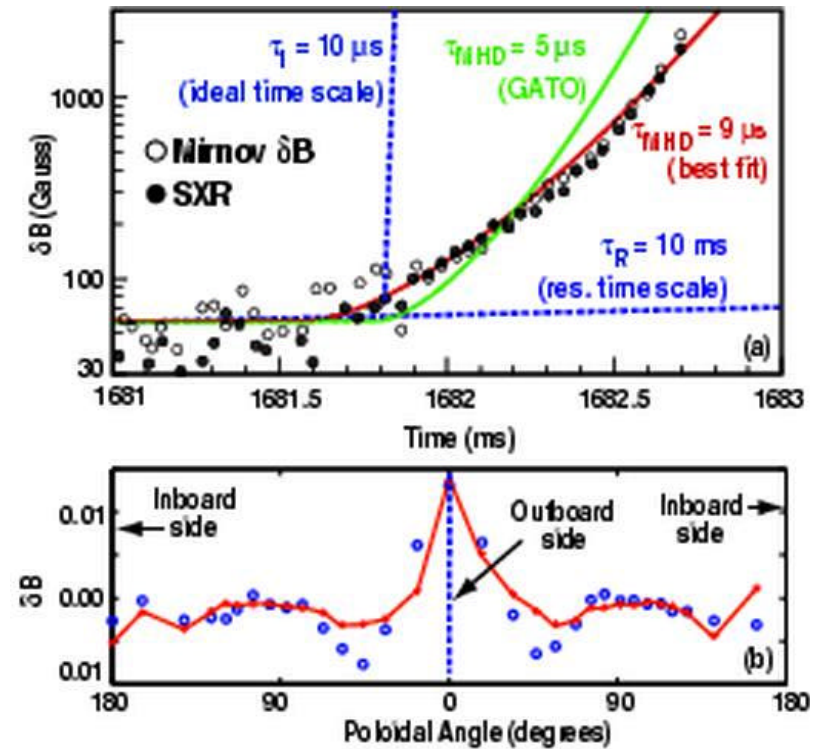


DIII-D SHOT #87009

- High- b disruption when heated slowly through critical b_N



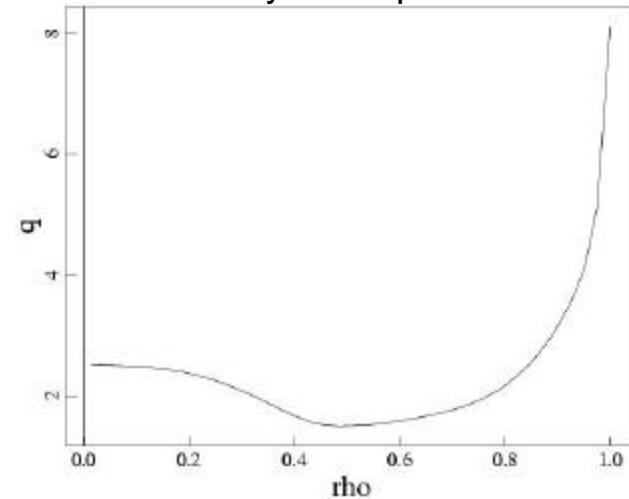
- Growth is faster than simple exponential



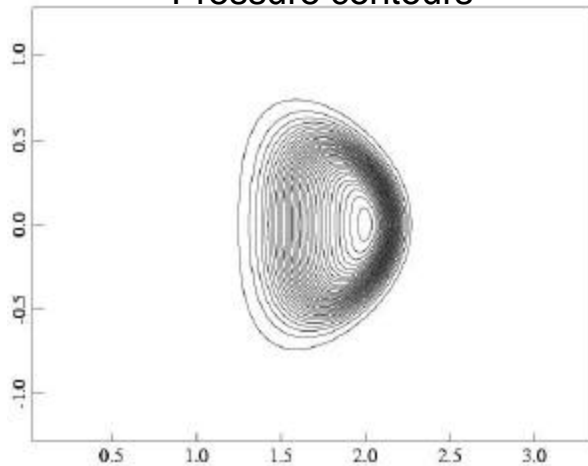
EQUILIBRIUM AT $t = 1681.7$ msec

- Equilibrium reconstruction from experimental data
- Negative central shear
- Gridding based on equilibrium flux surfaces
 - Packed at rational surfaces
 - Bi-cubic finite elements

