

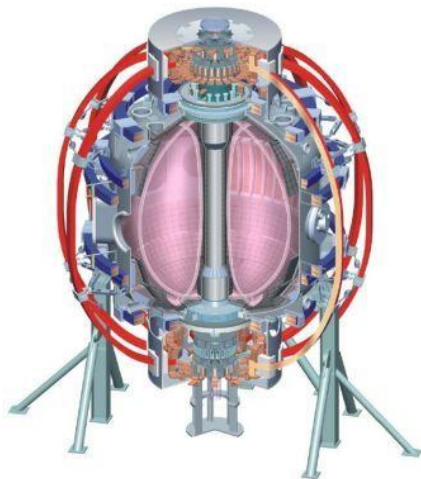
# Simulation of microtearing turbulence in NSTX and scaling with collisionality

Walter Guttenfelder<sup>1</sup>

J. Candy<sup>2</sup>, S.M. Kaye<sup>1</sup>, W.M. Nevins<sup>3</sup>, E. Wang<sup>3</sup>, J. Zhang<sup>4</sup>,  
 R.E. Bell<sup>1</sup>, N. Crocker<sup>4</sup>, B.P. LeBlanc<sup>1</sup>, G.W. Hammett<sup>1</sup>,  
 D.R. Mikkelsen<sup>1</sup>, Y. Ren<sup>1</sup>, H. Yuh<sup>5</sup>

<sup>1</sup>PPPL, <sup>2</sup>General Atomics, <sup>3</sup>LLNL, <sup>4</sup>UCLA, <sup>5</sup>Nova Photonics Inc.

**APS-DPP, Salt Lake City**  
**Nov. 14-18, 2011**



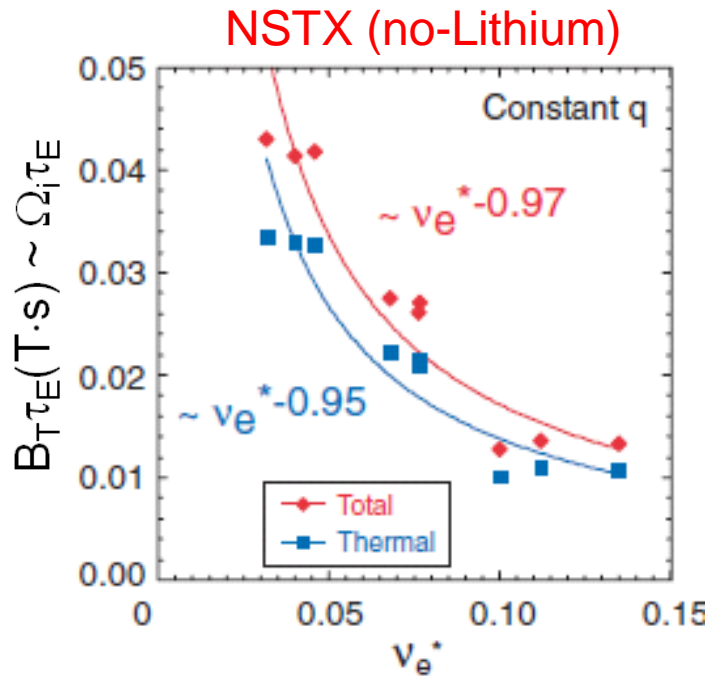
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 NFRI  
 KAIST  
 POSTECH  
 ASIPP  
 ENEA, Frascati  
 CEA, Cadarache  
 IPP, Jülich  
 IPP, Garching  
 ASCR, Czech Rep

# Overview

- Experimental motivation - favourable  $\Omega_i \tau_{E,th} \sim v_*^{-(0.8-0.95)}$  dependence in STs
  - Microtearing modes found to be unstable in experimental  $v_*$  scans,  $\gamma_{lin} \sim v_e$
- Linear microtearing properties for high- $\beta$  NSTX discharges
  - Electromagnetic, electron drift mode with narrow resonant current layer,  $\Delta_j < \rho_s$
  - Non-monotonic dependence on  $v_e^{e_i/\omega}$ , threshold in  $\beta_e$  and  $\nabla T_e$
  - $Z_{eff}$  (and  $s/q$ ) scaling distinct from ETG
- Non-linear simulations
  - Necessary to “resolve” (distinguish) each simulated rational surface
  - Transport is experimentally significant and dominated by magnetic “flutter”
- Scaling of non-linear transport
  - Predicted  $\chi_{e,sim} \sim v_e^{1.1}$  close to experimental trend
  - “Stiff” with  $\nabla T_e$  but suppressible by experimental levels of  $E \times B$  shear

# Experimental motivation - strong collisionality scaling in STs



← NSTX\* (Kaye et al., Nucl. Fusion 2007)

$$\Omega \tau_E^{\text{th}} \sim \nu_{*e}^{-0.95}$$

MAST (Valovič et al., Nucl. Fusion 2011)

$$\Omega \tau_E^{\text{th}} \sim \nu_{*e}^{-0.82}$$

ITER (PIPB, Doyle et al., Nucl. Fusion 2007)

$$\Omega \tau_E^{\text{th},04(2)} \sim \nu_{*e}^{-0.2}$$

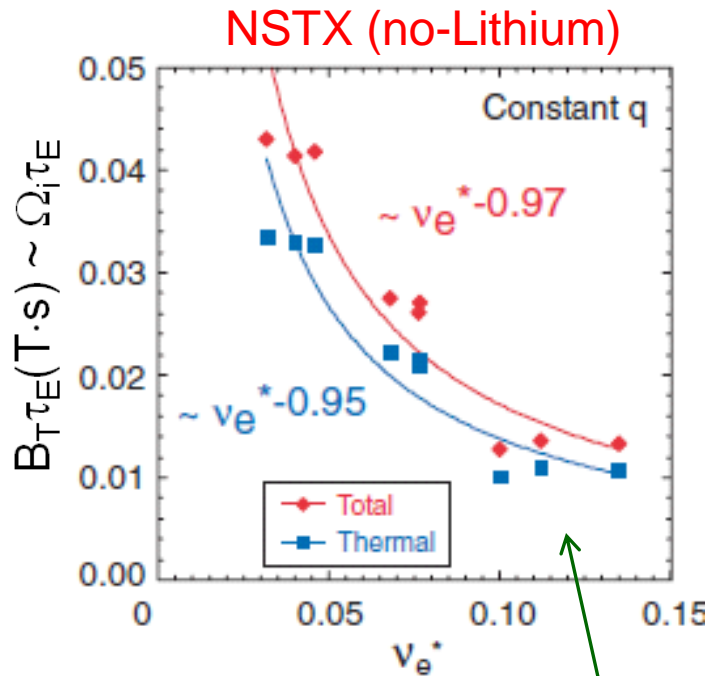
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S. Kaye, PP-9.30 (Thurs. PM)

S. Gerhardt, YI-2.2 (Fri. AM)

- Ion transport is neoclassical, consistent with strong toroidal flow and flow shear
- What is the cause of anomalous electron thermal transport?
- Will favorable  $\tau_E$  scaling hold at lower  $\nu_{*e}$  envisioned for next generation ST (high heat flux, CTF, ...)?

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**Following simulations based on a single NSTX high- $v_{*}$  discharge**  
 $B_T=0.35\text{T}$ ,  $I_p=0.7\text{ MA}$ ,  $P_{\text{NBI}}=4\text{ MW}$ ,  $n_e \approx 6 \times 10^{19}\text{ m}^{-3}$ ,  $T_e(0) \sim 1\text{ keV}$

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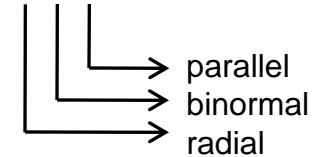
# GYRO\* used for gyrokinetic simulations

- Eulerian solver of gyrokinetic-Maxwell equations, evolving  $\delta f(E, \lambda, r, \alpha, \theta)$

$$\delta f \sim \delta \hat{f}(r, \theta) e^{-in\alpha} \quad \alpha = \phi + v(r, \theta) \approx \phi - q(r)\theta$$

$$\underline{k_\theta \doteq \frac{nq}{r}}$$

High-aspect ratio,  
low  $\beta$  limit



- Kinetic ions (D+C) and electrons, general equilibrium
- Fully collisional & electromagnetic ( $\delta A_{||}$ ,  $\delta B_{||}$ ) (both important in NBI heated ST)
- Freedom to include toroidal flow and flow shear (important in NBI heated ST)
- Can use experimental profile variations,  $T(r)$ ,  $n(r)$ ,  $q(r)$ , etc... (likely important in ST,  $\rho_s/a \sim 1/100$ ,  $\rho_s/L \sim 1/40$ )

\*J. Candy & R.E. Waltz, Phys. Rev. Lett. **91**, 045001 (2003); J. Comp. Physics **186**, 545 (2003); <https://fusion.gat.com/theory/Gyro>  
J. Candy, Phys. Plasmas Control. Fusion **51**, 105009 (2009); E.A. Belli & J. Candy, Phys. Plasmas **17**, 112314 (2010).

