



Studies of Improved Confinement in Heliotron J

Presented by
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for Heliotron J Team

Studies of Improved Confinement in Heliotron J

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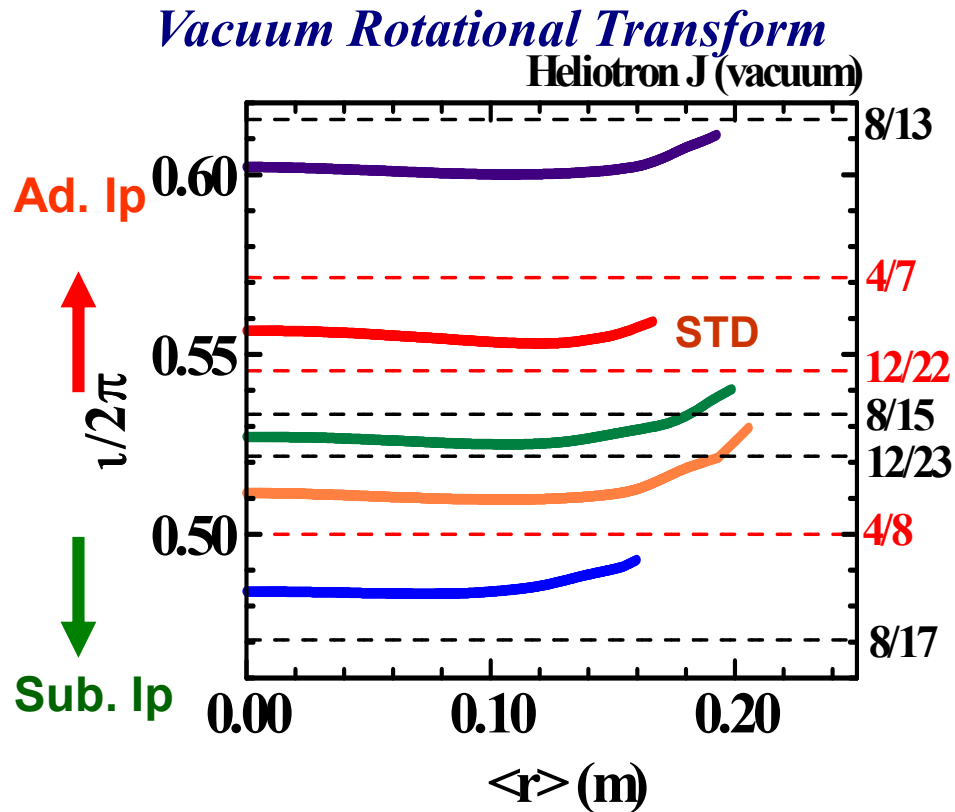
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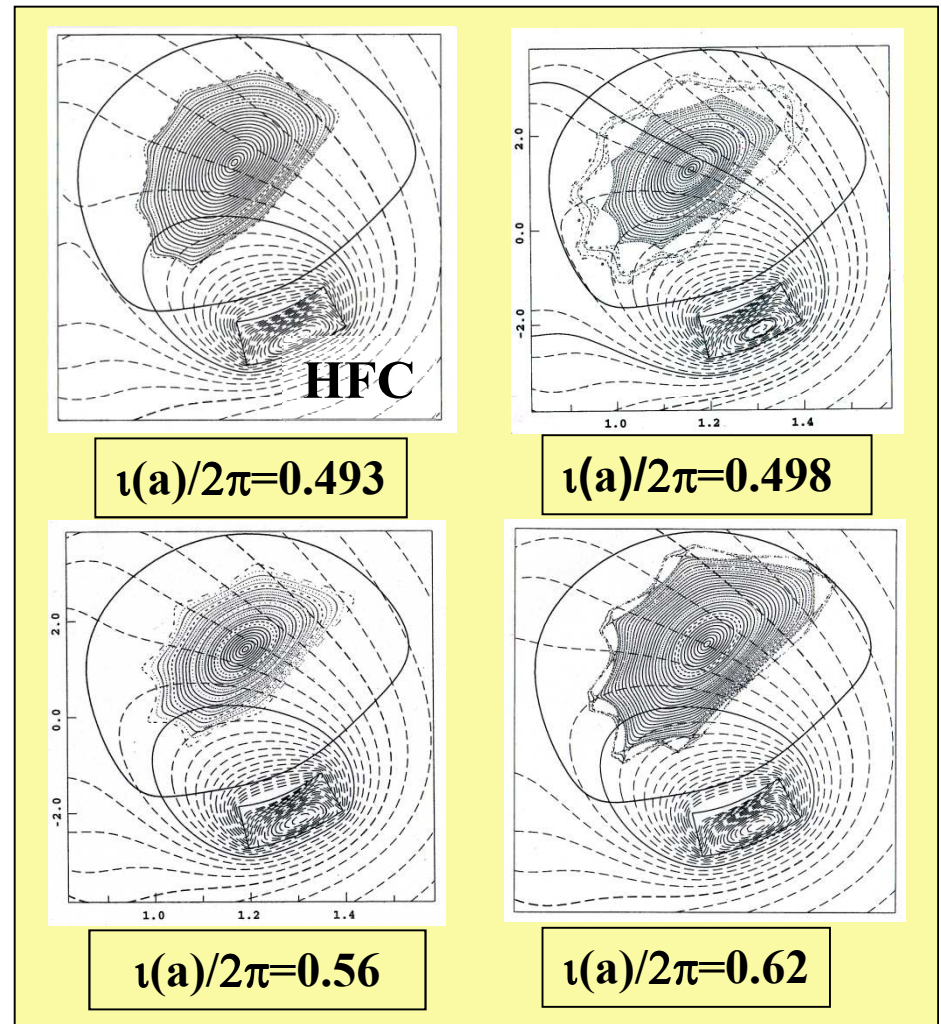
- **Introduction**
- **Transition to Improved Confinement**
 - Configuration Effects
 - Effects of Heating Scenario
- **Fueling Control**
- **Summary**

*Due to the time limitation, this talk will focused on “Transition”.
For other topics, please contact me or refer the proceeding*

The edge magnetic surfaces are strongly modified by the resonant conditions.



The effects of rational surfaces have been studied is not only from the MHD activities but also from the appearance of transport barriers viewpoints.



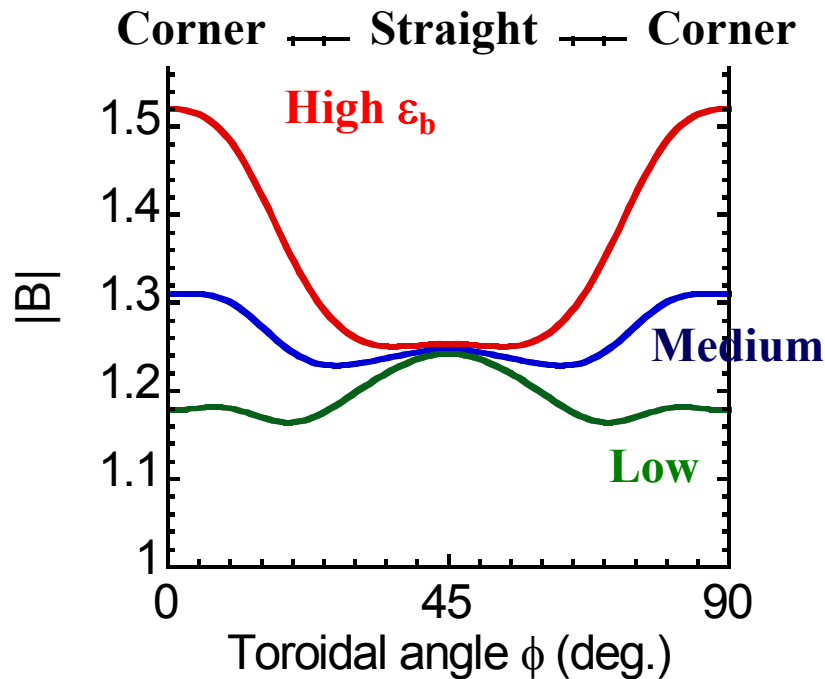
Bumpiness Control

∇B drift can be suppressed with bumpiness control.

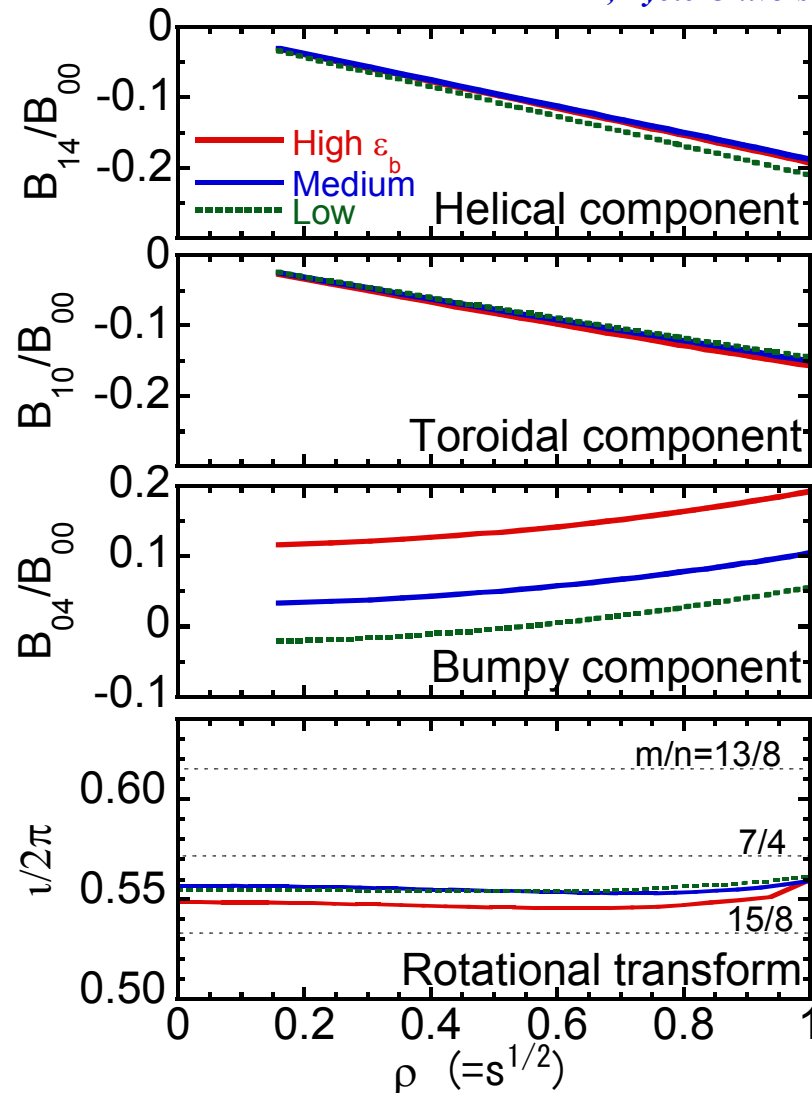
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Magnetic field strength at plasma axis