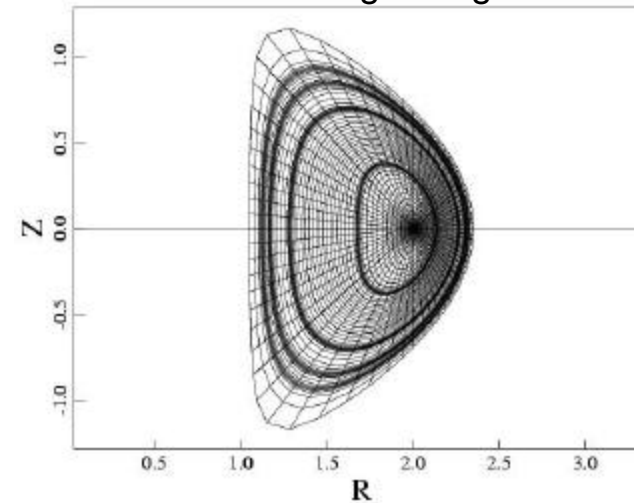
Safety factor profile



Pressure contours

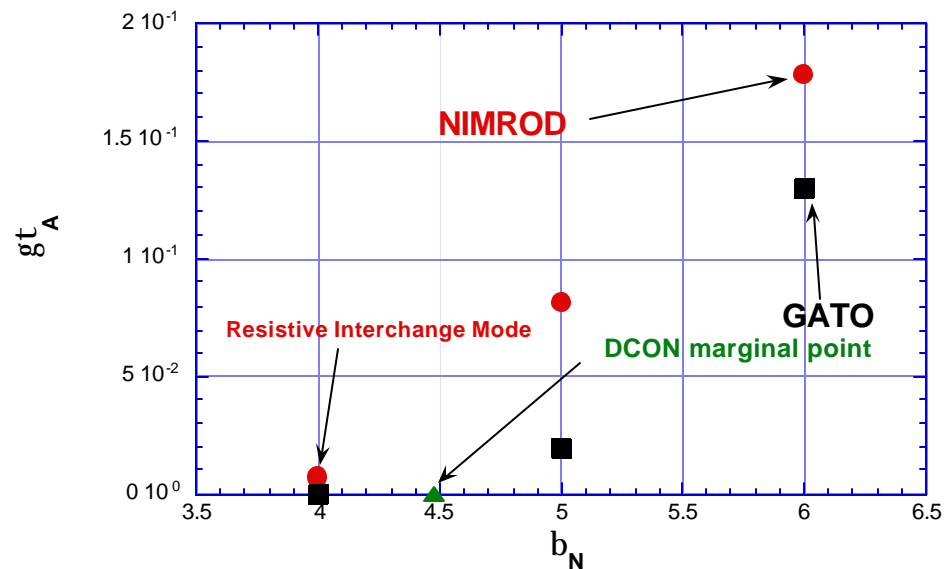


Poloidal gridding



LINEAR STABILITY

- Present version of NIMROD requires conducting wall at outer flux surface
 - Critical b_N for ideal instability larger than in experiment
- NIMROD gives slightly larger ideal growth rate than GATO
- NIMROD finds resistive interchange mode below ideal stability boundary



THEORY OF SUPER-EXPONENTIAL GROWTH

- In experiment mode grows faster than exponential
- Theory of ideal growth in response to slow heating (*Callen, Hegna, Rice, Strait, and Turnbull, Phys. Plasmas 6, 2963 (1999)*):

Heat slowly through critical b : $b = b_c(1 + g_h t)$

Ideal MHD: $w^2 = -\hat{g}_{MHD}^2 (b/b_c - 1) \rightarrow g(t) = \hat{g}_{MHD} \sqrt{g_h t}$

Perturbation growth:

$$\frac{dx}{dt} = g(t)x \quad \rightarrow \quad x = x_0 \exp[(t/t)^{3/2}], \quad t = (3/2)^{2/3} \hat{g}_{MHD}^{-2/3} g_h^{-1/3}$$

- Good agreement with experimental data

NONLINEAR SIMULATION WITH NIMROD

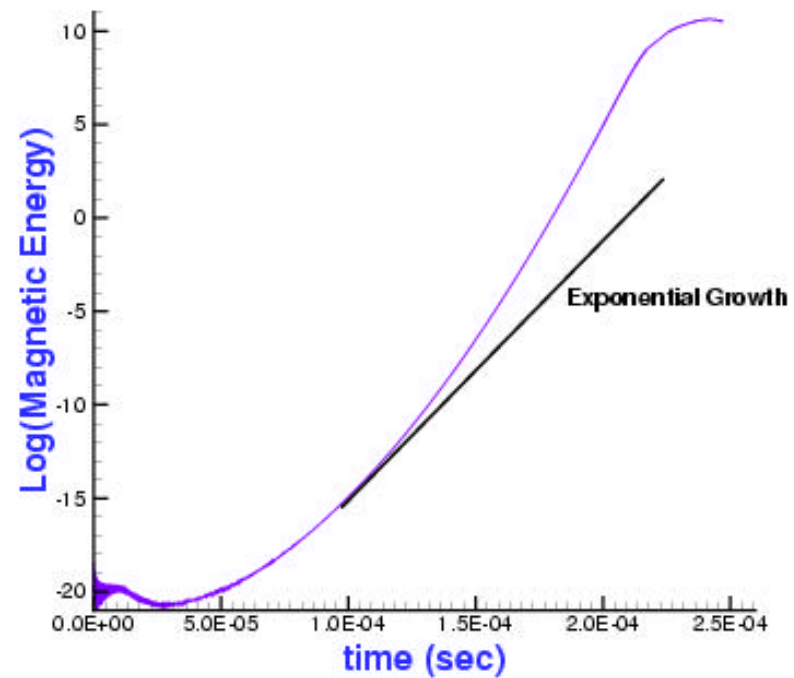
- Initial condition: equilibrium below ideal marginal b_N
- Use resistive MHD
- Impose heating source proportional to equilibrium pressure profile

$$\frac{\partial P}{\partial t} = \dots + g_H P_{eq}$$

$$\rightarrow b_N = b_{Nc}(1 + g_H t)$$

- Follow nonlinear evolution through heating, destabilization, and saturation

Log of magnetic energy in $n = 1$ mode vs. time
 $S = 10^6$ $Pr = 200$ $g_H = 10^3 \text{ sec}^{-1}$