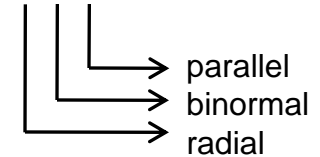
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- All following linear calculations performed in the local flux-tube limit (periodic BC's) without toroidal flow & shear*

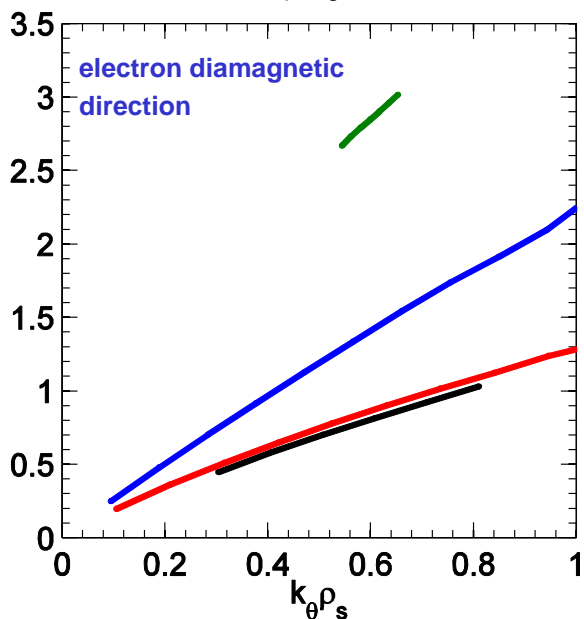
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# Microtearing modes found to be unstable in many high $v_*$ discharges

- Microtearing dominates over  $r/a=0.5-0.8$ ,  $k_\theta \rho_s < 1$  ( $n \approx 5-70$ )
- Real frequencies in electron diamagnetic direction,  $\omega \approx \omega_{*e} = (k_\theta \rho_s) \cdot (a/L_n + a/L_{Te}) \cdot (c_s/a)$
- ETG mostly stable due to larger  $Z_{\text{eff}} \approx 3$ ,  $(R/L_{Te})_{\text{crit,ETG}} \sim (1 + Z_{\text{eff}} T_e/T_i)$

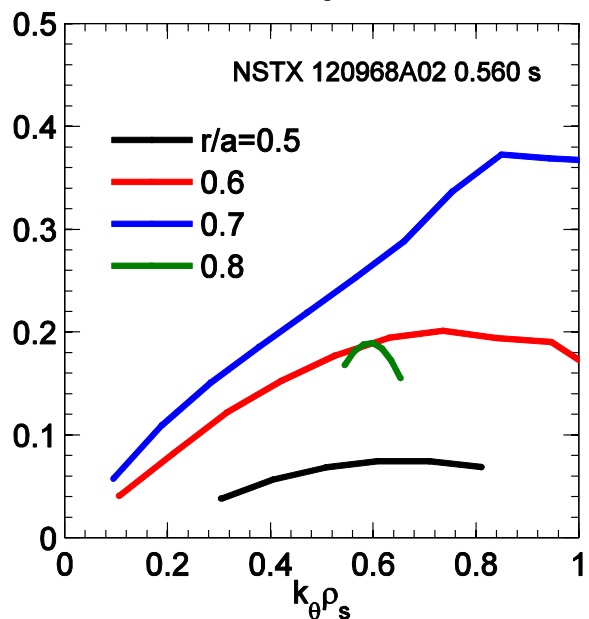
real frequencies

$$\omega_r (c_s/a)$$

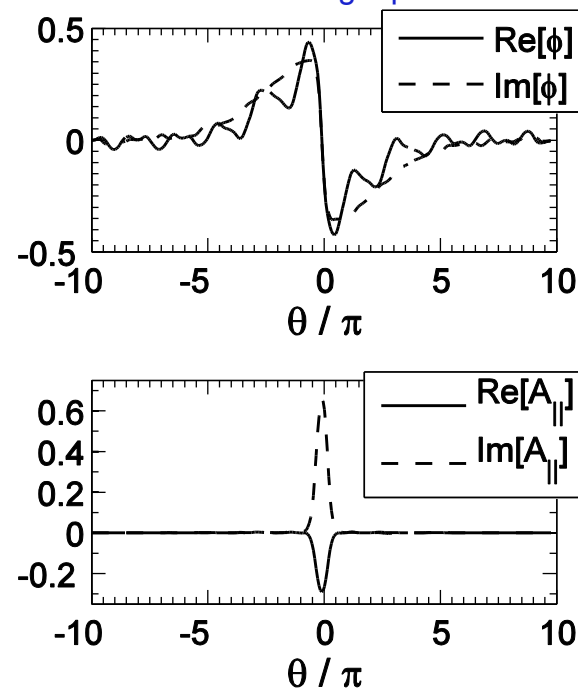


growth rates

$$\gamma (c_s/a)$$



eigenfunctions in "ballooning" space

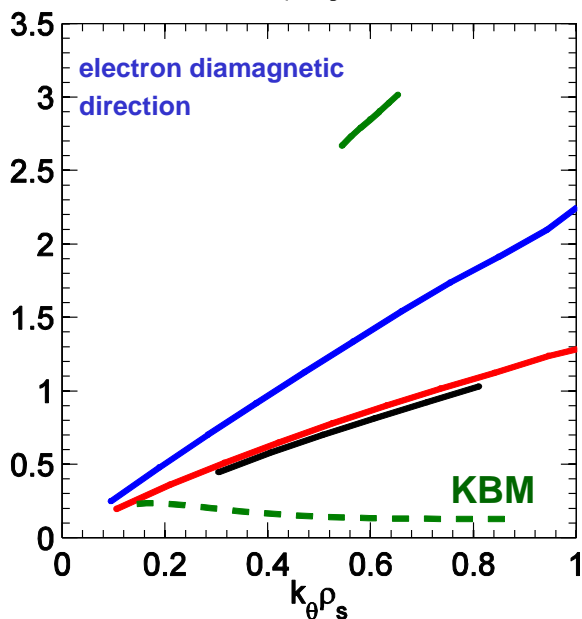


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- KBM competes farther out ( $r/a \geq 0.8$ ) where  $\alpha_{\text{MHD}} = -q^2 R \beta'$  much larger (larger  $q$ ,  $a/L_n$ )

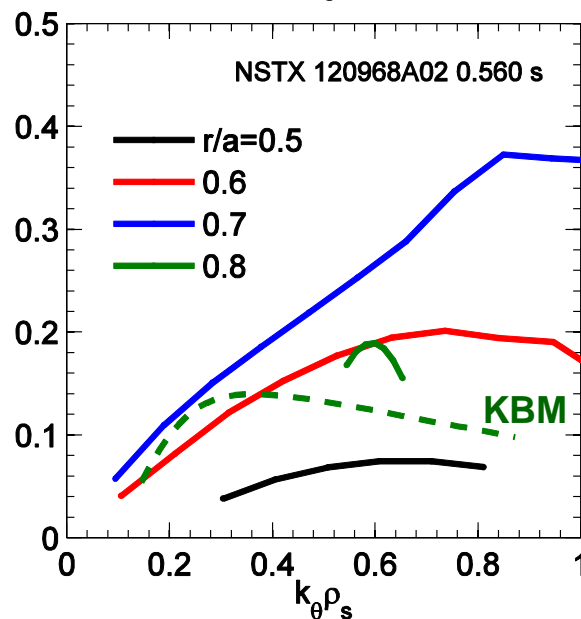
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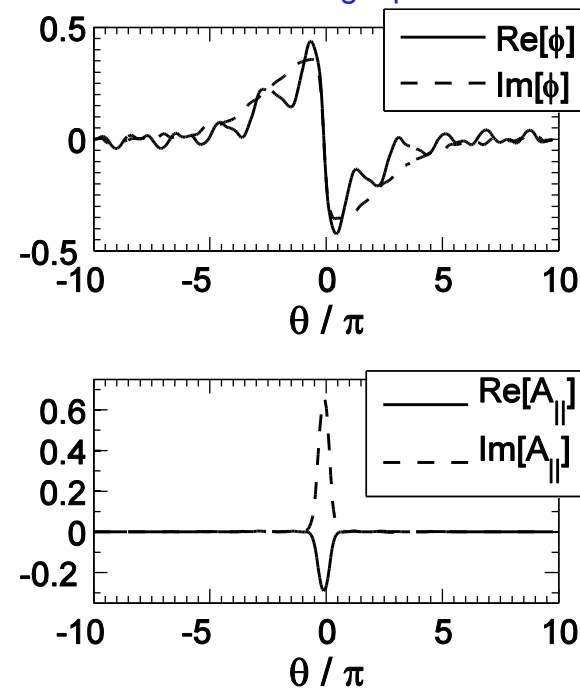


growth rates

$$\gamma (c_s/a)$$



eigenfunctions in "ballooning" space



Following calculations mostly for  $r/a=0.6$



# Linear microtearing instability

- High- $m$  tearing mode around a rational  $q(r_0)=m/n$  surface ( $k_{\parallel}(r_0)=0$ )  
(Classical tearing mode stable for large  $m$ ,  $\Delta' \approx -2m/r < 0$ )
- Driven by  $\nabla T_e$  with time-dependent parallel thermal force\*  $\Rightarrow$  requires e-i collisions

## Conceptual linear picture

- Imagine helically resonant ( $q=m/n$ )  $\delta B_r$  perturbation  $\delta B_r \sim \cos(m\theta - n\phi)$
- $\delta B_r$  leads to radially perturbed field line, finite island width  $w = 4 \left( \frac{\delta B_r}{B} \frac{rR}{n\hat{s}} \right)^{1/2}$
- $\nabla T_e$  projected onto field line gives parallel gradient  $\nabla_{\parallel} T_{e0} = \frac{\vec{B} \cdot \nabla T_{e0}}{B} = \frac{\delta B_r}{B} \nabla T_{e0}$
- Parallel thermal force ( $R_{T\parallel} \sim -\alpha(\omega) n_e \nabla_{\parallel} T_e$ ) drives parallel electron current that reinforces  $\delta B_r$  via Amperes's law  $k_{\perp}^2 \rho_s^2 \hat{A}_{\parallel} = \frac{\beta_e}{2} \hat{j}_{\parallel}$ ,  $B_r = ik_{\theta} A_{\parallel}$
- **Instability requires sufficient  $\nabla T_e$ ,  $\beta_e$ ,  $v_e$ , and time dependence ( $\omega$ ) important**

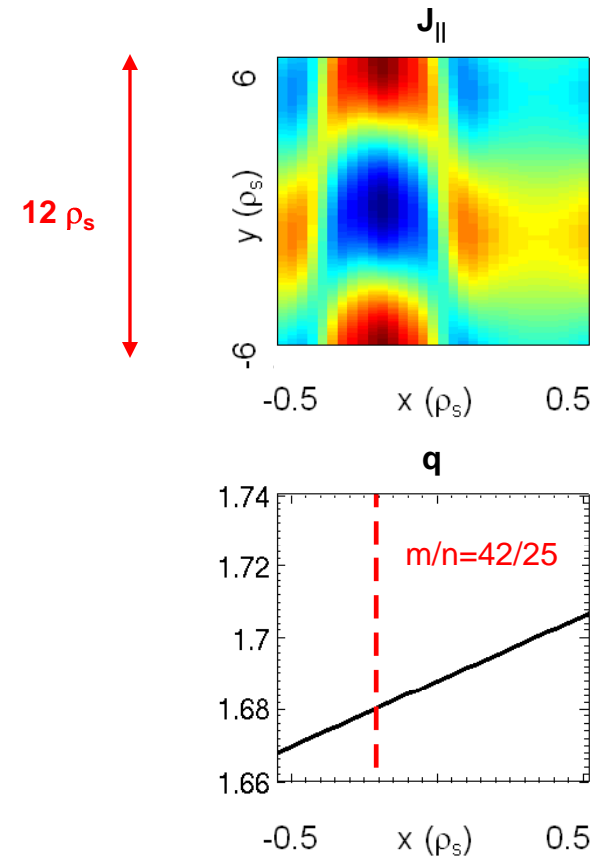
\*e.g. Hazeltine et al., Phys. Fluids 18, 1778 (1975); Gladd et al., Phys. Fluids 23, 1182 (1980);

D'Ippolito et al., Phys. Fluids 23, 771 (1980); M. Rosenberg et al., Phys. Fluids 23, 2022 (1980).

