

# Linear and nonlinear stability of a toroidal plasma and M3D results

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# Topics

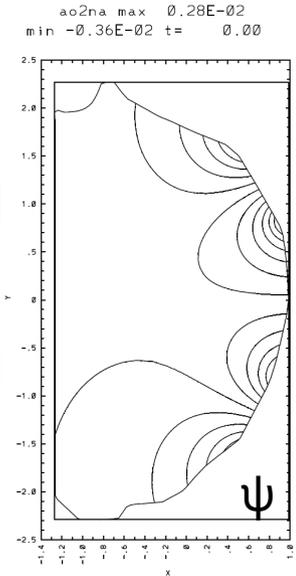
- Summary of current M3D simulation work
- Linear and nonlinear stability of ELM edge instability
  - Differences
  - Magnetic tangle?
- New projects/ideas

# Summary of recent edge simulation work

- DIII-D Type I ELM (119690, 126006)
  - Nonlinear compared to linear stability and growth rates
  - New form for I-coil fields for M3D ( $B_\phi$  and  $A_\phi$ ) with high toroidal resolution + other fields (C-coil, bus, error): D. Orlov, UC-SD
  - Full toroidal spectrum of  $n=3$  RMP I-coil field – Fourier aliasing
  - TBD: Add all non-axisymmetric fields to ELM (126006)
  - Two-fluid + toroidal rotation effect on growth rates - important!
    - Combination has NL stabilizing effect, stronger than either alone
- NSTX lithium divertor ELM suppression
  - ELM seen in both pre-lithium and lithium 'stabilized' case (MHD, at high  $\eta$ . Numerical stability worse in ST geometry. Better grid, higher resolution.)
- CMOD EDA regime with QCM edge oscillation
  - Diamagnetic-profile equilibrium is MHD stable at 10x actual resistivity
  - Waiting for kinetic profile equilibrium, part of Joint Milestone 2011.

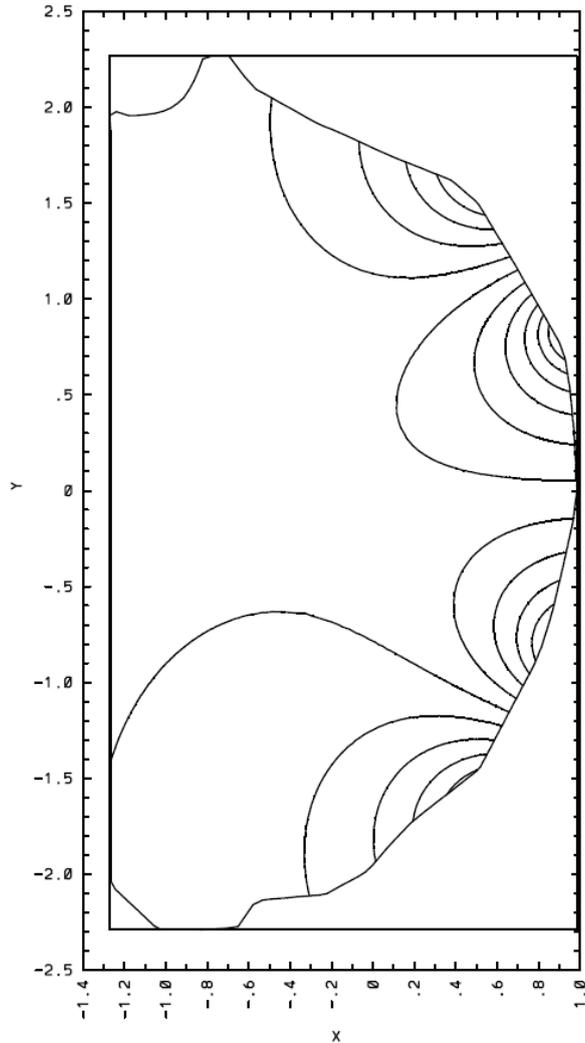
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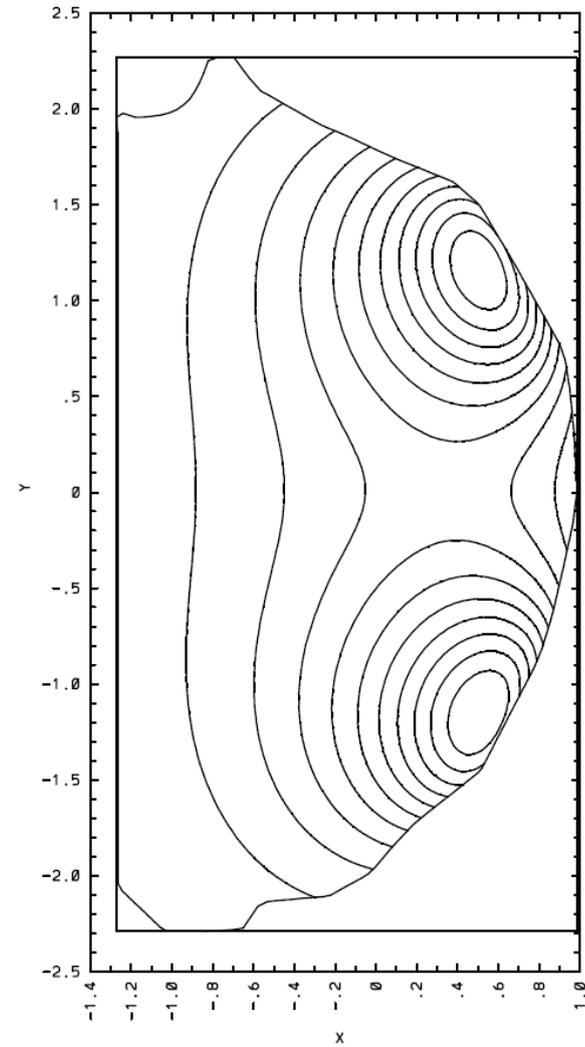
# RMP (I-coil) field

ao2na max 0.28E-02  
min -0.36E-02 t= 0.00



$(R/R_0)\psi$

rm $\psi$  max 0.77E-02  
min 0.30E-03 t= 0.00



$RB_\phi$

# Summary of edge simulations -2-

- DIII-D EHO 128542
  - Ran original equilibrium in MHD/ MHD+rot/ 2F/ 2F+rot – similar type of instability with different growth rates; low  $n=2,3$  as in experiment
  - Running new equilibrium reconstruction with  $q_0 > 1.0$ ; strange.
  - TBD: add error-field and correction (I-coil).
- MAST ELM – new case; testing.
- New wall-load and divertor diagnostic (VisIt), including fluxes  $v_n \cdot X$ 
  - ELM divertor traces qualitatively similar to experimental measurements
  - Harmonics to compare to experiment
- Still missing cases for V&V:
  - Pre/post RMP comparison for density pumpout in H-mode. (Can't use 126006. DIII-D cases identified, waiting for good data.)
  - ELM crash with fast time data for detailed comparison to expt.
- NERSC Cray XT-6 now allows NL sim at  $(2x)^3$  resolution of 2009/10 with good turn around!

