High Beta Issues in a Helical System

- with emphasis on achievement and prospec of LHD experiments

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High beta is a critical mission as well as a challenging physical issue

Fusion gain

 P_{\cdot}

$$\Gamma_{usion} / P_{loss} \propto nT\tau_E \propto \beta B^2 \tau_E \propto \beta^2 B^4 V$$

High beta → Compact & Efficient reactor
← Restriction by heat and neutron loads

"High beta" is a vital issue in helical system

Reason

Critical issue in a helical reactor: to secure sufficient space for blanket and shield > 1 m

Solution 1 : Large machine Solution 2 : Thin magnet → Low B → High beta





High beta is a critical mission as well as a challenging physical issue

MHD instabilities to limit
beta in tokamak are
inevitably 3-D
⇔ share commonalty
in methodology

Resistive wall mode (RWM)







Neoclassical Tearing Mode (NTM)

However, characteristics and effect of MHD instability are different between tokamak and helical system

← Bias (help) by vacuum confining magnetic field is essential

← Commonalty in physics is very challenging

Development of β in LHD experiment is very encouraging

5 % of β has been achieved without any disastrous instabilities while MHD activity is certainly seen as the theory predicts.



High β ranging 5 % is maintained in steady state for longer than 100 \times τ_{E}

Achievable β is limited by the available heating power without a hard β -limit.







- 1. Introduction
- 2. Gaps in non-dimensional parameters
- Characteristics of growth of MHD instabilities
 ✓ Usually benign but violent under specific condition
- 4. Enhancement of stochastic field and its effect
 ✓ 3-D effect
- 5. Effect of resonant magnetic perturbation
 ✓ Ideal mode violated by reduction of magnetic shear
 ✓ Resistive interchange mode
- 6. Transport associated with high- β equilibrium and resistive interchange mode
- 7. Summary

Establishment of integrated physics model is required to bridge over the gap

Normalized gyro radius

 $ho^{st} \propto T^{1/2}$ / aB

 $\frac{\text{Collisionality}}{v_h^* \propto n / T^2}$



Gyro-Bohm nature in confinement is preferable $\chi \propto \rho^*$

Collisionless (Long mean free path) high β is an important issue

ex. effect of stochasiticity

Establishment of integrated physics model is required to bridge over the gap

Magnetic Reynolds Number : S



$$S = \tau_{R} / \tau_{A} \propto \frac{aBT_{e}^{3/2}}{ZA^{1/2}n_{i}^{1/2}} \propto \frac{\beta^{1/2}}{\nu_{b} * \rho *^{2}}$$

Role of resistivity→ Resistive interchange mode

Growth rate of resistive mode (B.Coppi, NF (1966))

$$\gamma \propto S^{-1/2}$$

Pay careful attention to extrapolation γ is reduced by 1/6 ?

Note: LHD cannot realize a plasma with full dimensional similarity

Integrated physics model with reliable predictability which bridges over the gap of non-dimensional parameters



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LHD has large flexibility of magnetic configuration

Poloidal field coil

- ➔ Position of magnetic axis
- → Elongation



NIFS-PE1479

Control of current center of helical coil

HC-C

HC-I

- ➔ Aspect ratio
 - ➔ Rotational transform

Plasma

Local Island Divertor coil

- ➔ Generates m/n=1/1 Resonant Magnetic Perturbation (RMP)
 - ➔ Magnetic island



Mercier criterion can be violated and quantification of low-m mode stability is an issue

- 12 ρ=0.5 (+~1/2) <mark>⊖ dβ_{kin}/d</mark>ρ Mercier 8 dβ/dp (%) currentless m/n=2/1 unstable γ/ω_^=1<u>0⁻²</u> 0 10² ρ**=0.5 (ε**=0.5) $b_{RMS}^{~~MS}/B_{ heta}$ (10⁻⁵) 10¹ 10⁰ 10⁻¹ n/m = 1/210⁻² disappear 10⁻³ 3 2 ()> (%) <β_{dia}
- R_{ax} =3.6m : inward shifted
- Mercier unstable region does not inhibit access to higher β and no destructing event is observed.
 - ➔ unlikely in tokamaks
- ✓ Low-*n* interchange mode unstable region looks prohibited
 - inhibits access to higher-β regime in the case with further inward shift of R_{ax} (<3.6m), where stability is violated seriously.
- Excitation and disappearance of MHD modes corresponds to prediction by Mercier criterion
- These MHD modes usually saturate at harmless level



Sawtooth-like activity is repeated in unstable regime for ideal interchange mode



- ✓ With pellet injection, pressure profile is peaked compared with gas-puff-fueled plasma.
- ✓ When the pressure gradient exceeds the low-*n* unstable condition, sawtooth-like MHD activities are destabilized, but pressure gradient continues to evolve. 11/28



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Core Density Collapse limits the central beta



✓ IDB/SDC plasmas, peaking of the plasma is limited by so-called CDC

✓ Half of particles are expelled from the high density core in sub ms.



3-D MHD equilibrium without assumption of nested flux surfaces

-0.4

2.5

MHD equilibrium



3.0

3.5

4.0

R (m)

 \rightarrow Significant T_{e} gradient exists in the edge stochastic area

Y.Suzuki, I.34 on Fri.

3%

4.5



Results from HINT well describes deformation of magnetic surfaces

Z (m)

E

3-D equilibrium consistent with experimental observation, i.e., Shafranov shift, pressure profile, etc.