# Linear mode structure in perpendicular plane illustrates key microtearing mode features

- Narrow resonant current channel ( $\approx 0.3\rho_s \approx 1.4$  mm) centered on rational surface

x-y perpendicular plane ( $\theta=0$ )

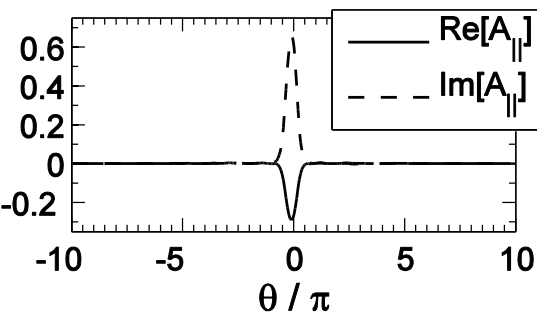


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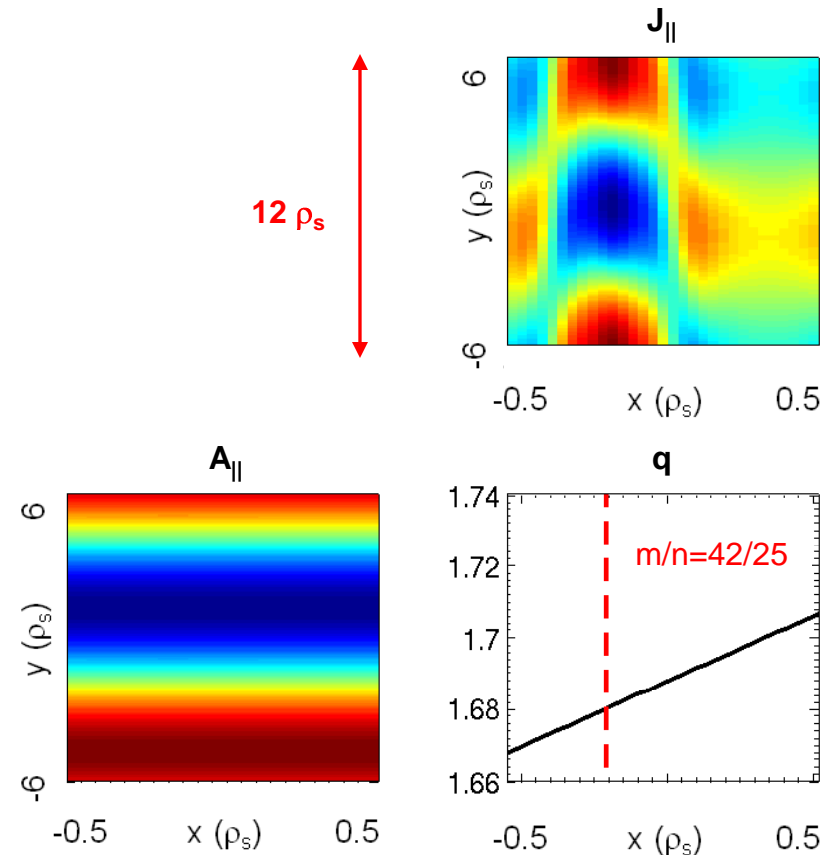
- Narrow resonant current channel ( $\approx 0.3\rho_s \approx 1.4$  mm) centered on rational surface
- Finite  $\langle A_{\parallel} \rangle_{\theta}$  (resonant tearing parity), strongly ballooning

“ballooning” space

$$k_{\perp}(\theta) = \hat{s} k_{\theta}(\theta - \theta_0)$$



x-y perpendicular plane ( $\theta=0$ )

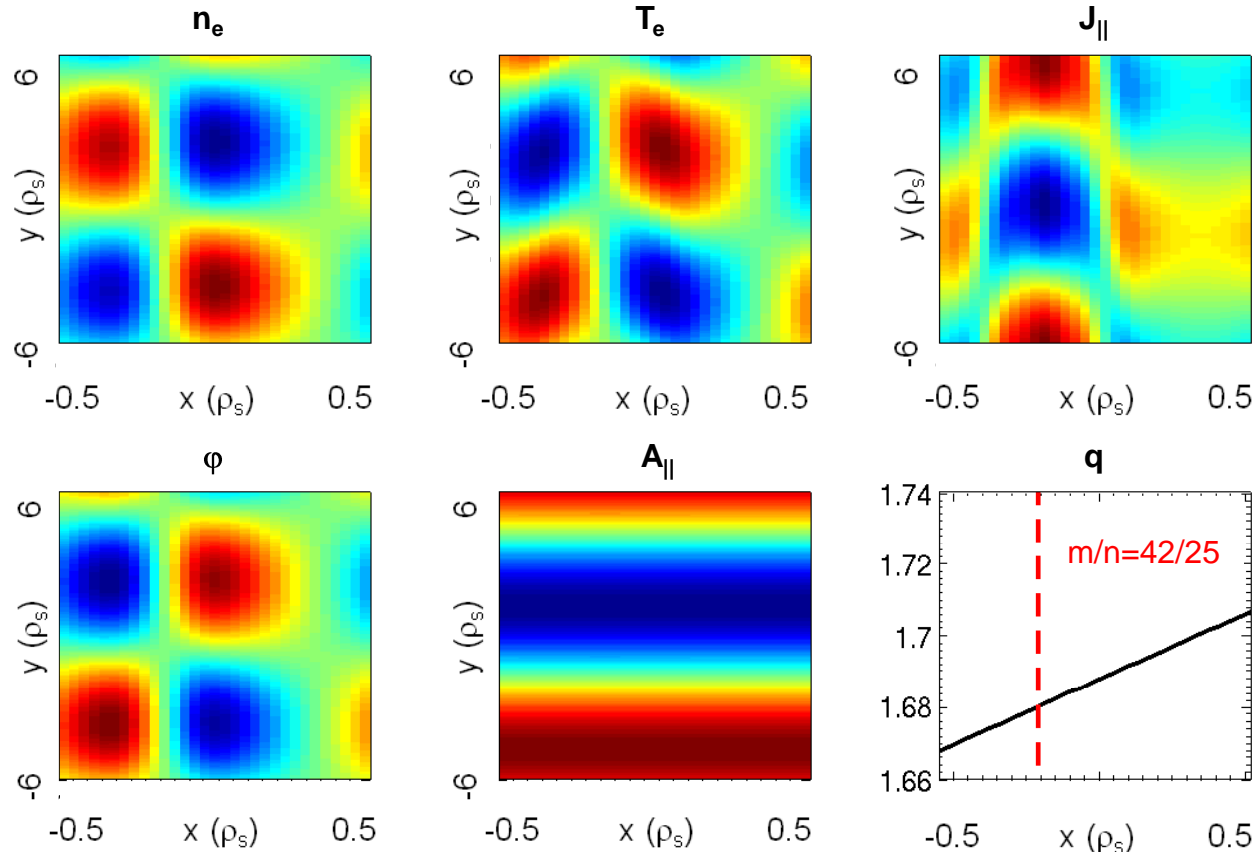
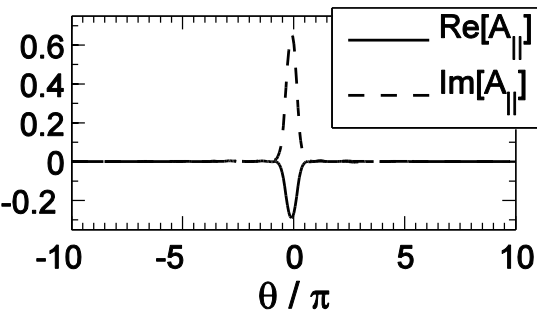
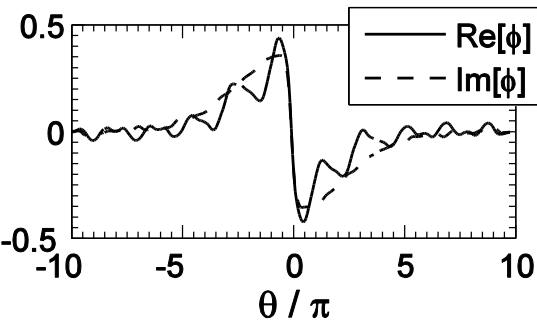


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- Finite  $\langle A_{\parallel} \rangle_{\theta}$  (resonant tearing parity), strongly ballooning
- Narrow  $n_e$  &  $T_e$  perturbations
- Nearly unmagnetized/adiabatic ion response  $\Rightarrow \frac{\tilde{n}}{n_0} \approx -Z_{\text{eff}} \left( \frac{e\tilde{\phi}}{T_i} \right)$   
x-y perpendicular plane ( $\theta=0$ )