# Linear vs Nonlinear ELM instability

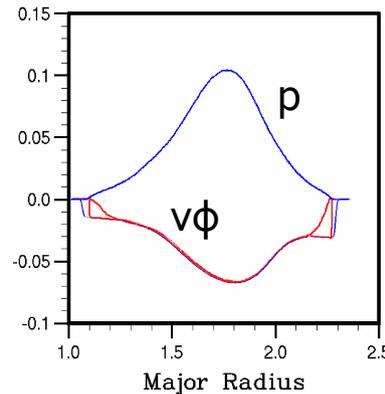
- What role does the magnetic tangle play in stability?
- Linear perturbation theory excludes full magnetic tangle
  - Small magnitude of perturbation,  $|\tilde{p}| \ll |p_o|$
  - Single toroidal harmonic  $\leftrightarrow$  linearization drops nominally smaller terms
- In theory, a magnetic tangle results from any small enough transverse perturbation; should be biggest near X-point, away from ballooning-type instability driving term at midplane
  - Not linearized, not single harmonic; stochastic
  - Asymptotic field line splitting in different directions; Field splitting can be obtained by linear superposition of equilibrium + single-n perturbing field
  - Not flux tube boundary conditions
  - X-point system does not preserve energy since X-lines intersect domain boundary (Only small exterior effect for simulation?)
- Other nonlinear effects are important in ELM
  - NL harmonic interaction leads to low mode-number and n=1 effects

# Linear vs Nonlinear: Growth rates for DIII-D ELM

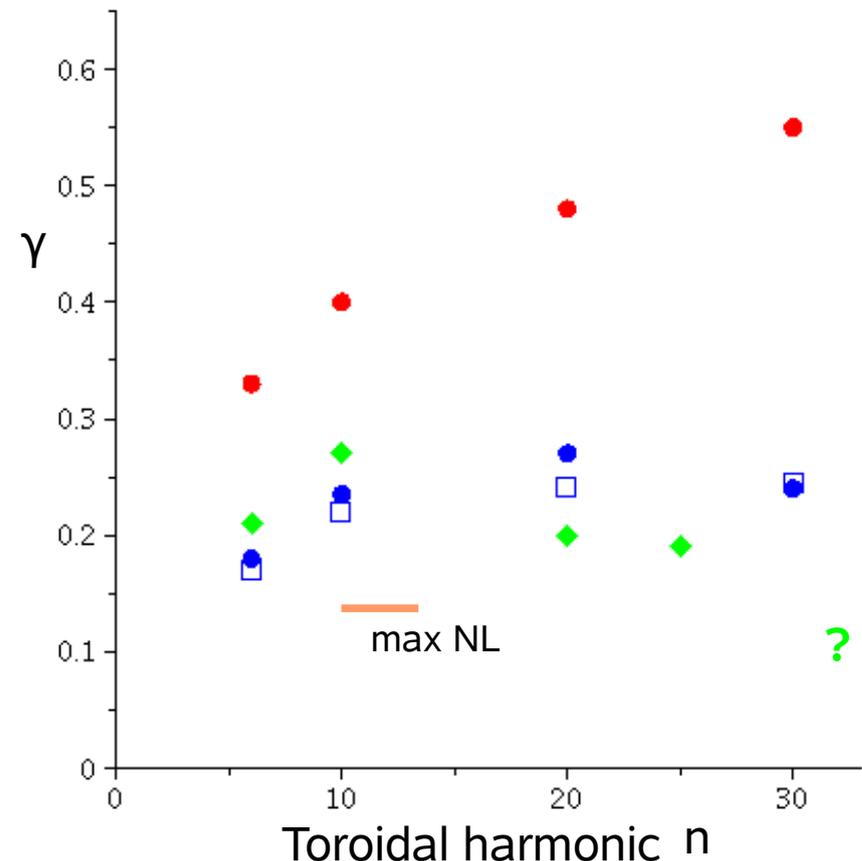
- Compute linear growth rates for the DIII-D Type I ELM 126006 case
  - Match nonlinear simulation conditions, for comparison
- **Linear and nonlinear growth rates are different in MHD.** Strong nonlinear effects affect ELM at finite, but nonlinearly small, size.
  - Linear rates show expected MHD ballooning dependence; inc. with  $n$
  - Nonlinear, dominant harmonics are moderate  $n=10,13$ ;  $m$
  - Maximum NL  $\gamma$  is smaller (0.13 vs. 0.5+ for  $n=23$  or 0.35 for  $n=10$ ).
  - Linear  $\gamma$  reduced by toroidal rotation, maximum NL  $\gamma$  increased.
- Not exactly same models
  - Linear pert has no  $\partial n/\partial t$ ; NL evolves density.
- More accurate linear perturbation results should use higher resolution, especially higher harmonics with strong rotation shear.

