

Progress on Nonlinear Resistive MHD Code Verification Problems with M3D-C1

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Outline

1. Linear resistive MHD test problems for JA-2
2. Nonlinear resistive MHD test problems for JA-2
3. Nonlinear resistive MHD test problem with resistive wall and error fields

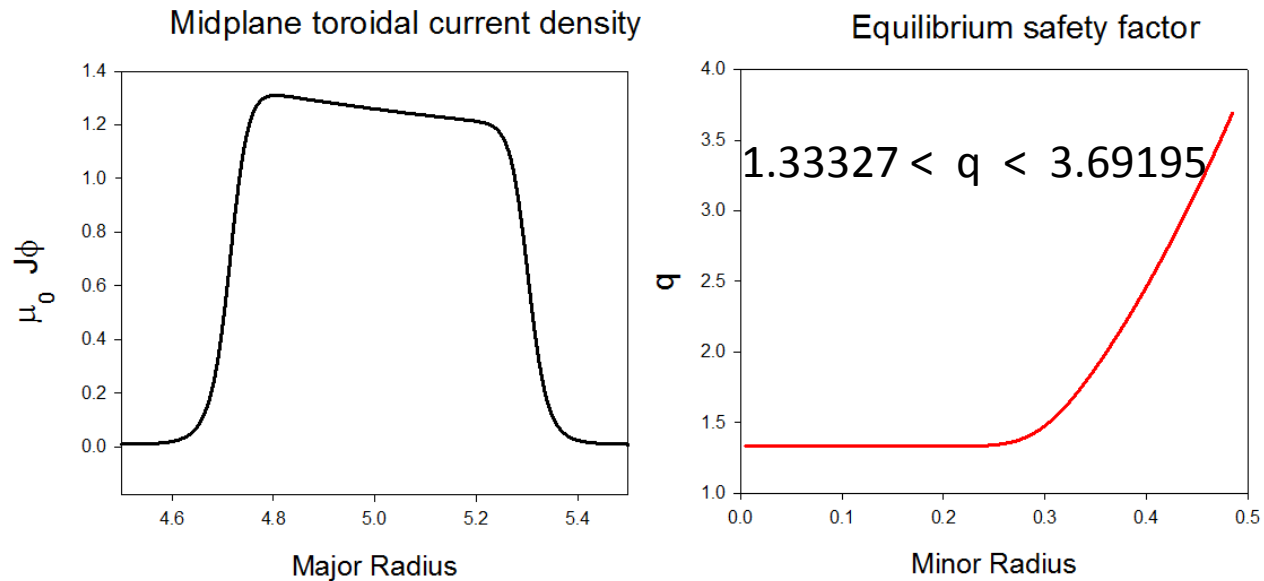
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Progress on ITPA JA-2 benchmark with M3D-C1

We started with a cylindrical equilibrium with a given $q(r)$ profile. This was then converted to a toroidal $A=10$ equilibrium with the same $q(\psi)$. This was unstable to both $n=1$ and $n=2$ modes.

$$q(r) = 1.33 \left[1 + (r^2 / .354)^4 \right]^{1/4}$$

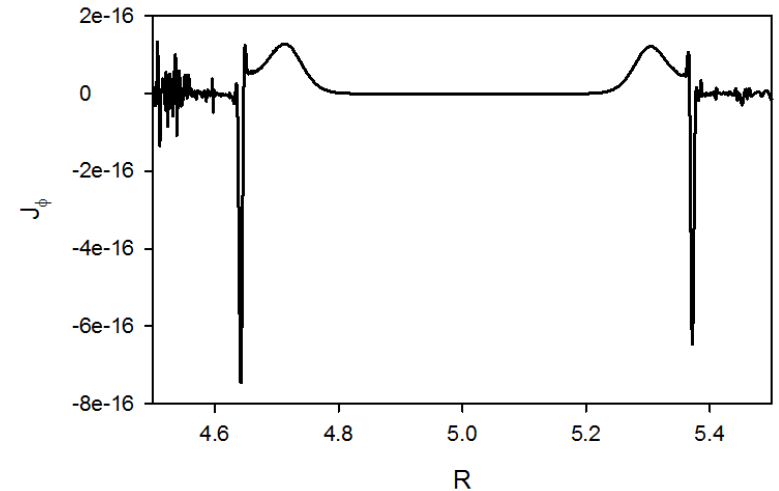
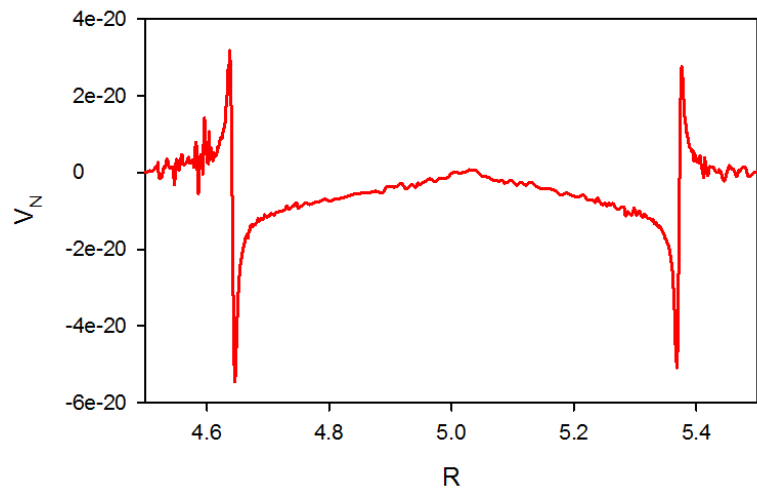
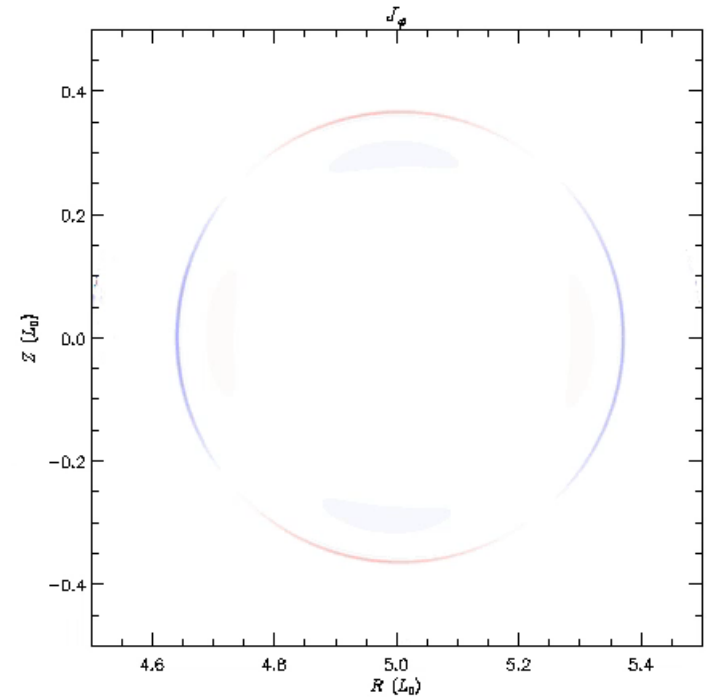
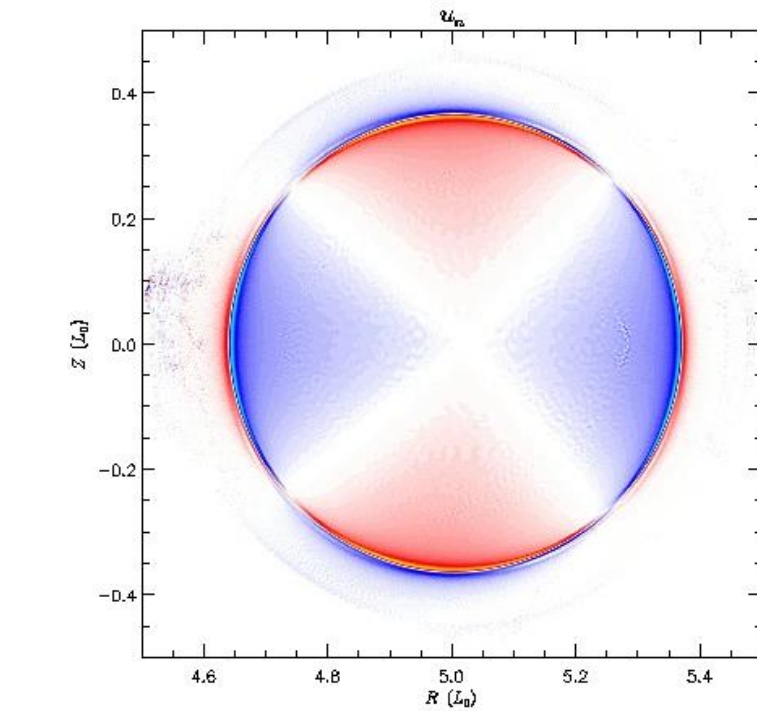