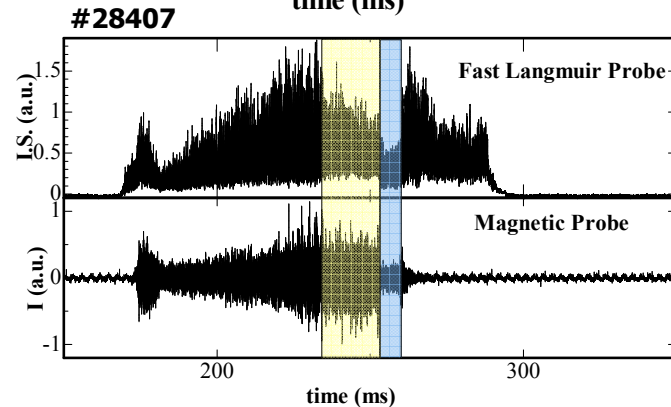
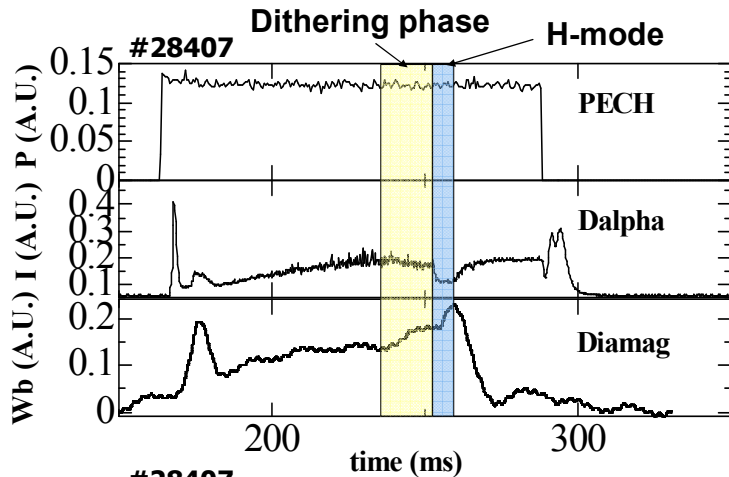
- High ϵ_b : Higher mirror ripple
 Low ϵ_b : Mirror field reversal (at axis)
 Constant parameters :
 - $R_{ax}/\langle a_p \rangle = 1.2 \text{ m}/0.17 \text{ m}$, $V_p \sim 0.7 \text{ m}^3$
 - Edge rotational transform
 - Magnetic well in entire region



The Heliotron J experiments have revealed the existence of the spontaneous transition to improved confinement mode.

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ECH-only plasma

- Observation with fast cameras revealed the existence of filament structure parallel to B in the peripheral plasma turbulence.
- Phase image analysis suggests
 - ✓ The CCW poloidal rotation of the filament structure during the L-mode.
 - ✓ During the L-H transition, this rotation speed decreases.
 - ✓ After the transition, it restarts the movement but in CW direction.

Ref. Nishino, C09 (Tue)

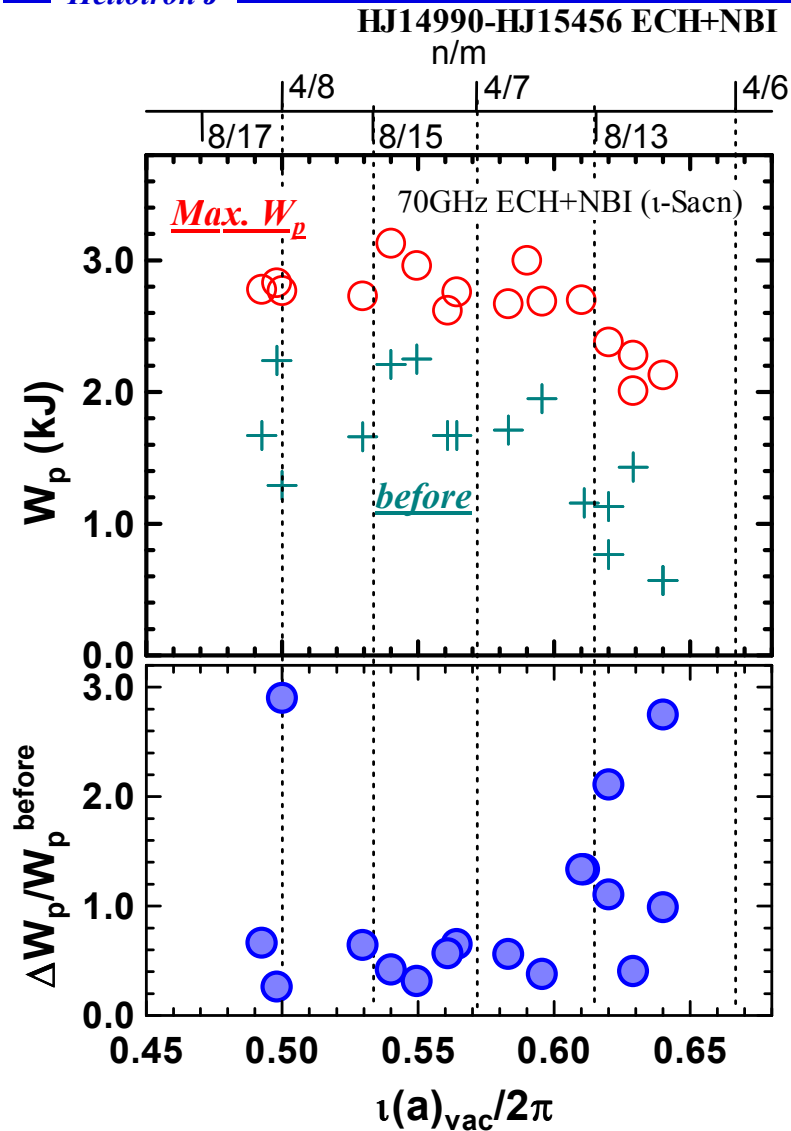


L-mode

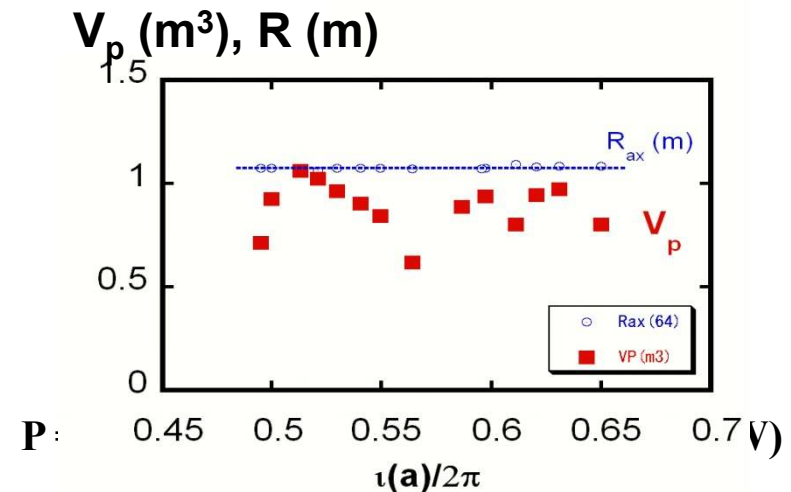
L-H transition

H-mode

Before the transition, W_p strongly depends on $\iota(a)/2\pi|_{vac}$, but it is moderated after the transition.

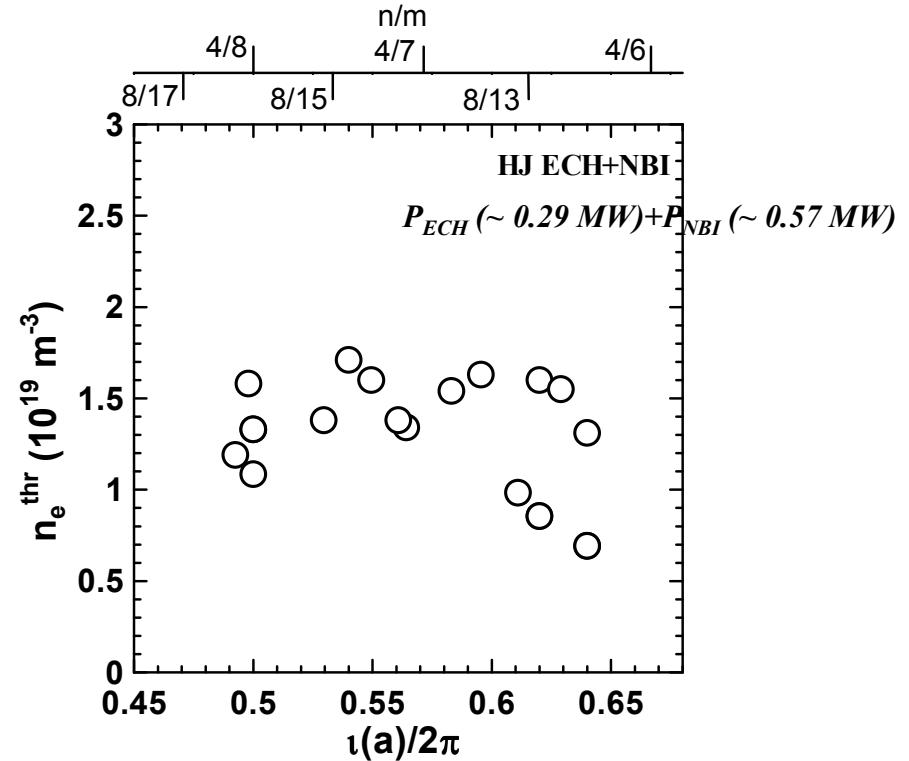
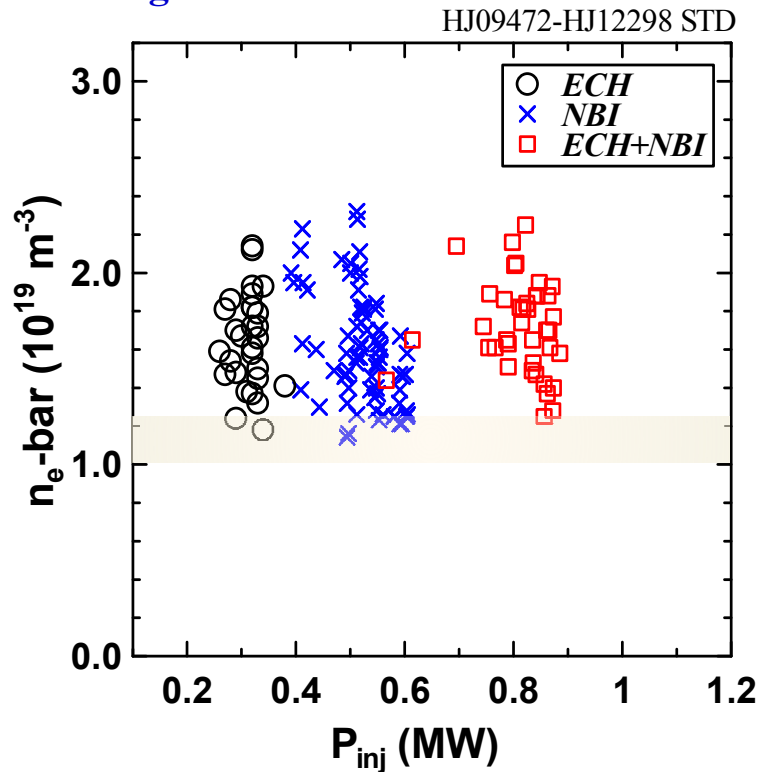