# Linear growth rates reduced by toroidal rotation

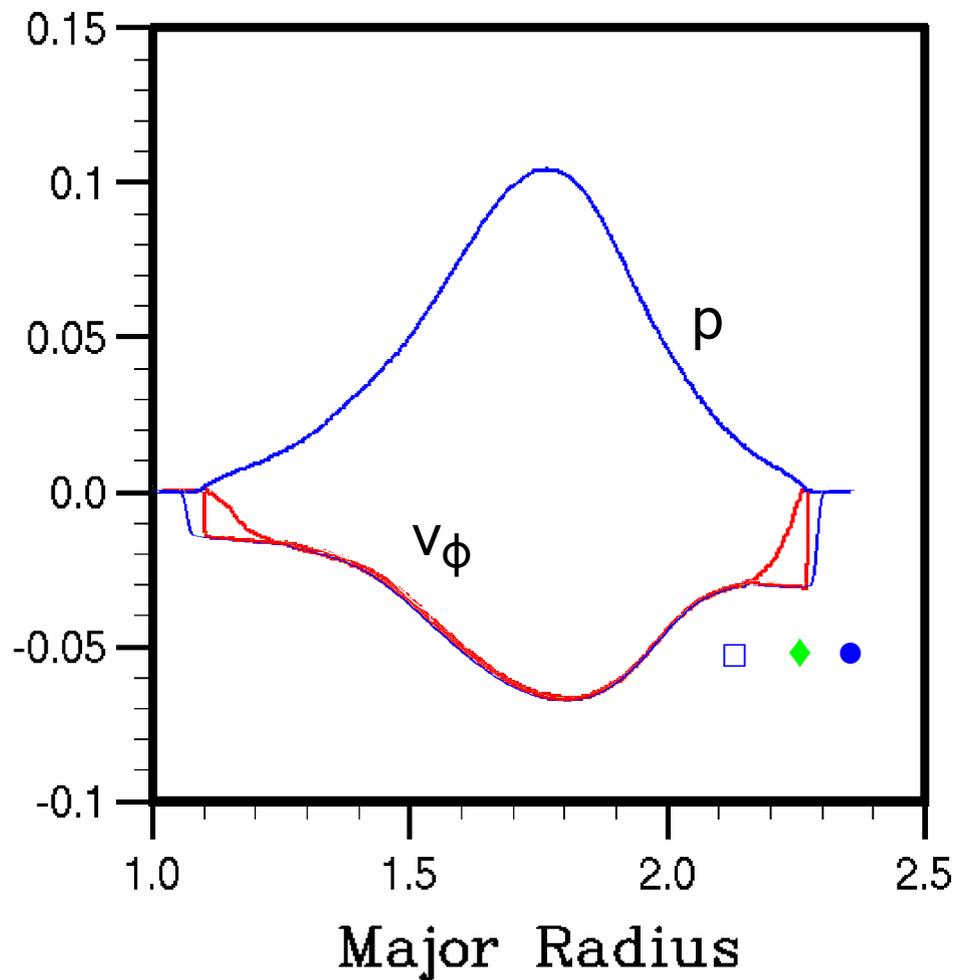
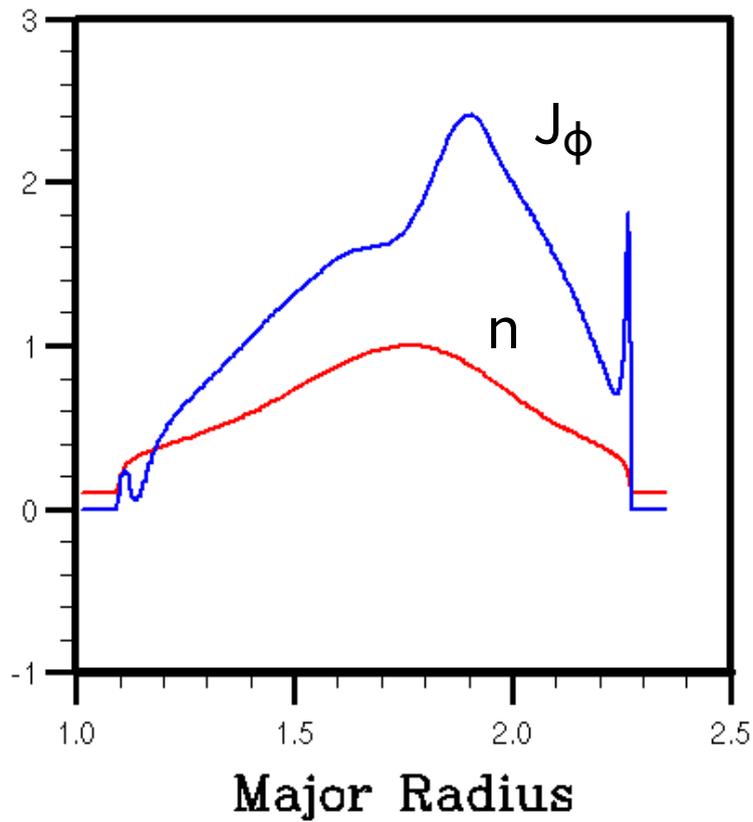
- DIII-D ELM 126006
- MHD without rotation (●) has expected ballooning behavior
- MHD with toroidal rotation, varying edge rotational shear
  - Experimental profile (□)
  - Modified to constant  $\Omega$  over  $0.7 < \psi < 1$  near plasma edge, chopped to  $\Omega=0$  at separatrix (◆) ( $n=30$  not converged)
  - Const over edge, but  $\Omega \rightarrow 0$  smoothly (tanh) starting outside separatrix but well inside wall (●)
- Rotation is stabilizing
- Rotational shear effects weak, unless shear is very strong



- MHD, no rotation
- Rotation, exptl profile
- ◆ Rot, const over edge, zero at separatrix
- Rot, const over edge, smoothly to zero outside



# Midplane profiles, expanded



DIII-D 126006 ELM  
 $\eta = 3 \times 10^{-8}$ ,  $\mu = 6 \times 10^{-6}$

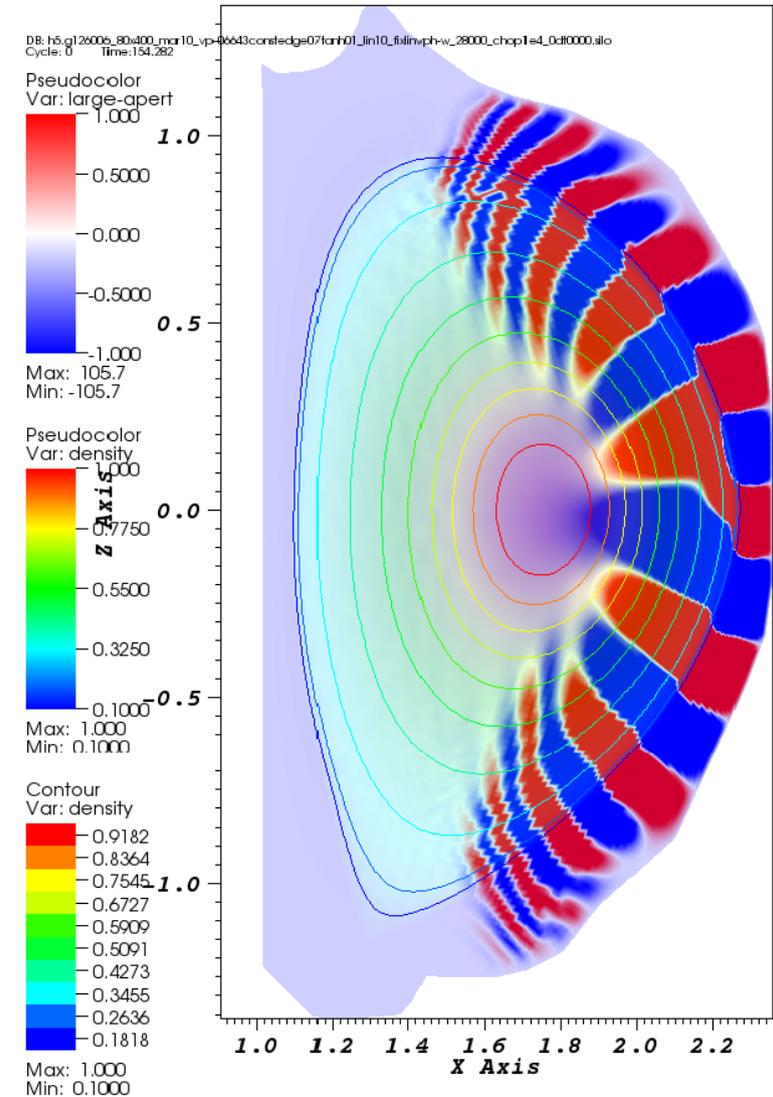
# Linearization excludes the formation of a magnetic tangle