Equilibrium parameters (SI units): $R=5$, $a=0.5$, $B_T = 4.2$, $n_0 = 10^{20}$

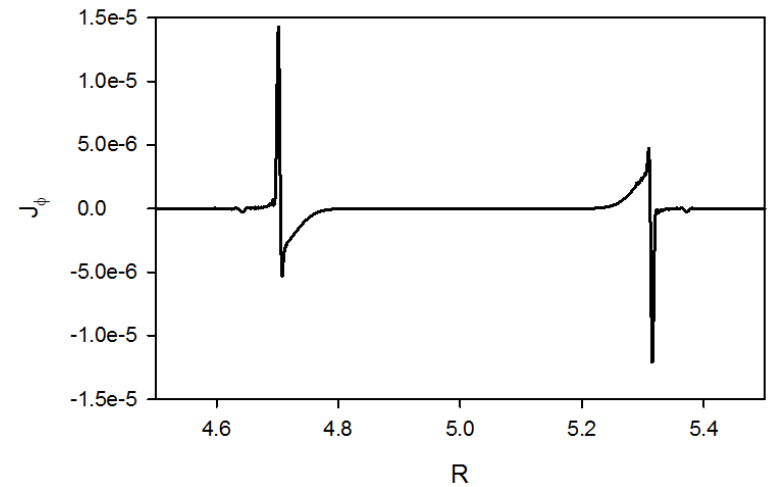
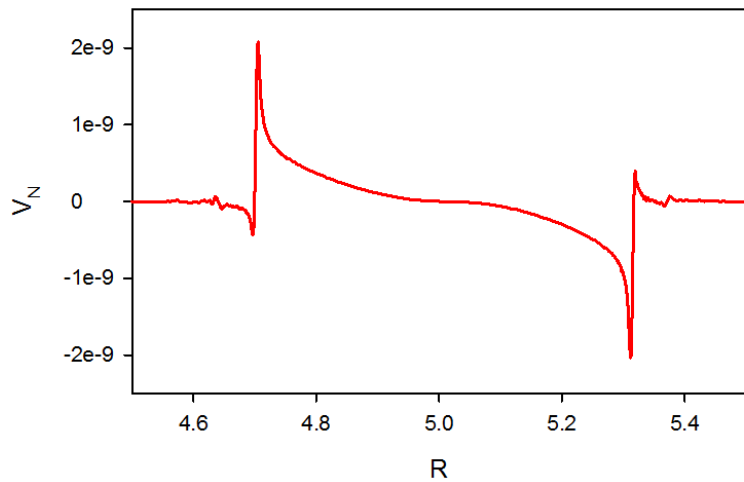
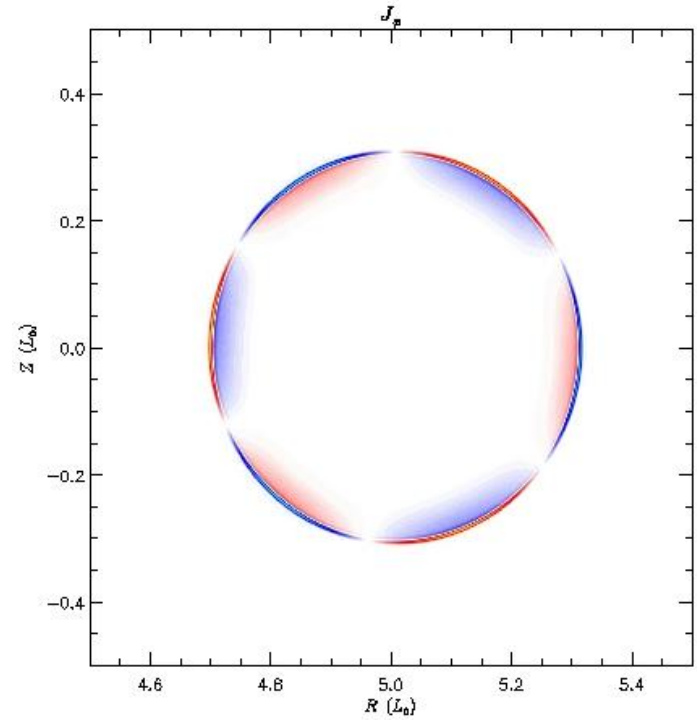
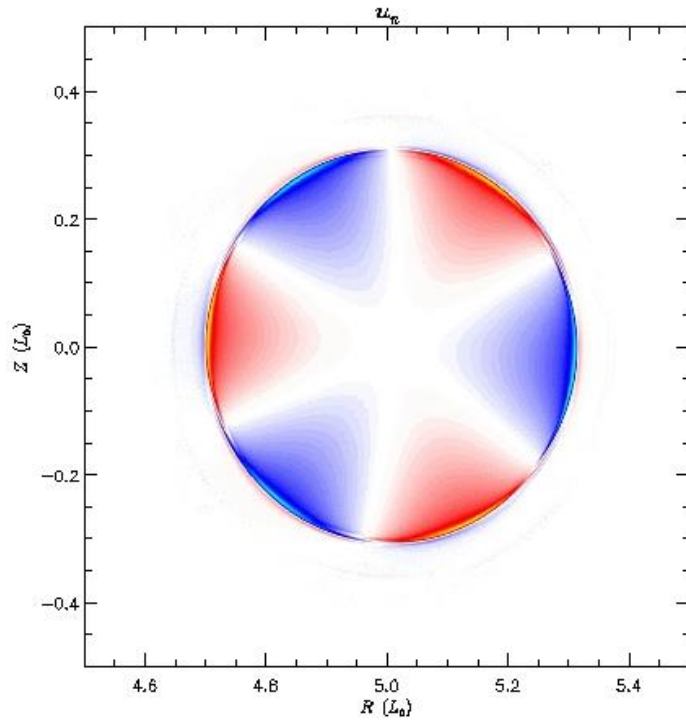
$$\eta = 0.21 \times 10^{-6} \eta_0, \quad \nu = 10^{-8} \nu_0, \quad t_0 = [\mu_0 n_0 M_i]^{1/2} = 4.58 \times 10^{-7}$$

$$\eta_0 = \left[\frac{\mu_0}{n_0 M_i} \right]^{1/2} = 7.52, \quad \nu_0 = \left[\frac{n_0 M_i}{\mu_0} \right]^{1/2} = 0.132, \quad S \equiv \frac{a^2 B_T}{R} \frac{\eta_0}{\eta} = 10^6$$

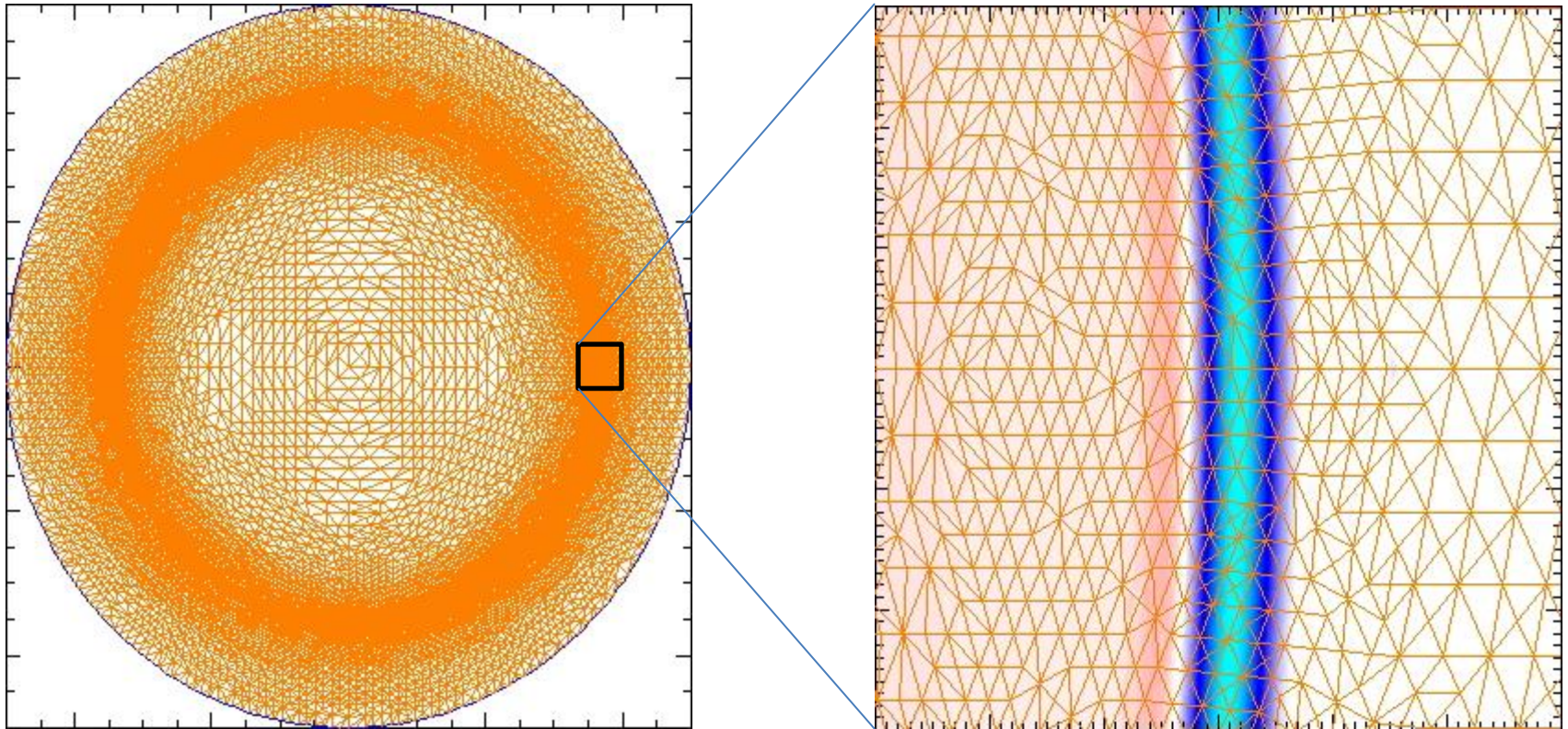
2. Normal displacement and perturbed current eigenfunctions for $n=1$ mode



