

a. Thomson Scattering cross-section

nonrelativistic equation of motion: $\ddot{\vec{\beta}} = \frac{-e}{mc} \vec{E}_i$ for an electron

$$\text{Radiated field: } \vec{E}_s = \frac{-e}{4\pi\epsilon_0} \left[\frac{1}{Rc} (\hat{s} \times (\hat{s} \times \ddot{\vec{\beta}})) \right]_{\text{ret}}$$

$$\Rightarrow \vec{E}_s = \frac{e^2}{4\pi\epsilon_0 mc^2} \left[\frac{1}{R} \hat{s} \times (\hat{s} \times \ddot{\vec{E}}_i) \right]_{\text{ret}}$$

$$\text{classical electron radius: } r_e = \frac{e^2}{4\pi\epsilon_0 mc^2} = 2.8 \cdot 10^{-15} \text{ m}$$

$$\vec{E}_s = \left[\frac{r_e}{R} \hat{s} \times (\hat{s} \times \ddot{\vec{E}}_i) \right]_{\text{ret}}$$

$$\frac{dP}{d\Omega} = R^2 c \epsilon_0 |\vec{E}_s|^2 = r_e^2 \sin^2 \theta c \epsilon_0 |\vec{E}_i|^2$$

where θ is the angle between \vec{E}_i and the observation point

Define the differential scattering cross section as the ratio of $\frac{dP}{d\Omega}$ to the incident power per unit area, $c\epsilon_0 |\vec{E}_i|^2$

$$\Rightarrow \frac{d\sigma}{d\Omega} = r_e^2 \sin^2 \theta$$

$$\text{Integrated over solid angle, } \sigma = \frac{8\pi}{3} r_e^2$$

b. $l = 1 \text{ cm}$ $n_e = 10^{20} \text{ m}^{-3}$ $d\Omega = 0.01$, 90° scattering

$$\begin{aligned} \text{Total fraction of power scattered into detector} &= 0.01 \times n_e r_e^2 l \\ &= 7.8 \cdot 10^{-14} \end{aligned}$$