- Linearized magnetic perturbation can only be large where plasma perturbation is large, i.e.,  $\tilde{\mathbf{J}}$  and  $\tilde{\psi}$  must match  $\tilde{p}$ :

$$\tilde{\mathbf{J}} \times \mathbf{B}_o + \mathbf{J}_o \times \tilde{\mathbf{B}} = \nabla \tilde{p} + \rho_o \left( \frac{\partial \tilde{\mathbf{v}}}{\partial t} + (\mathbf{v}_{\phi o} \cdot \nabla) \tilde{\mathbf{v}} \right) + \dots$$

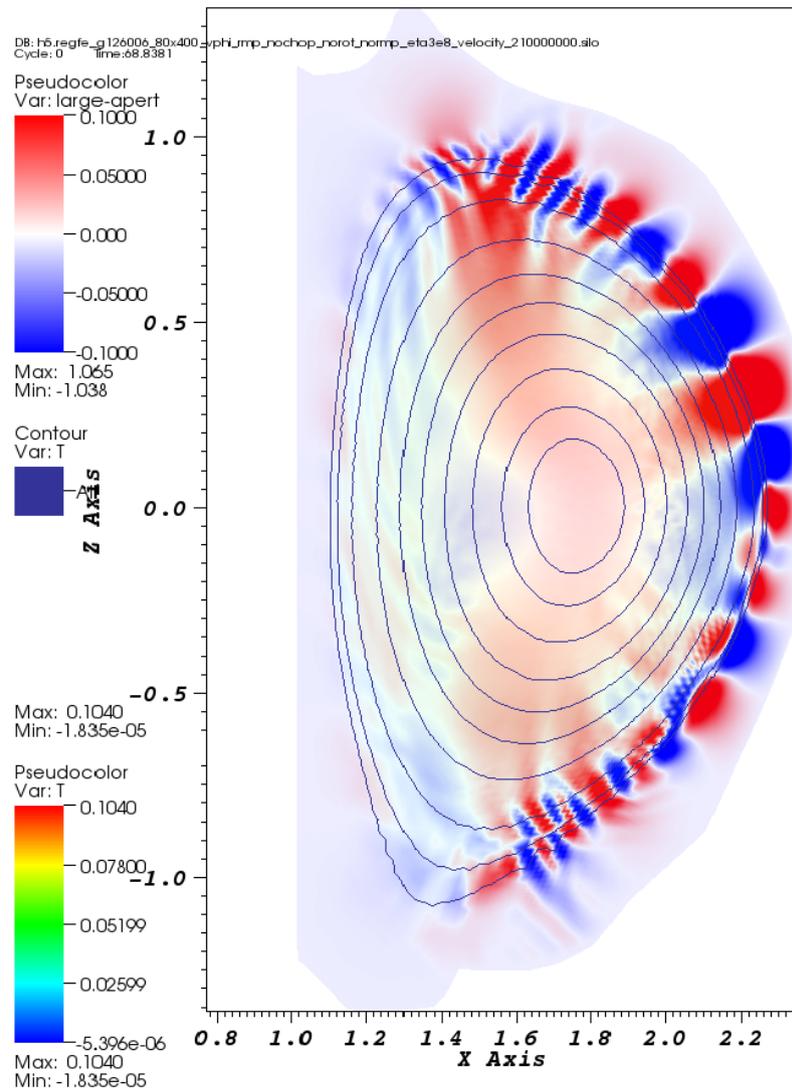
- Test: Calculate linear perturbation in standard manner\*, then multiply it by very large factor and plot →

- X-point regions and inboard side have  $|\tilde{\psi}|/|\tilde{\psi}_{max}|$  less than  $10^{-8}$ . Poloidal extent is approximately the same for all variables.
- $\tilde{\psi}$  plotted on density, n=10 mode, rotation (●)
- \*Linear calculation solves almost fully NL equations. After each  $\Delta t$  time step, resets n=0 part and filters to given harmonic n, also controls perturbation magnitude.

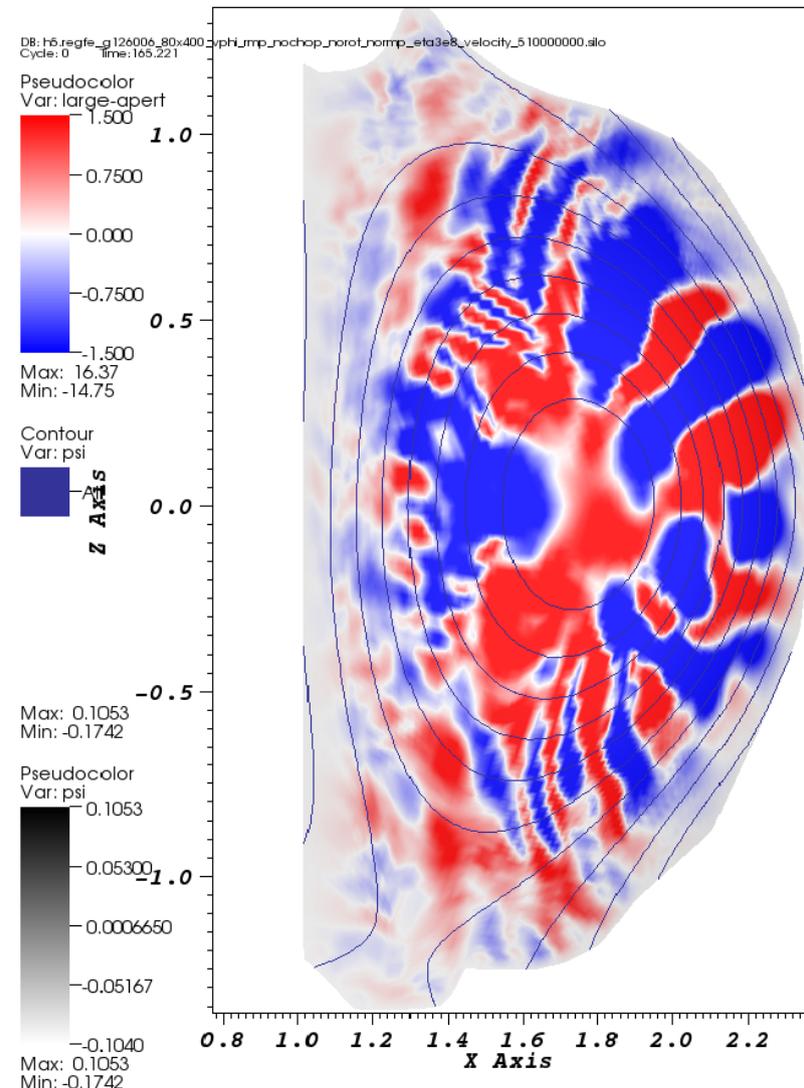


$\tilde{\psi}$  on density, n=10

- Nonlinear ELM forms a tangle early (inboard and near-X fraction of  $\tilde{\psi}$  is 1/20-1/200 of maximum instead of less than  $10^{-8}$ ; fraction declines later, but absolute magnitude increases



