Implicit Heat Conduction in M3D and Recent Scaling Results

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Background

The parallel heat transport model in M3D has historically been the "artificial sound wave": the MHD-like energy equation

$$\frac{\partial p}{\partial t} + \mathbf{v} \cdot \nabla p + \gamma p \nabla \cdot \mathbf{v} = \frac{2}{3} \nabla_{\perp} \cdot \left[n \chi_{\perp} \nabla_{\perp} \left(\frac{p}{\rho} \right) \right] + \text{heating terms}$$

where

$$\nabla_{\perp} \equiv \hat{R} \frac{\partial}{\partial R} + \hat{z} \frac{\partial}{\partial z}$$

is supplemented by a pair of wave equations for temperature and an auxiliary variable *u*:

$$\frac{\partial T}{\partial t} = s \frac{\mathbf{B} \cdot \nabla u}{n}$$
$$\frac{\partial u}{\partial t} = s \mathbf{B} \cdot \nabla T + v \nabla_{\perp}^{2} u$$

for sound speed s and artificial viscosity v.

Drawbacks of the Artificial Sound Approach

- This model is not a part of standard MHD, and not directly comparable to results from codes that implement the standard model (unclear what the equivalent χ_{\parallel} should be).
- Because the perpendicular operator is perpendicular to the φ direction rather than **B**, parallel and perpendicular transport are not cleanly separated.
- The implementation of the parallel operator is explicit in time, unlike the perpendicular one, and so can restrict the time step when *s* is large.

Parallel Heat Diffusion

The MHD energy equation with anisotropic heat conduction can be written ∂p ______

$$\frac{\partial p}{\partial t} + \mathbf{v} \cdot \nabla p + \gamma p \nabla \cdot \mathbf{v} = -\frac{2}{3} \nabla \cdot \mathbf{q} + \text{heating terms}$$

where the heat flux \mathbf{q} is given by

$$\mathbf{q} = -n \Big[\Big(\chi_{||} - \chi_{\perp} \Big) \hat{\mathbf{b}} \hat{\mathbf{b}} \cdot \nabla T + \chi_{\perp} \nabla T \Big].$$

Straightforward solution of this equation using finite differences or loworder finite elements is susceptible to numerical pollution of perpendicular heat transport by parallel when $\chi_{\parallel}/\chi_{\perp} \gg 1$, as is typical.

Günter, *et al.*, have a 2nd-order scheme based on centered differences that largely avoids this problem¹.

¹S. Günter, K. Lackner, and C. Tichmann, J. Comp. Phys. 226 (2007) 2306-2316.

Günter's Method for Finite Elements

The Galerkin integral of the heat conduction term against a test function is

$$-\int \lambda_i \nabla \cdot \mathbf{q} \, dV = \int \mathbf{q} \cdot \nabla \lambda_i \, dV + \text{surface term}$$

Substituting in the expression for \mathbf{q} , we rewrite the integral on the RHS as

$$-\int n\left(\chi_{||}-\chi_{\perp}\right)\tilde{q}_{||}\hat{\mathbf{b}}\cdot\nabla\lambda_{i}dV-\int n\chi_{\perp}\nabla T\cdot\nabla\lambda_{i}dV$$

The key procedure is the replacement of the factor $\hat{b} \cdot \nabla T$ in the first integral with

$$\widetilde{q}_{||m} \equiv rac{\int_{\Delta_m} \hat{\mathbf{b}} \cdot \nabla T dV}{\int_{\Delta_m} dV},$$

reducing the order of the representation by one, in this case from a C^0 bilinear function to a C^{-1} piecewise constant function over element *m*.

M3D Implementation

New routine advances

$$\frac{\partial \tilde{T}}{\partial t} = \nabla \cdot \left\{ n \left[\left(\chi_{||} - \chi_{\perp} \right) \hat{\mathbf{b}} \, \hat{\mathbf{b}} \cdot \nabla \, \tilde{T} + \chi_{\perp} \nabla \, \tilde{T} \right] \right\}$$

where

$$\tilde{T} = T - T_0$$

with arbitrary source term T_0 ; performs 3D numerical Galerkin integration (3×7-point quadrature) using trilinear test functions to construct global mass matrix **M** and heat conduction operator matrix **Q**; solves

Vertex values affect function over 12 adjoining triangular prisms

$$\left(\overline{\mathbf{M}} - \theta \delta t \overline{\mathbf{Q}}\right) \cdot \widetilde{\mathbf{T}}^{n+1} = \left[\overline{\mathbf{M}} + (1 - \theta) \delta t \overline{\mathbf{Q}}\right] \cdot \widetilde{\mathbf{T}}^{n}$$

using GMRES, $\theta = 1/2$ for 2nd-order accurate implicit time advance.

Qualitative Comparison with Artificial Sound Wave Method

- Use circular cross-section equilibrium, aspect ratio 3; 1.33 < q < 4.79.
 - Peak $T = 3.1 \times 10^{-7}$ near axis.
- Add temperature perturbation $\delta T = 3 \times 10^{-5} \exp\left\{-\left(\frac{\varphi - \pi}{0.07709}\right)^2 - \left[\frac{(R - 3.53)^2 + (z - 0.53)^2}{(0.04)^2}\right]\right\}$

• Freeze density and **B** and **v** fields, evolve *T* only; $\chi_{\perp}=10^{-50}$ and either $s=6, v=10^{-3}$ or $\chi_{\parallel}=60$



Artificial Sound Results

t = 0.025



t = 1.5



Implicit χ_{\parallel} Results

t = 0.025



t = 0.15



Late Time States

Implicit $\chi_{||}$, t = 15.0

Artificial sound, t = 9.0



Numerical instability

Quantitative Test of Accuracy

The figure of merit is the normalized RMS parallel temperature gradient, defined as



2nd Test: Find steady state using misaligned grid with inhomogeneous bcs

- Start with circular cross-section equilibrium, aspect ratio three. ۲
- Add 1,1 helical perturbation to poloidal flux to shift surfaces away from ٠ mesh.
- Use boundary condition T=0 on outer surface; T=1 on inner surface defined ۲ by ψ =-0.235.
- Run to steady state using 32 planes, 141 radial zones.





Detail of misaligned mesh, surfaces





Summary

- True implicit parallel and perpendicular heat conduction have been implemented in M3D.
- Use of Günter's method appears to allow high accuracy with linear finite elements even when tensor is highly anisotropic and mesh is misaligned.
- The new operator is an improvement on the old one in some respects.
- Further verification is needed. Because cases with analytic solutions are scarce, benchmarking with other codes would be useful.

M3D Scales to 12k+ Franklin Processors



• Algebraic multigrid preconditioner used for GMRES solves.

• Reverse Cuthill-McKee (RCM) matrix re-ordering used to reduce fill-in in ILU preconditioning for other linear solves.

• Base case is nonlinear C-Mod sawtooth with 24,000 vertices/CPU.

• ~200 toroidal CPUs, ~50 poloidal CPUS for largest cases.

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