



Spectrum of multi-region-relaxed magneto-hydrodynamic modes in slab geometry or **Putting the D in MRMHD** a prescription for all that ails ideal MHD!

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If it ain't broke, why fix it?

- Ideal MHD is overconstrained
 - No heat transport along field lines
 - No reconnection so islands or chaos cannot form
 - Thus inapplicable to hot and 3D plasmas!
- Fix by removing the bad constraints and keeping the good, doing more with less!





MRxMHD: M stands for **Multi-region** (aka waterbag) Rx stands for **Relaxed**; ...D stands for **Dynamics**





Fundamental postulates of new general reformulation of MHD:

□ ∃ transport interfaces, \mathcal{I}_i or $\Gamma_{i,j}$, or $\partial \Omega_{i,j}$ (e.g. nested tori or island separatrices), that act like sheets of ideal-MHD plasma

 Plasma relaxes (in some generalized Taylor sense) in regions *P_i* (or *Ω_i*) bounded by the interfaces
Only a *subset* of ideal-MHD invariants apply



SPEC (currently) uses MRxMHS, not MRxMHD:

MRxMHS = Multi-region Relaxed MagnetohydroStatics (i.e. equilibrium theory)

- Taylor relaxation *energy* principle
- constant pressure in each region
- MRxMHD = Multi-region Relaxed Magnetohydro*Dynamics*

New approach:: ^{*} use *Hamilton's Principle* — stationarity of time-integrated *Lagrangian*

⇔ constant *temperature* in each region

supports sound waves within relaxation regions as well as radially compressible and Alfvén modes + *tearing*

- can treat *development* of resonant current sheets
- can add equilibrium flow to SPEC and will be basis for a new time-evolution waterbag code

Ref. Stuart Hudson's talk yesterday



MRxMHD Lagrangian is *kinetic energy* minus MHD *potential energy* + constraint terms:

• MHD Lagrangian density in region *i*

$$\mathcal{L}^{\text{MHD}} = \rho \frac{v^2}{2} - \frac{p}{\gamma - 1} - \frac{B^2}{2\mu_0}$$

• Constrained Lagrangian in region *i*

$$L_i = \int_{\Omega_i} \mathcal{L}^{\text{MHD}} dV + \tau_i (S_i - S_{i0}) + \mu_i (K_i - K_{i0})$$

• Helicity and entropy macroscopic invariants

$$K_i \equiv \int_{\Omega_i} \frac{\mathbf{A} \cdot \mathbf{B}}{2\mu_0} \, dV \qquad S_i \equiv \int_{\Omega_i} \frac{\rho}{\gamma - 1} \ln\left(\kappa \frac{p}{\rho^{\gamma}}\right) dV$$



In varying action, ρ is constrained *holonomically* to the displacement ξ of each fluid element:

• Mass conserved *microscopically*, i.e. pointwise

$$\delta \rho = -\nabla \cdot (\rho \boldsymbol{\xi}) \text{ in } \Omega_i$$

- Helicity and entropy constrained *macroscopically*, throughout Ω_i, using Lagrange multipliers μ_i and τ_i, while *p* and **A** are free fields
- Including vacuum field energy, total Lagrangian is

$$L = \sum_{i} L_{i} - \int_{\Omega_{v}} \frac{\mathbf{B} \cdot \mathbf{B}}{2\mu_{0}} \, dV$$

• Setting variation of action to 0 gives EL equations: $\delta \int Ldt = 0$



Equations within Ω_i

• Mass conservation (microscopic constraint)

$$\frac{\partial \rho}{\partial t} = -\boldsymbol{\nabla} \boldsymbol{\cdot} (\rho \mathbf{v})$$

• $\delta p \Rightarrow$ Isothermal equation of state

$$p = \tau_i \rho$$
 (N.B. $\tau_i = C_{\mathrm{s}i}^2$)

δA ⇒ Beltrami equation

$$\nabla \times \mathbf{B} = \mu_i \mathbf{B} \quad (\text{N.B.} \Rightarrow \mathbf{j} \times \mathbf{B} = 0)$$

• $\boldsymbol{\xi} \Rightarrow$ Momentum equation (Euler fluid) $\rho\left(\frac{\partial \mathbf{v}}{\partial t} + \mathbf{v} \cdot \boldsymbol{\nabla} \mathbf{v}\right) = -\boldsymbol{\nabla} p$



Equations on interface $\Gamma_{i,j}$

ξ ⇒ Force balance

$$\left[p + \frac{B^2}{2\mu_0}\right]_{i,j} = 0$$

Surface constraints

 $\mathbf{n}_i \cdot \mathbf{B} = 0 \quad \text{on } \partial \Omega_i$ $\mathbf{n}_i \cdot [\![\mathbf{v}]\!]_{i,j} = 0 \quad \text{on } \partial \Omega_{i,j}$

• Complete set of equations, consistent because derived from single scalar function *L*



Proving the MRxMHD pudding:

- Q1) What is the MRxMHD spectrum and what are the effects of field-line curvature and <u>equilibrium mass flow</u> on stability?
- Q2) When are the current sheets topologically stable towards internal plasmoid formation (reconnection)?
- Q3) When do unstable modes saturate at a low level or develop nonlinearly into explosive events?



What happens in static limit?

•
$$\partial_t \to \mathbf{0} \Rightarrow \nabla \cdot (\rho \mathbf{v}) = 0 \quad \mathbf{v} \cdot \nabla \mathbf{v} = -\tau_i \nabla \ln \rho$$

only solutions valid for any flowline configuration, from nested surfaces to arbitrarily chaotic, are

$$\rho = \rho_{0i} \exp\left(-\frac{v^2}{2\tau_i}\right) \diamondsuit p = p_{0i} \exp\left(-\frac{v^2}{2\tau_i}\right)$$

(N.B. *incompressible* in limit $v/C_s \rightarrow 0$) and

$$\nabla \times \mathbf{v} = \alpha_{0i} \exp\left(-\frac{v^2}{2\tau_i}\right) \mathbf{v}$$

Almost isomorphous to **B** equation: should be implementable in SPEC. Derivable variationally?



Slow limit? Switch-on slab boundary ripple:*



Ripple amplitude: $\alpha = 0.003$ Current density exhibits sign reversal

2-region MRxMHD Hahm-Kulsrud model: mirror-image ripple top and bottom excites modulated current sheet at x = 0





Full *t*-dependence: linear modes in slab





Loading all mass on interfaces causes problems:



- Growth rate goes to ∞ when Newcomb node goes through interface
- Growth rate zero if wall or **k.B**=0 is at interface



Conclusion

- Action-based MRxMHD shows great promise
 - Very simple
 - Includes reconnection and flow in natural way
- We need to check physical reasonability of predictions in simple models
- Need both to extend SPEC and build a new time-evolution code