Just before crash, t=68.3



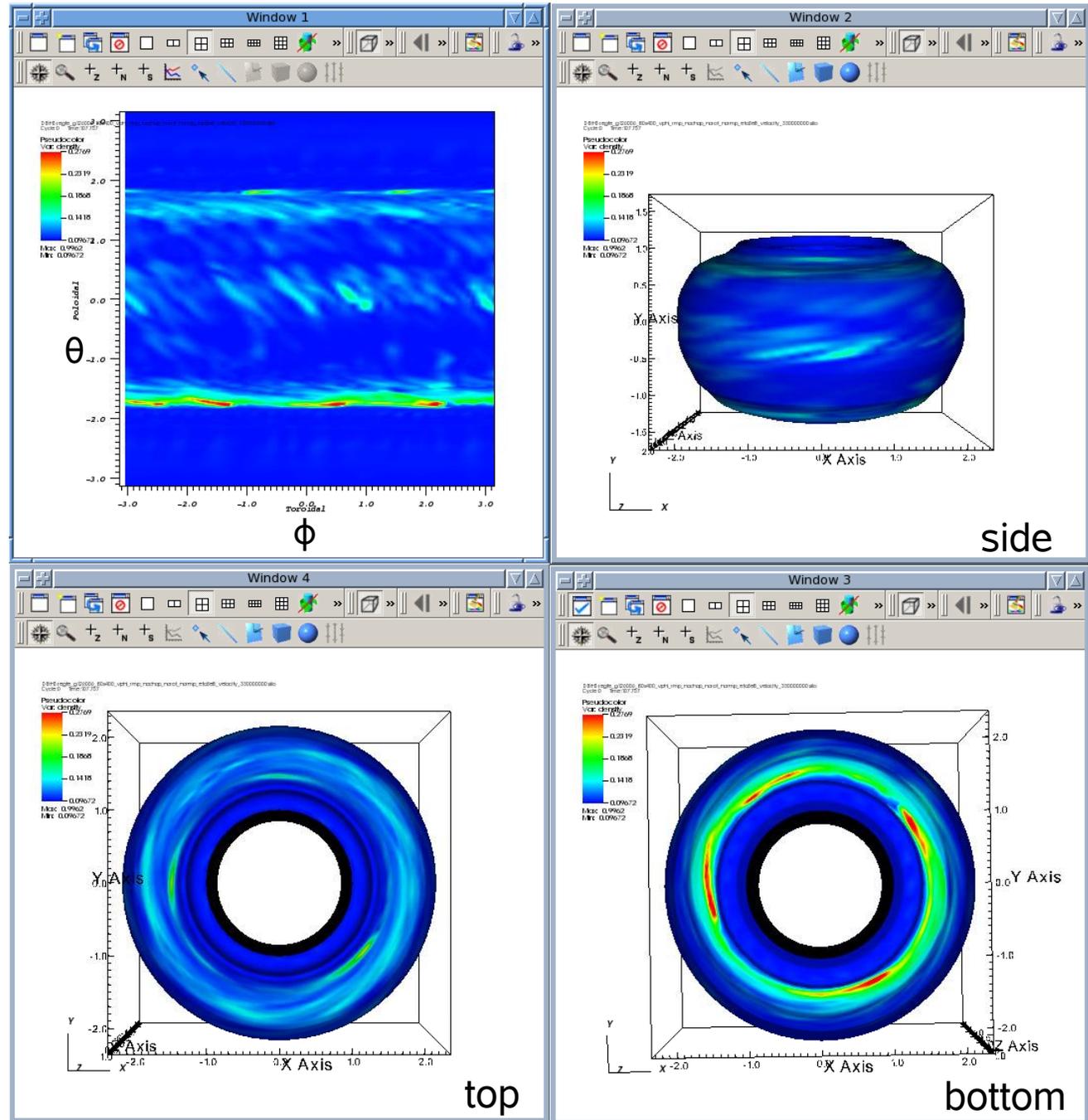
After initial outburst, t=165  
max  $\psi$  x16 vs earlier time

# Comments

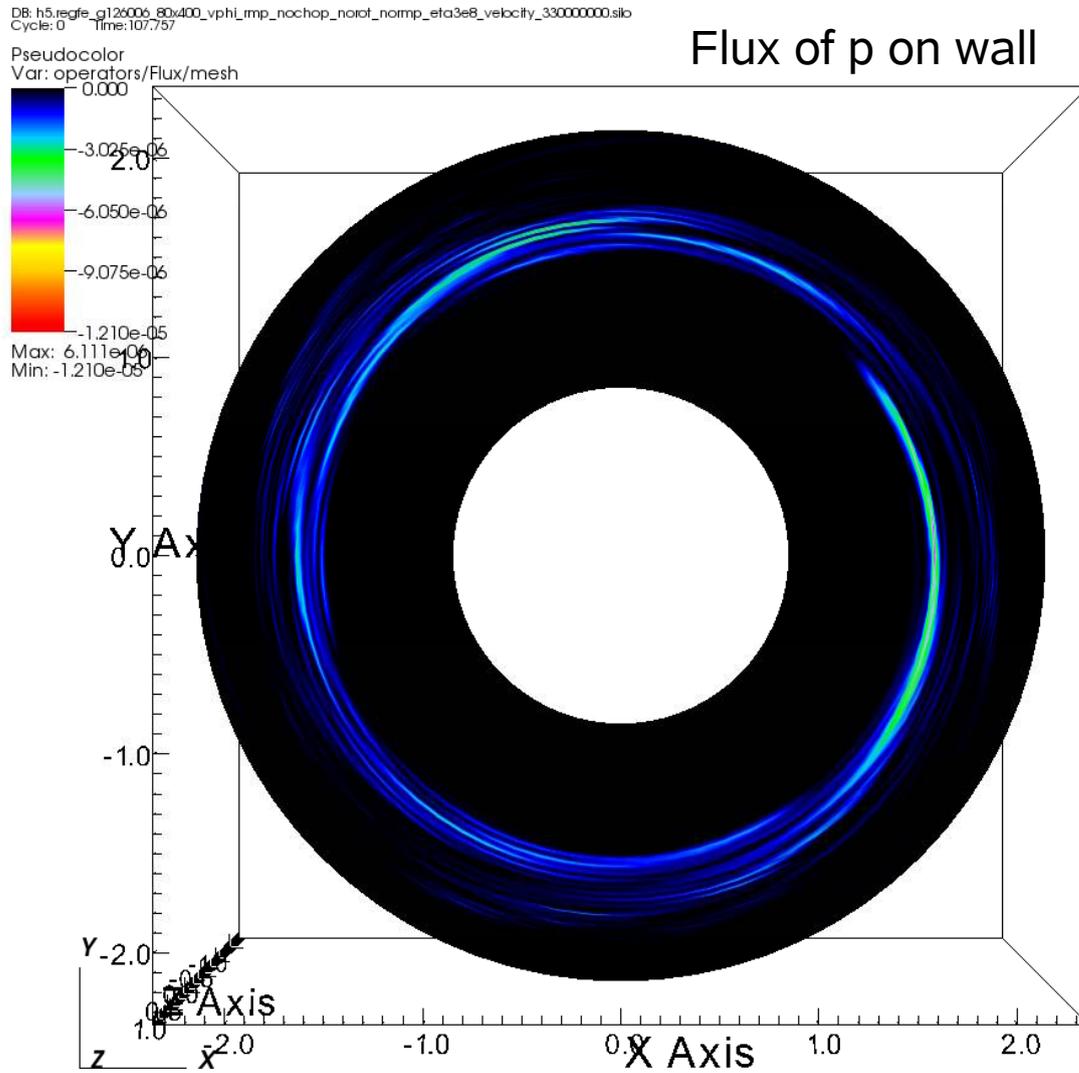
- Absence of the magnetic tangle in linear simulation is unlikely to result from numerical reasons (e.g., not enough resolution near X-point), but cannot be completely ruled.
- Linear theory does not predict tangle; major restrictions agree with those in simulation, so conclude '**No linearized tangle**'.
  - Tangle requires propagation along B; too slow to grow as  $e^{\gamma t}$
- Magnetic tangle should be stabilizing nonlinearly, since requires additional work to drive a field perturbation away from main plasma instability
  - Seen in an indirect test of nonlinear evolution (Sugiyama, PoP 2010)
- Some other results suggest that the density evolution may have strong stabilizing effects on linear edge perturbations (not part of standard MHD linear model)
  - Here, no  $\partial n/\partial t$  in linear case, since very steep edge density gradient in ELM case needs to be better resolved for linear convergence.