3. Normal displacement and perturbed current eigenfunctions for n=2 mode

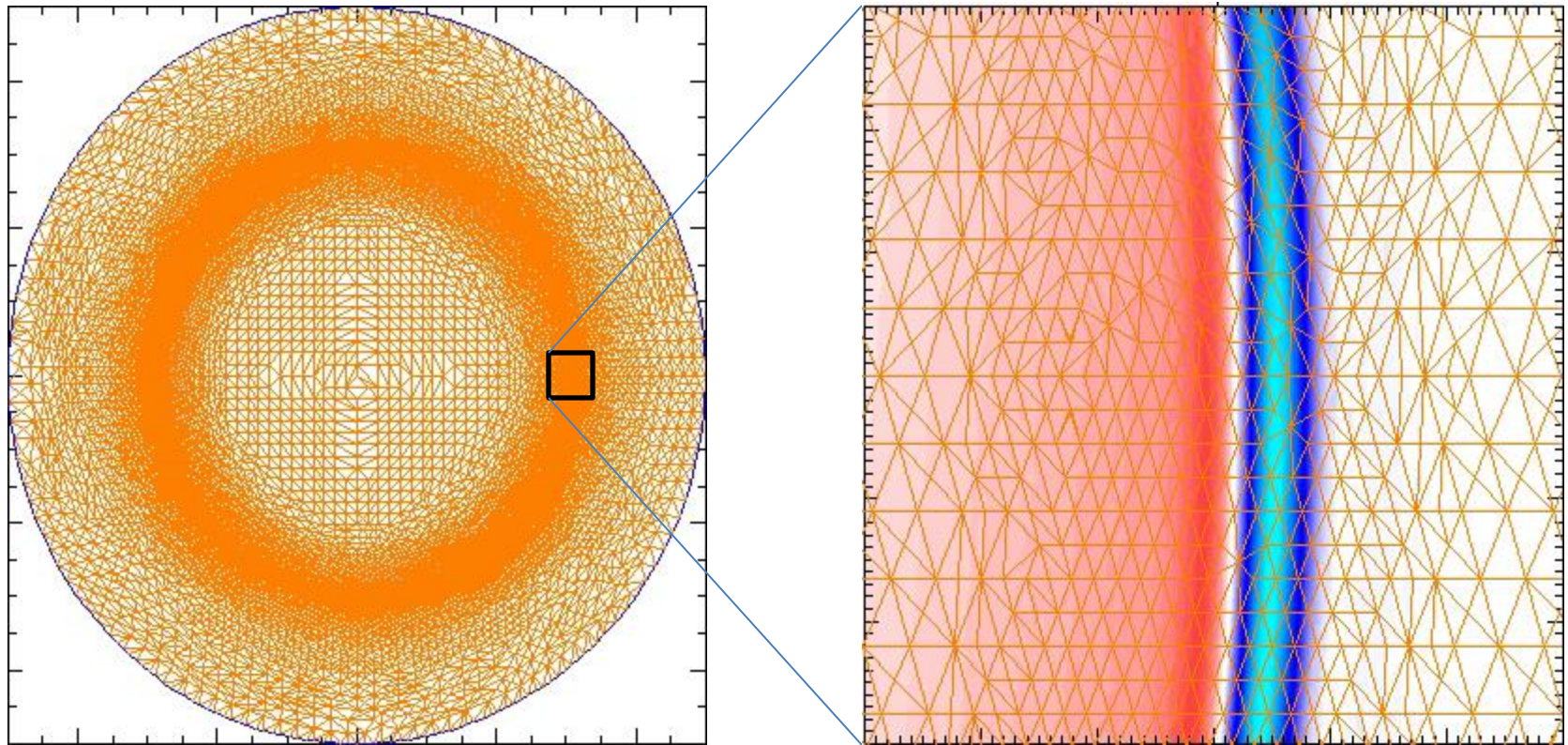


4. Adapted mesh used for $n=1$ mode



Mesh size varies from 0.002 to 0.03

5. Adapted mesh used for $n=2$ mode



Mesh size varies from 0.002 to 0.03

6. Convergence study for the $S=10^6$, $\beta=0$, $A=10$ equilibrium

DT	ΔX_{\min}	$\gamma_{n=1} t_H$	$\gamma_{n=2} t_H$
$5 \tau_A$	0.002	1.64 E-3	7.19 E-3
$5 \tau_A$	0.004	1.70 E-3	7.87 E-3
$10 \tau_A$	0.002	1.65 E-3	7.29 E-3

$$\tau_H \equiv \frac{R}{B_T} [\mu_0 n_0 M_i]^{1/2}$$

7. Dependence on geometry and beta

n=1		
geom	β	$\gamma \tau_H$
A=10	0	1.64 E-3
A=10	10^{-6}	1.13 E-3
Cyl	0	1.49 E-3
Cyl	10^{-6}	1.49 E-3

n=2		
geom	β	$\gamma \tau_H$
A=10	0	7.19 E-3
A=10	10^{-6}	5.00 E-3
Cyl	0	2.89 E-3
Cyl	10^{-6}	2.00E-3

- n=2 is always more unstable than n=1
- Finite pressure is stabilizing in torus, not so much in cylinder
- Toroidal geometry is more unstable than cylinder, especially for n=2

8. Test of growth rate with different numerical options:

Cylindrical test case:

$\rho_0 = 1.E-4$, $dt = 2.0$, $\mu = 1.e-5$, $\eta = 1.e-5$ ($S = 2.1 \times 10^4$), uniform mesh $\Delta x \sim 0.03$

isplitstep=1							isplitstep=0
	impmod=0			impmod=1			
	16	32	linear	16	32	linear	linear
n=1	.0063	.0063	.0063	.0070	.0070	.0071	.0072
n=2	.0048	.0048	.0048	.0064	.0064	.0064	.0066

Conclusions:

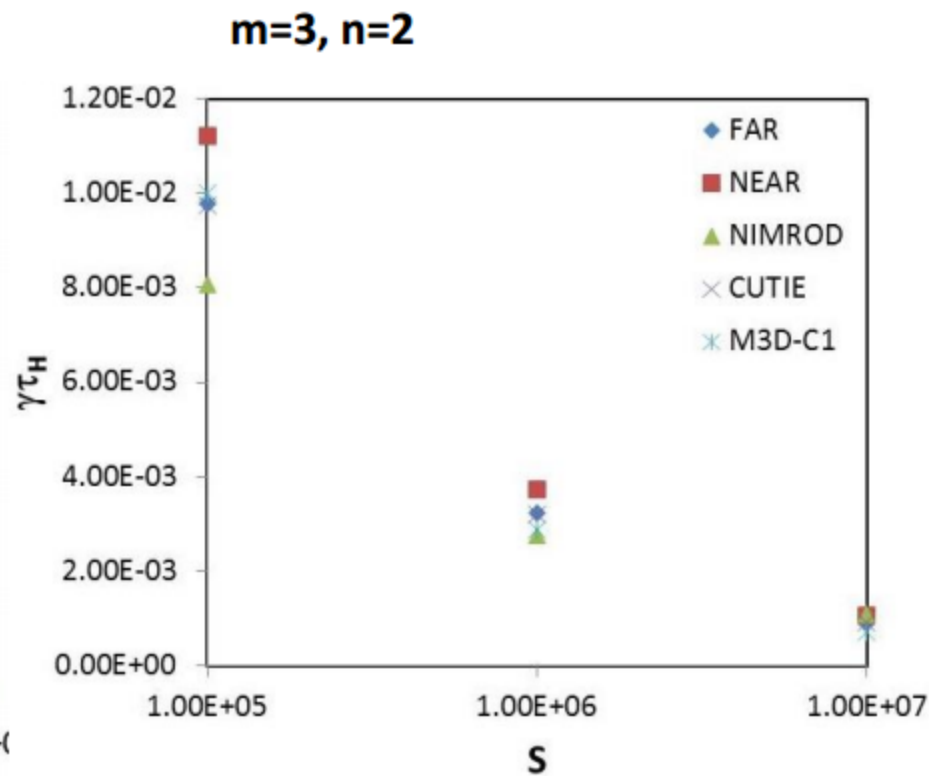
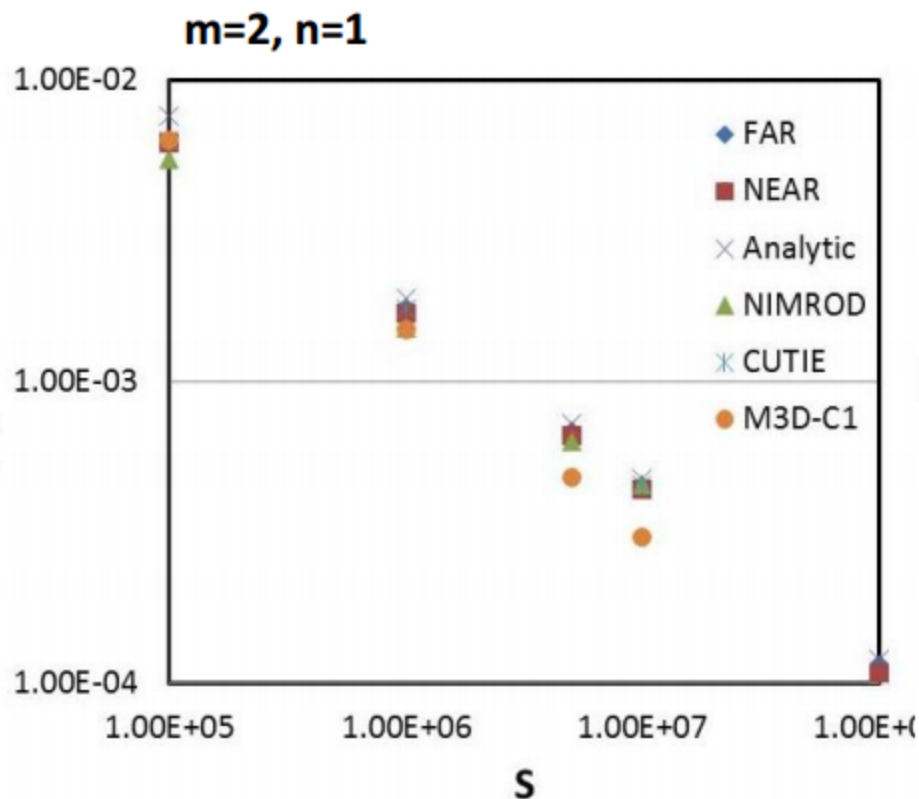
1. isplitstep=0 is most accurate, but only viable for linear runs
2. impmod = 1 (Caramana advance) more accurate than impmod=0
3. Good convergence in toroidal mode number for low-n modes (1 and 2)