- The transition phenomena were observed almost all $\iota(a)/2\pi$ cases for the input power range in the experiment (ECH+NBI).
- Some configurations show a remarkable improvement but the others show a slight improvement.



One robust condition to the transition is the core density.
The value of the threshold line-averaged density is **not sensitive to the heating method, P_{inj} level, $\iota(a)/2\pi|_{vac}$.**

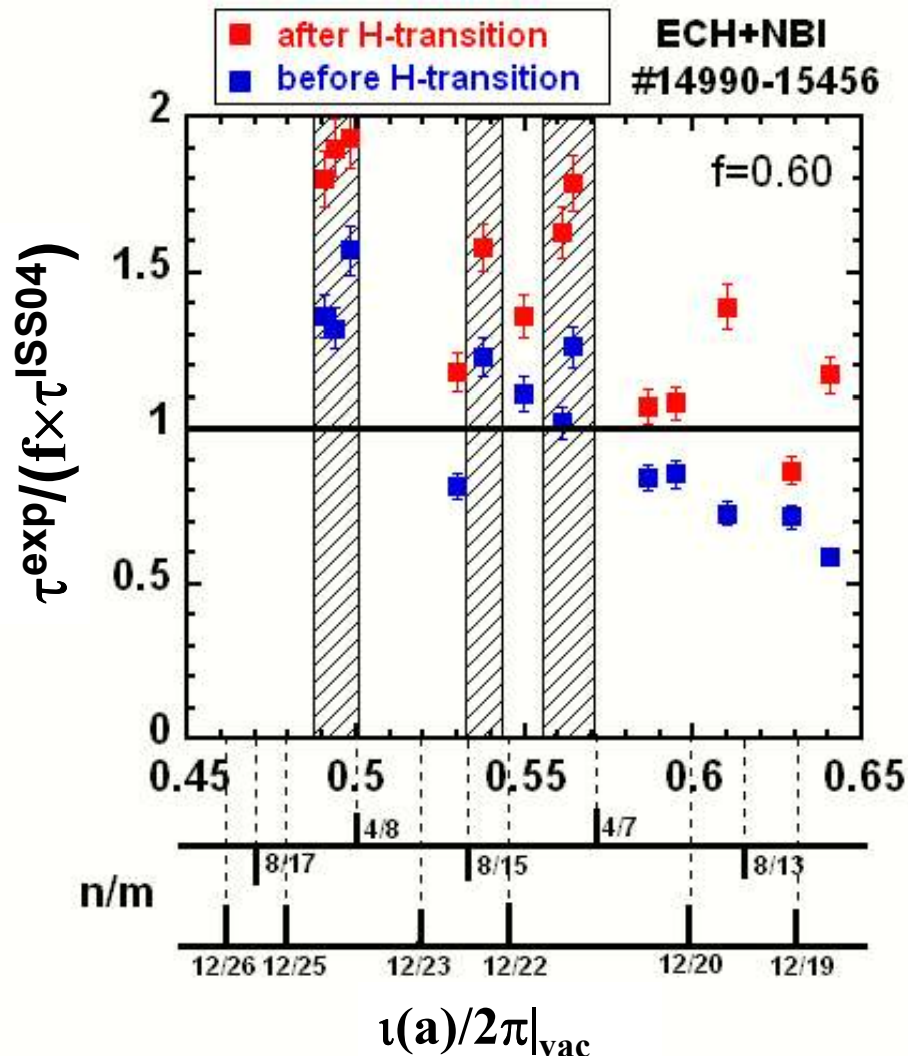
STD configuration



- The threshold density is a robust condition of the transition.
- The threshold density (lower density limit for the transition) is $\sim 1\text{-}2 \times 10^{19} \text{ m}^{-3}$.
- Only for higher values of $\iota/2\pi$, some exceptional low threshold density cases are observed. \Leftrightarrow MHD activities?

the “quality of improvement” based on ISS04 seems to depend on the configuration.

F. Sano, et al., Nucl. Fusion 45 (2005) 1557.



- Iota-windows for high quality H-mode close to the low-mode rationals
- All configurations basically have “stellarator shear” (vacuum) \Rightarrow hard to get edge flux surfaces on the “right-side” of the rationals.
- Difference in the poloidal viscous damping among the configurations can cause the difference in the “quality”?
 - Still open question.
 - Bias experiment will give us some information.

Electrode Bias Exp. in Heliotron J under a Low Field Condition

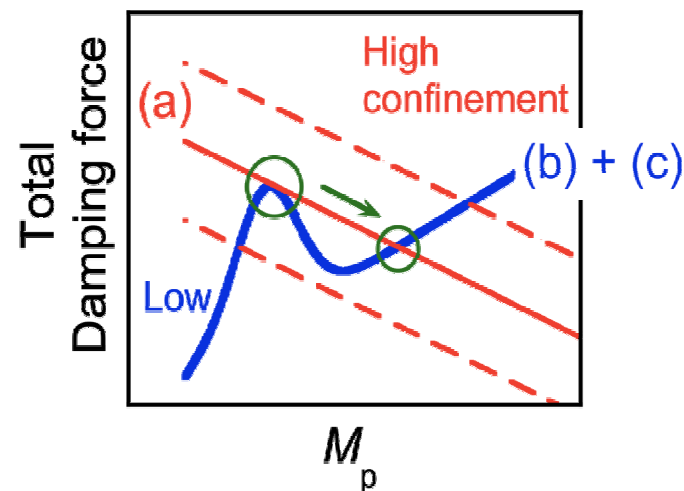
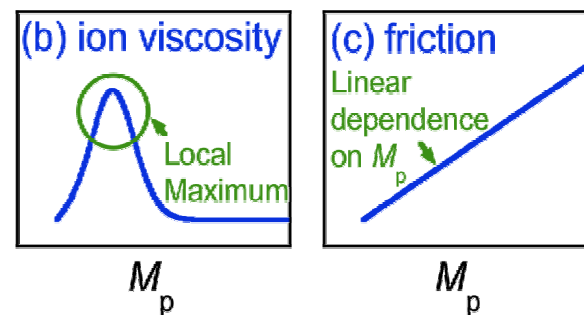
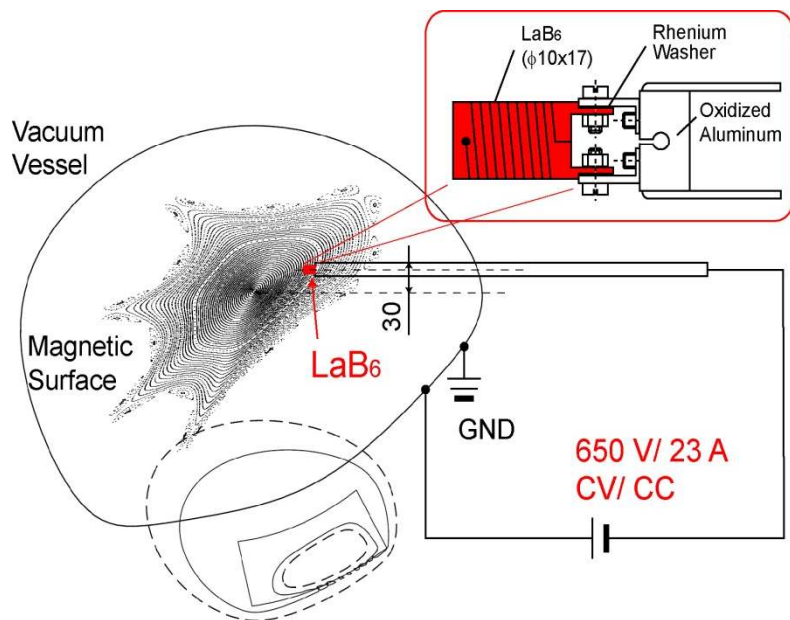
S. Kitajima, et al., Tohoku Univ.

Externally controlled torque could give us **a lot of (but NOT sufficient) information** on the viscous damping.

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Electrode of LaB₆ (Hot Cathode)



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Note: This experiment was performed under a special configuration for 2.45GHz ECH discharge.