# ELM wall loads are strongly asymmetric

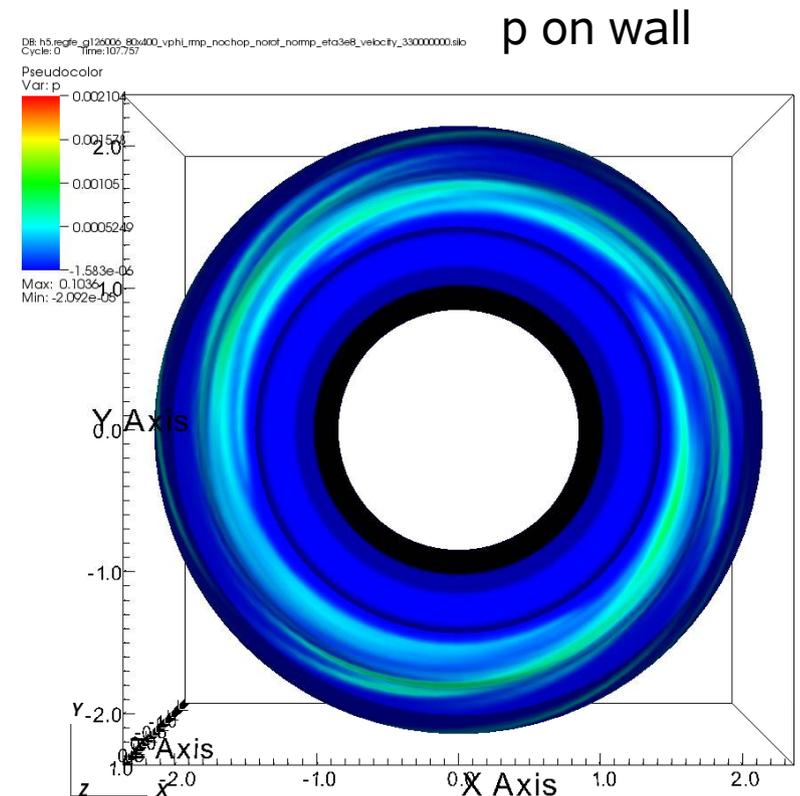
- **New wall diagnostic in VisIt for M3D** (LBL Vis group, H. Childs)
- Strong asymmetry in divertor (when density first hits lower outer divertor)  
Concentrated points during crash!
- Helical stripes on top and bottom divertors follow field lines, overall
- Locations and magnitudes change on fast MHD time scale



