

A chaotic collection of thoughts on stochastic transport

what are the issues that M3D must consider to accurately determine heat transport

which analytical and numerical methods may complement the M3D algorithm

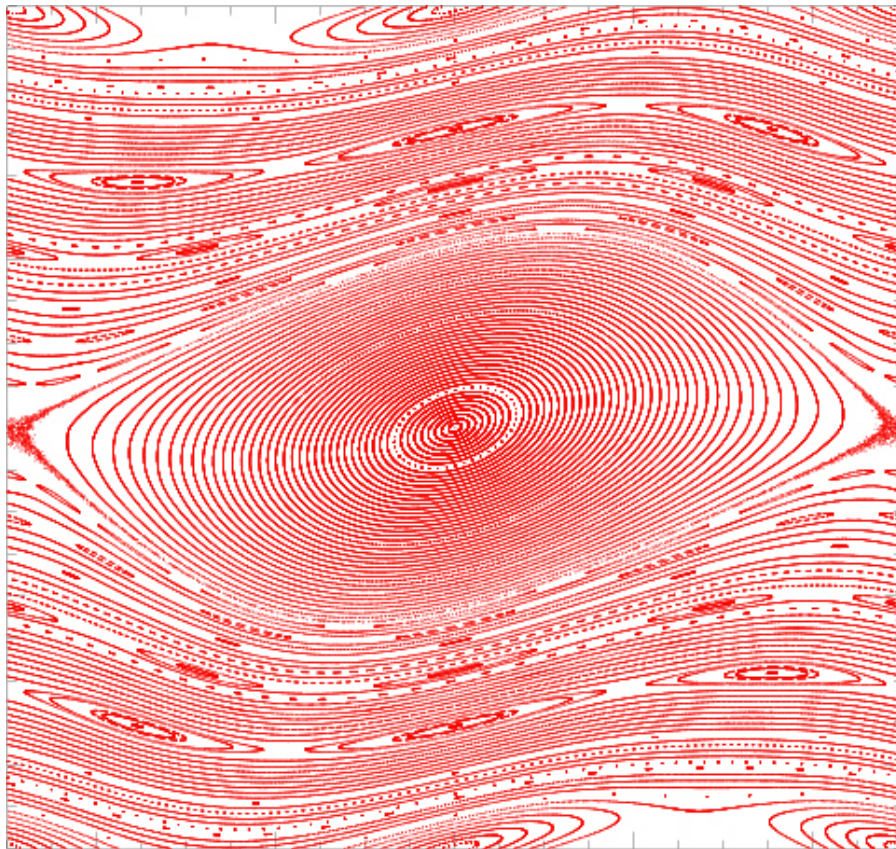
Stuart Hudson

thanks for interesting discussions with Josh Breslau, Ravi Samtaney, Roscoe White, Allen Boozer

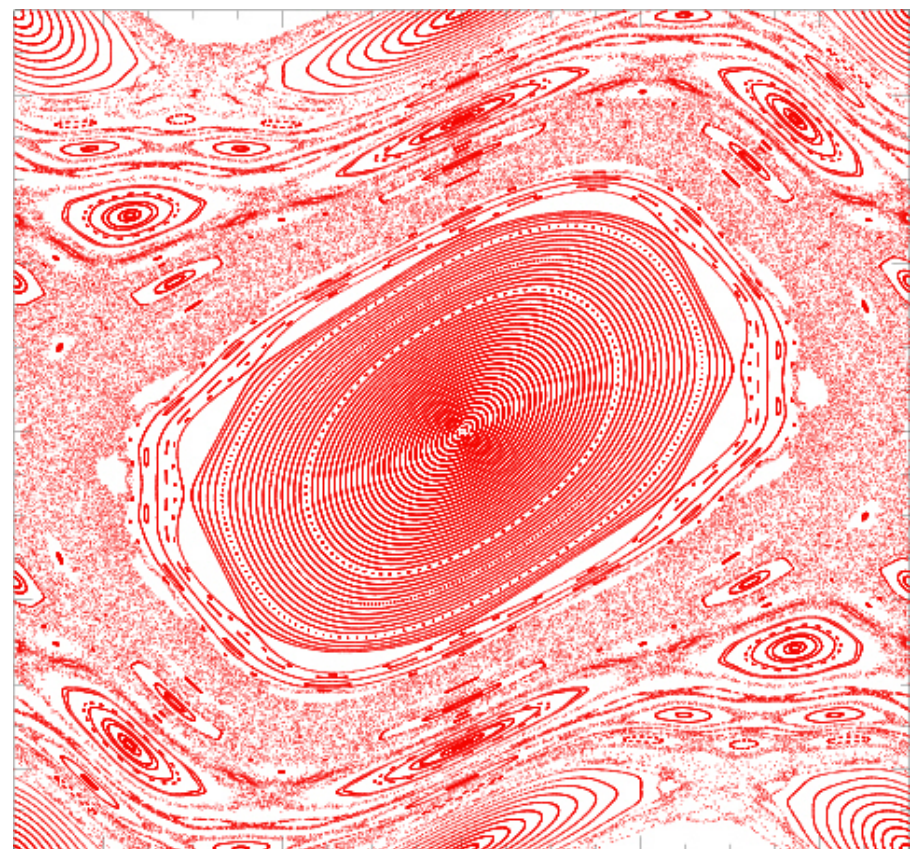
How is heat transported across islands and chaotic regions ?

$$\mathbf{q} = -\kappa_{\parallel} \nabla_{\parallel} T - \kappa_{\perp} \nabla_{\perp} T, \quad \nabla \cdot \mathbf{q} = S, \quad \kappa_{\parallel} \gg \kappa_{\perp}$$

heat flux \mathbf{q}



heat source S



A spectrum of analytical and numerical approaches exist . .

1) Analytical

- 1) tearing mode theory following Fitzpatrick (island but no chaos)
- 2) Rechester-Rosenbluth (strong chaos)

2) Monte Carlo

- 1) particle pushing (statistical)

3) Point Gaussian model

- 1) pushing local gaussians

4) Magnetic Coordinate approach

- 1) locate KAM surfaces

5) Markov model of transport through cantori

- 1) following MacKay, Meiss, Percival

6) Full Numerical Simulation

- 1) finite elements, M3D, NIMROD,

*I will discuss
each of these
approaches*



Fitzpatrick [PoP 2(3), 825, 1995] applied tearing mode analysis to determine scale island width

Balancing $\kappa_{\parallel}(\mathbf{b} \cdot \nabla)^2 \sim \kappa_{\perp}(\nabla_{\perp})^2$
gives scale island width

