





# Magnetic islands and singular currents at rational surfaces: from MRXMHD to ideal MHD

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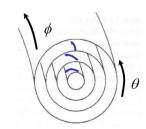
Stuart Hudson, Amitava Bhattacharjee, Per Helander

# Ideal MHD predicts singular currents

$$\mathbf{j}_{\perp} = (\mathbf{B} \times \nabla p)/B^2$$
$$\mathbf{B} \cdot \nabla u = -\nabla \cdot \mathbf{j}_{\perp}$$

Magnetic coordinates  $(\psi, \theta, \phi)$ 

$$\stackrel{\text{deal MHD}}{\Longrightarrow} \sqrt{g} \ \mathbf{B} \cdot \nabla \equiv \mathbf{t} \partial_{\theta} + \partial_{\phi}$$



Fourier decomposition

$$u = \sum_{m,n} u_{mn} e^{i(m\theta - n\phi)}$$

$$(tm-n)u_{mn} = i(\sqrt{g} \nabla \cdot \mathbf{j}_{\perp})_{mn}$$

Equation type

$$xf(x) = h(x)$$

$$x \equiv tm - n, \ h(x) \sim p'$$

$$u_{mn}(x) = h(x)/x + \hat{j}_{mn}\delta(x)$$

Pfirsch-Schlutter

Delta current

# How to calculate 3D MHD equilibrium?

### **EITHER** <u>Insist on nested flux surfaces (ideal MHD)</u>

- > VMEC
  - based on energy functional with ideal constraints
  - cannot resolve rational surfaces

### OR Relax assumption of flux surfaces

- Resistive MHD
  - initial value calculation, not based on energy minimum
  - islands develop and break rational surfaces
- ➤ Relaxed MHD
  - based on energy functional with fewer constraints
  - islands develop at rational surfaces

# Multiregion RelaXed MHD



Taylor's theory

Fewer constraints

$$F = W + \frac{\mu}{2} \left( \underbrace{\int_{V} \mathbf{A} \cdot \mathbf{B} \ dV}_{H} - H_{0} \right)$$

Topology:  $\mathbf{B} \cdot \nabla \psi \big|_{\partial V} = 0$ 

Given  $p, \Delta \psi, H_0$ 

$$\delta F = 0 \Longrightarrow \nabla \times \mathbf{B} = \mu \mathbf{B}$$

[Taylor, 1974]

MRXMH

 $F = W + \frac{\mu}{2} \left( \int_{V} \mathbf{A} \cdot \mathbf{B} \ dV - H_{0} \right) F = \sum_{l=1}^{N} \left[ W_{l} + \frac{\mu_{l}}{2} \left( H_{l} - H_{l0} \right) \right] W = \int_{V} \left( \frac{p}{\gamma - 1} + \frac{B^{2}}{2} \right) dV$ 

Topology:  $\mathbf{B} \cdot \nabla \psi \big|_{\partial V_i} = 0$ 

Given  $p_l, \Delta \psi_l, \Delta \psi_{p,l}, H_{l0}$ 

$$\delta F = 0 \Longrightarrow \begin{bmatrix} \nabla \times \mathbf{B} = \mu_l \mathbf{B} \\ [[p + B^2/2]] = 0 \end{bmatrix} \delta W = 0 \Longrightarrow \mathbf{j} \times \mathbf{B} = \nabla p$$

[Dewar, Hole, Hudson, 2006]

**Ideal MHD** 

More constraints

$$W = \int_{V} \left( \frac{p}{\gamma - 1} + \frac{B^2}{2} \right) dV$$

Topology:  $\mathbf{B} \cdot \nabla \psi = 0$ 

Given  $p(\psi), \psi_p(\psi)$ 

$$\delta W = 0 \implies \mathbf{j} \times \mathbf{B} = \nabla p$$

[Kruskal and Kulsrud, 1958]

### Existence of stepped-pressure equilibria?

# Existence of Three-Dimensional Toroidal MHD **Equilibria with Nonconstant Pressure**

OSCAR P. BRUNO

PETER LAURENCE

California Institute of Technology Universita di Roma "La Sapienza"