# Flux of pressure ( $p \cdot v_n$ ) shows multiple striations in divertor

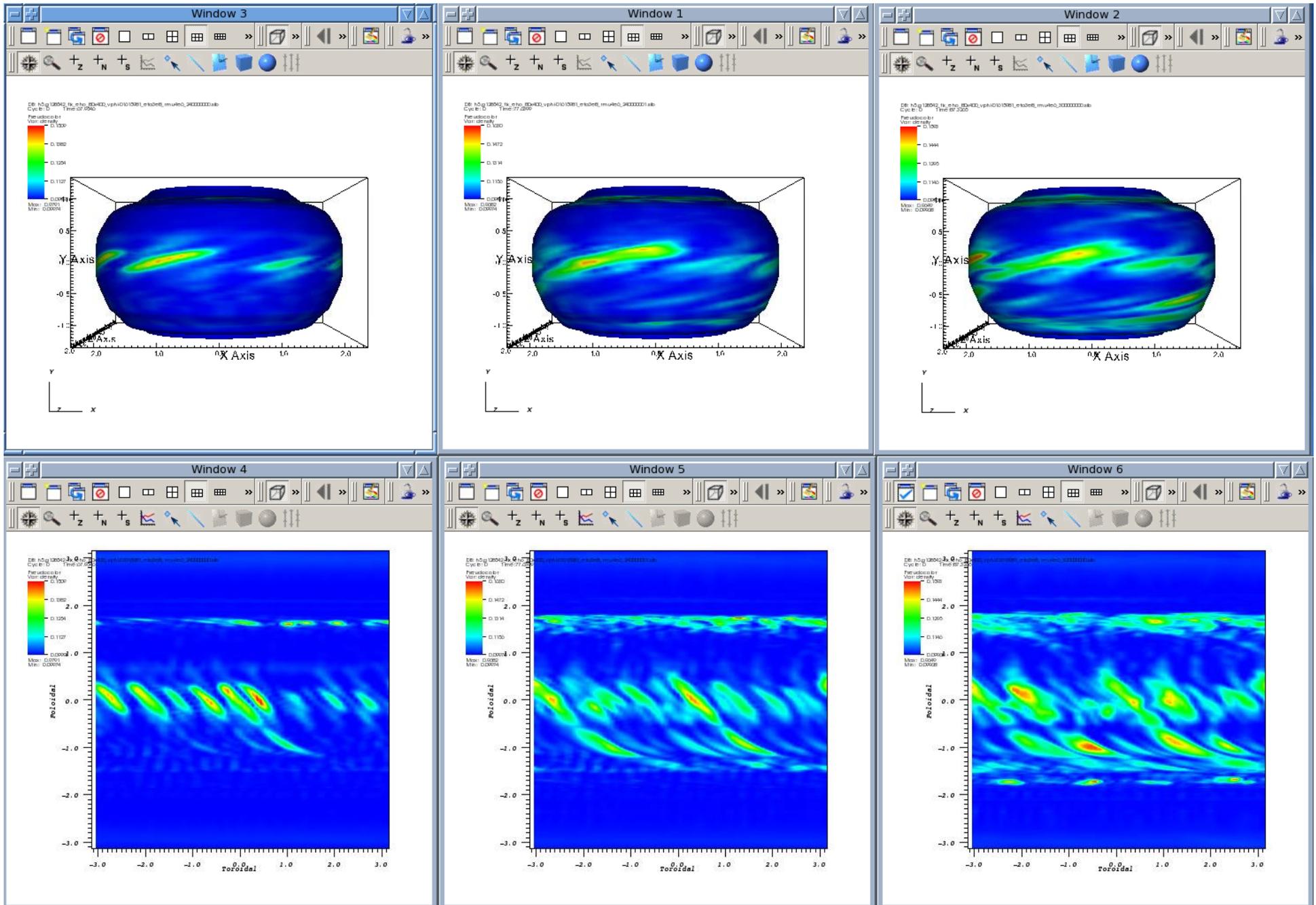


- Raw variable p (or n) is broader, smoother



- Need time integrals to match experimental observations

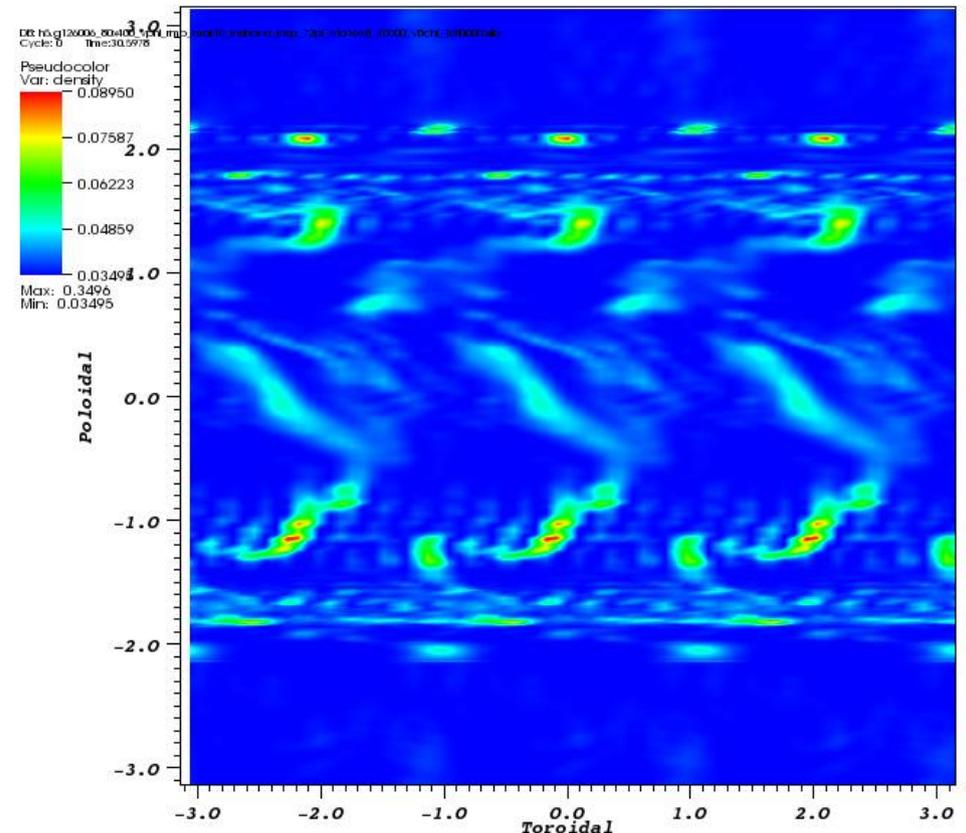
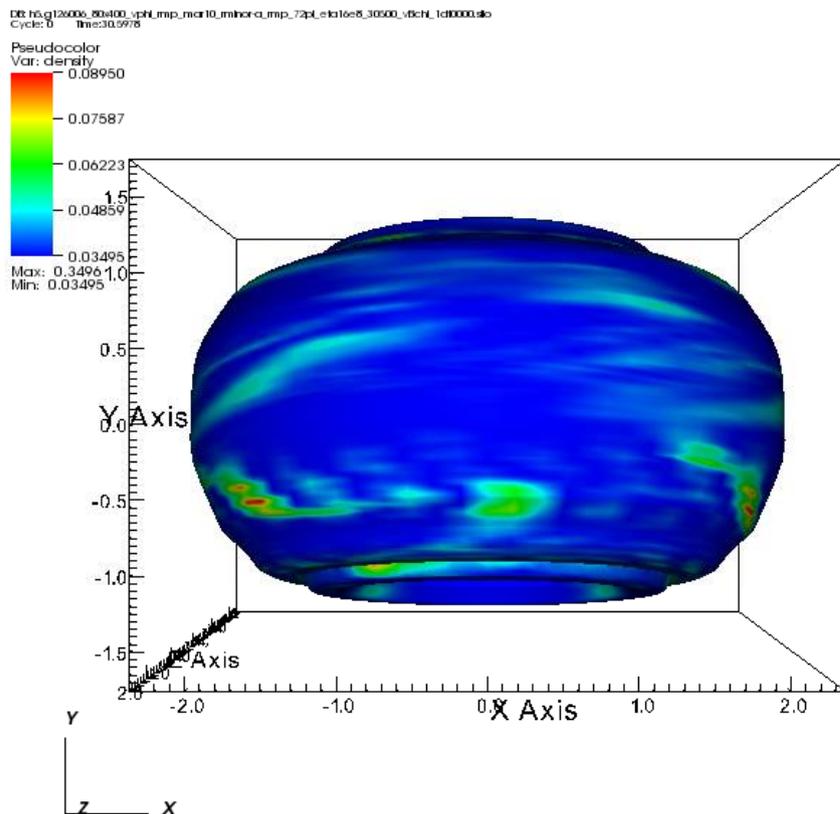
# DIII-D 128542 EHO instability: Density on wall. (Experiment has dominant n=2 mode.)



time →

# DIII-D RMP fields from I-coil

- Full I-coil field with new algorithm for M3D by D. Orlov:  $B_\phi$  and  $A_\phi$  (magnetic vector potential, from J. Hansen (2002) algorithm)
- Analysis of spectrum shows  $n=3,9, \dots$ ,  $B_\phi$  has many harmonics near coils
- Preliminary example: density pump-out to wall with  $n=3$  (*old* RMP!)