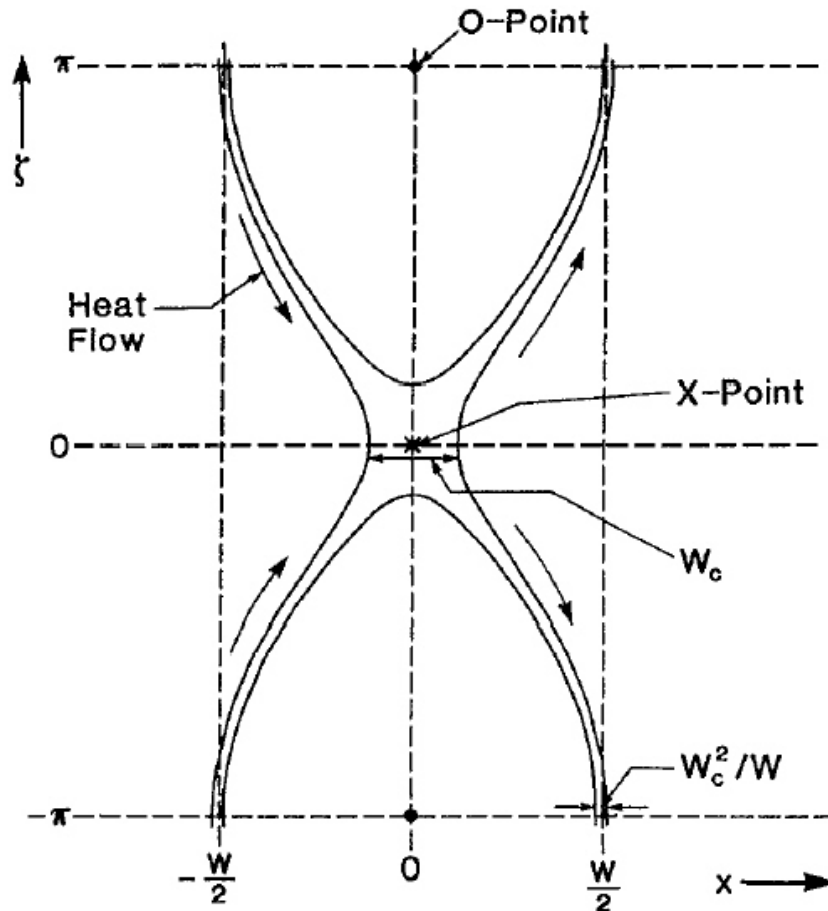
$$\frac{W}{r} \propto \left(\frac{\kappa_{\parallel}}{\kappa_{\perp}} \right)^{1/4}$$