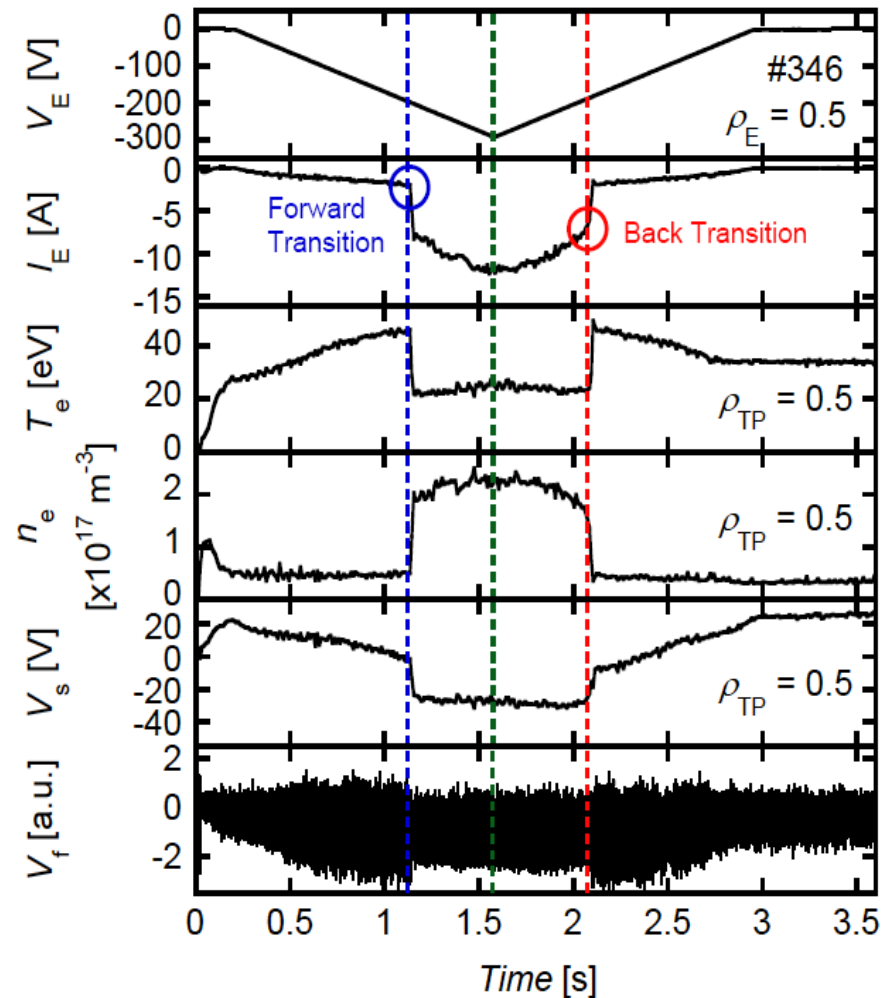
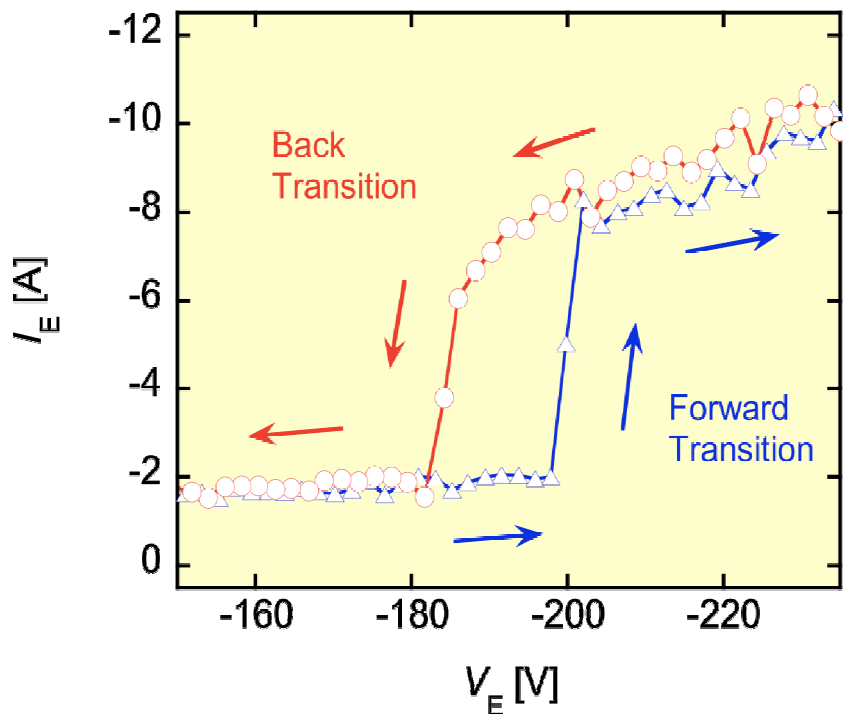
Time Evolution of a Biased Discharge

Triangular voltage was applied in order to give rise to forward/ back transition.

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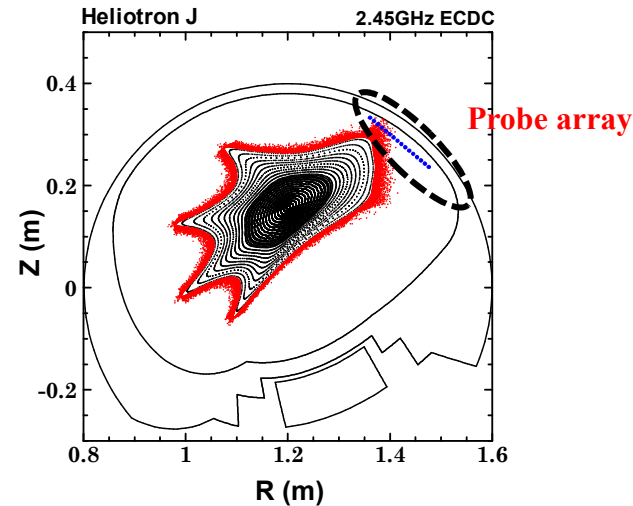
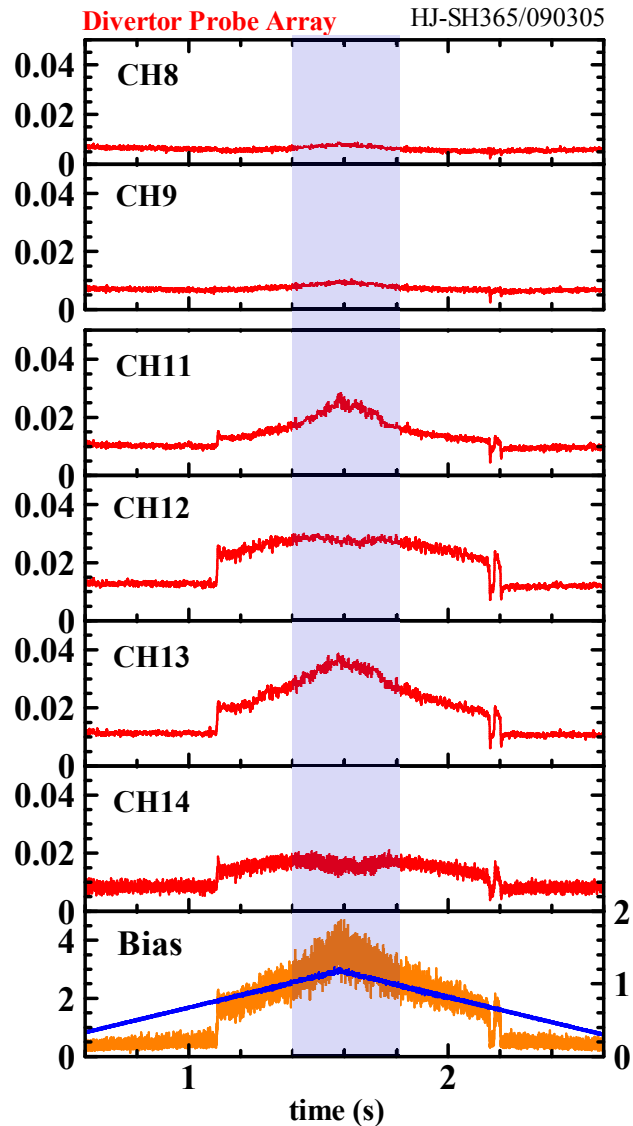
- **Rapid increase of I_E and n_e .**
- **V_s changed from positive to negative.**
- **Drop in the fluctuation level.**
- **Clear hysteresis of I_E .**



Bias Exp. suggests plasma rotation can modify the edge field topology.

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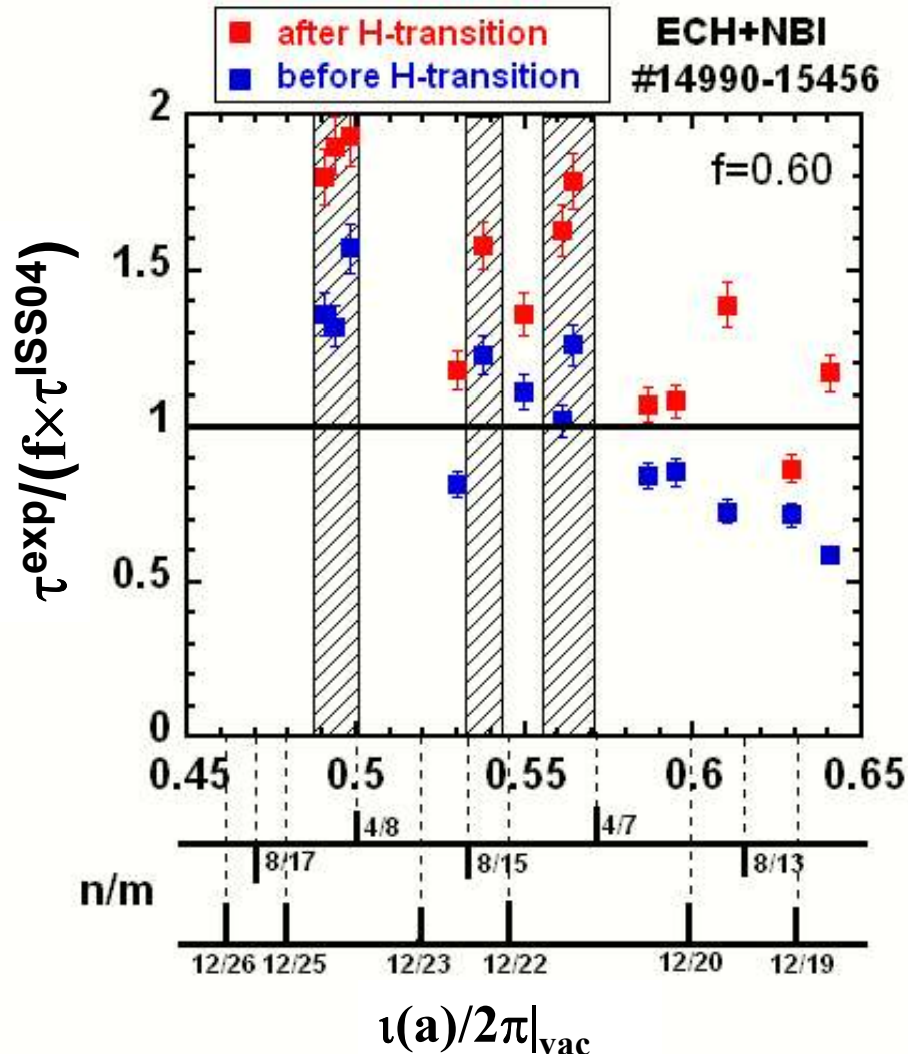
**Time traces of ion-saturation currents at divertor probe array during biasing:
Increase for some channels,
but decrease for other channels.
→ Modification of edge field topology!**

The “quality of improvement” based on ISS04 seems to depend on the configuration.