We establish an existence result for the three-dimensional MHD equations

$$(\nabla \times \mathbf{B}) \times \mathbf{B} = \nabla p$$
$$\nabla \cdot \mathbf{B} = 0$$
$$\mathbf{B} \cdot n|_{\partial T} = 0$$

with  $p \neq const$  in tori T without symmetry. More precisely, our theorems insure the existence of sharp boundary solutions for tori whose departure from axisymmetry is sufficiently small; they allow for solutions to be constructed with an arbitrary number of pressure jumps. © 1996 John Wiley & Sons, Inc.

Communications on Pure and Applied Mathematics, Vol. XLIX, 717–764 (1996)

# Stepped-Pressure Equilibrium Code (SPEC)

- Implementation of MRXMHD
- Finds minimum of

$$F = \sum_{l=1}^{N} \left[ W_l + \frac{\mu_l}{2} \left( H_l - H_{l0} \right) \right]$$

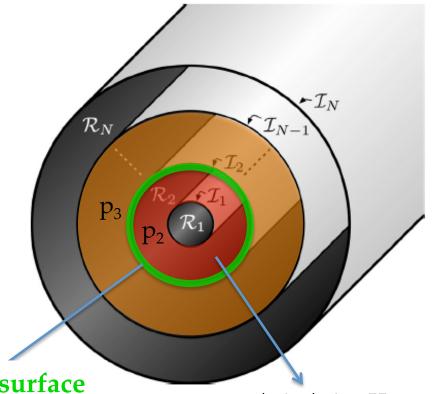
Topology:  $\mathbf{B} \cdot \nabla \psi \big|_{\partial V_l} = 0$ 

Given  $p_l, \Delta \psi_l, \Delta \psi_{p,l}, H_{l0}$ 



 $\mathcal{R}_l$ :  $\nabla \times \mathbf{B} = \mu_l \mathbf{B}$  KAM surface

$$\mathcal{I}_l: [[p+B^2/2]] = 0$$



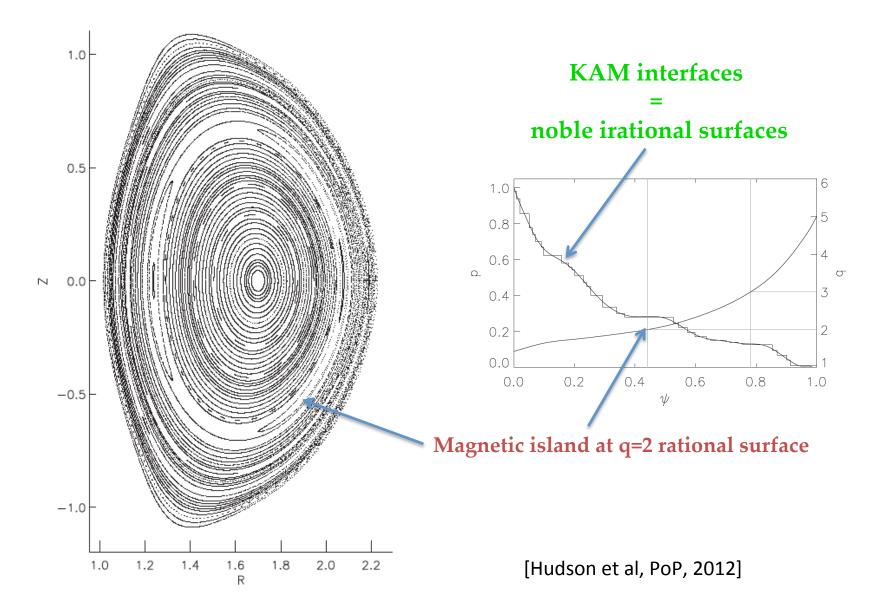
 $\Delta \psi, \Delta \psi_p, H_0$ 

or  $\Delta \psi, \Delta \psi_p, \mu$ 

or  $\Delta \psi, \epsilon^-, \epsilon^+$ 

[Hudson et al, PoP, 2012]

