## **QAS Design of the DEMO Reactor**

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The NSTAB code solves differential equations in conservation form, and the TRAN test particle code tracks guiding center orbits in a fixed background, to provide simulations of equilibrium, stability and transport in tokamaks and stellarators. These codes are well correlated with experimental observations and have been validated by convergence studies. Bifurcated 3D solutions of the 2D tokamak problem have been calculated that show persistent disruptions and ELMs crashes occurring in ITER, which does not pass the NSTAB simulation test. Therefore we have designed a QAS stellarator with similar proportions as a candidate for the DEMO fusion reactor [1]. Our configuration has two field periods and an exceptionally accurate 2D symmetry that furnishes excellent thermal confinement and good control of the prompt loss of alpha particles. Robust coils are found from a filtered form of the Biot-Savart law based on a distribution of current over a control surface for the coils and the current in the plasma defined by the equilibrium calculation. Computational science has settled these issues of equilibrium, stability, and transport, so what remains to be developed is an effective plan to construct the coils and build a divertor.

[1] Garabedian, P.R. and McFadden, G.B., J. Research NIST 114 (2009) 229-236; online at http://www.nist.gov/jres.

## **Quasi-Axisymmetric Stellarator Design**



Magnetic fusion of hydrogen forms helium and emits energetic neutrons from a torus of plasma shaped to optimize confinement. Modular coils generate a field keeping the ions from hitting walls.

### **QAS Design**



This configuration is a good candidate for DEMO because its geometry is similar to that of the ITER tokamak. There is enough flexibility in the shaping of the plasma to control the stellarator contribution to  $\iota(s)$  so as to compensate for unforeseen complications with the bootstrap current.

## **QAS Flux Surfaces**





Flux surfaces at four cross sections over the two field periods of a stellarator equilibrium at  $\beta = 0.04$ .

## **QAS Pressure and Rotational**



#### **Transform Profiles**



For toroidal flux in the interval 0 < s < 1 we have pressure  $p(s) = 0.039(1 - s^2)^{1.5}$  and rotational transform  $l(s) = 0.55 - 0.40(0.68 - s)^2$ .

# **QAS Outer Boundary**



## In cylindrical coordinates (r, v, z):

$$r(u,v) + iz(u,v) = e^{iu} \Sigma \Delta_{mn} e^{-imu + inv}$$

**Table 1.** Fourier coefficients  $\Delta_{mn}$  defining the QAS stellarator in Fig. 2 for values m = -1, 0, 1, 2, 3, 4 of the poloidal index and values n = -1, 0, 1, 2, 3, 4 of the toroidal index. This is a configuration that depends on bootstrap current for good performance

$m \setminus n$	-1	0	1	2	3	4
-1	0.200	0.140	0.000	0.000	0.000	0.000
0	0.000	1.000	0.006	0.007	0.000	0.000
1	0.042	2.500	0.084	0.006	-0.002	0.000
2	0.010	-0.100	-0.350	-0.058	-0.010	0.000
3	0.000	0.000	0.030	0.050	0.017	0.004
4	0.000	0.000	0.025	-0.020	0.000	0.000

# **QAS Spectrum B**<sub>mn</sub>(s)



$$B = \nabla S \times \nabla \theta = \nabla \phi + \varsigma \nabla S$$

$$\frac{1}{|B|^2} = \sum B_{mn}(s) \cos[m\theta - (n - \iota m)\phi]$$

$$\frac{J\mathbf{g}\mathbf{B}}{|\mathbf{B}|^2} = p'(s)\sum \frac{mB_{mn}(s)}{[n-\iota(s)m]}\cos[m\theta - (n-\iota(s)m)\phi]$$

$$\delta B_{mn}(s) \sim \delta \Delta_{mn}$$

B <sub>mn</sub> (s)>0.005							
m	n	0.188	s 0.535	0.966			
0	0	0.913	0.863	0.789			
0	1	0.012	-0.007	-0.016			
1	-1	0.003	-0.003	-0.009			
1	0	0.222	0.348	0.405			
1	1	0.000	-0.002	0.017			
1	3	0.002	0.003	0.005			
2	0	0.027	0.082	0.128			
2	1	-0.000	0.000	-0.001			
2	2	0.004	0.005	-0.001			
3	0	0.003	0.019	0.047			
3	1	0.001	0.006	-0.008			
3	2	0.001	0.004	0.006			
3	3	-0.001	0.000	0.005			
4	0	-0.001	0.004	0.018			
5	0	-0.001	0.001	0.007			







Three distinct coils of the QAS stellarator with just two field periods. When the torus of plasma is removed the coils are seen to be well enough separated so they can be constructed on the length scale of a reactor.

### **Bifurcated Tokamak Equilibrium**



Left: cross section of a bifurcated tokamak equilibrium with  $\beta = 0.035$  displaying a magnetic island at  $\iota = 2/3$ . The 2D solution (right) is stable for small perturbations of the plasma, but becomes unstable for larger 3D displacements. This behavior is symptomatic of neoclassical tearing modes (NTMs) and edge localized modes (ELMs).



For toroidal flux in the interval 0 < s < 1 we have pressure  $p(s) = 0.037(1 - s^{1.8})^{1.8}$  and rotational transform  $\iota(s) = 0.880 - 0.322$  s.



Convergence of the discrete 3D solution of a 2D tokamak problem to the level of roundoff error after 10<sup>7</sup> iteration cycles. This NSTAB simulation establishes the risk of disruptions and ELMs crashes in the ITER project.

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### Thank You