F. Sano, et al., Nucl. Fusion 45 (2005) 1557.

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- All configurations basically have “stellarator shear” (vacuum)
 \Rightarrow hard to get edge flux surfaces on the “right-side” of the rationals.
- Difference in the poloidal viscous damping among the configurations?
- **But, is that all?**
 - *How about the “plasma effects”, especially effects of (non-induction) plasma current.*
 - *Experiments suggest the change of the iota and/or the size of the confinement region.*

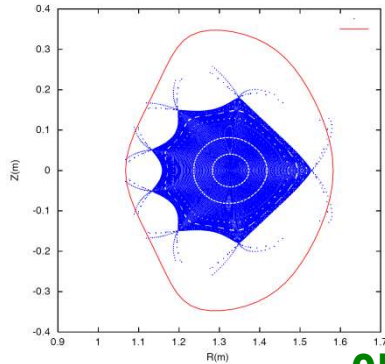
HINT2 predicts the change of the field topology caused by plasma current:

$$\beta(s) = 0.5 \cdot (1-s)^2 [\%], \quad j_p = j_{p0} \cdot (1-s)^2$$

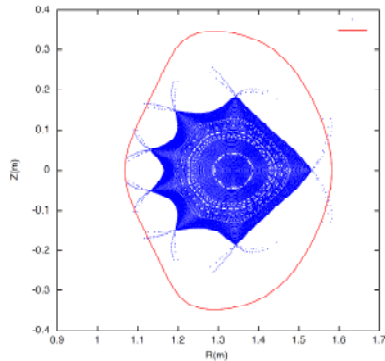
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- “Additive” plasma current can modify not only the $1/2\pi$ -profile but also the “**shape**” of magnetic surfaces.
- Asymmetric effect of plasma current due to the “proximity” to a low-mode rational number.

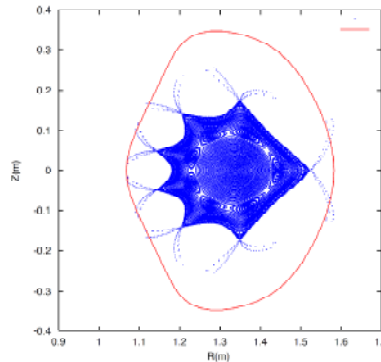
vacuum



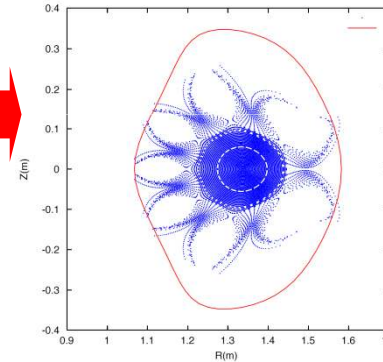
0kA



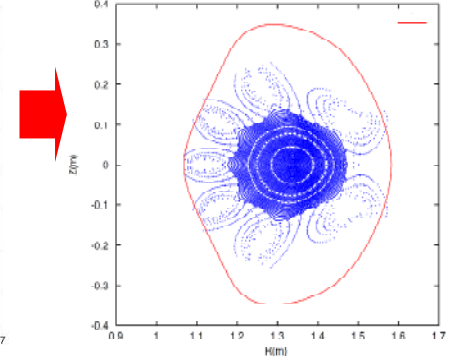
ad. 1kA



ad. 2kA

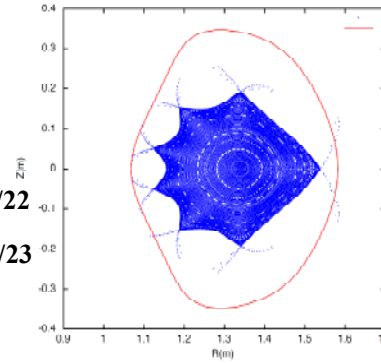
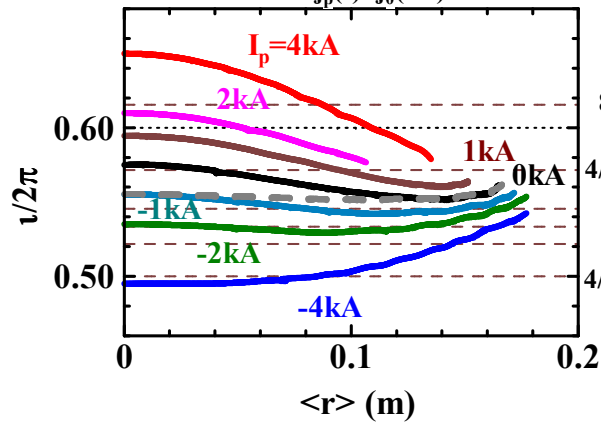


ad. 4kA

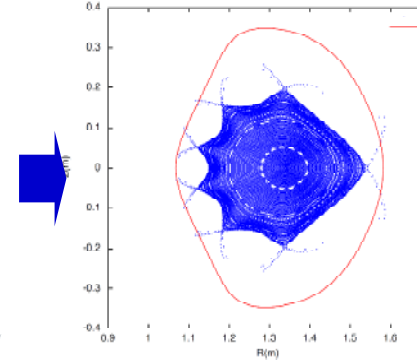


$$\beta(s) = b_0(1-s)^2, \quad \beta_0 = 0.5 \% \\ j_p(s) = j_0(1-s)^2$$

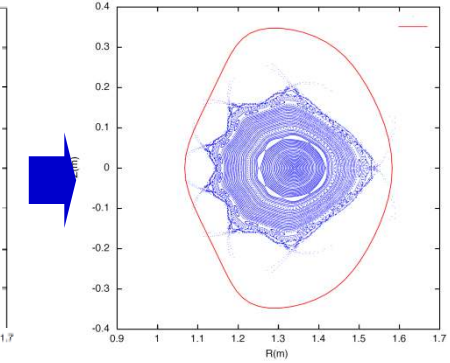
HINT2



sub. 1kA



sub. 2kA



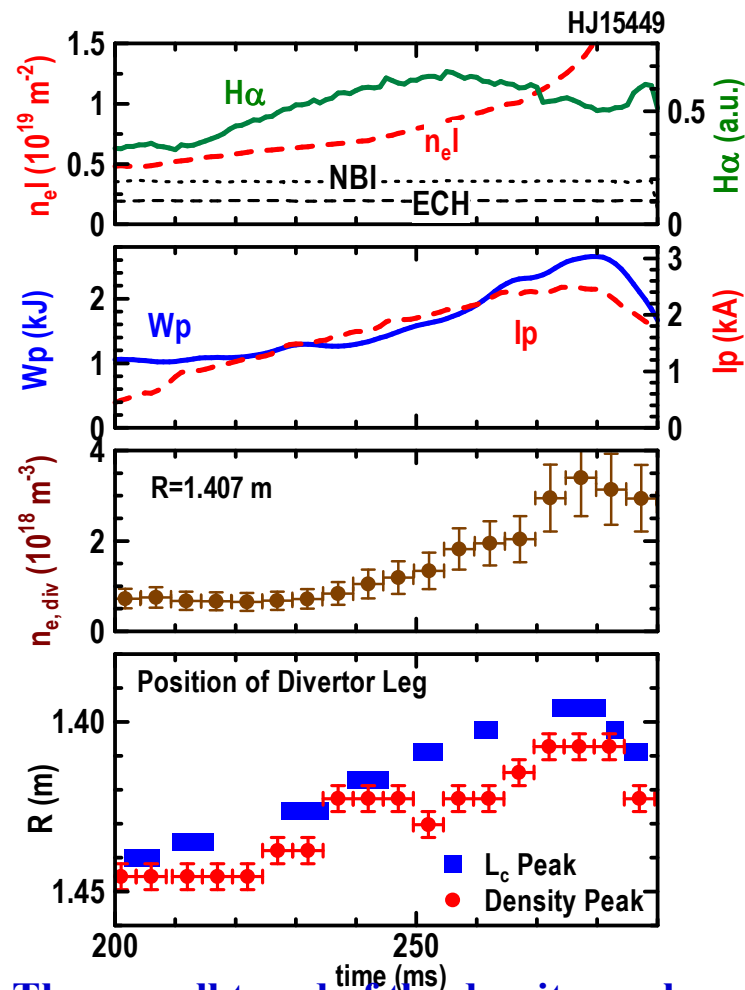
sub. 4kA

Experiments suggest the change of the edge field topology during a plasma shot.

ECH (≈ 0.3 MW) + NBI (≈ 0.7 MW, co-injection) @ STD

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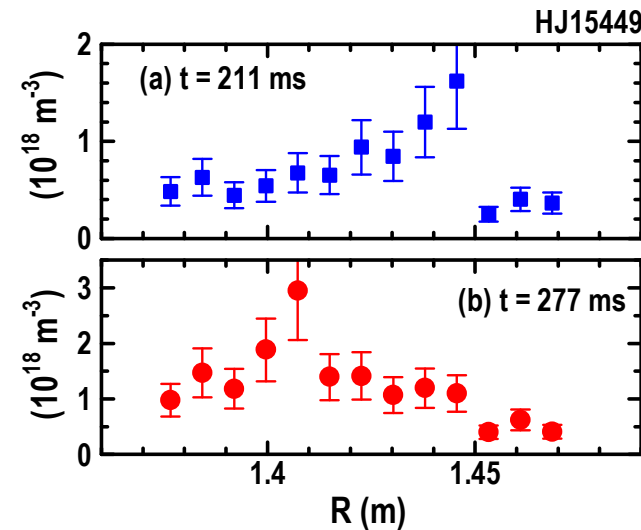
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- The overall trend of the density-peak-position shift seems to be well synchronized with the change of W_p or I_p .

- The observed non-inductive toroidal plasma current is gradually increases as increase of the stored energy.