For explanation of impmod=0,1, see: Ferraro and Jardin, JCP 228 (2009) 7742-7770

Linear Benchmarking



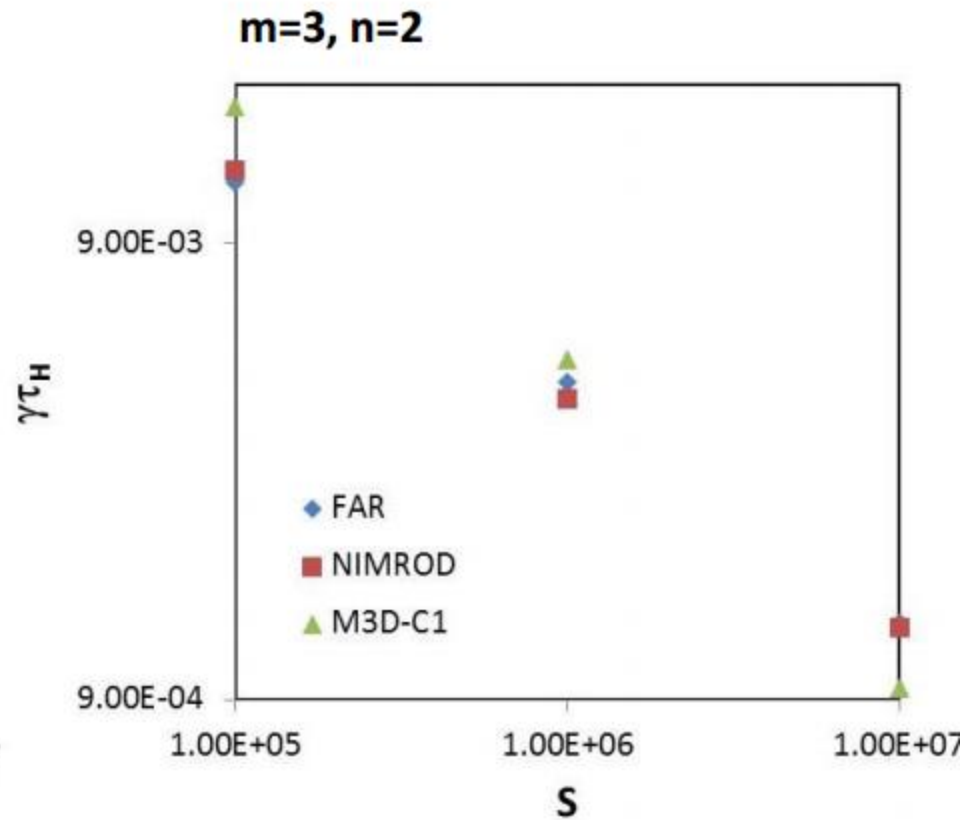
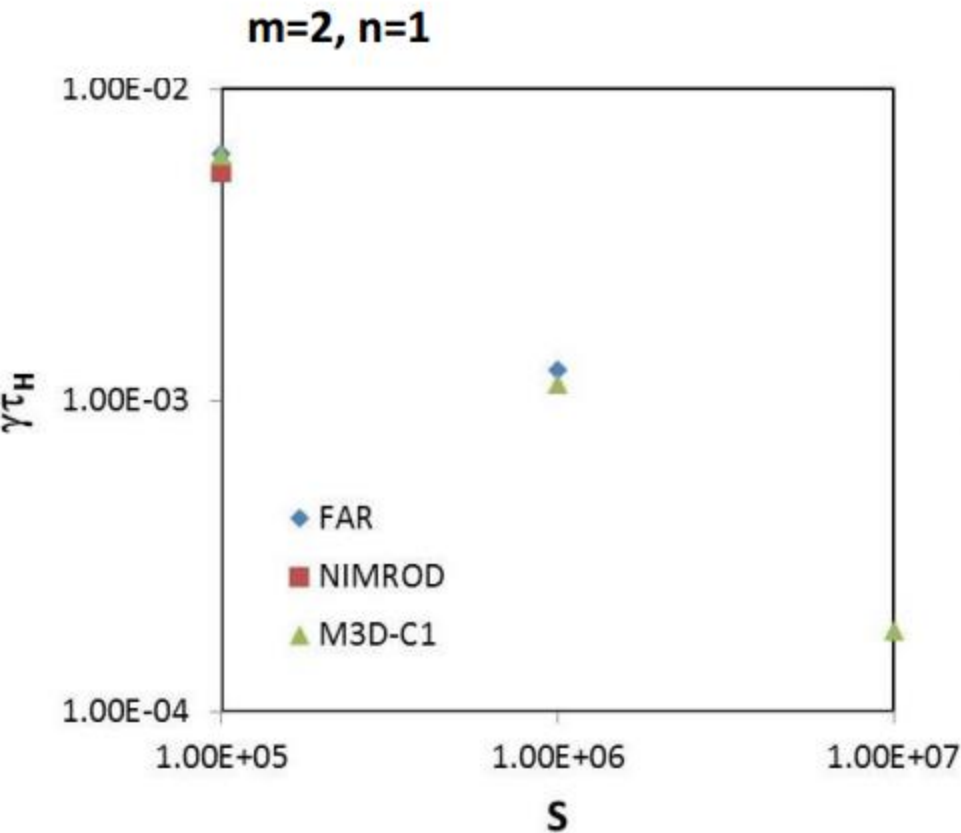
Cylindrical $q=1.33(1 + (r/0.595)^8)^{0.25}$ and $\beta=0$



Linear Benchmarking



$$R/a=10, \quad q=1.33(1 + (r/0.595)^8)^{0.25} \quad \text{and} \quad \beta_0=1.1 \times 10^{-7}$$



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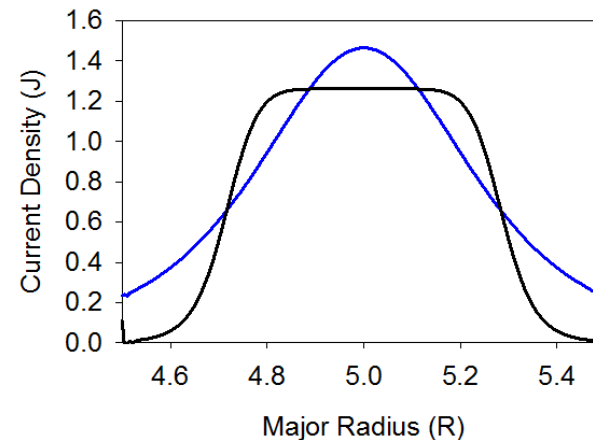
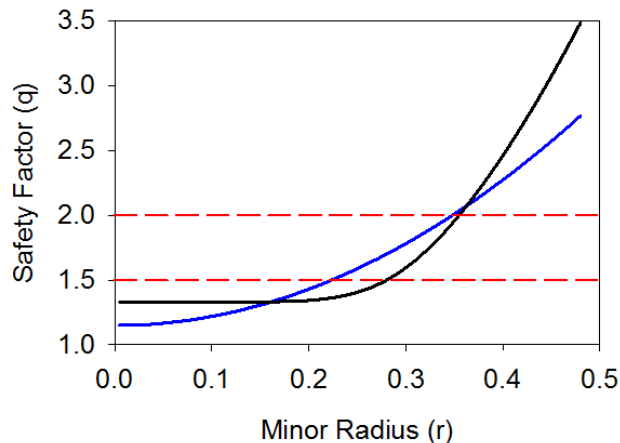
Non-Linear Runs looked at Two Equilibria

$$q(r) = 1.33 \left[1 + \left(\frac{r^2}{.354} \right)^4 \right]^{1/4}$$

- **Two-Mode** Case (same as linear)
- n=1 and n=2 unstable
- Cylinder, A=3, A=10 torus
- Linear and Nonlinear

$$q(r) = 1.15 \left[1 + \left(\frac{r^2}{.6561} \right) \right]$$

- **One-Mode** Case
- Only n=1 unstable
- Cylinder, A=3, A=10 torus
- Linear and Nonlinear