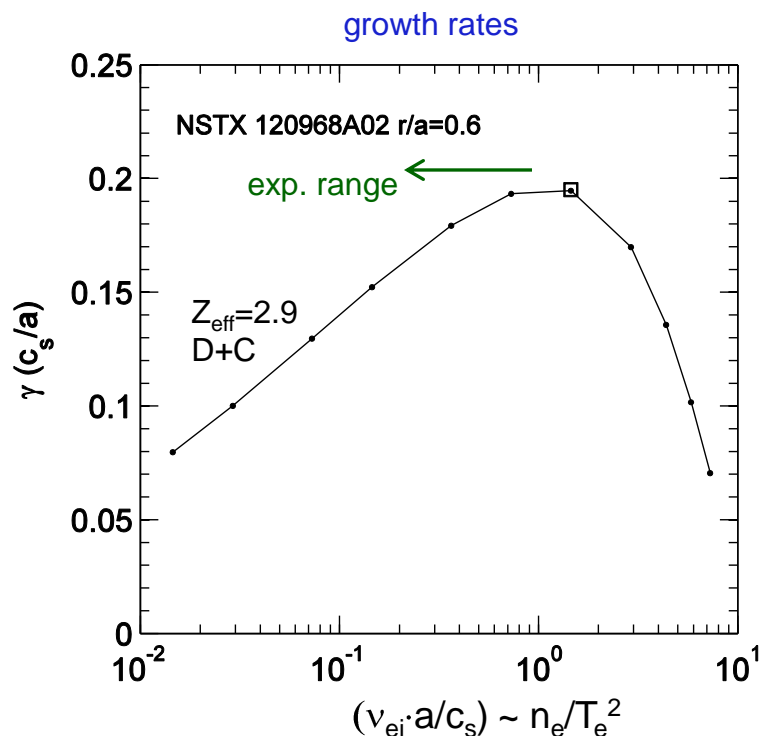
“ballooning” space

$$k_r(\theta) = \hat{s}k_{\theta}(\theta - \theta_0)$$



# A distinguishing feature of the microtearing mode is the non-monotonic dependence on $v^{e/i}/\omega$

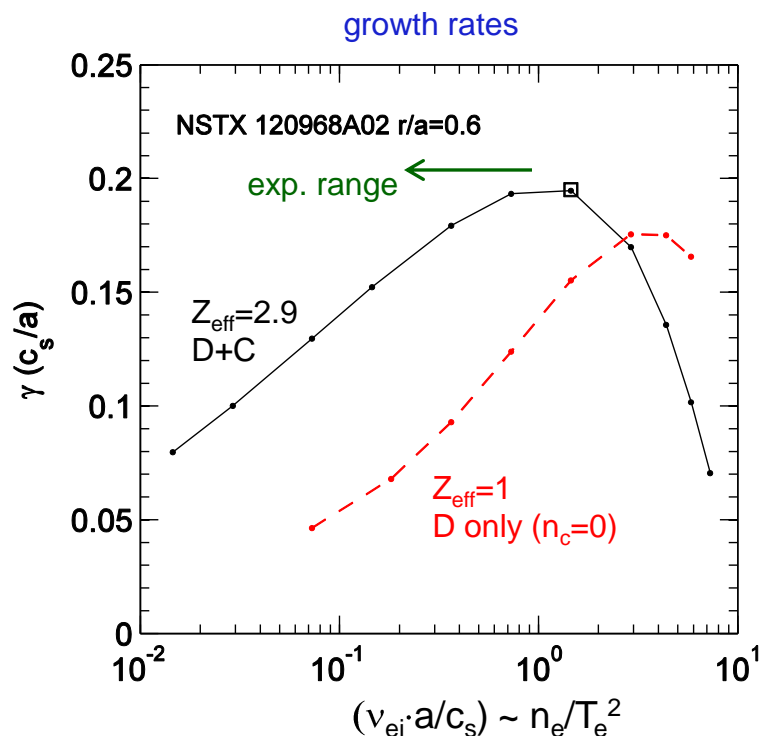
- Peak  $\gamma$  occurs for  $v^{e/i}/\omega = Z_{\text{eff}} \cdot v_{ei}/\omega \sim 1-6$ , similar to slab calculations (Gladd et al., 1980)
- $\gamma$  decreases with  $v_e$  in experimental range, qualitatively consistent with confinement scaling



$$v^{e/i} = Z_{\text{eff}} v_{ei} \propto Z_{\text{eff}} \frac{n_e}{T_e^{3/2}}$$

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- $\gamma$  decreases with  $v_e$  in experimental range, qualitatively consistent with confinement scaling
- In addition to shifting peak in  $v^{e/i}/\omega$ ,  $Z_{\text{eff}}$  can enhance instability through shielding potential (from adiabatic ion response,  $\delta n_i \sim -Z_{\text{eff}} \delta \phi / T_i$ )



$$v^{e/i} = Z_{\text{eff}} v_{ei} \propto Z_{\text{eff}} \frac{n_e}{T_e^{3/2}}$$

[Jenko et al. \(2001\)](#)

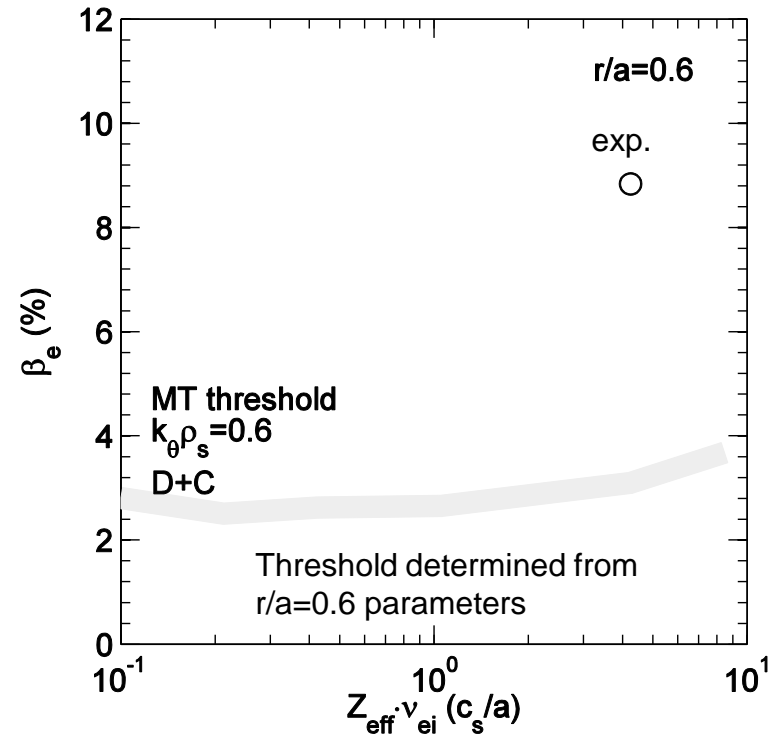
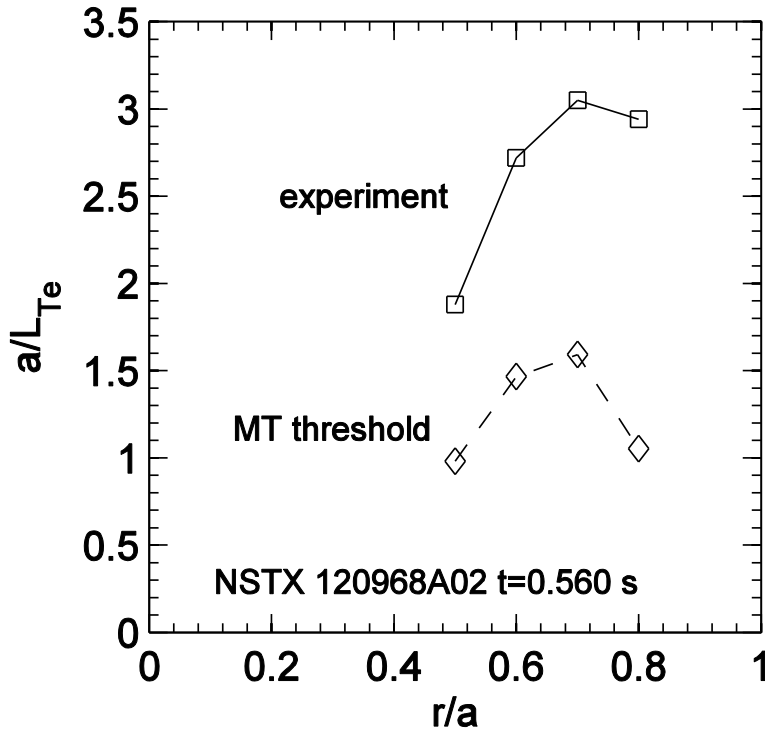
- $Z_{\text{eff}}$  (and  $s/q$ )\* dependence opposite to ETG expectations

$$\left( \frac{R}{L_{Te}} \right)_{\text{crit}}^{\text{ETG}} \sim \left( 1 + Z_{\text{eff}} \frac{T_e}{T_i} \right) \left( 1.3 + 1.9 \frac{s}{q} \right) (\dots)$$

\* Guttenfelder et al., *Scaling of linear microtearing stability for a high collisionality NSTX discharge*, submitted to Phys. Plasmas (Oct, 2011)