for $w \ll W$, κ_{\perp} dominates

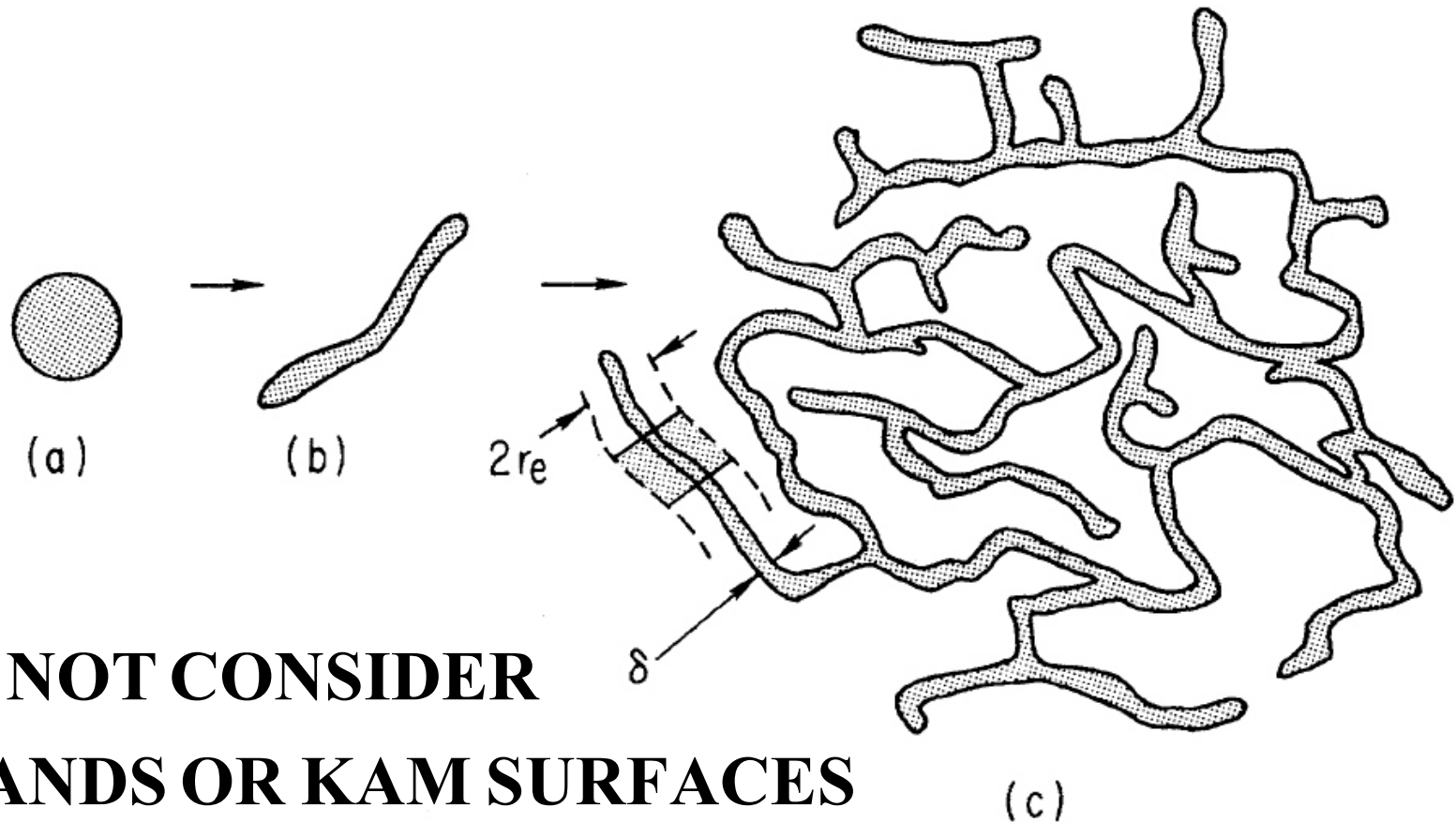
for $w \gg W$, κ_{\parallel} dominates

heat transported along boundary layer

DID NOT CONSIDER CHAOS

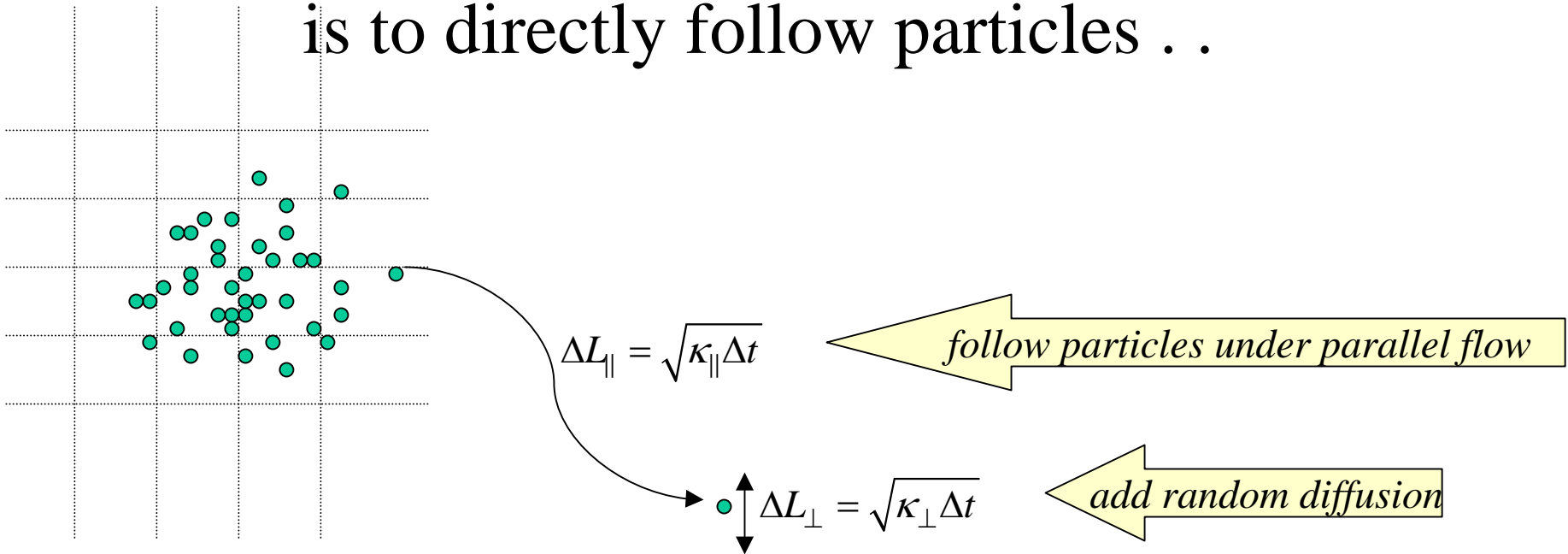


Rechester & Rosenbluth [PRL 40(1), 38, 1978]
considered the enhancement to particle
diffusion due to the stochastic field



**DID NOT CONSIDER
ISLANDS OR KAM SURFACES**

Monte Carlo : A potentially exact approach
for determining particle transport,
is to directly follow particles . .



- 1) *can handle islands, chaos, KAM surfaces . . .*
- 2) *need lots of particles to give accurate statistics*

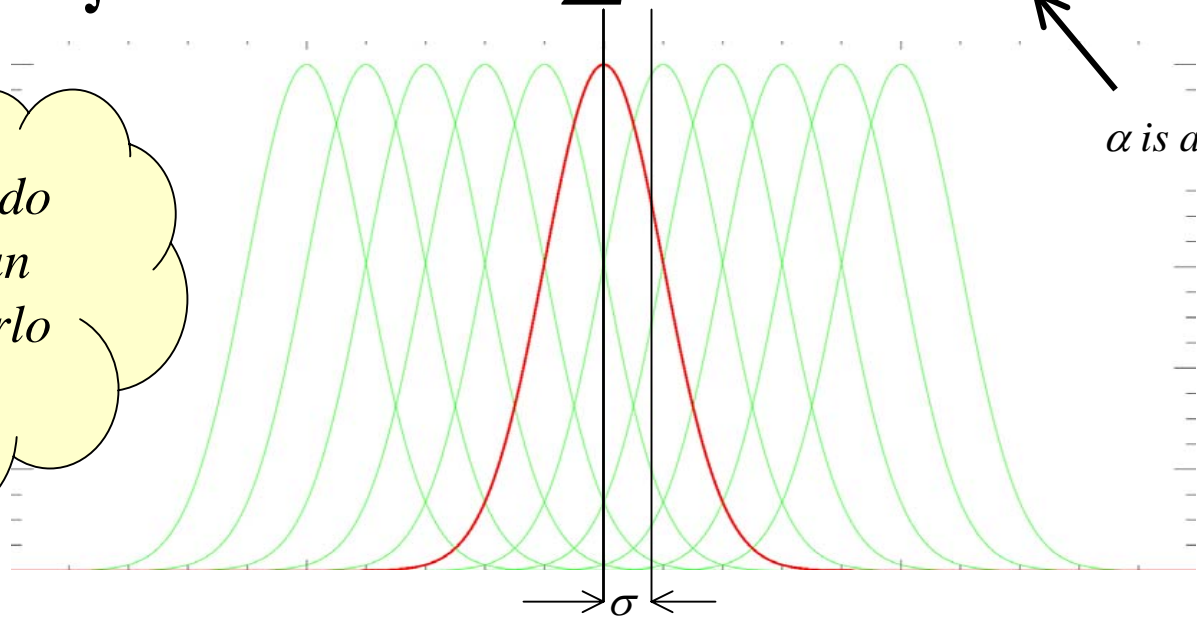
The following few slides will introduce the Gaussian model

The gaussian plays a special role in diffusion

$$g_\sigma(x) = \frac{1}{\sigma} \exp\left(\frac{-x^2}{2\sigma^2}\right) \text{ solves } \partial_t T = D\partial_{xx}^2 T \text{ with } \sigma = \sqrt{2Dt + \sigma_0^2}.$$

Model : approximate T by summation of gaussians . . .

$$\text{Consider } f(x) = \int dx' f(x') \delta(x-x') \rightarrow \sum \Delta x f_i g_\sigma(x-x_i) \alpha$$



α is an overlap constant

hope is to do better than Monte Carlo method

By representing T as sum of gaussians, effect of diffusion is approximated.

$$\mathbf{x} = (x, y), \quad \Gamma = \begin{pmatrix} \sigma_x & 0 \\ 0 & \sigma_y \end{pmatrix}$$

In 2D, $T(\mathbf{x}) = \sum_{i,j} T_{i,j} \exp \left[-\frac{1}{2} (\mathbf{x} - \mathbf{x}_{i,j})^T \cdot \Gamma^2 \cdot (\mathbf{x} - \mathbf{x}_{i,j}) \right] \alpha.$

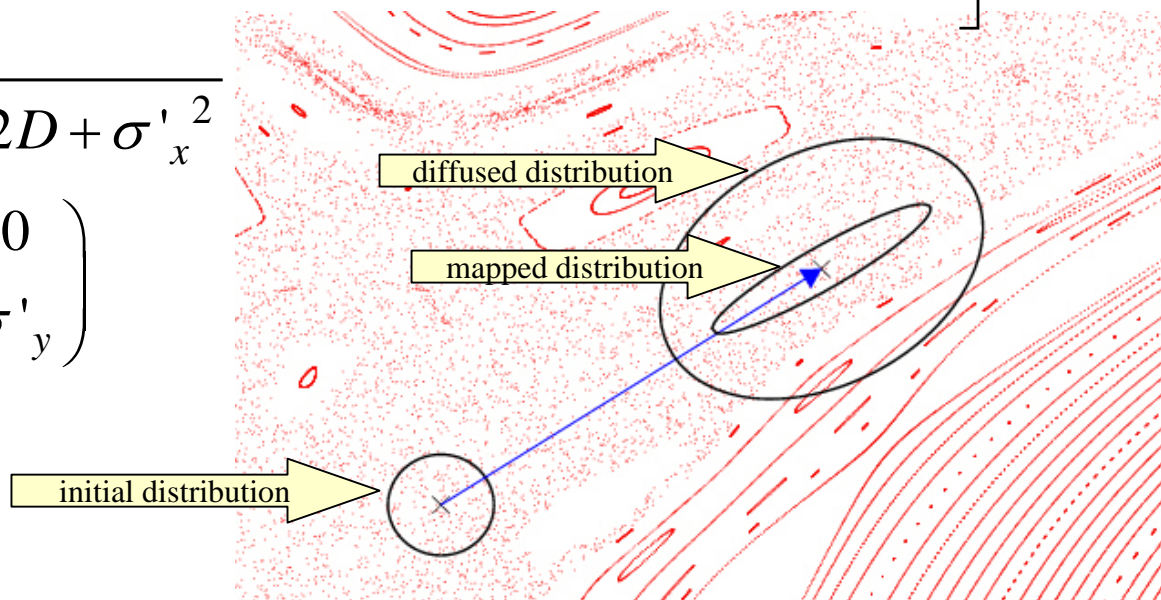
Under map, $\mathbf{x}' = \mathbf{F}\mathbf{x}$, gaussian is rotated & stretched