Non-Linear Results

1. Two-Mode case: Ran Nonlinear A=10 Torus numvar=3

- $S = 2 \times 10^4$ $\nu/\eta = 1$
- Ran both 32 and 16 toroidal planes (with hermit cubic elements)

In both cases, islands continue to grow and becomes totally stochastic by Time Slice 30. ($t=1500$) Very similar results for $N=16$ and $N=32$ plane cases. Error $\propto N^4$

2. One-Mode Case: Ran Cylindrical geometry case to saturation with reduced MHD

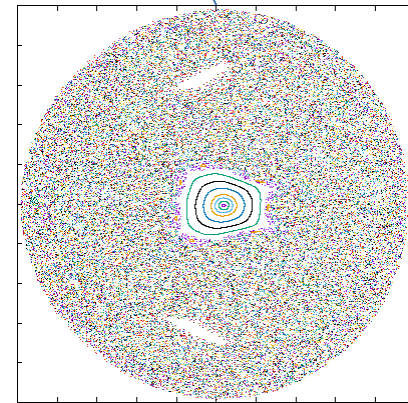
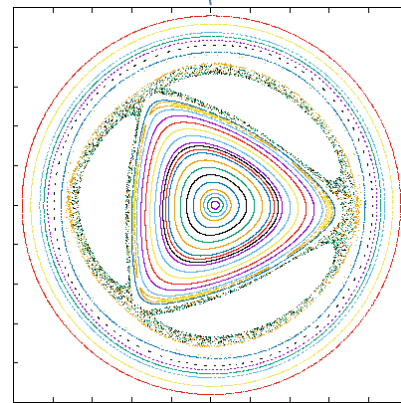
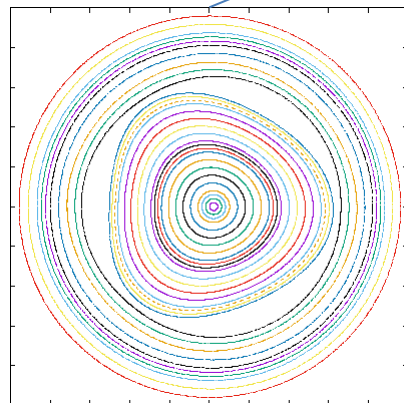
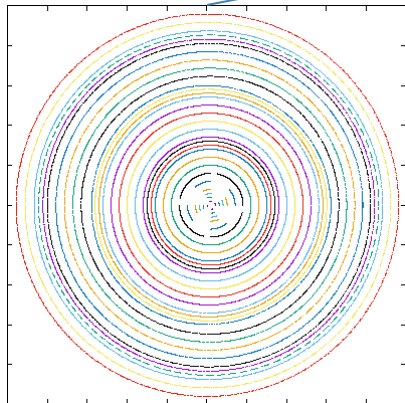
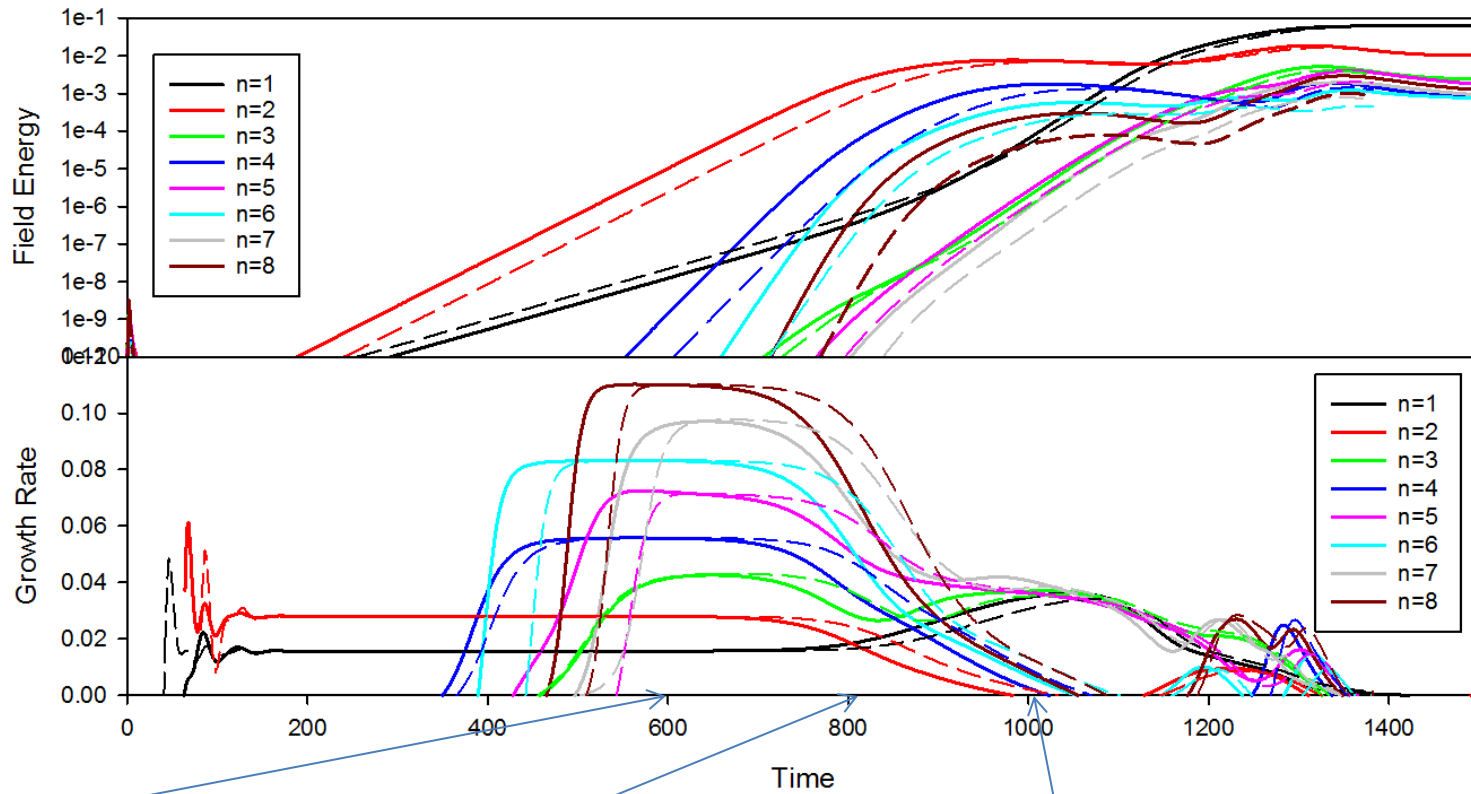
- $S = 10^5$, $\nu/\eta = .047$

Much slower island growth than above. Island saturates at about $W/a = .14$ at $t=50,000$

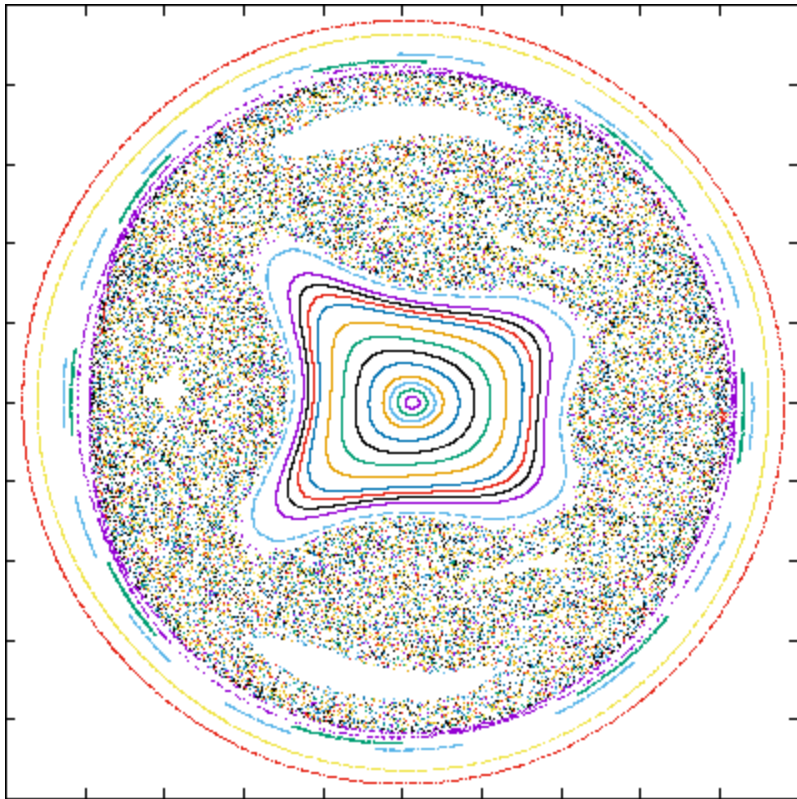
- Now running A=10 toroidal cases with full MHD