SCALING WITH HEATING RATE

- NIMROD simulations also display super-exponential growth
- Simulation results with different heating rates are well fit by $x \sim \exp[(t-t_0)/t]^{3/2}$

- Time constant scales as

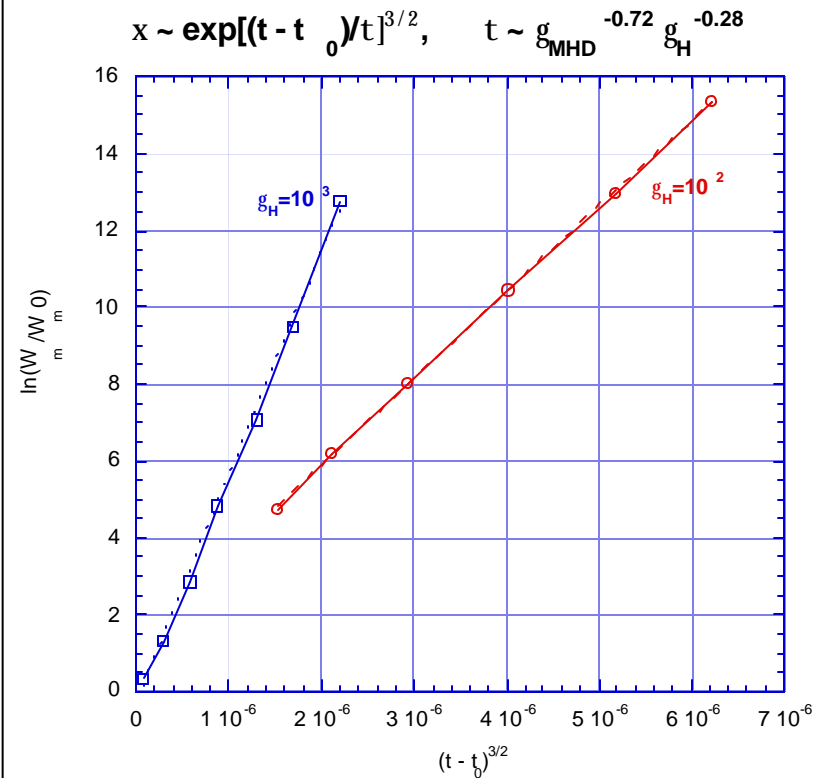
$$t \sim g_{MHD}^{-0.72} g_H^{-0.28}$$

- Compare with theory:

$$t = (3/2)^{2/3} \hat{g}_{MHD}^{-2/3} g_h^{-1/3}$$

- Discrepancy possibly due to non-ideal effects

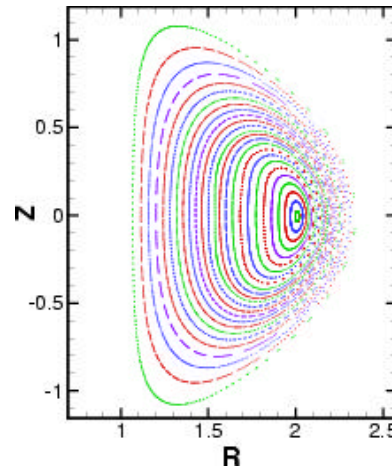
Log of magnetic energy vs. $(t - t_0)^{3/2}$
for 2 different heating rates



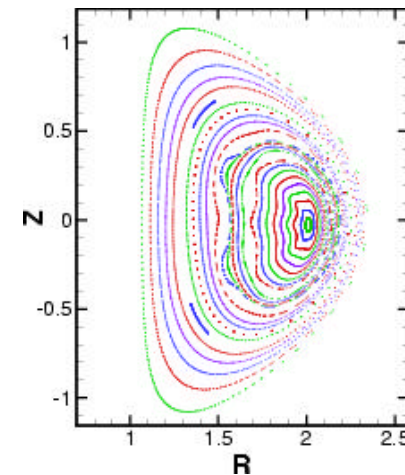
EVOLUTION OF MAGNETIC FIELD LINES

- Simulation with small but finite resistivity
- Ideal mode yields stochastic field lines in late nonlinear stage
- Implications for degraded confinement
- Disruption?