## Example: DIII-D equilibrium with perturbed boundary



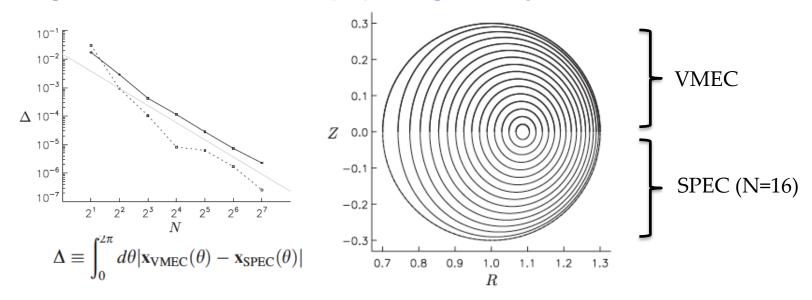
### Ideal MHD is a limit of MRXMHD

### The infinite interface limit of multiple-region relaxed magnetohydrodynamics

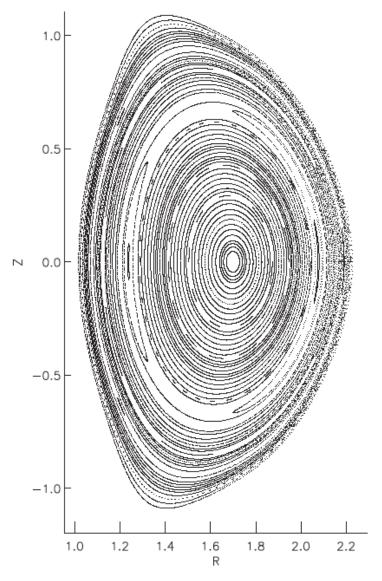
G. R. Dennis, <sup>1,a)</sup> S. R. Hudson, <sup>2</sup> R. L. Dewar, <sup>1</sup> and M. J. Hole <sup>1</sup> <sup>1</sup>Research School of Physics and Engineering, Australian National University, ACT 0200, Australia <sup>2</sup>Princeton Plasma Physics Laboratory, P.O. Box 451, Princeton, New Jersey 08543, USA

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We show the stepped-pressure equilibria that are obtained from a generalization of Taylor relaxation known as multi-region, relaxed magnetohydrodynamics (MRXMHD) are also generalizations of ideal magnetohydrodynamics (ideal MHD). We show this by proving that as the number of plasma regions becomes infinite, MRXMHD reduces to ideal MHD. Numerical convergence studies illustrating this limit are presented. © 2013 American Institute of Physics. [http://dx.doi.org/10.1063/1.4795739]



# Singular currents at rational surfaces should arise in MRXMHD with sufficiently high N



- ➤ MRXMHD allows **B** discontinuities and therefore delta currents
- ➤ Ideal MHD delta currents represent the shielding of magnetic islands

$$\mathbf{j} \equiv u\mathbf{B} + \mathbf{j}_{\perp}$$
$$u_{mn}(x) = h(x)/x + \hat{j}_{mn}\delta(x)$$

**Approach 1:** interface at rational surface

**Approach 2:** squeeze island with irationals

## A simple analytical model can be built

- ➤ Non-axisymmetric, shaped, toroidal geometry is too complicated
- > Simplify to reveal fundamental physics and build understanding
  - ➤ Geometry: cartesian, topologically toroidal plane
  - ➤ Boundary: m=1, n=0, small perturbation (resonance n/m=0)
  - Number of volumes: N=1 (Taylor), N=2, N=3
  - ➤ Constant pressure: [[B²]]=0 at interfaces

$$s = +1 \qquad \qquad R = R_{l,0} + R_{l,1} \cos \theta$$

$$\mathbf{x} = \theta \hat{\mathbf{i}} + \zeta \hat{\mathbf{j}} + R(s, \theta, \zeta) \hat{\mathbf{k}}$$