Mode Growth vs time for **Two-mode case**

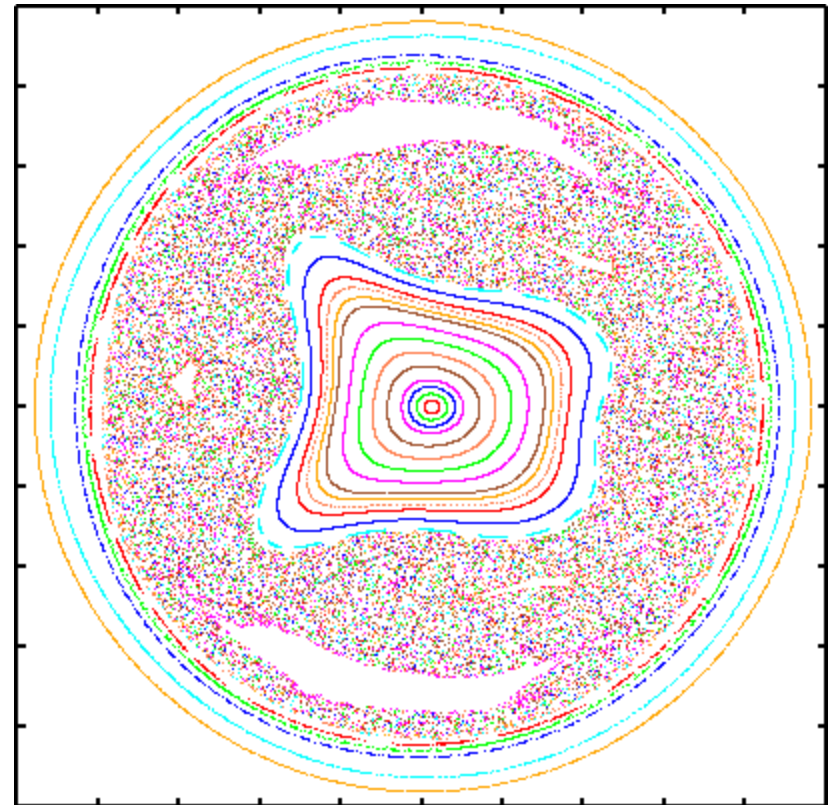
Comparison of 16 planes (solid) and 32 planes (dashed)



Comparison of surfaces at time $t=1200$ for **Two-Mode case** with 16 and 32 Planes

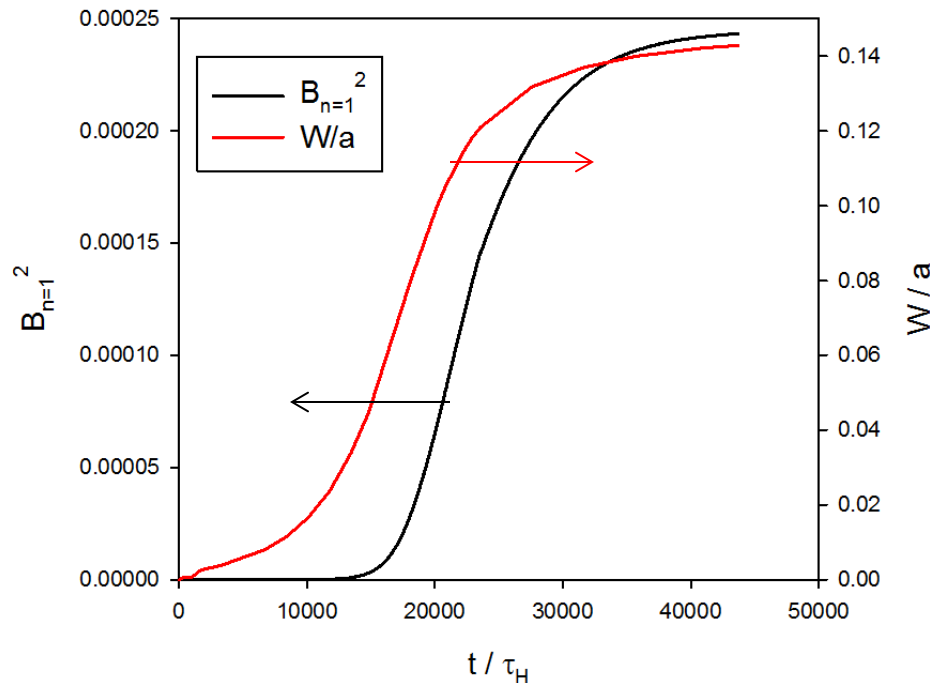


16 Planes

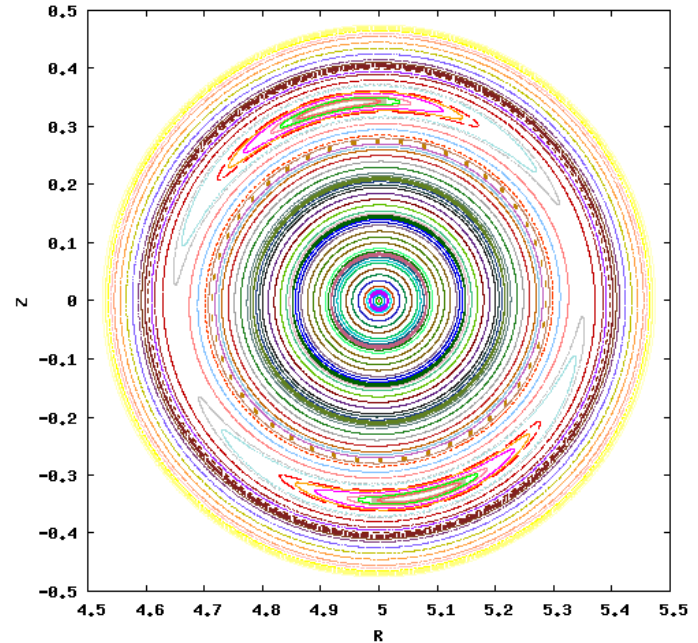


32 Planes

Single Unstable Mode Case in a cylinder



Island width and magnetic energy in $n=1$ harmonic vs time.



Poincaré plot at final time.

$$q(r) = 1.15 \left[1 + \left(\frac{r}{0.81} \right)^2 \right]$$

SI Units: $R=5, a=0.5, B_T = 4.2, n_0 = 10^{20}$

$$S = \frac{a^2 B_T}{\eta R} \left[\frac{\mu_0}{n_0 M_i} \right]^{1/2} = 10^5, \quad \tau_H \equiv \frac{R}{B_T} [\mu_0 n_0 M_i]^{1/2}, \quad \nu/\eta = .047$$

$N=0$ (axisymmetric) equilibrium not advanced in time.

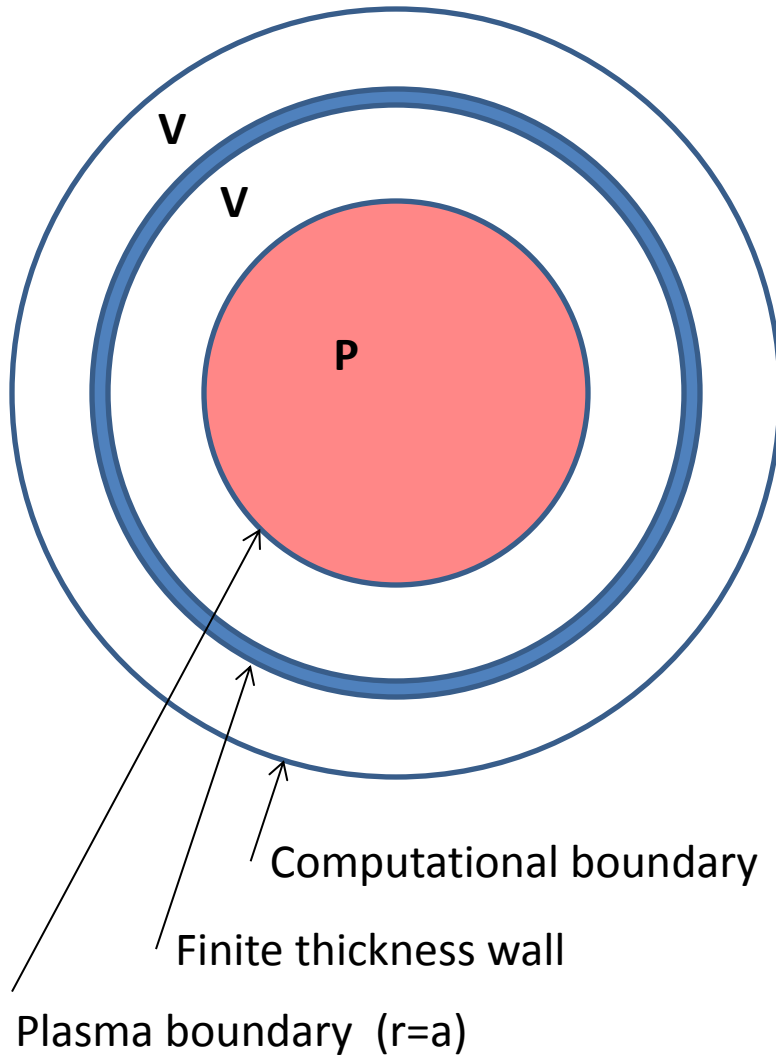
Low $\beta < 10^{-7}$

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Model Cylindrical Equilibrium

- Poloidal flux low order polynomial in r^2
- Current density vanishes at a
- $q_0 < q < q_a$



$$\Psi(r) = \begin{cases} \frac{B_T r^2}{2q_0 R} \left[1 - (r/a)^2 \left(1 - \frac{3q_0}{2q_a} \right) + \frac{1}{3} (r/a)^4 \left(1 - 2\frac{q_0}{q_a} \right) \right] & r < a \\ \frac{a^2 B_T}{6q_0 R} \left[1 + \frac{5q_0}{2q_a} + 6\frac{q_0}{q_a} \ln(r/a) \right] & r \geq a \end{cases}$$

$$J(r) = \frac{B_T}{6q_0 R} \left[1 - 4(r/a)^2 \left(1 - \frac{3q_0}{2q_a} \right) + 3(r/a)^4 \left(1 - 2\frac{q_0}{q_a} \right) \right]$$

$$q(r) = \begin{cases} q_0 \left[1 - 2(r/a)^2 \left(1 - \frac{3q_0}{2q_a} \right) + (r/a)^4 \left(1 - 2\frac{q_0}{q_a} \right) \right]^{-1} & r < a \\ q_a (r/a)^2 & r \geq a \end{cases}$$