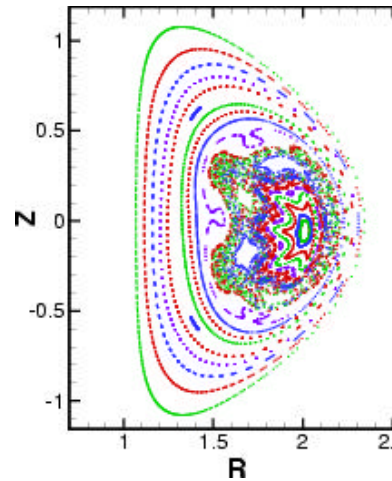
$t = 1.99 \times 10^{-4}$ sec.



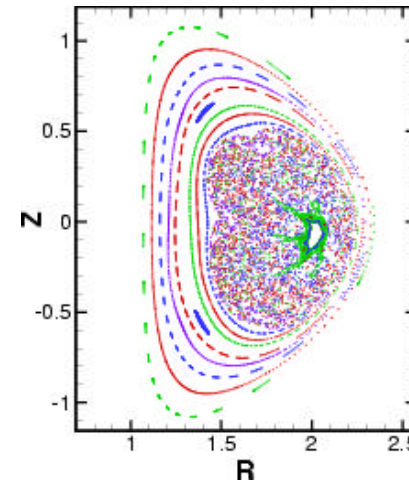
$t = 2.102 \times 10^{-4}$ sec.



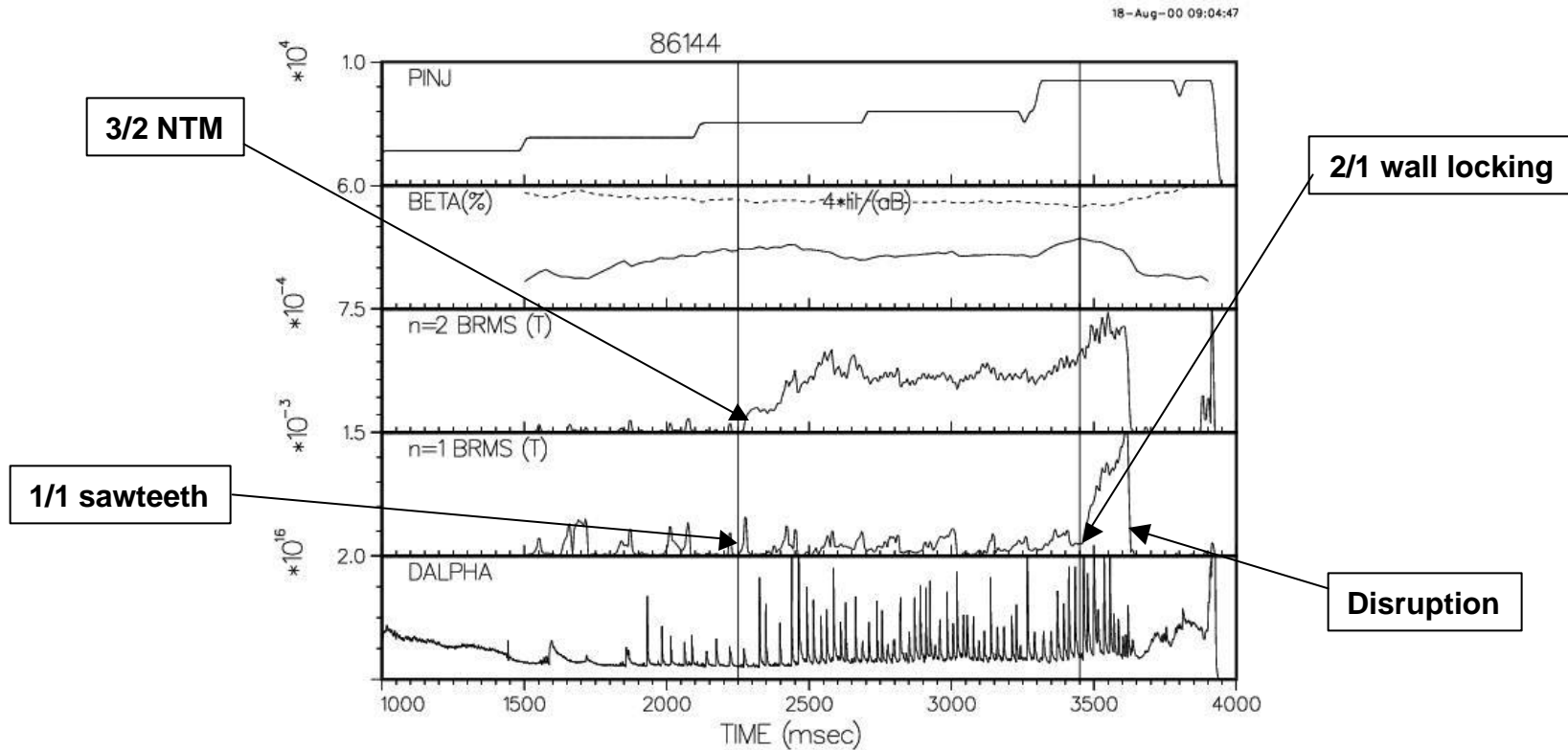
$t = 2.177 \times 10^{-4}$ sec.