# Microtearing instability exhibits thresholds in electron temperature gradient and beta

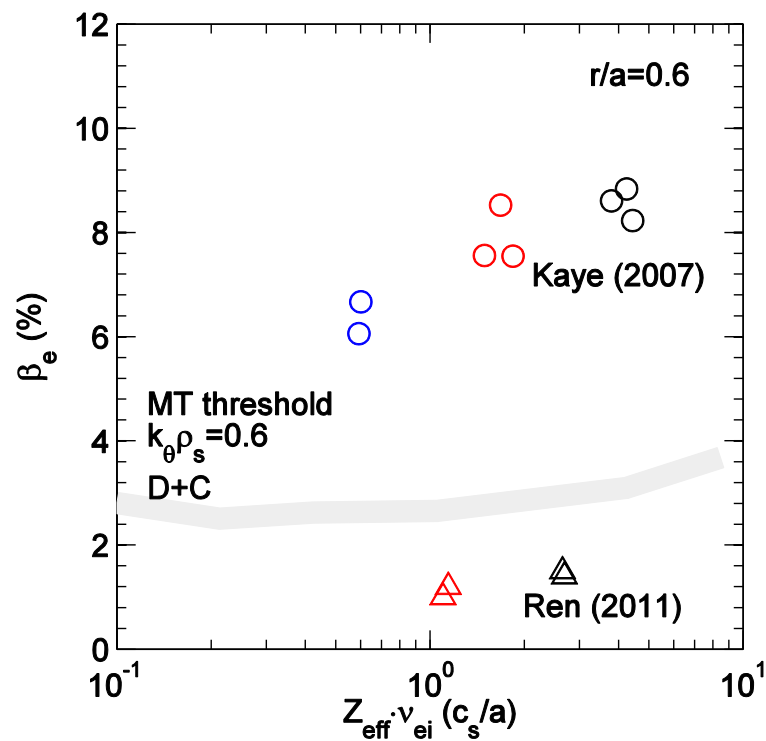
- In this high- $v_*$  discharge,  $a/L_{Te}$  and  $\beta_e$  are 2-3 $\times$  larger than linear thresholds



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# NSTX has studied electron transport for a range of beta and collisionality

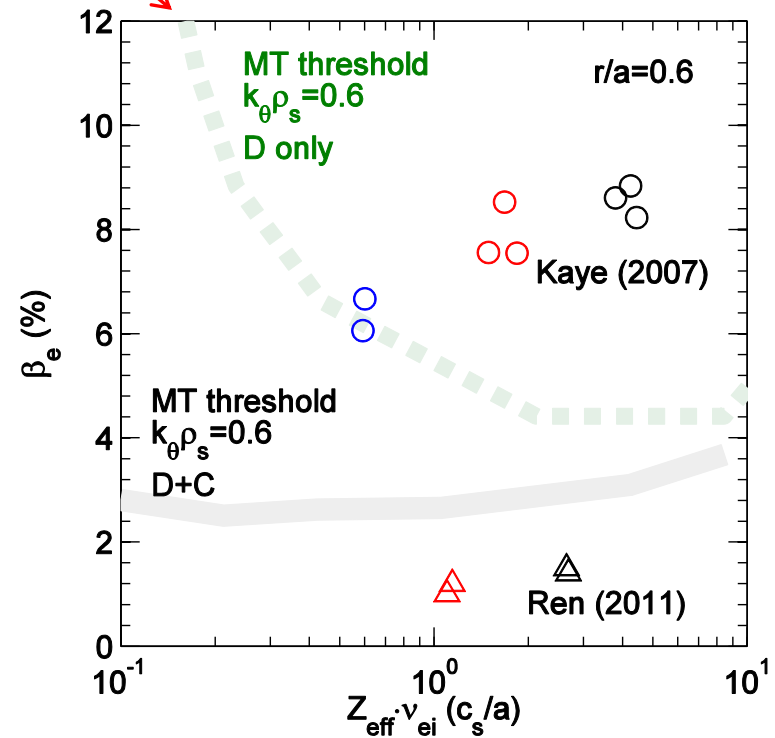
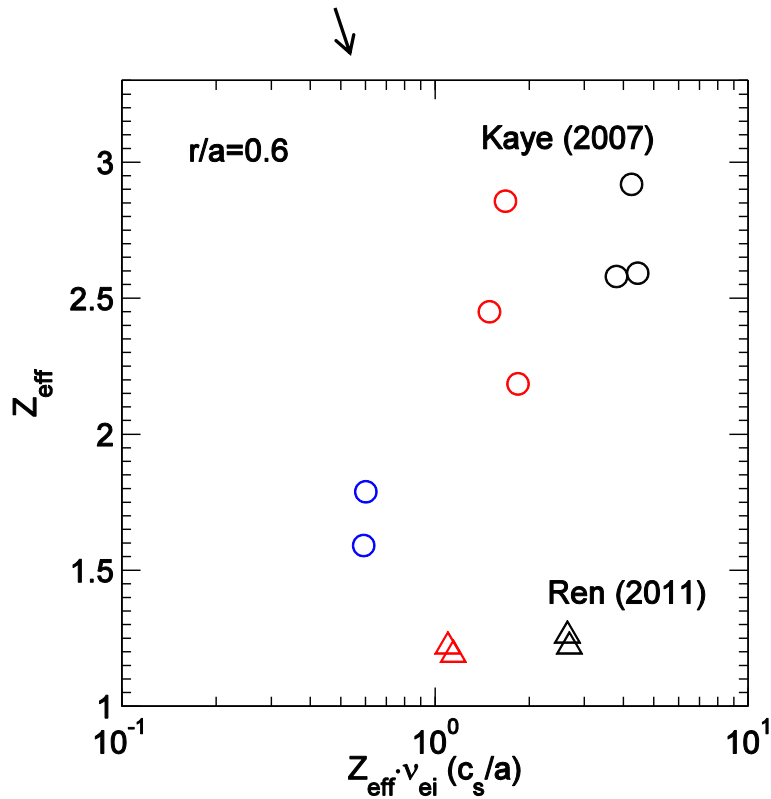
- Ren (2011)  $v_*$  experiment (invited talk TI-2.2) performed at lower  $\beta_e$  compared to Kaye (2007) (lower density and NBI power)





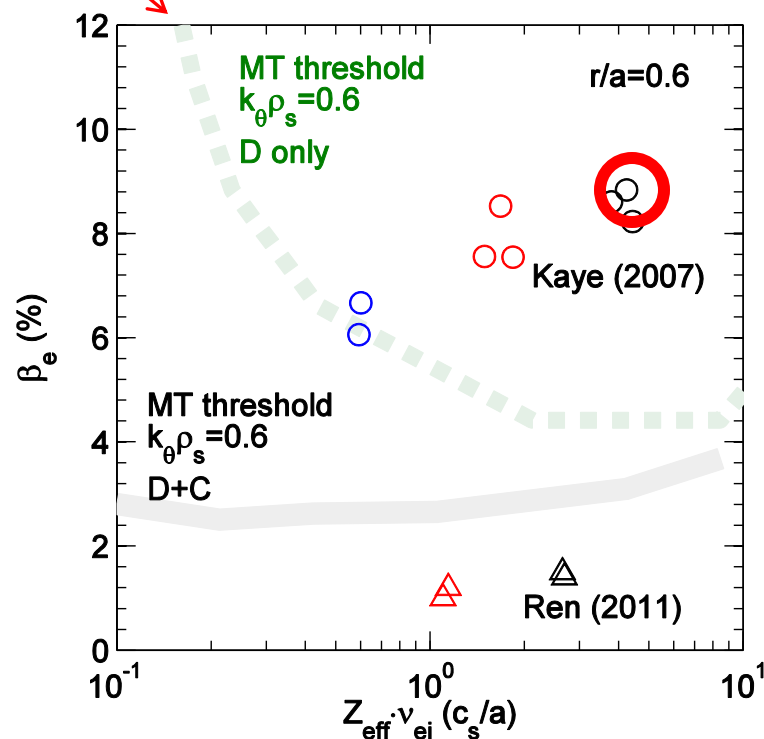
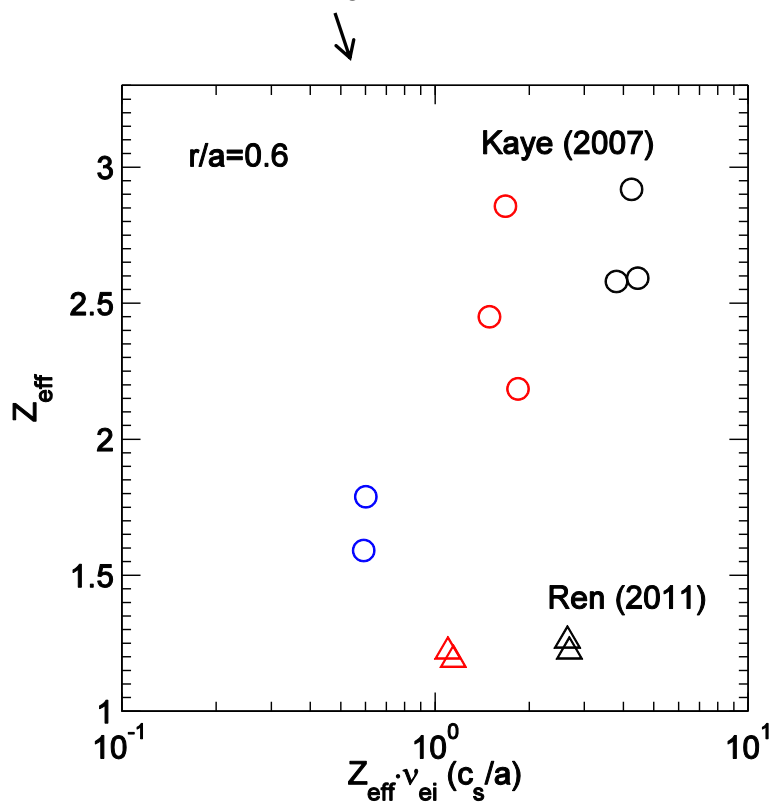
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- Also at lower  $Z_{\text{eff}}$  – **increase in MT threshold, but smaller ETG threshold**



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- **Nonlinear simulations run for high- $\beta_e$ , high- $v_e$  where only microtearing unstable**

# First nonlinear microtearing simulations in NSTX\*

- Simulations where only microtearing unstable, no ETG (NSTX 120968,  $r/a=0.6$ )
  - Electromagnetic ( $\phi, A_{\parallel}$ ) and collisional ( $v_e$ )
  - Varying  $E \times B$  shear (mostly  $\gamma_E=0$ )
  - Deuterium only (but  $Z_{\text{eff}}$  in collision operator)
  - “Local” → no profile variation in equilibrium quantities
  - $k_{\theta}\rho_s=[0,0.105,0.21,\dots]$ , same as  $n=[0,5,10,\dots]$