Stationary State Equilibrium equations

Uniform density: $T \sim p$; Spitzer resistivity: $\eta \sim p^{-3/2}$

Constant loop voltage: $\eta J \sim (\eta_0 p_0^{3/2}) p^{-3/2} J \sim V_L$

$$\implies p = p_0 \left[1 - (r/a)^2 \left(4 - 6 \frac{q_0}{q_a} \right) + (r/a)^4 \left(3 - 6 \frac{q_0}{q_a} \right) \right]^{2/3} \quad V_L = B_T \eta_0 p_0^{3/2} / 6 q_0 R$$

$$\nabla p = \mathbf{J} \times \mathbf{B} \quad \implies \frac{1}{2} \frac{d}{dr^2} B_z^2 = \frac{-\frac{dp}{dr^2} + \frac{B_z^2}{(Rq)^2} \left(1 - \frac{r^2}{q} \frac{dq}{dr^2} \right)}{1 + r^2 / (Rq)^2} \quad B_z(a) = B_T$$

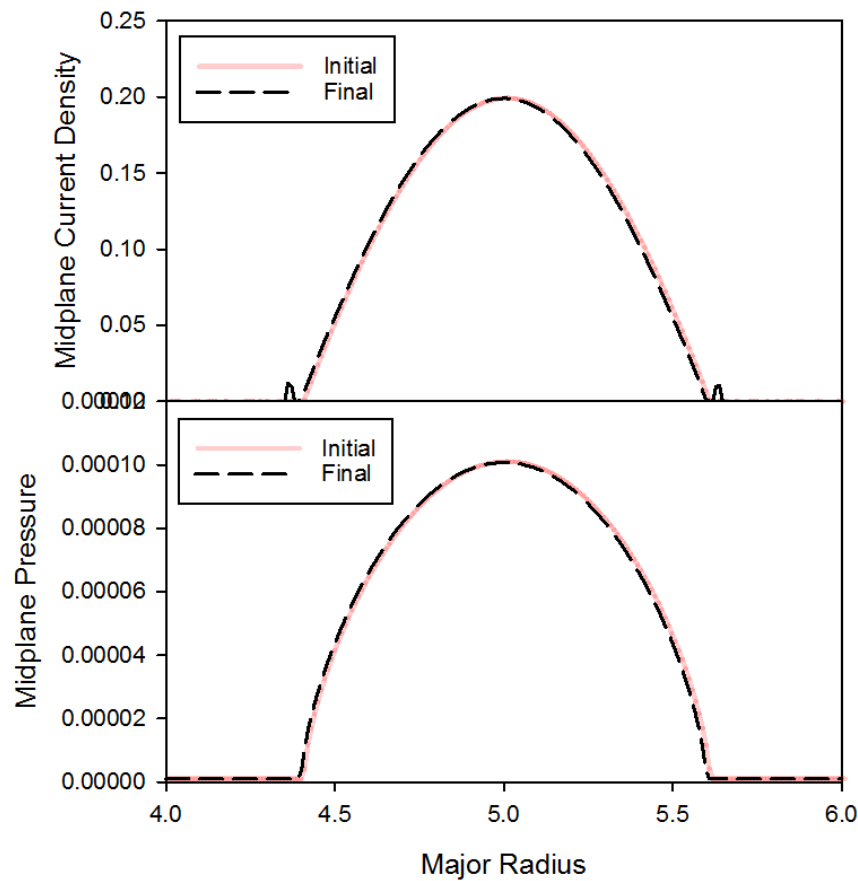
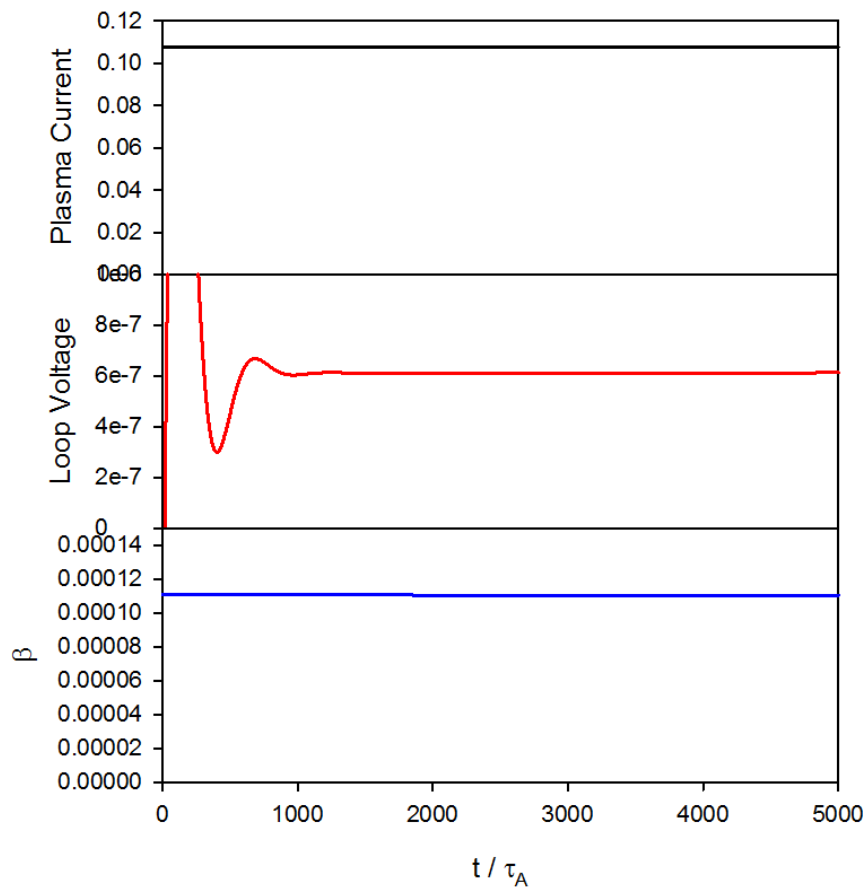
Stationary pressure profile: $\frac{1}{r} \frac{d}{dr} \left(r \kappa \frac{dp}{dr} \right) = -\eta J^2 = -V_L J$

$$\implies \kappa \frac{dp}{dr} = -\frac{V_L B_T}{6 q_0 R} \frac{1}{r} \int_0^r r \left[1 + (r/a)^2 \left(-4 + 6 \frac{q_0}{q_a} \right) + (r/a)^4 \left(3 - 6 \frac{q_0}{q_a} \right) \right] dr$$

$$\kappa = \kappa_0 \frac{\left[1 - 2(r/a)^2 \left(1 - \frac{3}{2} \frac{q_0}{q_a} \right) + (r/a)^4 \left(1 - 2 \frac{q_0}{q_a} \right) \right] \left[1 - 4(r/a)^2 \left(1 - \frac{3}{2} \frac{q_0}{q_a} \right) + 3(r/a)^4 \left(1 - 2 \frac{q_0}{q_a} \right) \right]^{1/3}}{\left[\left(1 - \frac{3}{2} \frac{q_0}{q_a} \right) - \frac{3}{2} (r/a)^2 \left(1 - 2 \frac{q_0}{q_a} \right) \right]}$$

$$\kappa_0 = \frac{3}{192} V_L B_T a^2 / q_0 R p_0 \left(1 - \frac{3}{2} \frac{q_0}{q_a} \right) \quad q_0, q_a, B_T, a, R, \eta_0, p_0 \implies V_L, \kappa_0$$

2D Nonlinear Run reaches stationary state



Linear Resistive Stability of Model Cylindrical Equilibrium with no wall (free boundary)

$$q_0 \quad q(r) = \begin{cases} q_0 \left[1 - 2(r/a)^2 \left(1 - \frac{3q_0}{2q_a} \right) + (r/a)^4 \left(1 - 2\frac{q_0}{q_a} \right) \right]^{-1} & r < a \\ q_a (r/a)^2 & r \geq a \end{cases}$$

	1.01	1.11	1.33	1.67	1.80	2.01
2.4						
2.6						
2.8						
3.0						
3.2						
3.4						
3.6						
3.8						
4.0						
4.2						
4.4						
4.6						
4.8						
5.0						
5.2						
5.4						
5.6						
5.8						