$t = 2.2 \times 10^{-4}$ sec.



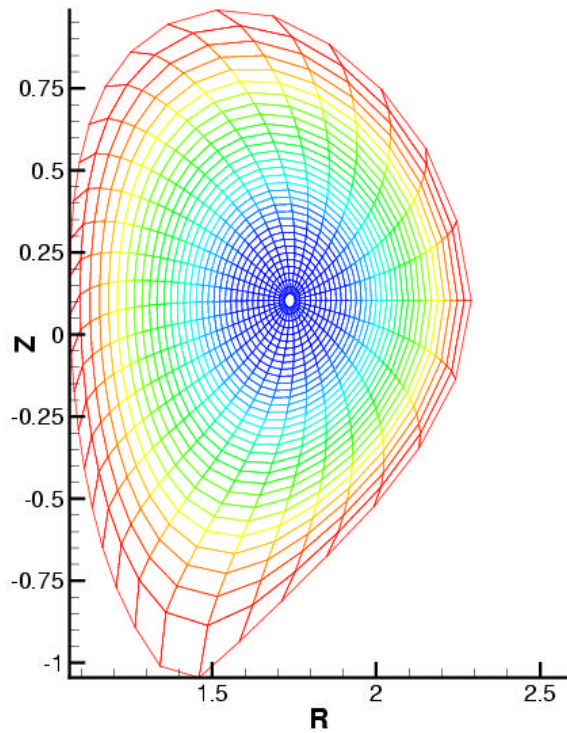
DIII-D SHOT #86144



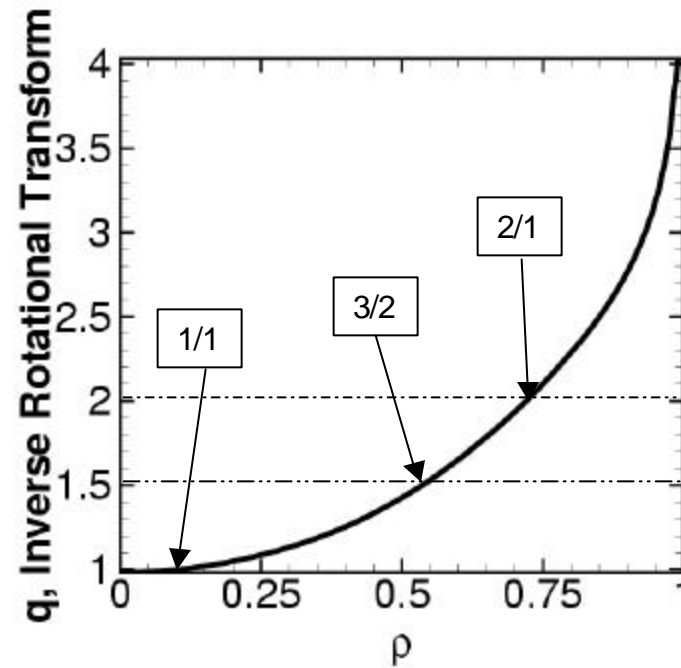
- Sawtooth discharge
- 3/2 NTM triggered at 2250 msec
- 2/1 locks to the wall

EQUILIBRIUM AT $t = 2250$ msec

Grid (Flux Surfaces)