$$s \in [-1, 1], \text{ and } \theta, \zeta \in [0, 2\pi]$$

$$\nabla \times \mathbf{B} = \mu_l \mathbf{B}$$

$$\hat{\mathbf{k}}$$

$$\hat{\mathbf{i}}$$

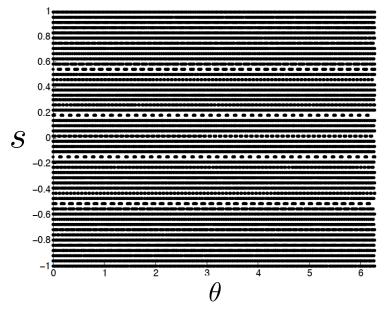
$$s = -1$$

$$R = R_{l-1,0} + R_{l-1,1} \cos \theta$$

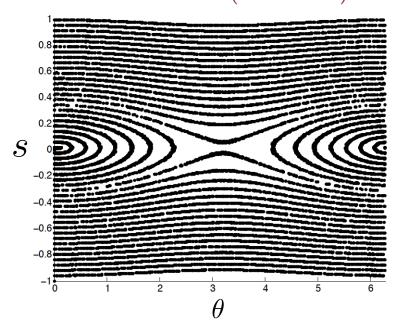
### Toroidally symmetric solution for small perturbation

- ightharpoonup Toroidal symmetry is assumed:  $\partial_{\zeta} = 0$
- Solution of the form:  $\mathbf{B} = (B_s, B_\theta, B_\zeta) = \mathbf{B}(s, \theta; \mu, B_0, \hat{B}_0, \Delta_l, \epsilon_{l-1}, \epsilon_l)$
- ightharpoonup Can relate  $(\mu, B_0, \hat{B}_0) \leftrightarrow (\Delta \Psi_l, \iota_l^+, \iota_l^-)$  then e.g. choose  $\iota_l^{\pm} = \pm 1.618033...$
- Poincaré section for the magnetic field trajectories

### **UNPERTURBED** ( $\varepsilon = 0$ )

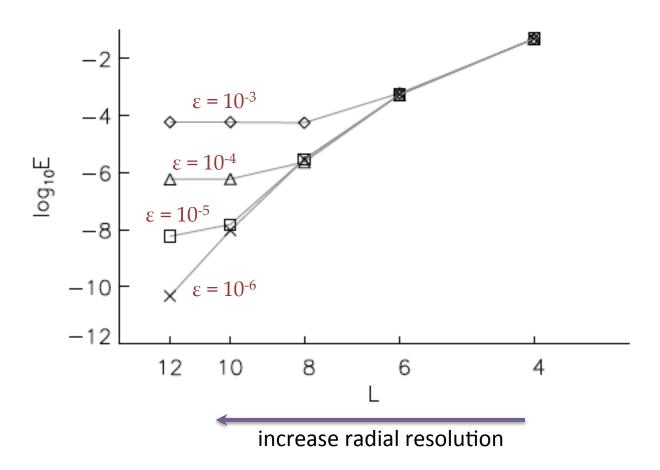


#### **PERTURBED** ( $\varepsilon = \pm 2 \times 10^{-2}$ )

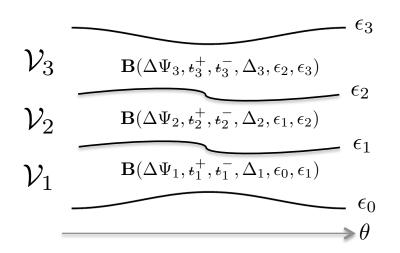


### SPEC converges to the analytical solution

Compare analytical solution with SPEC (for small perturbations)

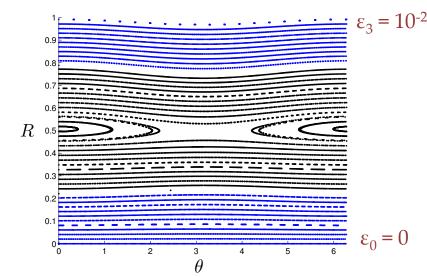


### Multiple relaxed volumes can be coupled



### Additional geometric freedom is constrained by

$$[[B^2]]=[[B^2]]_{m=0}+[[B^2]]_{m=1}\cos\theta=0$$
 at the internal interfaces provides  $\Delta_1,\Delta_2$  provides  $\epsilon_1,\epsilon_2$ 