\mathbf{U} is rotation matrix
 \mathbf{W} is elongation matrix

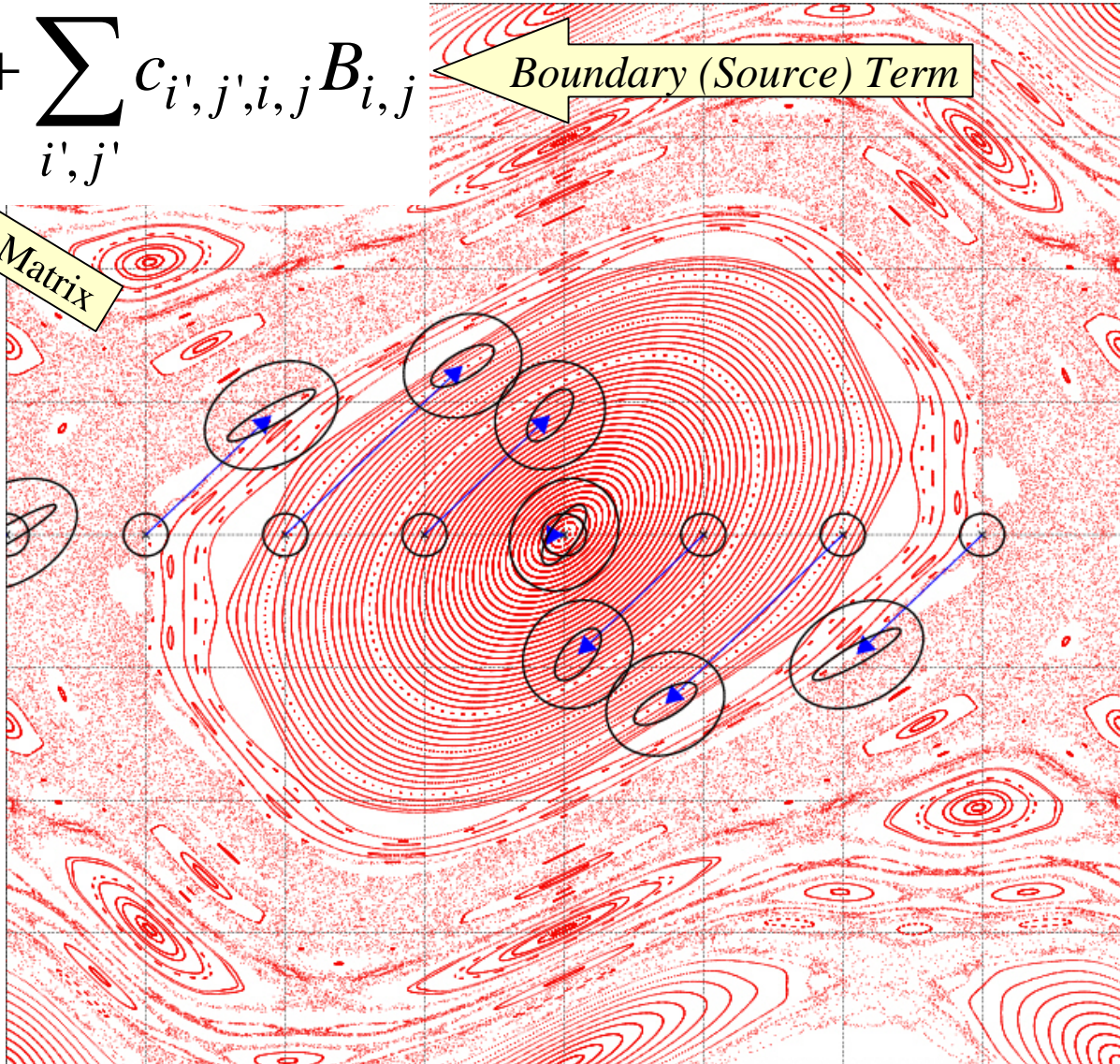
$$T(\mathbf{x}) = \sum_{i,j} f_{i,j} \exp \left[-\frac{1}{2} (\mathbf{x} - \mathbf{x}'_{i,j})^T \cdot \mathbf{U} \cdot \mathbf{W}^{-1} \cdot \Gamma^2 \cdot \mathbf{W}^{-1} \cdot \mathbf{U}^T \cdot (\mathbf{x} - \mathbf{x}'_{i,j}) \right] \alpha,$$

and then diffuses $\sigma'_x \rightarrow \sqrt{2D + \sigma'_x{}^2}$

where $\Gamma' = \Gamma \cdot \mathbf{W} = \begin{pmatrix} \sigma'_x & 0 \\ 0 & \sigma'_y \end{pmatrix}$



The model determines Temperature evolution

$$T_{i,j}^{n+1} = \sum_{i',j'} c_{i',j',i,j} T_{i',j'}^n + \sum_{i',j'} c_{i',j',i,j} B_{i,j}$$


Linear System

$$\mathbf{T}^{n+1} = \mathbf{C} \mathbf{T}^n + \mathbf{B}$$

Steady state solved directly

$$(\mathbf{I} - \mathbf{C}) \mathbf{T} = \mathbf{B} \quad (\text{use BiCGStab})$$