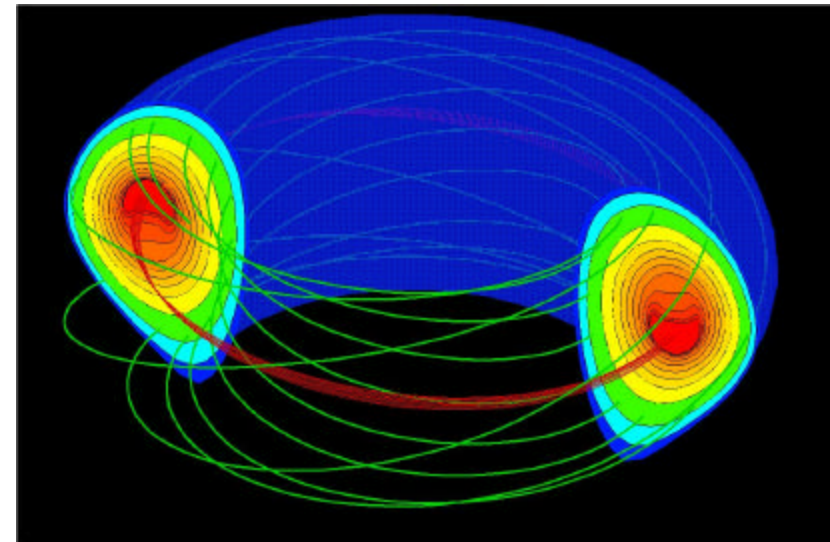
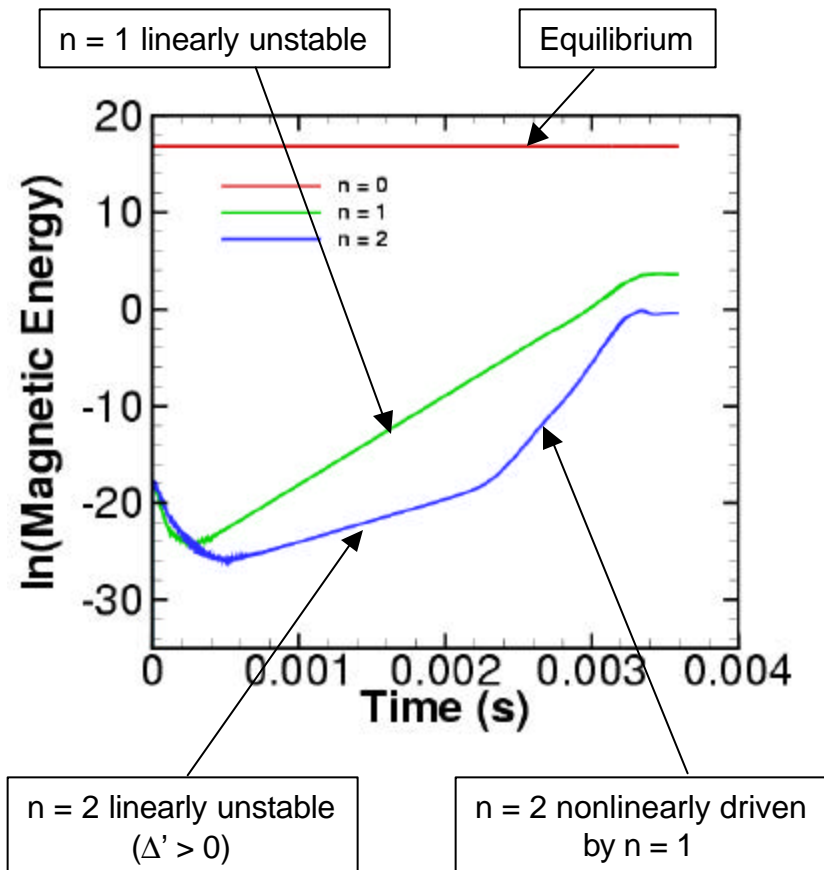


q - profile



- ITER-like discharge
- $q(0)$ slightly below 1

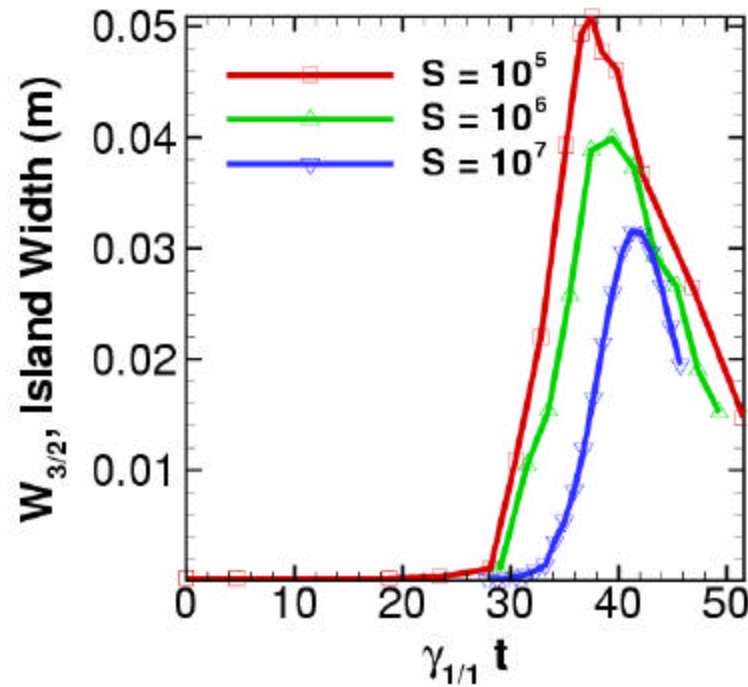
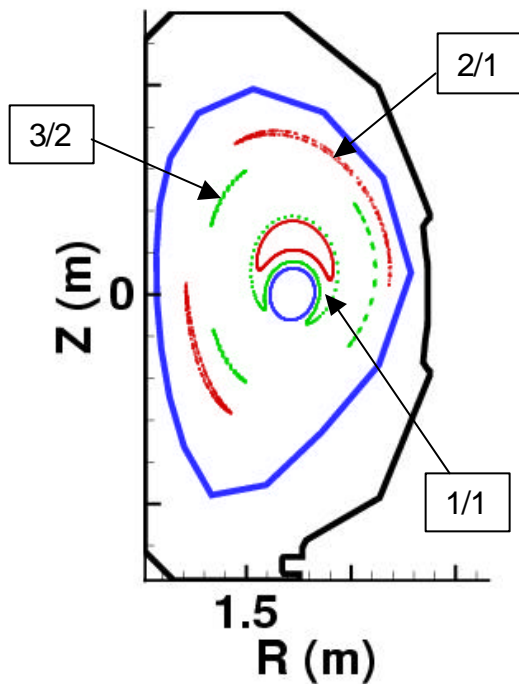
DISCHARGE IS UNSTABLE TO RESISTIVE MHD