$L_x \times L_y = 80 \times 60 \rho_s$
$n_x \times n_y = 400 \times 8$ ( $\Delta x = 0.2 \rho_s$ )
$n_{\theta} = 14$ (parallel mesh points)
$n_E = 8, n_{\lambda} = 12 \times 2$ (velocity space)

120968  $r/a=0.6$  surface

$a/L_{Te} = 2.73$        $a/L_n = -0.83$

$q = 1.69$        $s = 1.75$

$\kappa = 1.7$        $\delta = 0.13$

$T_e/T_i = 1.05$        $Z_{\text{eff}} = 2.9$

$\beta_e = 8.8\%$        $v_{ei} = 1.46 c_s/a$

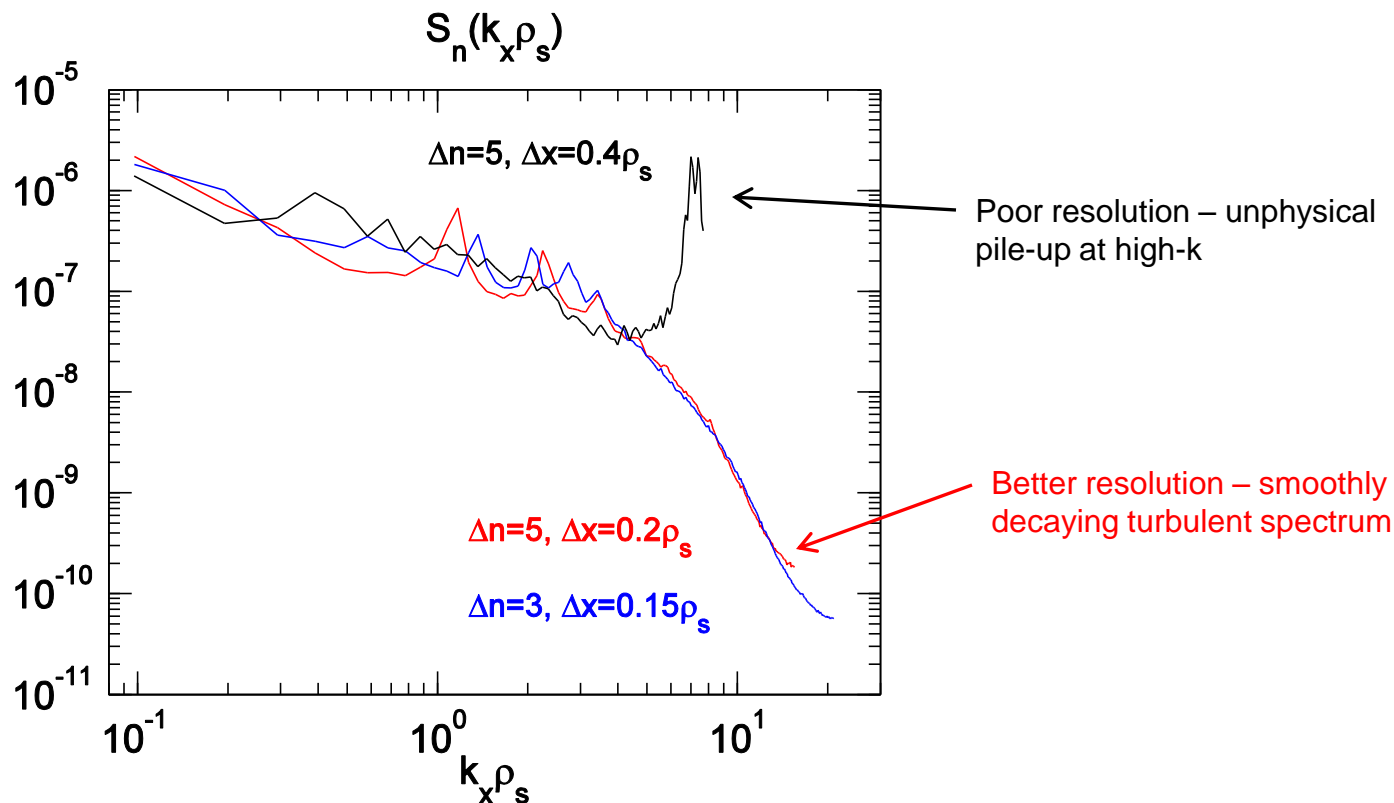
( $\beta_{e,\text{unit}} = 2.5\%$ )

## Acknowledgements: NERSC & OLCF (INCITE award FUS023)

\* W. Guttenfelder et al., Phys. Rev. Lett. **106**, 155004 (2011)

# Fine radial resolution required to obtain decaying nonlinear spectra

- Unphysical pile-up at high-k with insufficient resolution ( $\Delta x = 0.4 \rho_s$ )
- Smoothly decaying turbulent spectrum with better resolution ( $\Delta x = 0.2 \rho_s$ ,  $\Delta x = 0.15 \rho_s$ )



- Similar high-k pile-up observed in first careful attempts of GS2 MAST simulations – for more discussion see Applegate Ph. D. thesis (2007, Imperial College London)

# Fine radial resolution required to distinguish *linear* resonant layers of fastest growing mode

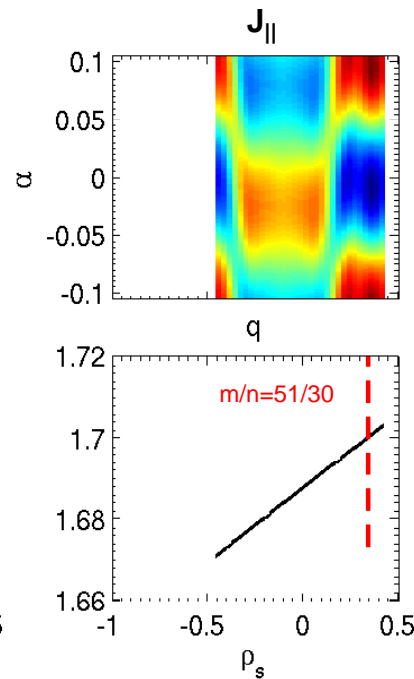
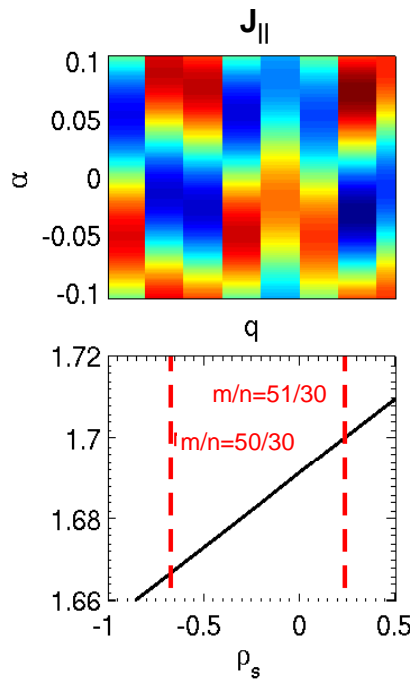
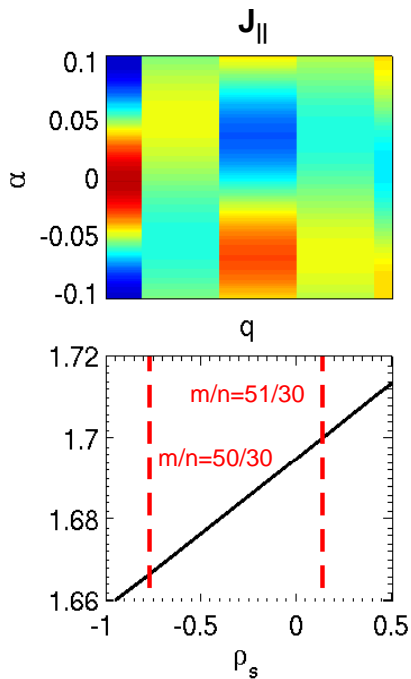
Linear calculation using box width and resolution of nonlinear simulations

“Typical” linear flux-tube calculation

$L_x = 80 \rho_s = 88 \Delta r_{\text{rat}}$   
 $n_x = 200$   
 $\Delta x = 0.4 \rho_s$

$L_x = 80 \rho_s = 88 \Delta r_{\text{rat}}$   
 $n_x = 400$   
 $\Delta x = 0.2 \rho_s$

$L_x = 0.9 \rho_s = \Delta r_{\text{rat}}$   
 $n_x = 32$   
 $\Delta x = 0.03 \rho_s$

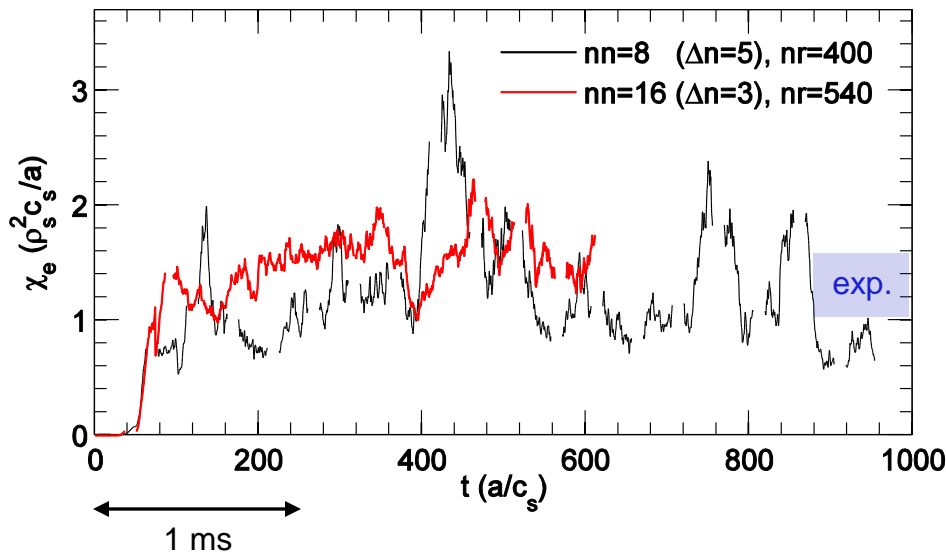


Rough rule-of-thumb:  $\Delta x \leq \min[\Delta r_{\text{rat}}]/4$      $\Delta x/\rho_s \leq 1/(4 \cdot \max[k_0 \rho_s] \cdot s)$