this is really advection-diffusion

$$\partial_t T = \mathbf{v} \cdot \nabla T + \mathbf{D} \nabla^2 T$$

& operator splitting

An example Temperature profile is shown

1) Steady state T is shown

boundary condition $T(r \leq 1) = 1$

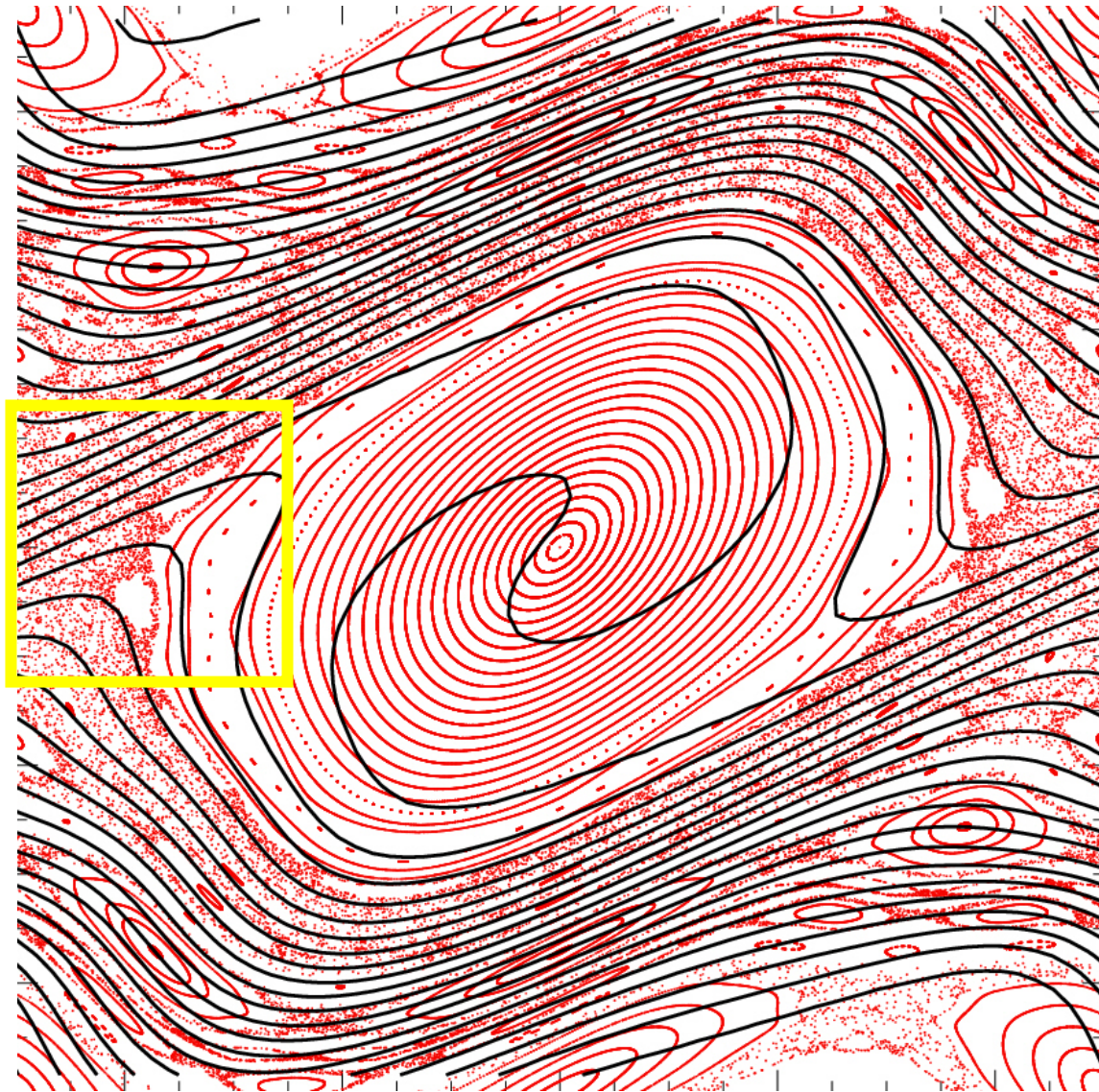
2) island width suff. large

$$\text{so } \kappa_{\parallel} (\mathbf{b} \cdot \nabla)^2 \sim \kappa_{\perp} (\nabla_{\perp})^2$$

unstable manifold

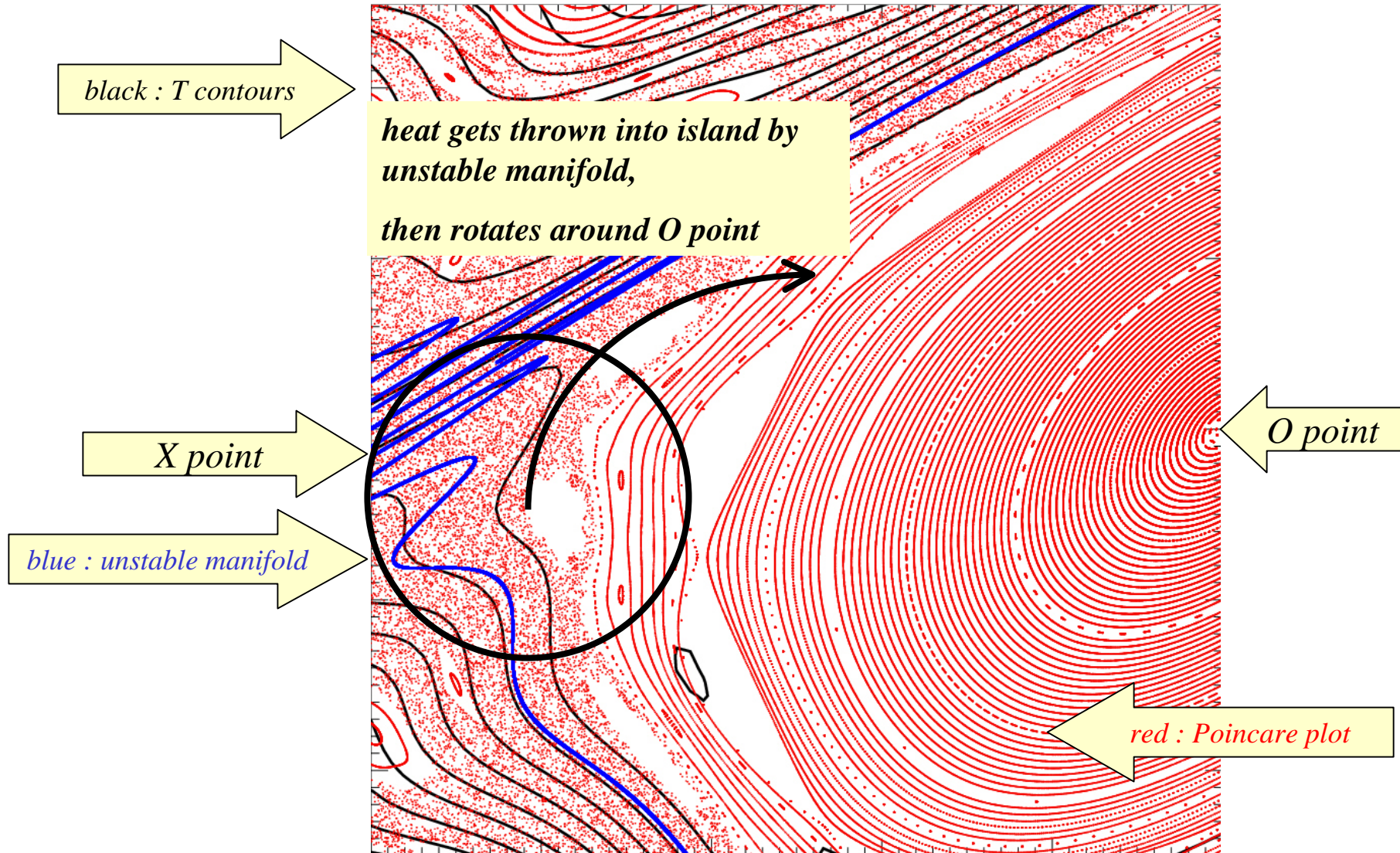
3) What is influencing T ?