Pressure and field lines

$$S = 10^7 \quad Pr = 10^3$$
$$g = 4.58 \times 10^3 / \text{sec} \quad g_{\text{exp}} \sim 1.68 \times 10^4 / \text{sec}$$

SECONDARY ISLANDS IN RESISTIVE MHD



- Secondary islands are small in resistive MHD
— $W_{\text{exp}} \sim 0.1$ m
- $3/2$ island width decreases with increasing S
- Need extended MHD to match experiment

NUMERICALLY TRACTABLE CLOSURES

- Resistive MHD is insufficient to explain DIII-D shot 86144
 - 3/2 magnetic island is too small
- Parallel variation of B leads to trapped particle effects
- Particle trapping causes neo-classical effects
 - Poloidal flow damping
 - Enhancement of polarization current
 - Bootstrap current
- Simplified model captures most neo-classical effects
(T. A. Gianakon, S. E. Kruger, C. C. Hegna, Phys. Plasmas (to appear) (2002))

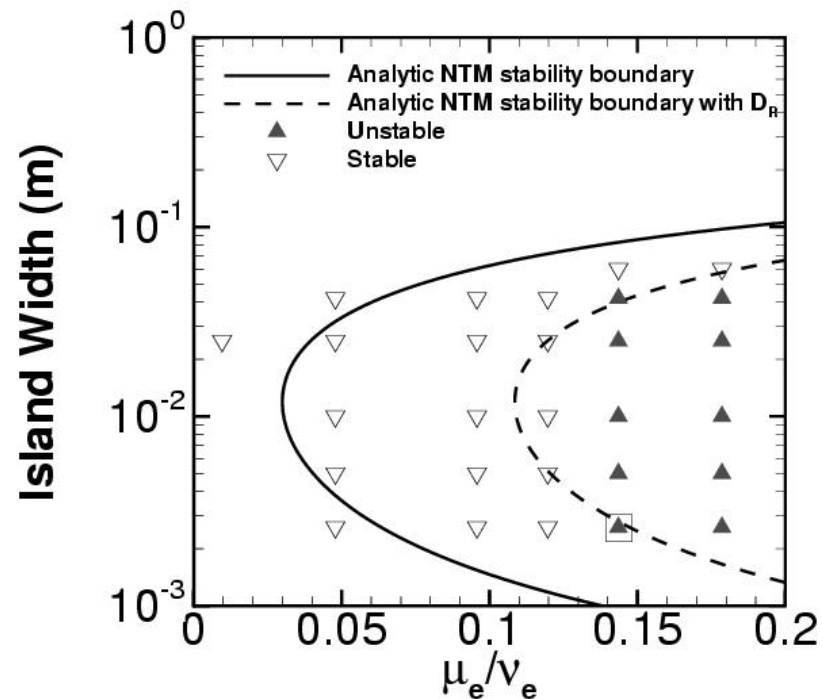
$$\nabla \cdot \Pi_a = m_a n_a m_a \langle B_0 \rangle^2 \frac{\mathbf{v}_a \cdot \nabla q}{(\mathbf{B}_0 \cdot \nabla q)^2} \nabla q$$

- For electrons, ideal MHD equilibrium yields bootstrap current

$$\nabla \cdot \Pi_e = - \frac{r_e m_e}{n_e} \frac{\langle B \rangle^2}{B^2} \frac{\mathbf{B}_0 \times \nabla p \cdot \nabla q}{(\mathbf{B}_0 \cdot \nabla q)^2} \nabla q$$