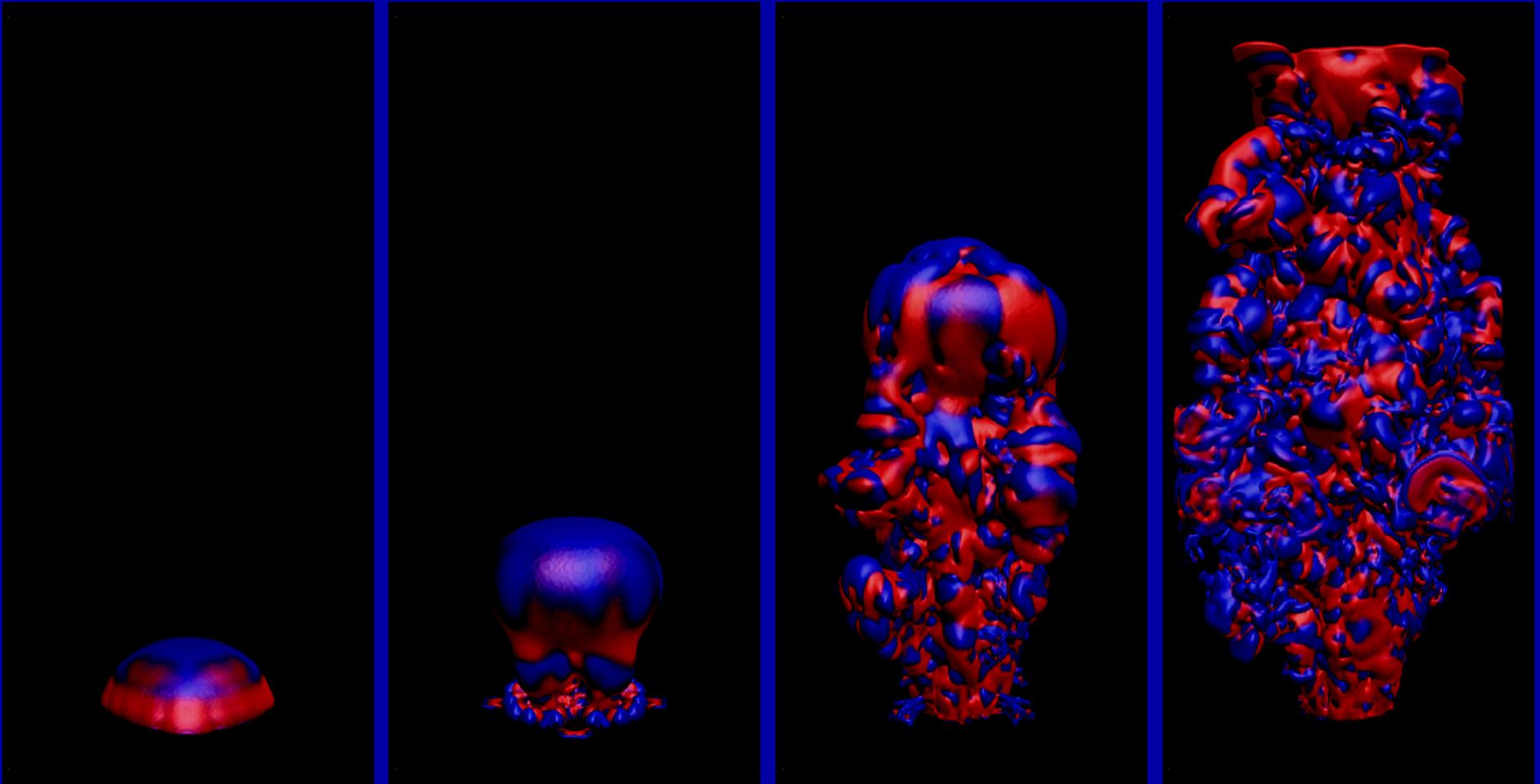
# Finite time Lyapunov exponents (FTLEs)

- How to characterize magnetic tangle and other superficially stochastic structures?
- Local definition of structures: Finite time Lyapunov exponents measure how fast local structures move apart or together.
  - Unlike regular Lyapunov exponents, which measure infinite-time growth or convergence
- New techniques to determine FTLEs for ordinary fluids recently developed and tested against experiments (mostly 2D, starting to go to 3D). Increasingly useful for real problems.
  - Haller (2001); Mathematical foundation (Shadden 2005)
- Older applications to plasmas targeted mainly homogeneous turbulence. Now, becoming practical to apply to instabilities with real structures.
- Bridge modern ideas of fluid turbulence/mixing and plasmas
- Study fundamental questions: Incompressible vs compressible MHD (also differences in GK and MHD magnetic evolution), num stability

# FTLEs for plasmas

- Visualization techniques to compute FTLEs for fluid velocity field  $\mathbf{v}$  are active area of research
  - Trace 'particle' paths in flow field, extract relative motion.
- Apply to plasma: not only  $\mathbf{v}$ , but  $\mathbf{B}$ 
  - Some extensions can be developed.
  - Incompressible MHD:  $\mathbf{v} \pm \mathbf{B}/\rho^{1/2}$ , vorticity  $w \pm J_\phi$ , etc.
- Nonlinear MHD simulations are a good test bed
  - H. Krishnan, LBL Vis Group (post-doc) working on M3D data
- Finding FTLEs is related to certain types of feature extraction (identify hills, valleys, level contours), so FTLE techniques can help analyze dynamic plasma structures, independent of the Lyapunov meaning
- Techniques will improve as computation capability improves
  - Next generation computing: highly parallelizable

# FTLE example: Fluid Jet



**Figure 4.** This image shows direct volume rendering of the time-varying Finite-Time Lyapunov Exponent fields (red indicates the forward-time exponent, blue shows the backward-time exponent) for four time steps, illustrating the formation of turbulence in a high-speed jet of entering a domain of stationary fluid. Individual turbulent structures and structure size and distribution can be observed directly from the volume rendering. Data set: C. Garth (UC Davis).

(VACETS SciDAC center (2007))

# Summary

- Linear vs nonlinear edge instability – important differences
  - Magnetic tangle is nonlinear
- MHD plasma edge stability/instability ongoing
  - Edge (ELM, ELM-free oscillations)
  - Edge + interior mode coupling
  - Initial two-fluid shows two-fluid+rotation is important
- Developing theoretical and practical tools to study questions raised: visualization with help from LBL/NERSC Vis group
  - Wall load and wall-flux diagnostic (working)
  - Finite time Lyapunov exponents
    - MHD structures and evolution: identification, local stability
    - Develop extensions from fluids to plasmas
- Other areas not discussed here...