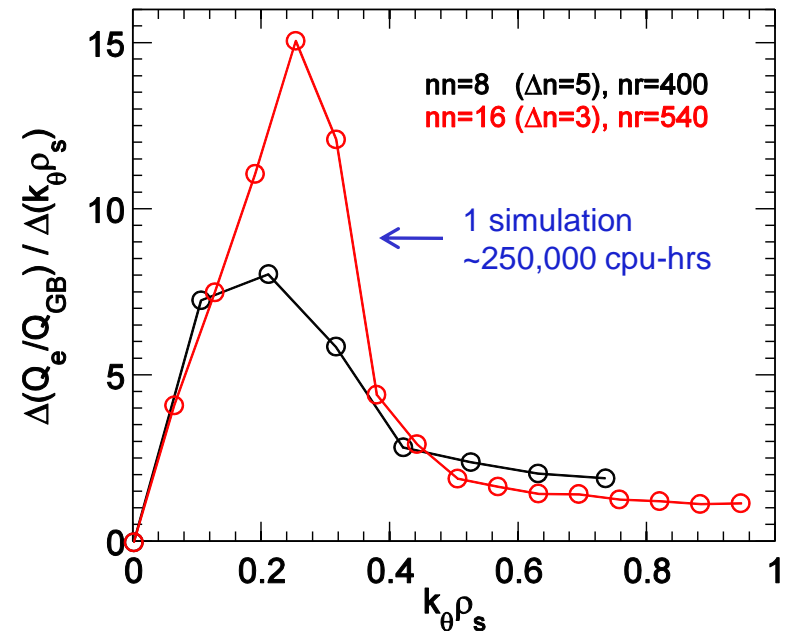
# Predicted electron thermal transport comparable to experiment

- Simulated transport ( $1.2 \rho_s^2 c_s/a$ ,  $6 \text{ m}^2/\text{s}$ ) comparable to experimental transport ( $1.0\text{-}1.6 \rho_s^2 c_s/a$ )
- Well defined peak in transport spectra ( $k_\theta \rho_s \approx 0.2$ ), downshifted from maximum  $\gamma_{\text{lin}}$  ( $k_\theta \rho_s \approx 0.6$ )
- Slowly decaying tail - predicted transport increases  $\sim 25\%$  with higher resolution

Transport time series



Fractional transport spectra

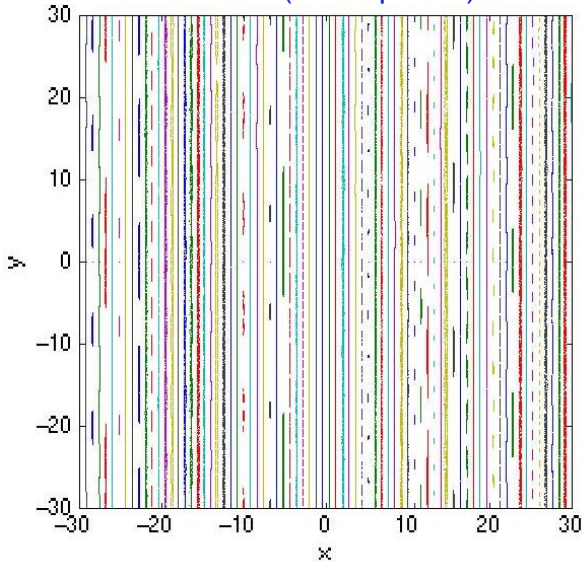


- Negligible particle, momentum, or ion thermal transport

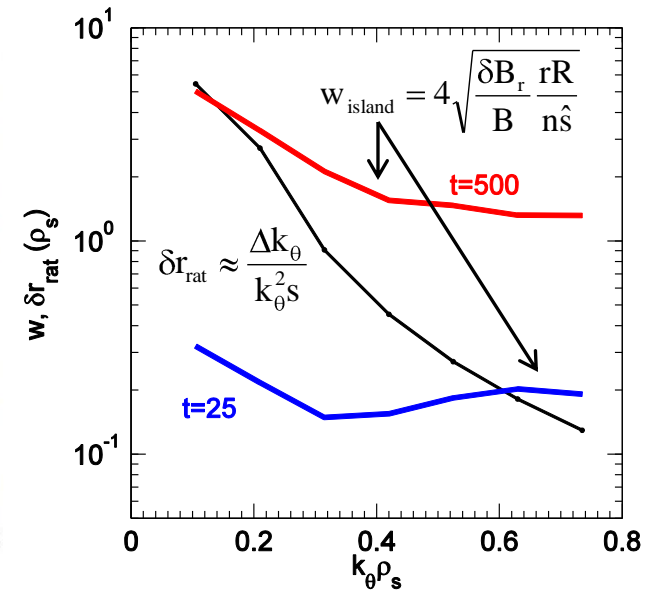
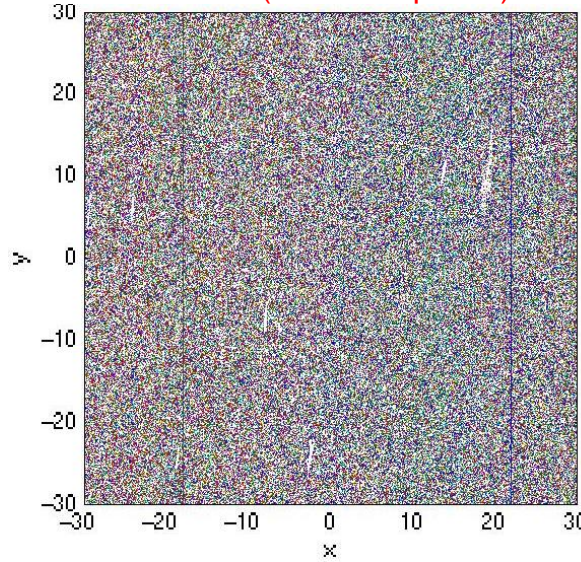
# ~98% of transport due to magnetic “flutter” contribution

- Flux surfaces become distorted in linear phase (t=25)
- Globally stochastic\* in saturated phase, complete island overlap  $w_{\text{island}}(n) > \delta r_{\text{rat}}(n)$

t=25 (linear phase)



t=500 (saturated phase)



- $\chi_{e,EM}$  close to *collisionless* Rechester-Rosenbluth\* ( $\lambda_{\text{mfp}}=12 \text{ m}$ ,  $L_c \approx 2.5 \text{ m}$ )

$$D_{\text{st}} = \lim_{s \rightarrow \infty} \frac{\langle [r_i(s) - r_i(0)]^2 \rangle}{2s}$$

$$\chi^{\text{RR}} \approx 2 \left( \frac{2}{\pi} \right)^{1/2} D_{\text{st}} v_{\text{Te}} f_p \approx 0.9 \left( \frac{\rho_s^2 c_s}{a} \right)$$

$f_p \approx 63\%$  passing particles

\*Wang et al., Phys. Plasmas (2011); Nevins et al., Phys. Rev. Lett. (2011);

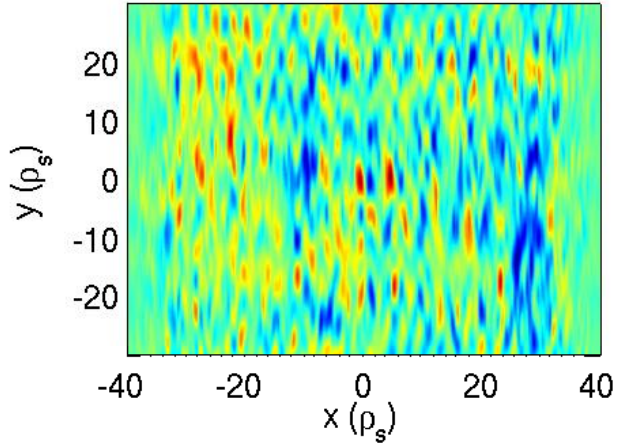
Rechester & Rosenbluth, Phys. Rev. Lett. (1978); Harvey et al., Phys. Rev. Lett. (1981)

# Narrow density perturbations remain in nonlinear simulations

- Narrow radial  $n$ ,  $\phi$ ,  $j_{\parallel}$  structures need to be resolved but  $A_{\parallel}$  very broad
- $\delta B_r/B \sim 0.15\% \sim \rho_e/L_{Te} = 0.065\%$
- $\delta B_r/B \sim \rho_e/L_{Te}$  analytic approximation from Drake et al. PRL 1980; used for NSTX in Wong et al. PRL 2007

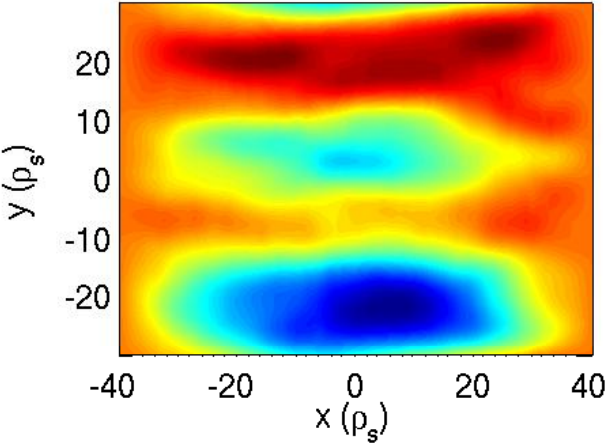
$\delta n/n \approx 0.5\%$

$\delta n$



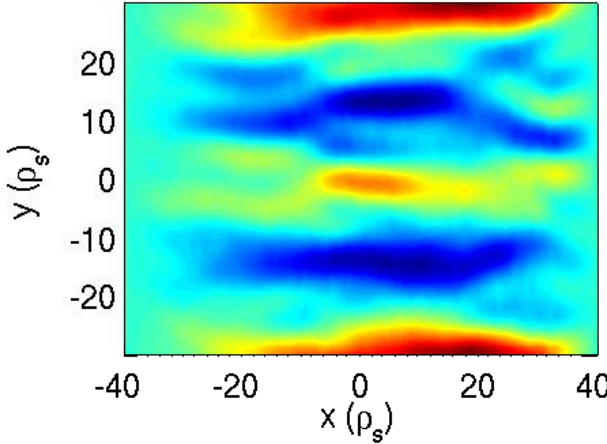
$\delta A_{\parallel}/\rho_s B \approx 0.8\%$