here standard map is used



The unstable manifold W^u determines Temp profile

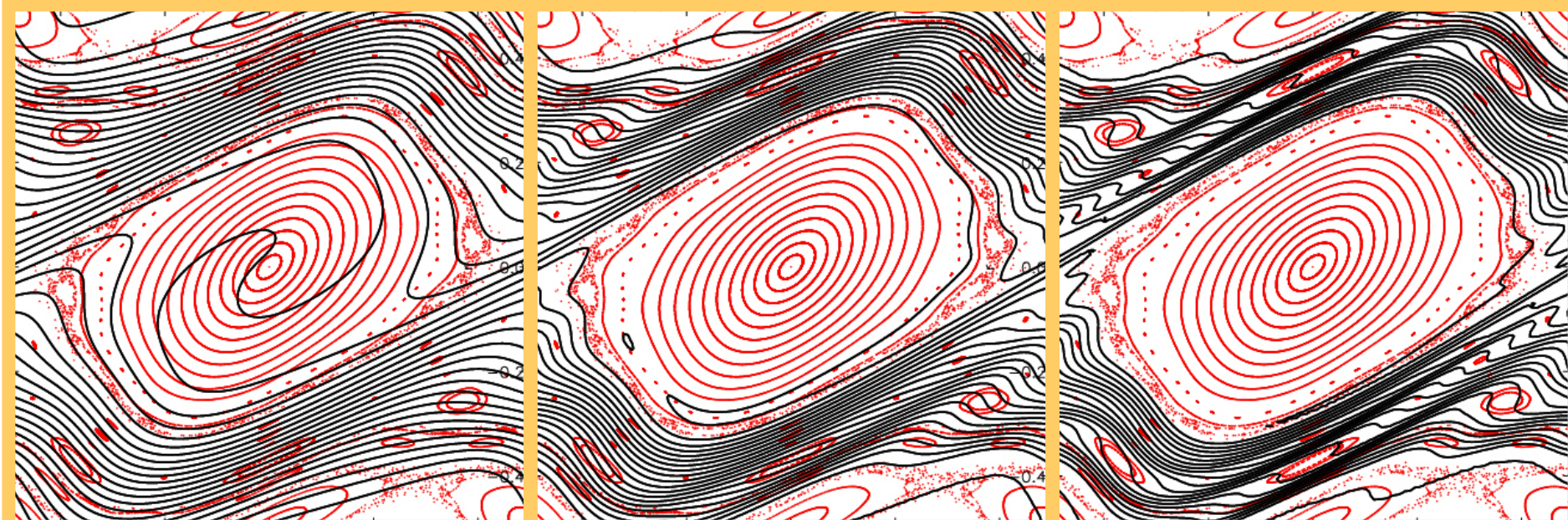
$$z \in W^u \Rightarrow F^j z \rightarrow z_0 \text{ as } j \rightarrow -\infty, \text{ where } z_0 \text{ is X point}$$



Decreasing diffusion κ_{\perp} leads to increased correlation with field

strong diffusion

weak diffusion



- 1) *I am still working through the details of this method . .*
- 2) *The method should be directly extendable to 3D chaotic fields with arbitrary diffusion . .*
- 3) *The parallel motion is solved exactly, so this method may reduce numerical pollution of perpendicular transport*
- 4) *The model is similar to both the Monte Carlo method and to direct matrix methods . .*

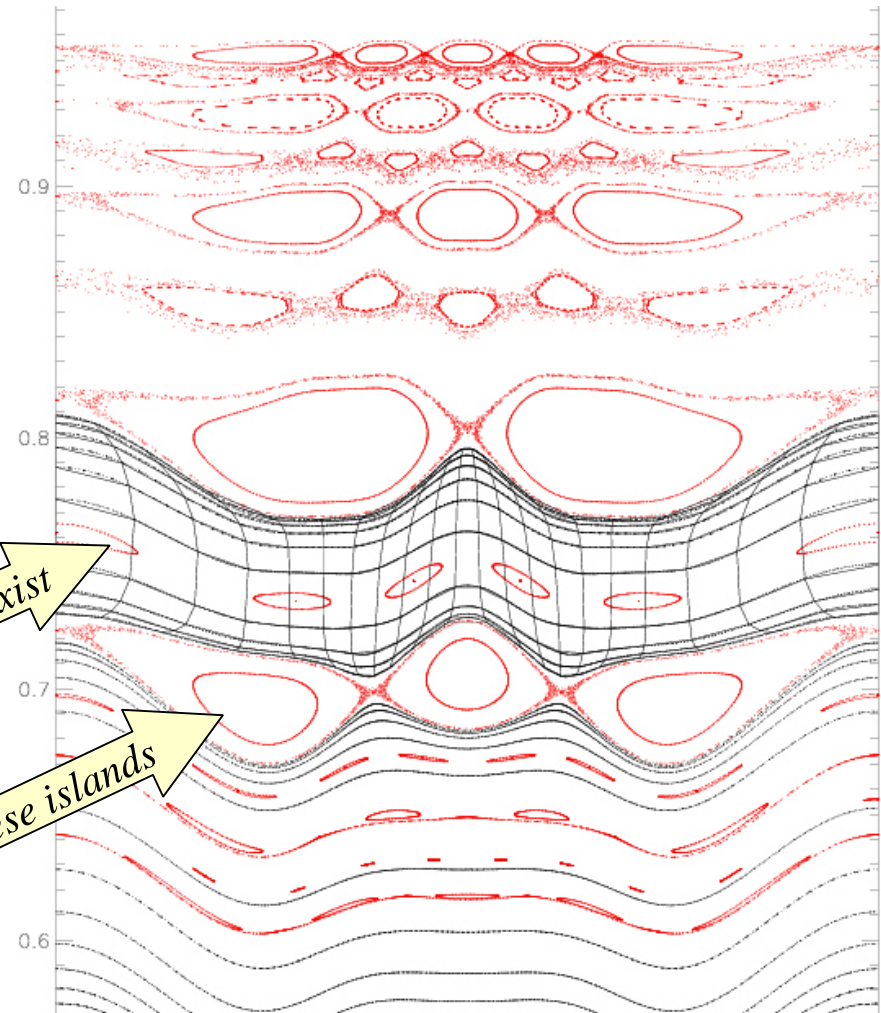
Now moving on to other approaches . . .

Can use discrete magnetic coordinates

- 1) *KAM surfaces are robust, and can construct magnetic coordinates between islands*
- 2) *The use of magnetic coordinates can reduce discretization errors*
- 3) *[Destruction of invariant surfaces and magnetic coordinates for perturbed magnetic fields, Hudson PoP 2004.](#)*
- 4) *Procedure implemented for M3D field*