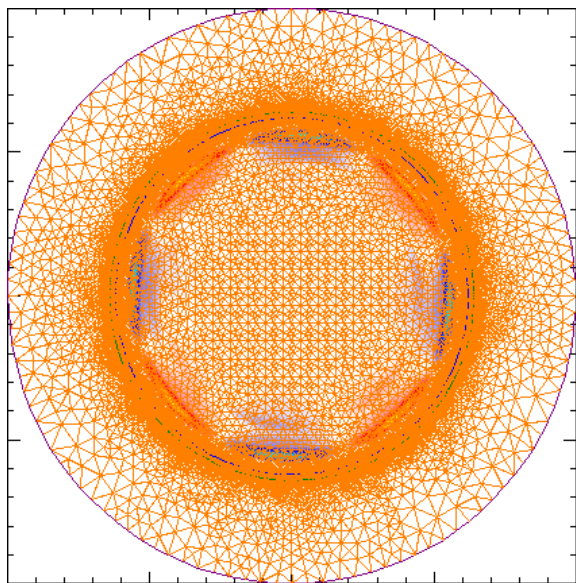
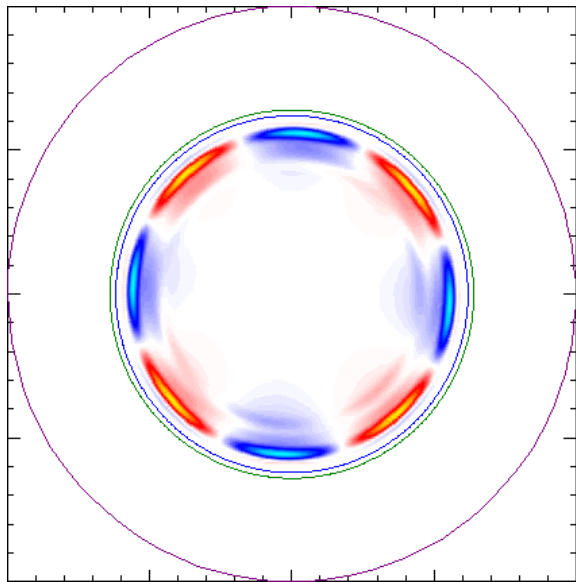
	2/1 most unstable
	3/1 most unstable
	4/1 most unstable
	Stable to all modes

monotonic solutions for:

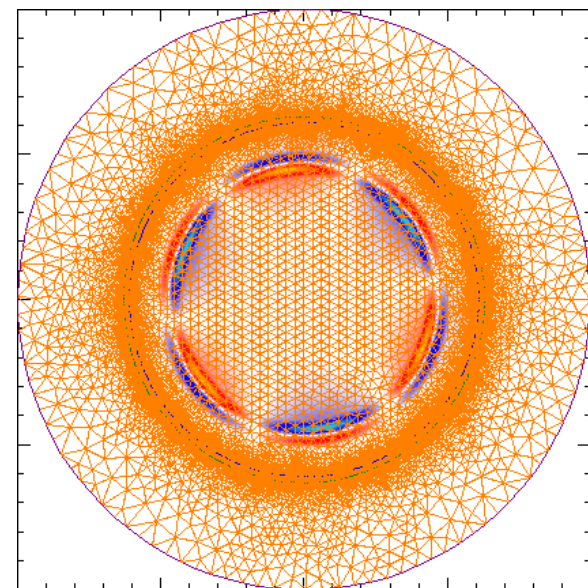
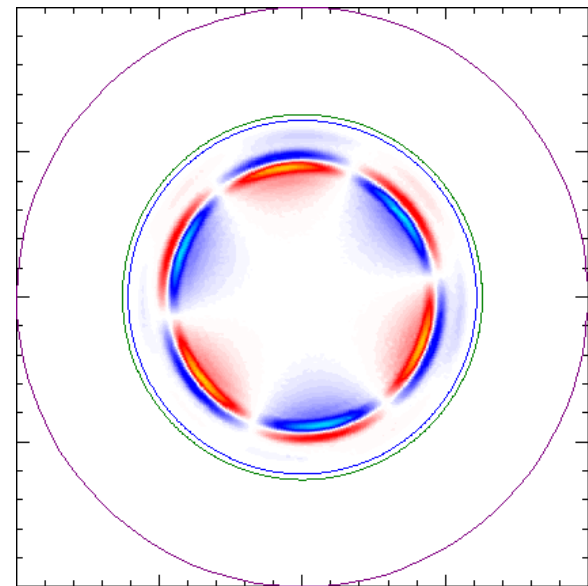
$$3q_0 > q_a > \frac{3}{2}q_0$$

Examine this case with a nearby wall

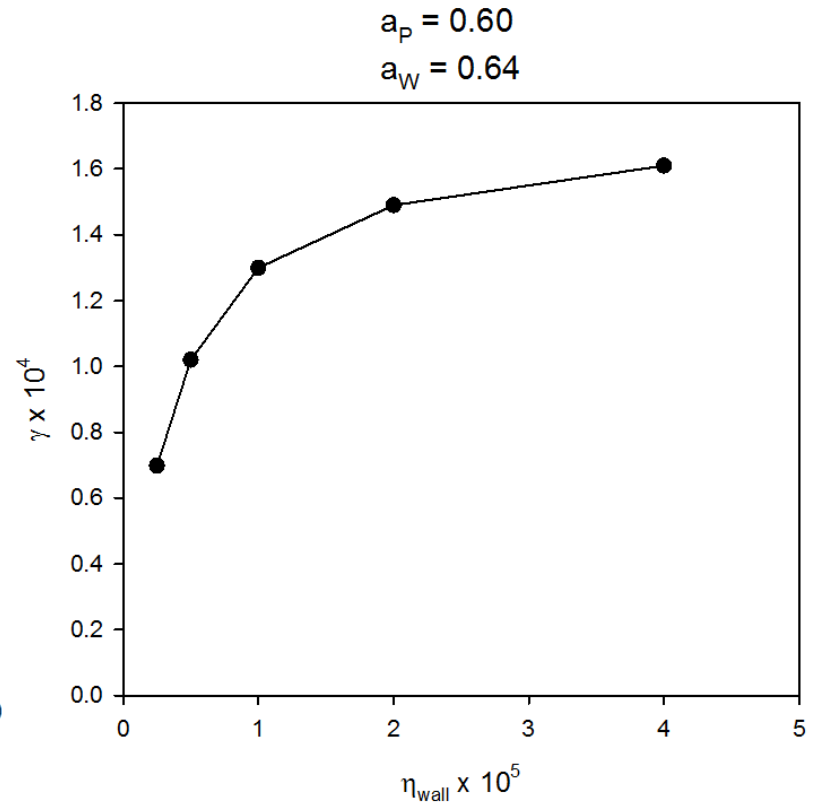
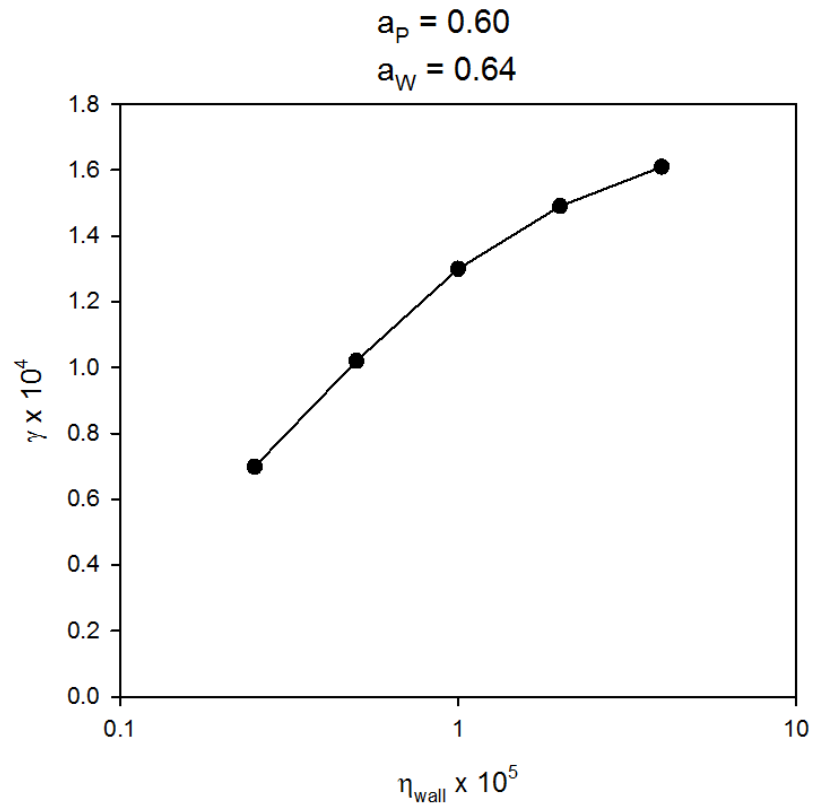
Wall at $r/a = 0.64$: unstable
(4/1) mode dominates



Wall at $r/a = 0.62$: stable!
(3/1) mode dominates



Growth rate of mode depends on wall resistivity



Next Steps

- Add sheared toroidal velocity with torque input (also to equilibrium equation)
- Add error field
- Study interaction of island with resistive wall and error field, and compare with theoretical results for mode locking
- Extend to Torus

References:

- [1] Fitzpatrick, R., Nuclear Fusion **33** (1993) p. 1049, Section (3)
- [2] Fitzpatrick, R., Nuclear Fusion **33** (1993) p. 1049, Section (6)
- [3] Waelbroeck, F. and Fitzpatrick, R. Phys. Rev. Lett. **78** (1997) p. 1703
- [4] Fitzpatrick, R., Plasma parameter scaling of the error-field penetration threshold in tokamaks", Phys. Plasmas, 10 (2003) , p. 1782
- [5] Fitzpatrick, R. and Waelbroeck, F. Phys. Plasma. 15, (2008) 012502
- [6] Fitzpatrick, R. Plasma Phys. Control. Fusion 54 ([2012](#)) [094002](#).
- [7] Reimerdes, H. et al. Nucl. Fusion 49 (2009) 115001
- [8] De Bock, M.F.M. et al Nucl. Fusion 48 (2008) 015007 .