– In the density range shown in the figure, the observed plasma current is considered to mainly consist of the bootstrap current and the NB induced current.



- The profiles of the diverted plasma density at $t = 211$ and 277 ms.

– At 211 ms, the density distribution is consistent with that expected from the vacuum field topology.

– At 277 ms, the density peak position shifts inward about 4 cm compared to the position at 211 ms.

Comparison between Co/CTR

NBI-only plasmas

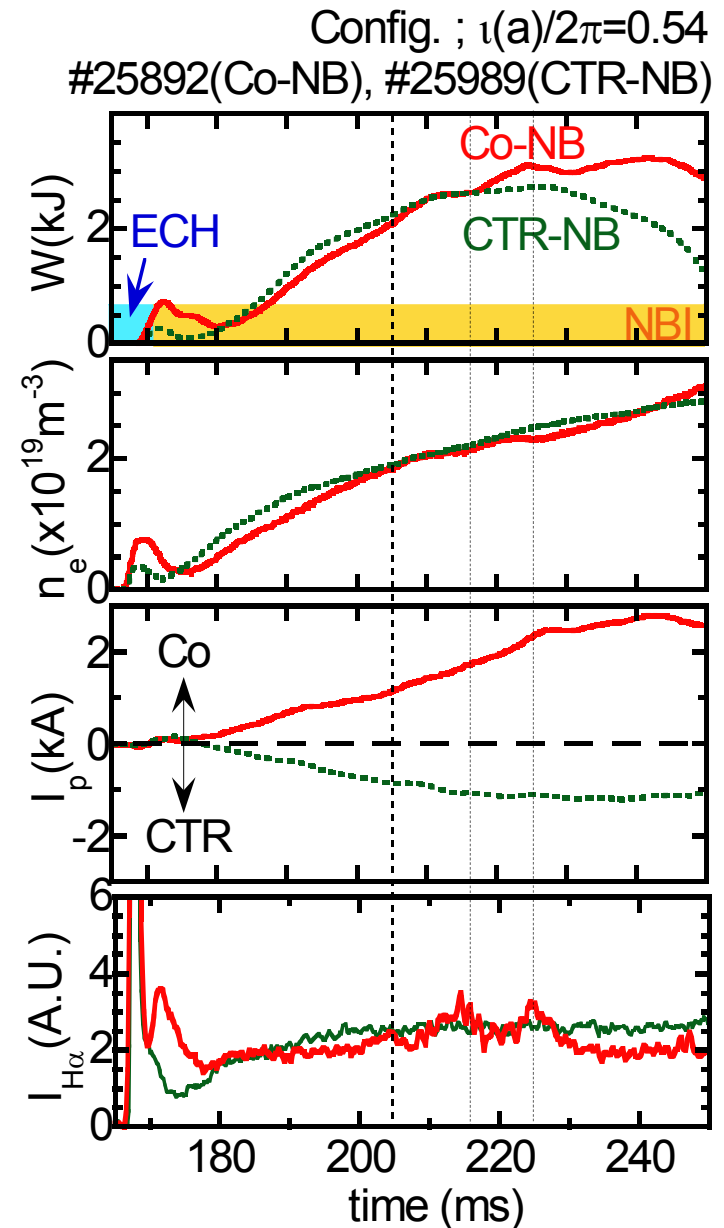
- No transition in CTR NBI -

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◆ In Co-NBI, the transition were observed, but no transition in CTR-NBI.

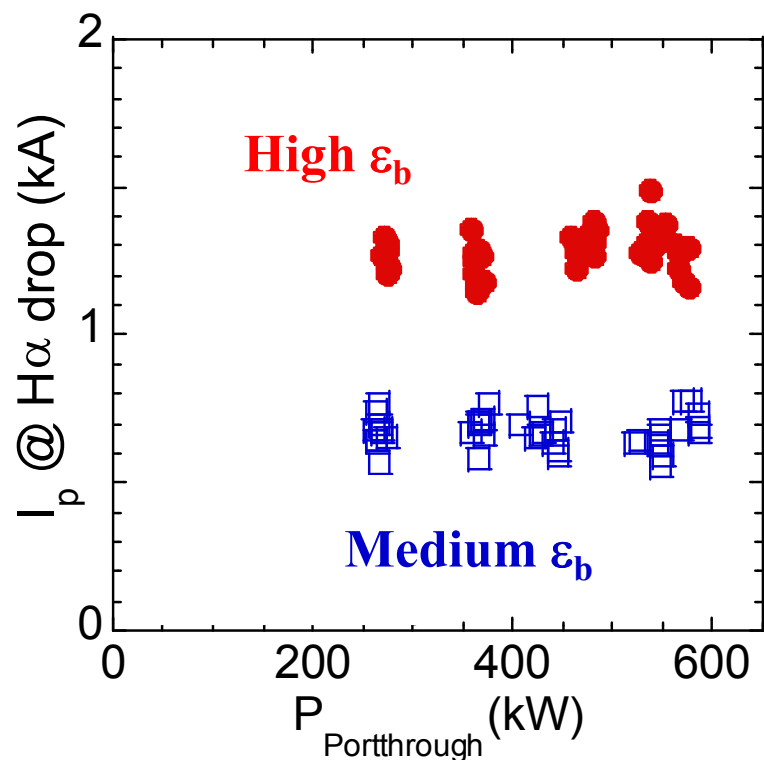
☞ Difference in the transition condition between these NBI-only plasmas due to

- Deformation of $\iota/2\pi$
 - Position of rational surface,
 - Shear,
 - Shape of LCFS,
- Direction of momentum input ?



The transition is observed when the plasma current reaches a critical value.

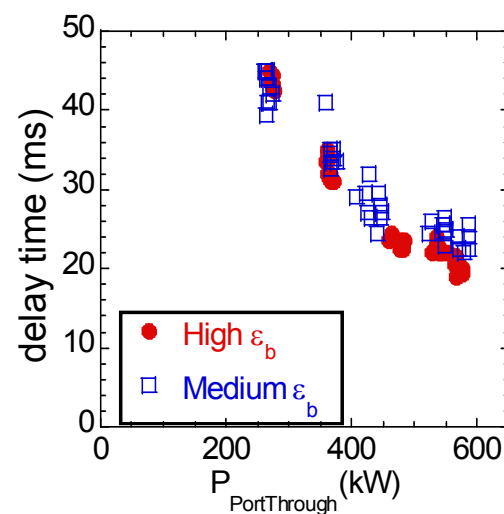
Dependence of I_p at transition on P_{NB}



– In high ϵ_b case, higher I_p was required.

- The critical value of the plasma current depends on the configuration;
 $0.7 \pm 0.1 \text{ kA}$ in middle ϵ_b
 $1.3 \pm 0.2 \text{ kA}$ in high ϵ_b cases.
- The observed time delay would be related to the growing-up time of the current.

Dependence of delay time on P_{NB}

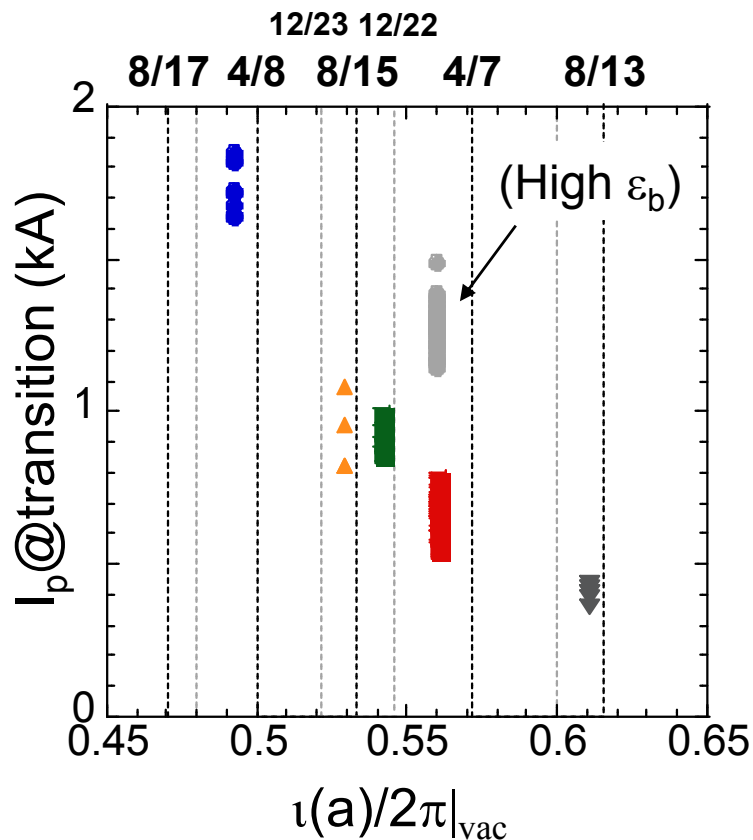


The effects of the plasma current on the field configuration should depend on the vacuum rotational transform. $\rightarrow i(a)/2\pi|_{vac}$ -scan experiment.

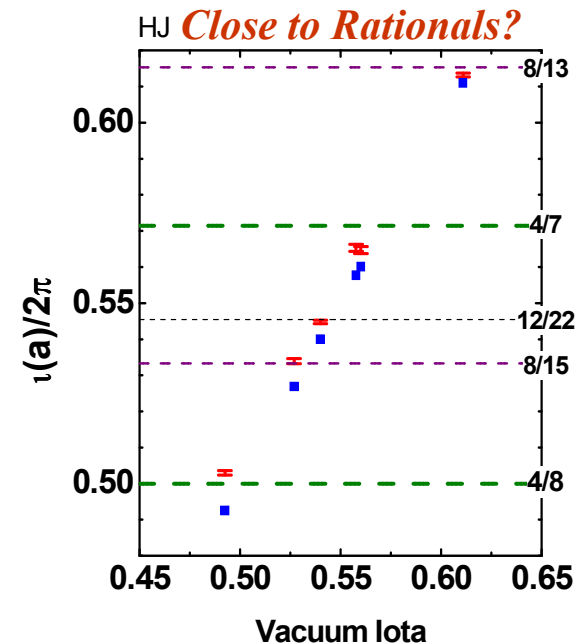
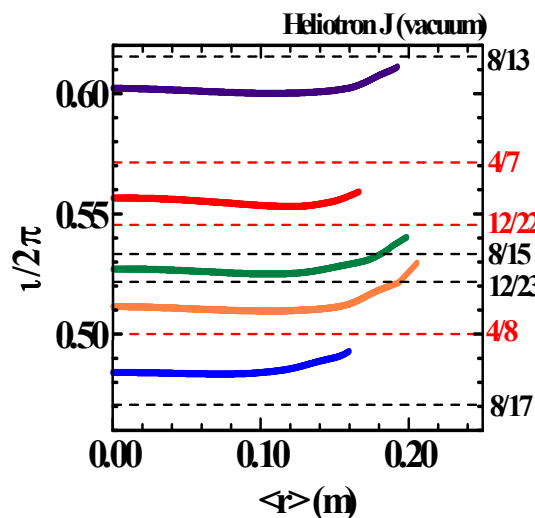
The critical current for *NBI-only plasma* strongly depends on $\iota/2\pi|_{\text{vac}}$.