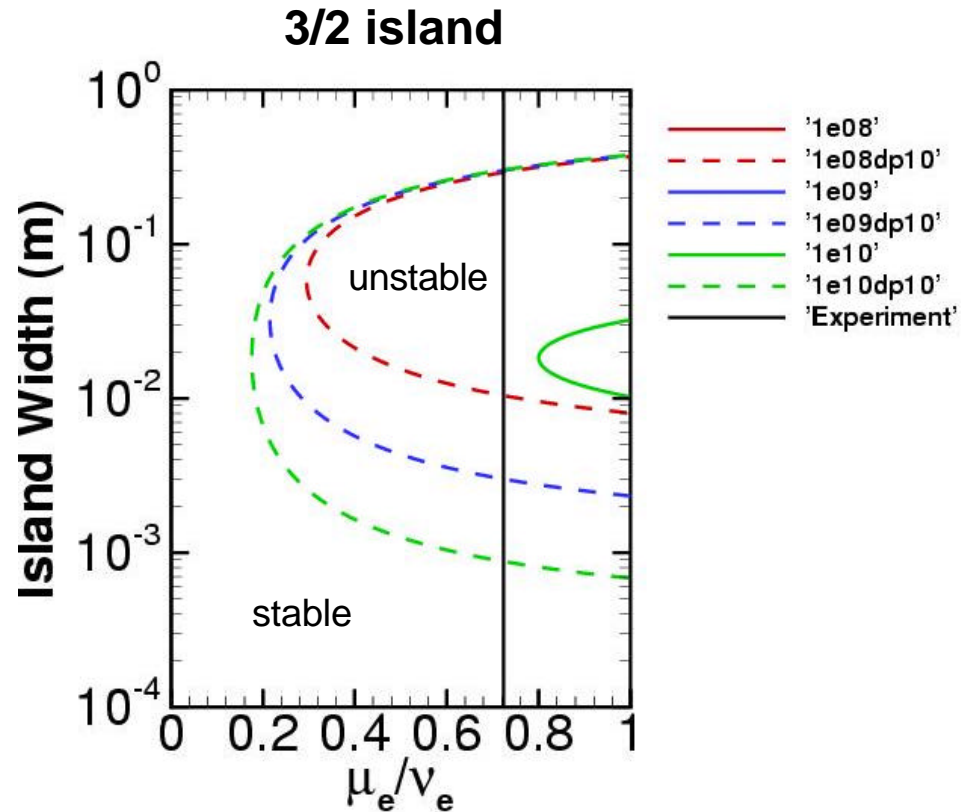
CLOSURES REPRODUCE NTM INSTABILITY

- TFTR-like equilibrium
- Comparison with modified Rutherford equation
- Initialize NIMROD with various seed island sizes
- Look for growth or damping
- Seek self-consistent seed and growth



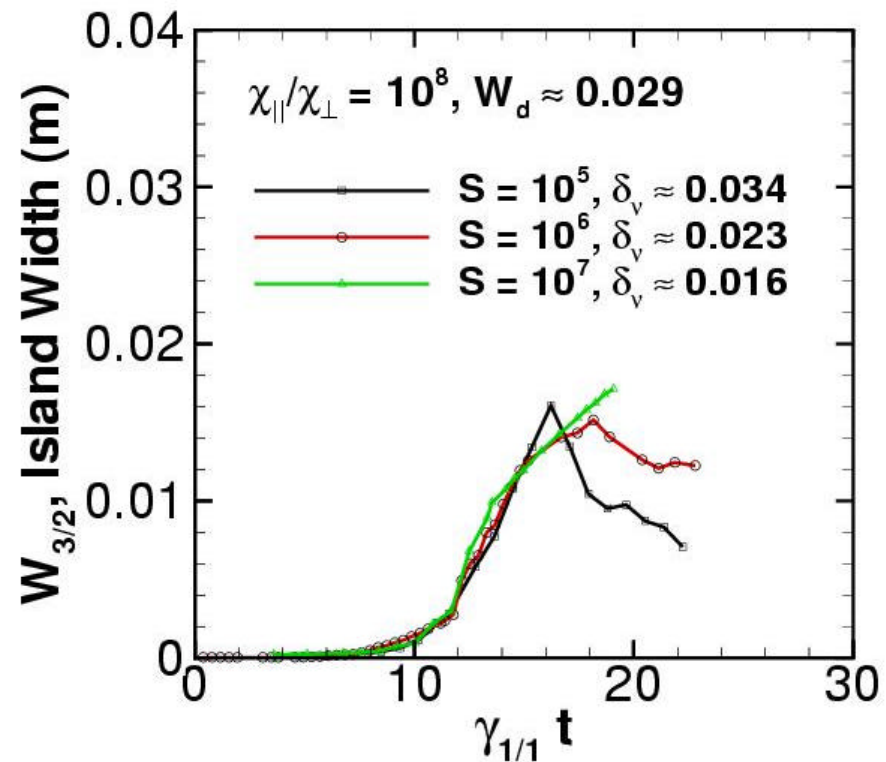
86144.2250 NTM STABILITY BOUNDARIES

- Use modified Rutherford equation
- 3 values of anisotropic heat flux
- 2 values of $D\zeta$
 - Vacuum
 - Reduced by factor of 10
- Experimental island width ~ 0.1 m



SELF-CONSISTENT NTM MAY REQUIRE HIGHER S

- $W_{3/2}$ versus time with NTM closure
- $W_{\text{exp}} \sim 0.1$ m
- Island width exceeds analytic threshold
- No NTM observed
- Island width less than visco-resistive layer width at low S
- Calculations at larger S are underway



Nonlinear NTM calculations are extremely challenging!

SUMMARY

- Nonlinear modeling of experimental discharges is possible, but extremely challenging
- *DIII-D shot #87009*
 - Heating through b limit
 - Super-exponential growth, in agreement with experiment and theory
 - Nonlinear state leads to stochastic fields
 - Calculations with anisotropic thermal transport underway
- *DIII-D shot #86144*
 - Secondary islands driven by sawtooth crash
 - Resistive MHD insufficient, requires neo-classical closures
 - Must go to large S ($\sim 10^7$) to get proper length scales
 - Calculations are underway