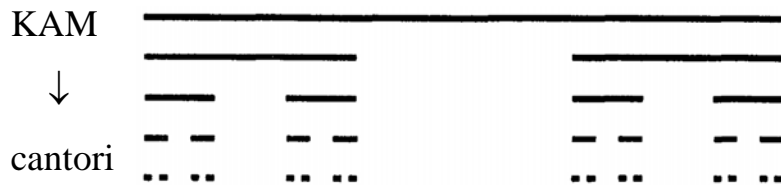
pretend these islands don't exist

flatten across these islands

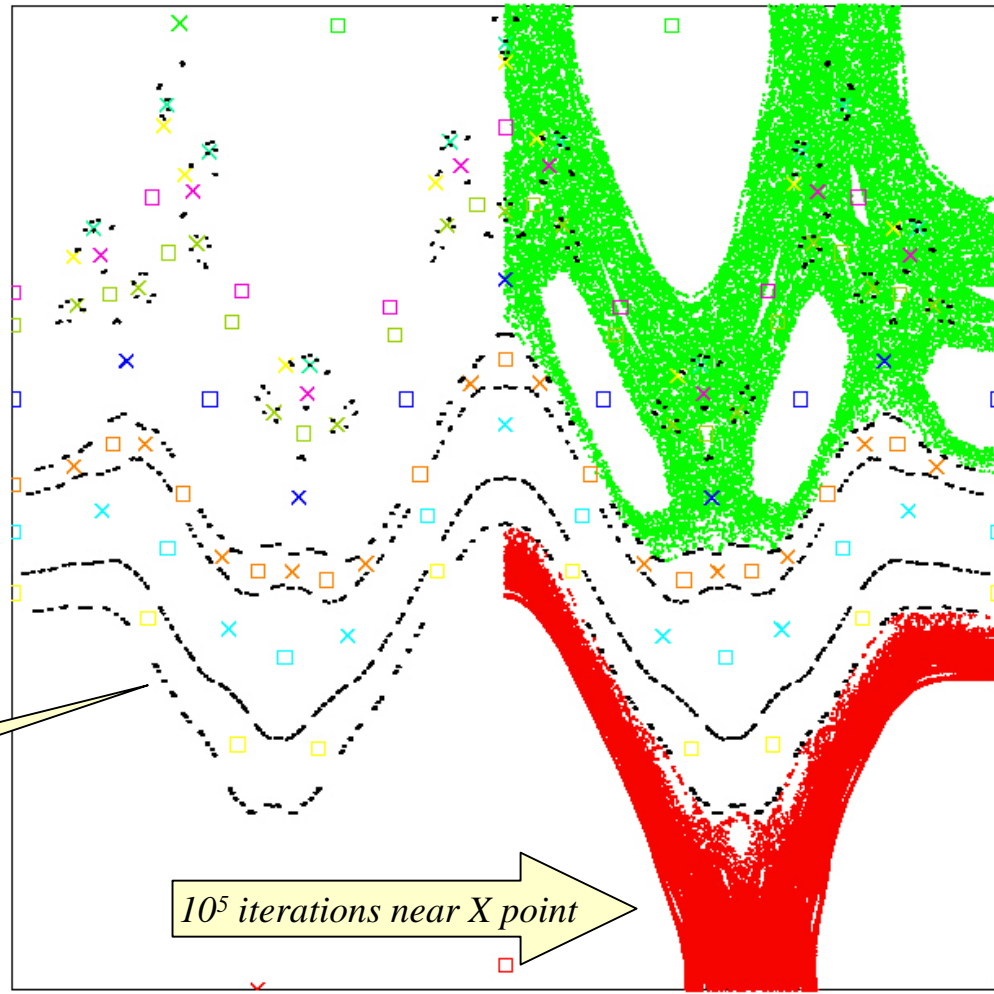


Can also use a Markov model for transport developed by MacKay, Meiss & Percival [PHYSICA D 13 (1-2): 55-81 1984]

- 1) *KAM surfaces disintegrate into leaky Cantor-set tori=cantori*



- 2) *Can develop a transport model based on probability of jumping from region to region $N_i' = \sum_j P_{i,j} N_j$*
- 3) *Cantori are approximated by high order periodic orbits eg. [Calculation of cantori for Hamiltonian flows, Hudson, Physical Review E, 2006.](#)*



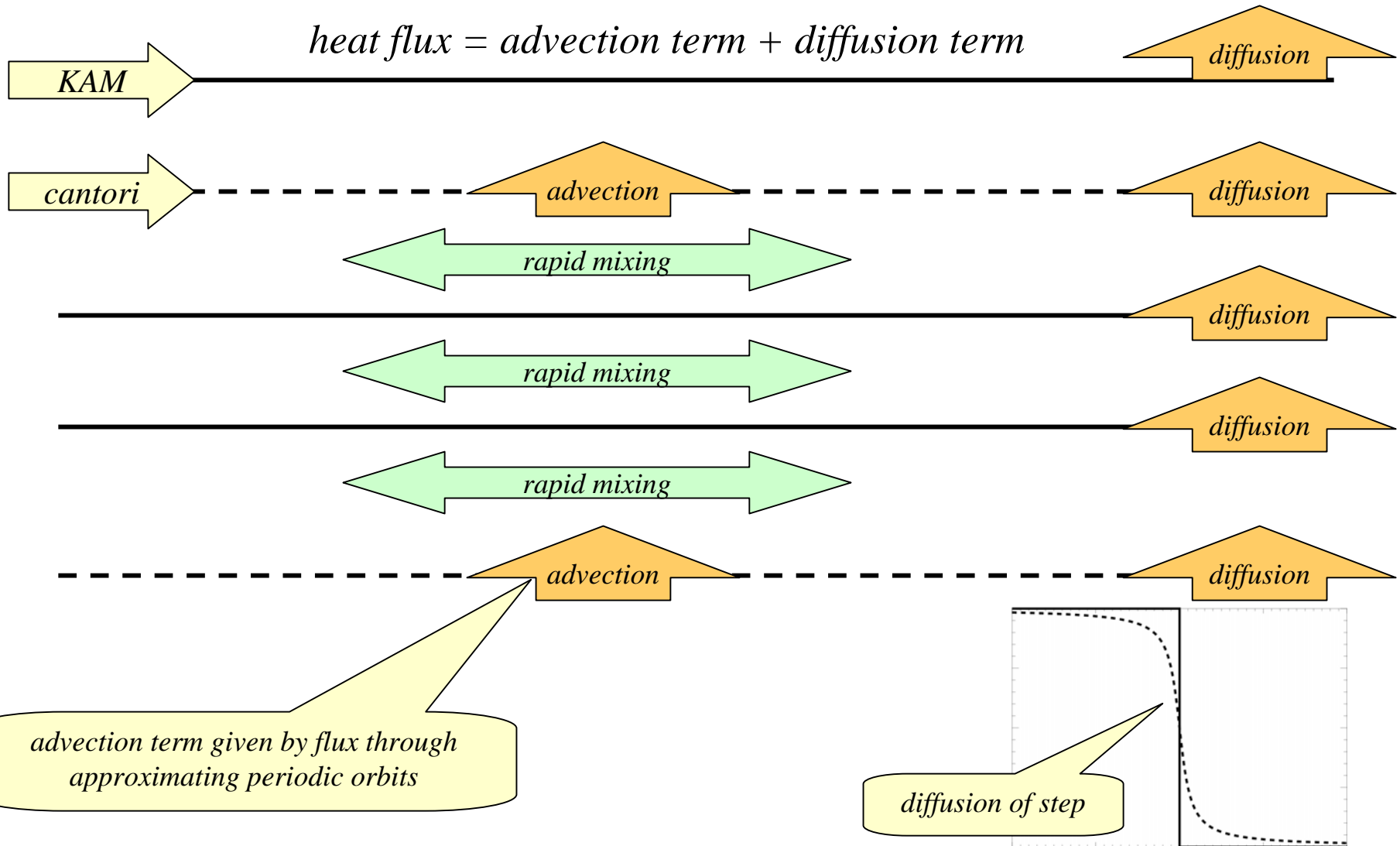
cantori can be very strong, partial barriers

10⁵ iterations near X point

A Markov model of advection-diffusion might look like this . .