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Radial profile of iota



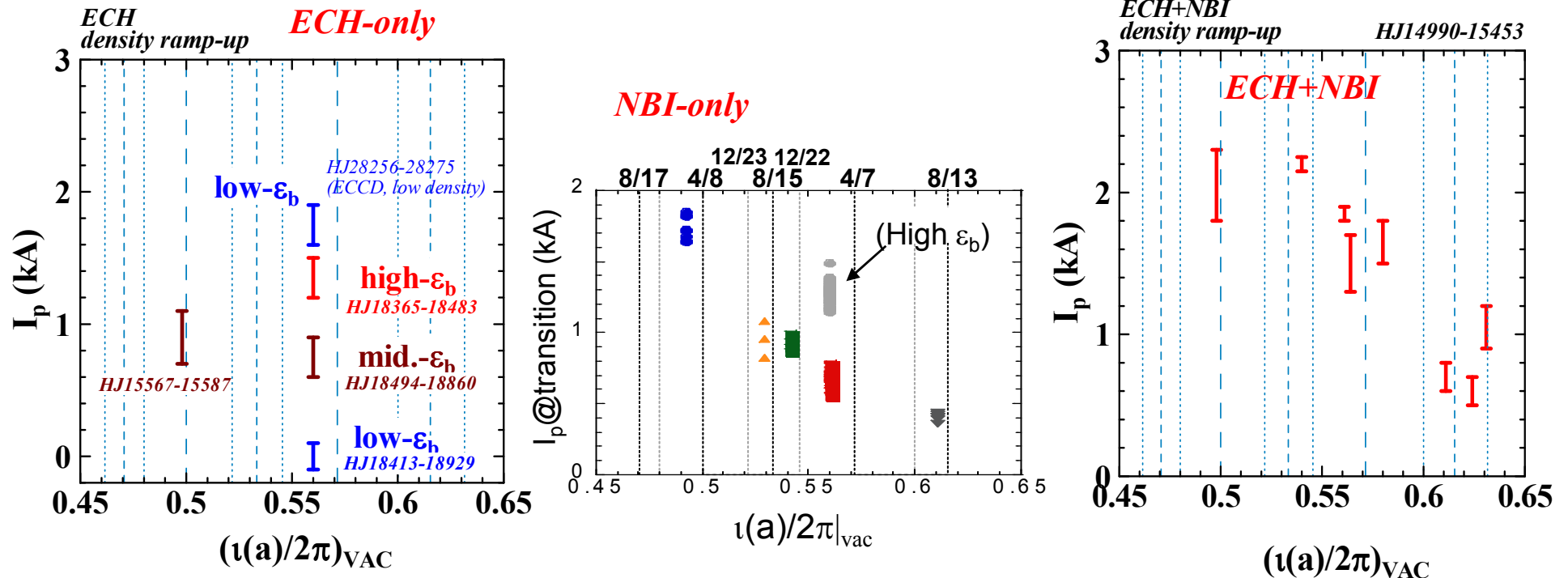
$0.25 \text{ MW} < P_{\text{inj}} < 0.6 \text{ MW},$
 $n_e \sim 1.5\text{-}2 \times 10^{19} \text{ m}^{-3}$

- ☞ So far, the transition was observed only in Co-NBI (ad. NBCD) plasmas ($P_{\text{inj}} < 0.6 \text{ MW}$).
- ☞ A “critical current” for the transition exists for all config.
- ☞ Does the change in the rotational transform caused by non-inductive plasma current trigger the transition? → Check other heating cases.

Comparison of the plasma current at the transition for *ECH-*, *NBI-*, and *ECH+NBI-*plasmas

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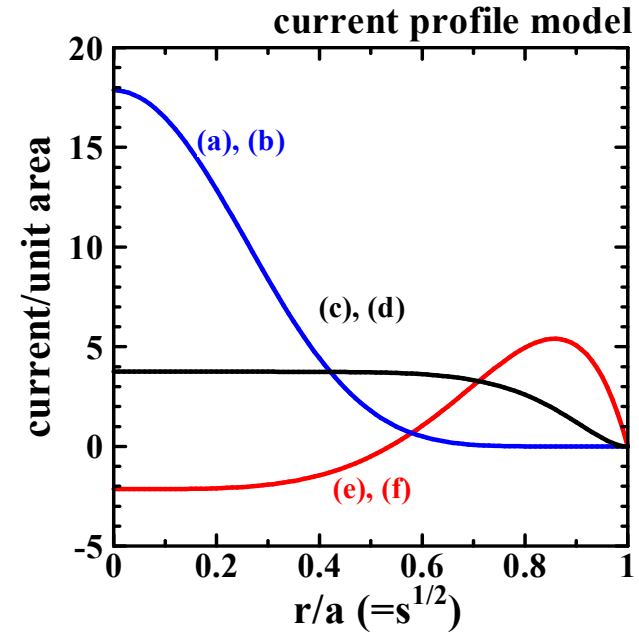
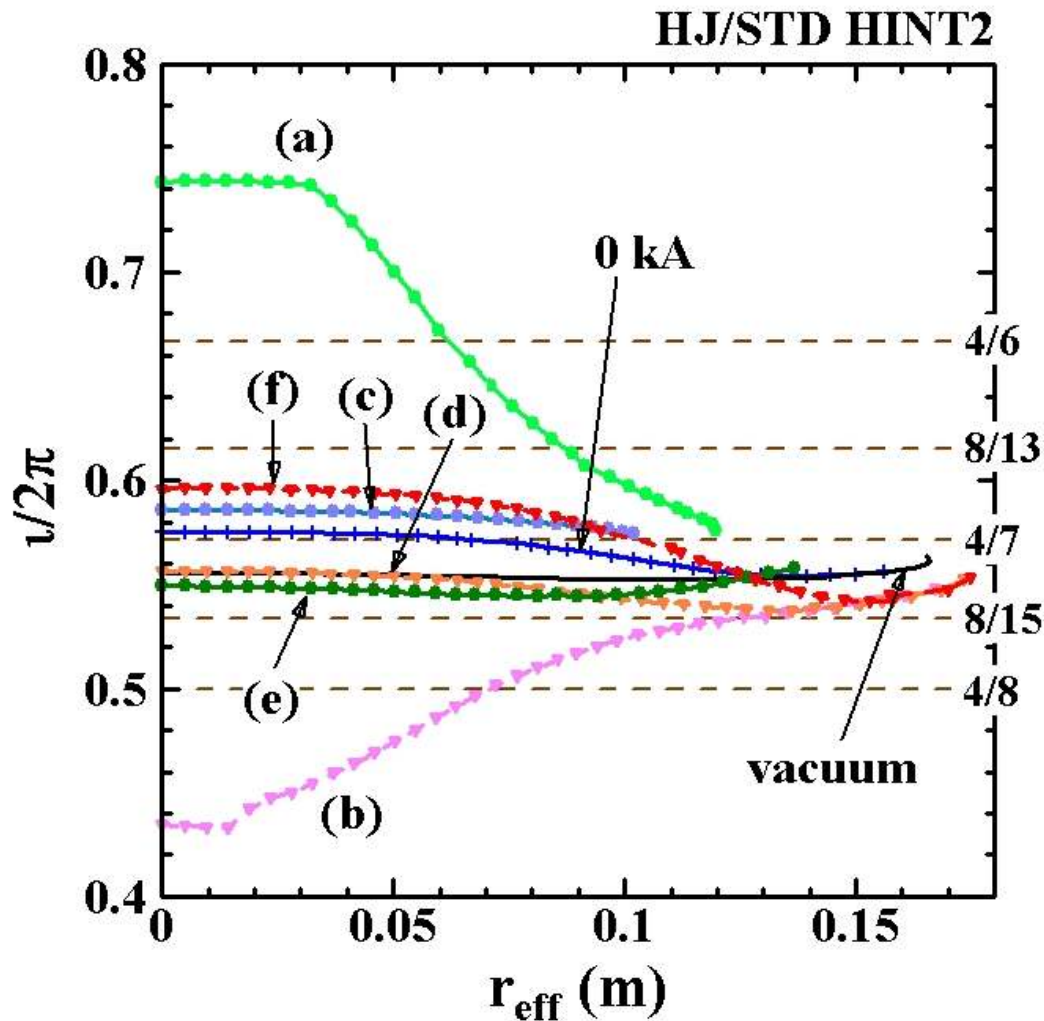
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- $I_{p,NET}$ itself seems not a key factor?
 → Need profile data!

Even for the same net current, the deformation of $\iota/2\pi$ and the shape of the magnetic surface strongly depend the current profile.