Significant pressure ($T_{\rm e}$) gradient exists in the edge stochastic area

Hypothesis

- 1) Plasma heals flux surfaces
- 2) Profile is consistent with characteristics of stochastic field
- 3) Somewhere between 1) & 2)

L_{C-TB} : connection length between the torus-top and - bottom

- ✓ L_C >> L_{C-TB}
 - → Pfirsch-Schlüter current is effective
 - ➔ Secure MHD equilibrium

✓ L_c >> MFP (even under a reactor condition)
 → Plasma is collisional enough to secure isotropic pressure



4. Enhancement of stochastic field and its effect

Electron temperature profiles in stochastic region at high beta operation

3D edge transport simulation (EMC3-EIRENE)

- \checkmark T_e profiles are sustained even at higher input power.
- ✓ Rechester-Rosenbruth model predicts larger radial transport at high T_{e} . $\chi_{\perp eff} \propto D_{FL} T_e^{0.5 \sim 2.5}$
- ✓ The 3D modelling indicates high T_e plasma sees different D_{FI}.



The parallel path becomes predominant in high T_e





m/n=1/1 mode is excited in case of high β , enhanced magnetic hill and loss of magnetic shear





a_{eff} (m)

5. Effect of resonant magnetic perturbation



Large scale MHD event is triggered in case of high β , enhanced magnetic hill and loss of magnetic shear

beta

1.0

0.5

<β> (%)

- ✓ Sawtooth-like amplitude correlates with plasma current



Unstable mode can be stabilized by RMP



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<β_{dia}> (%)

 $\bar{n}_{e} (10^{19} \text{ m}^{-3})$

 $R_{\rm ax} = 3.6 \text{ m}, B_{\rm t} = -0.425 \text{ T}$

NB

Fluctuation localized at the edge is predominant in regular high- β plasma

5

4

3

2

0

- ✓ $\beta \sim 4.8$ %, $\beta_0 \sim 9.6$ %, H_{ISS95} ~ 1.1
- Plasma was maintained for ~ $100\tau_{\rm F}$
- Shafranov shift $\Delta/a_{\rm eff} \sim 0.25$
- Edge MHD modes are dominantly observed.



is secured by magnetic shear \rightarrow Ideal mode is stable \rightarrow resistive mode

F.Watanbe, P03-09 on Thu.

Reynolds number dependence agree with resistive interchange mode

Magnetic Reynolds number is related with growth of resistive mode ($\gamma \propto \beta^{2/3} S^{-1/3}$), the change of topology (reconnection)

CHS: S=10³~10⁵ LHD: S=10⁶~10⁸



Comparison of MHD activities in between LHD and CHS



Saturation of peripheral MHD mode depends on S parameter. 22/28



- ✓ Pressure gradient near t = 1 surface gradually decreases with I_{LID} .
- ✓ Amplitude of the mode decreases with reduction of the gradient, and the mode disappears while finite gradient still remains.
 - ➔ below threshold of excitation



✓ Stabilization is due to reduction of D_R or other mechanism?

 Mode frequency slows down with decrease of the amplitude.



6. Transport associated with high- β equilibrium and resistive interchange mode

Energy confinement degrades moderately with increase of β 3 LHD H-Factor

Confinement dependence on R_{ax} in low β



Reconsideration of the normalization factor f_{ren} corresponding to the position of magnetic axis shifted due to β

- ➔ Degradation is mitigated
- ➔ But looks still remained



S.Sakakibara, H.Funaba, P03-07 on Thu. 24/28

6. Transport associated with high- β equilibrium and resistive interchange mode

Transport and density fluctuation increases with β







Degradation becomes large with enhancement of magnetic hill (large A_p).



Resistive g-mode turbulence

$$\chi_{GMT} \propto \beta^{4/3} S^{-4/3}$$

Confinement is anticipated to be improved in higher temperature operation



Summary



- 1. Achievement of high b of a reactor relevant 5 % in LHD is encouraging, however, integrated physics model is required to bridge over the gap to a reactor. Predictability should range over 10 for ρ^* , <u>60 for v_b^* and 200 for S</u>.
- 2. Documentation of <u>saturation of the unstable mode</u> is required to assess how much benign the interchange mode is.
- 3. <u>Mechanism of Core Density Collapse</u> is an open question.
- 4. <u>Advanced theory/model</u> to describe MHD stability in 3-D field, in particular, <u>with magnetic islands and/or RMP</u>, and <u>stochastic field</u> is demanding but necessary.
- 5. Confinement improvement is expected in large S regime.
- 6. Confinement of α particles on the way to high β and their driving instabilities (*K.Toi, PL04, on Fri.*)
- Discussed high beta issues are not specific to LHD/heliotron. Significant contribution is expected from smaller devices together with theory and simulation.





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On-going and Nearest Future Plan - for collisionless high-β plasmas -

- 1. Optimization of dynamic B_v control (S.Sakakibara, P03-08 on Wed.)
- 2. Upgrade in 2010
 ✓ NBI 5th beam line 7 MW, 60 keV
 → 30 MW in total
 ✓ Partial modification for closed divertor
 → 2/10 inboard section, proto-type



- 3. Closed helical divertor (S.Masuzaki, I.20 on Wed.)
- 4. Deuterium
 - ✓ Identification and documentation of isotope effect
 - ✓ Upgrade of NBI (32 MW in total)
- 5. Reactor design study FFHR : Force-Free Helical Reactor (T.Goto, C.13 on Wed.)