Perhaps a stepped temperature profile . . ?

$$\text{heat flux} = \text{advection term} + \text{diffusion term}$$



And, of course, there is the full numerical simulation

- 1) *Such numerical simulations are probably the most reliable approach, but lack insight and need to be supported by theory.*
- 2) *The ideas presented on previous slides may enable a better initial guess for iterative solutions . .*

eg. Gunter et al., JCP 209, 354, 2005

$$\left. \frac{\partial T}{\partial x} \right|_{i+\frac{1}{2},j+\frac{1}{2}} = \frac{1}{2(\Delta x)} ((T_{i+1,j+1} + T_{i+1,j}) - (T_{i,j+1} + T_{i,j})),$$

$$\left. \frac{\partial T}{\partial y} \right|_{i+\frac{1}{2},j+\frac{1}{2}} = \frac{1}{2(\Delta y)} ((T_{i+1,j+1} + T_{i,j+1}) - (T_{i+1,j} + T_{i,j})),$$

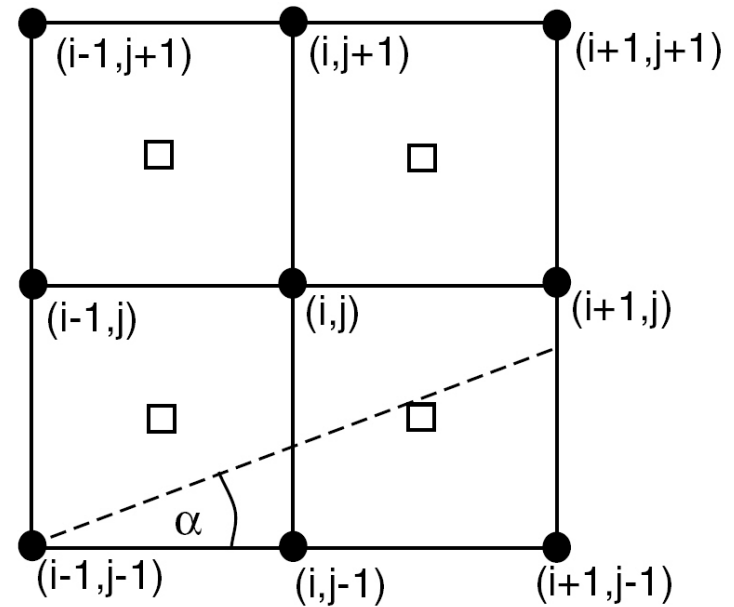
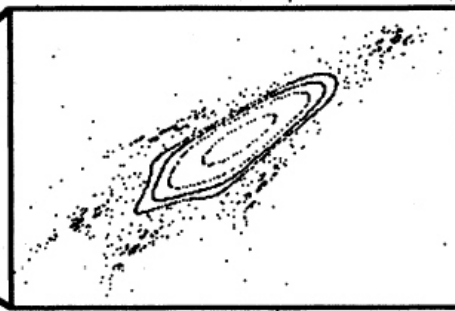
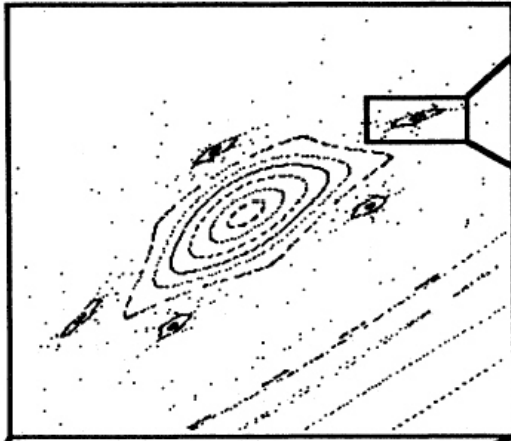


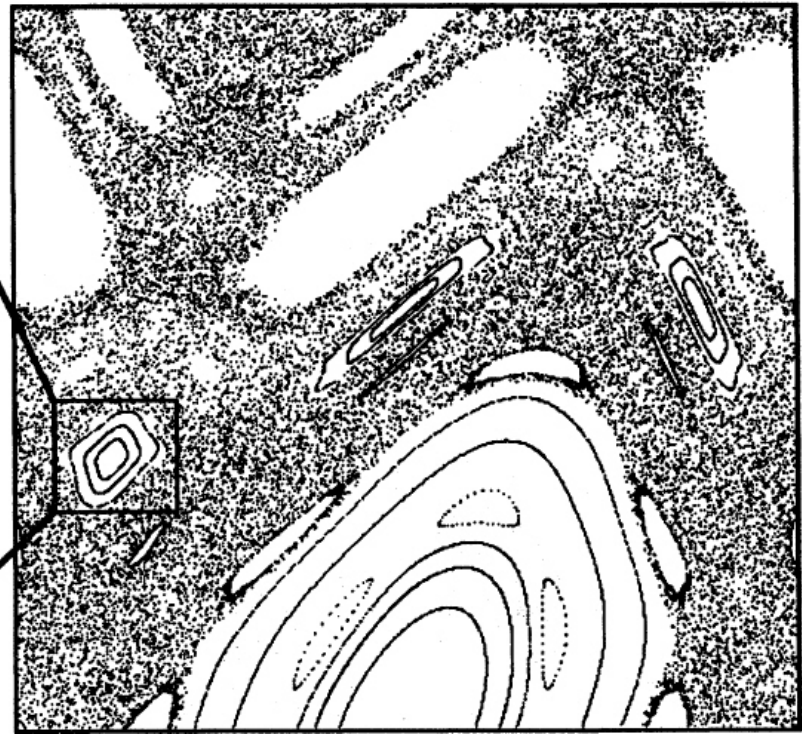
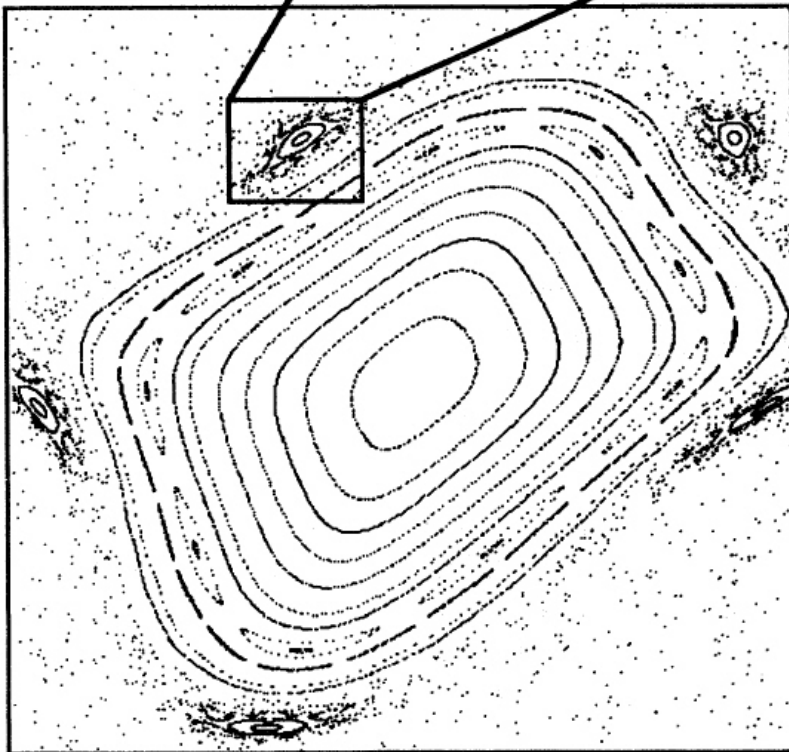
Fig. 1. Grid labelling and elementary cell.

as $\frac{\kappa_{\parallel}}{\kappa_{\perp}} \rightarrow \infty$, $\mathbf{B} \cdot \nabla T \rightarrow 0$
and the structure of \mathbf{B}
determines $T(\mathbf{x})$



Chaotic fields have infinite detail,
but what level of detail is required ?

Meiss, RMP 64(3), 795, 1992



The End