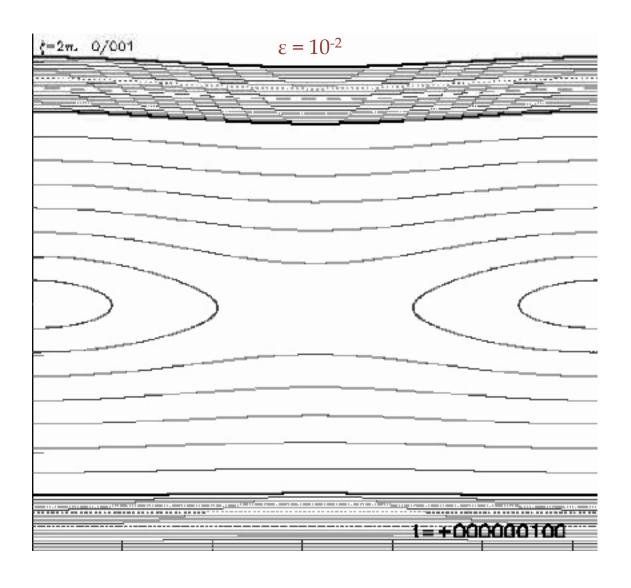


Can solve analytically: 
$$\mathcal{M} \begin{bmatrix} \epsilon_1 \\ \epsilon_2 \end{bmatrix} = \mathcal{D} \begin{bmatrix} \epsilon_0 \\ \epsilon_3 \end{bmatrix}$$

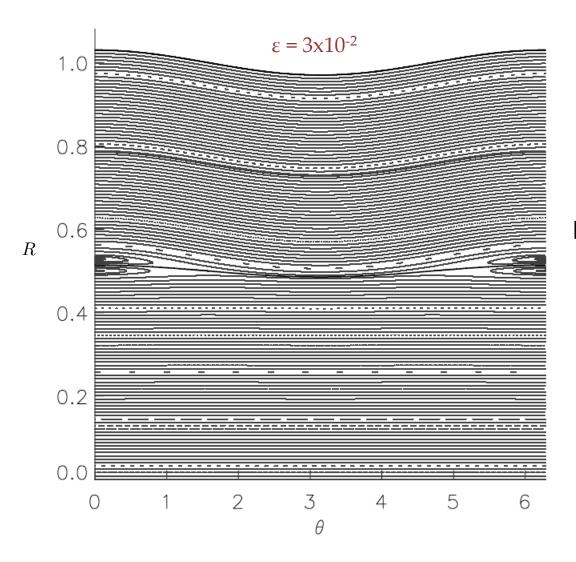
However, 
$$\lim_{\Delta\Psi_2 \to 0, \ \iota_2^{\pm} \to 0} \left[ \det \mathcal{M} \right] = 0$$

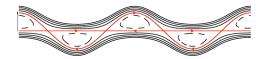
Solution becomes non-unique!

# Example: island squeezing with SPEC



### Example: interface at rational surface with SPEC





[Boozer and Pomphrey, PoP, 2010]

Suggested that delta current at the rational surface is not enough to completely shield the island

### Questions we need to answer

- ➤ We are developing a fundamental understanding of how MRXMHD can handle island shielding at rational surfaces
- ➤ Both theory and numerics suggest that magnetic island shielding may not be unique (delta current not unique).
- Is non-uniqueness possible in MRXMHD?
- Does magnetic island shielding require additional constraints?
- ➤ How does this connect to the Hahm-Kulsrud-Taylor model?
- ➤ How does this connect to the plasmoid solutions of Dewar et al?
- How does this connect to the Boozer-Pumphrey residual islands?