- MODEL CALCULATION -



$$\langle \beta \rangle \sim 0.32 \%, p = p_0(1-s^2)^2$$

$$(a) \quad j = j_0(1-s)^8, I_{net} = +2.0 \text{ kA}$$

$$(b) \quad j = j_0(1-s)^8, I_{net} = -2.0 \text{ kA}$$

$$(c) \quad j = j_0(1-s^4)^2, I_{net} = +2.0 \text{ kA}$$

$$(d) \quad j = j_0(1-s^4)^2, I_{net} = -2.0 \text{ kA}$$

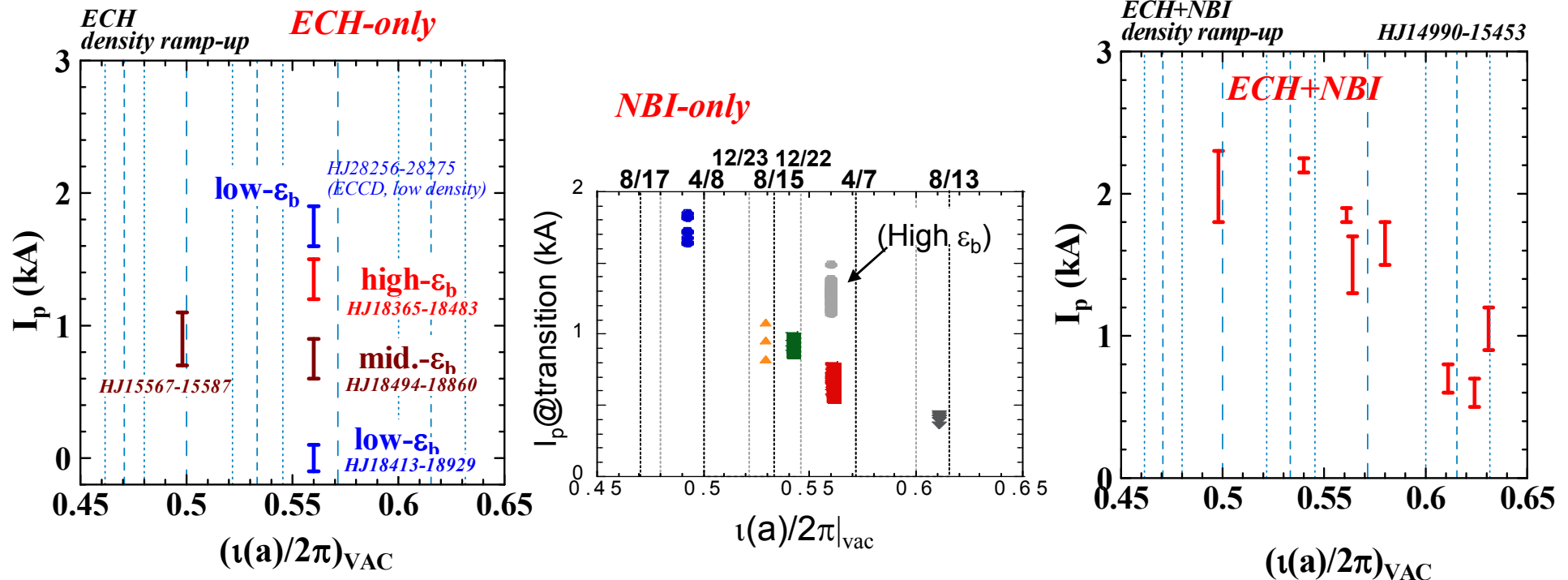
$$(e) \quad j = j_0(60(1-s^2)s^2 - 5(1-s^2)), I_{net} = +2.0 \text{ kA}$$

$$(f) \quad j = j_0(60(1-s^2)s^2 - 5(1-s^2)), I_{net} = -2.0 \text{ kA}$$

Comparison of the plasma current at the transition for *ECH-*, *NBI-*, and *ECH+NBI-plasmas*

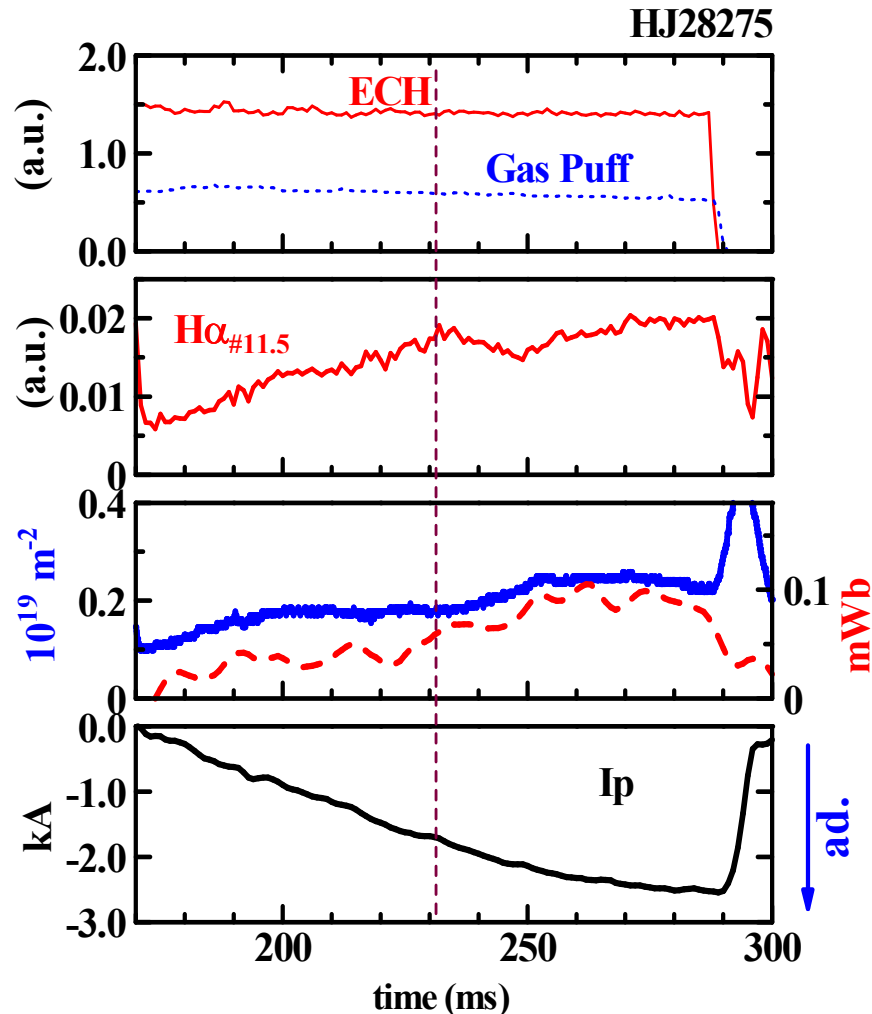
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- $I_{p,NET}$ itself seems not a key factor?
 → Need profile data!
 - Even in low density ECH-only plasma lower than the threshold density, the transition can be observed when the EC driven “additive” current is high enough.
 - Transition is observed in ECH+NBI plasmas even in the low ϵ_b case.

Transition to a better confinement mode was observed in a high ECCD current discharges.



- Low- ε_b , $1/2\pi(a) \sim 0.56$, ECH-only with $P_{\text{ECH}} > 0.3 \text{ MW}$
- The line-averaged density is **much lower than the critical low-density limit** observed in the previous experiments.
- *No enough experiments, from this point of view, in sub. high ECCD condition.*

Summary (1)

– *Transition to improved confinement* –

- No transition in CTR-NBI-only plasma
Existence of the “critical current” in Co-NBI-only plasma
→ **Change in the rotational transform** (and/or its radial profile) caused by plasma current and/or the direction of momentum input can **trigger the transition in NBI plasma.**
- No clear current dependence in ECH-only plasma
→ **Difference in the transition condition between NBI-only and ECH-only plasmas.**
→ ECH should have some effect on the transition condition in NBI plasma.
- **In low density** (lower than the threshold density) **ECH-only plasma, the transition can be observed** when the EC driven “additive” current is high enough.
→ Change in the rotational transform (and/or its radial profile) is effective even for ECH-only plasma.

Summary (2)

– Transition to improved confinement –

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for $n_e > n_{e,th}$

	ECH-only	NBI-only (Co)	NBI-only (CTR)	NBI (Balance) ECH+NBI (Balance)	ECH+NBI (Co)	ECH+NBI(CTR)
$i(a)/2\pi _{vac}$	X only for 0.493	OK	X	No Data	OK	No Data
ϵ_b^*	Δ for high- ϵ_b	X for low- ϵ_b	X	No Data	OK	No Data
$I_{p,th}^{**}$	No	Yes	–		Yes	
remarks	Transition is observed in low n_e , high I_p					

* Data for different $i(a)/2\pi|_{vac}$ are necessary.

** Data @ $I_p=0$ are necessary.

Thank you for your attention!

&

Thank you All Contributors for Nice Collaborations!

■ *Heliotron J*

■ *IAE, Kyoto University* ■ ■ ■

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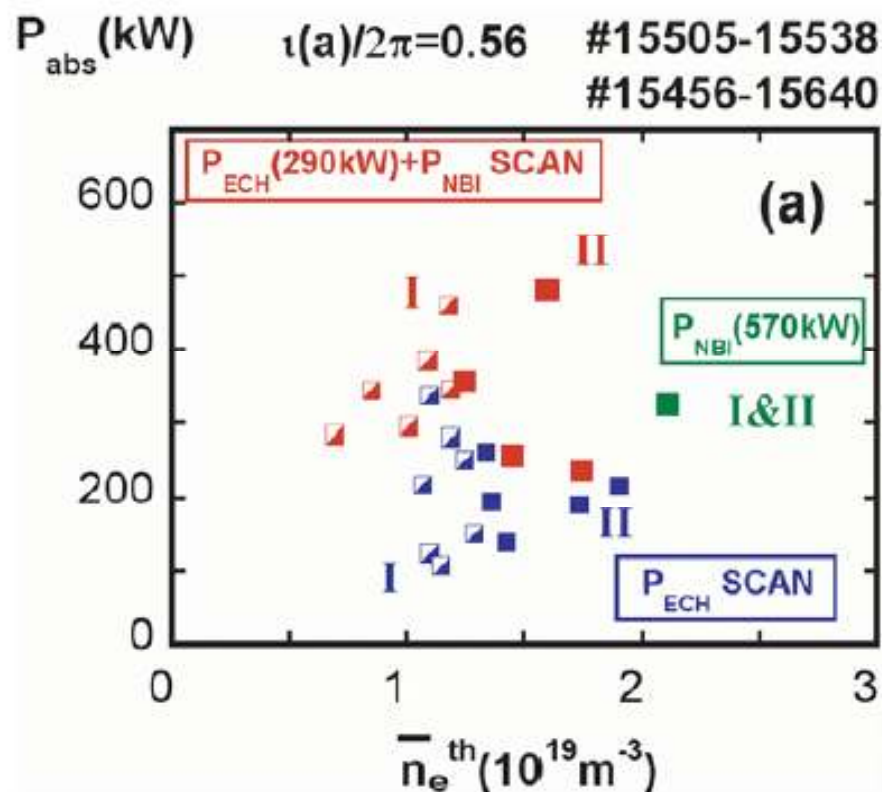
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Southwestern Institute of Physics, China

No clear threshold power for Phase-I



- No clear/simple relation between n_{th} and P_{abs} .
- At the $\iota(a)/2\pi \approx 0.493$, ECH-only discharges cannot make a transition. → configuration effects on the power threshold.
- Different power dependence for Phase-I and II?