$\delta A_{\parallel}$



$\delta B_r/B \approx 0.15\%$

$\delta B_r$

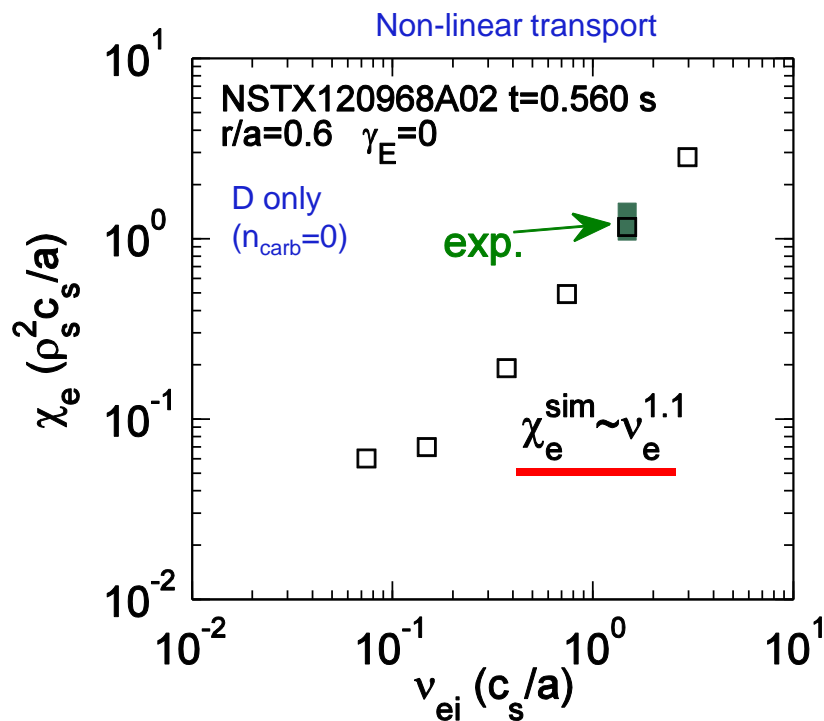


$\delta T_e/T_e \approx 2\%$

$\delta v_{e,\parallel}/c_s \approx 6\%$



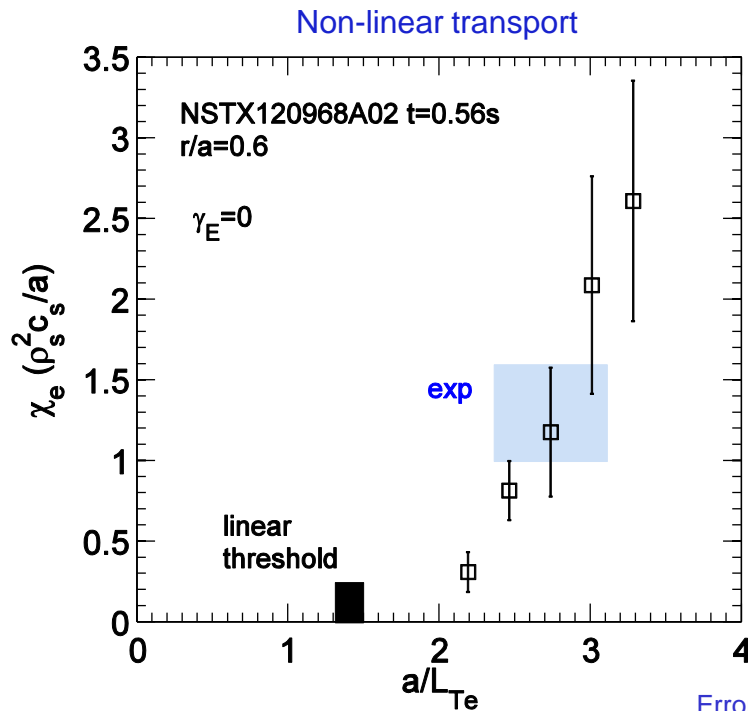
# Near linear scaling of transport with $v_e$ consistent with experimental scaling



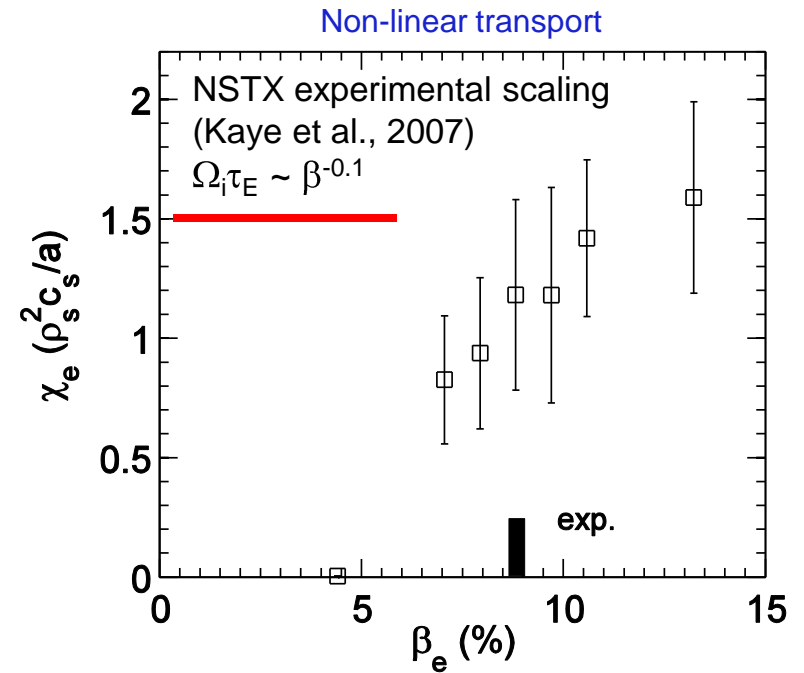
- As transport drops,  $a/L_{T_e}$  will increase (for fixed heat flux), at some point ETG (TEM?) should become important
- This transition likely to determine limit of “favorable”  $v_*$  scaling
- Also likely to depend on ionic charge ( $Z_{\text{eff}}$ ) – above with D only ( $n_c=0$ )

# Predicted transport “stiff” with $\nabla T_e$ , increases with $\beta_e$

- Complicates simple interpretation from  $\chi_{e,sim} \sim \nu_e^{1.1}$  scaling
- Useful to characterize scaling of threshold gradient (work in progress)



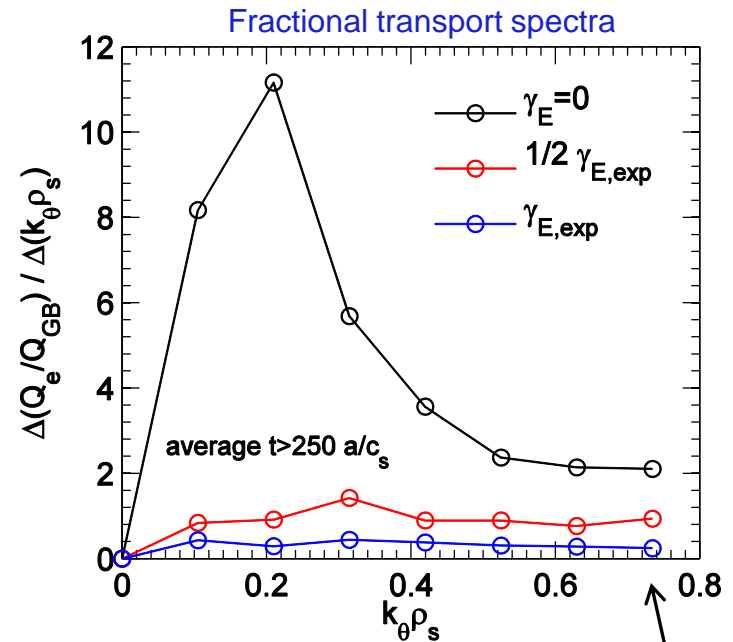
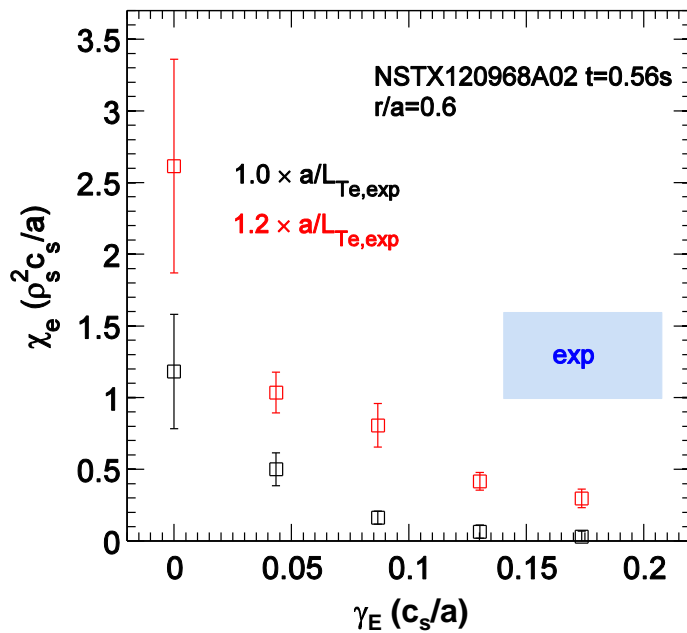
Error bars represent rms variation in transport



- Non-linear threshold  $\sim 40\%$  bigger than linear threshold -- possible influence from limited numerical resolution

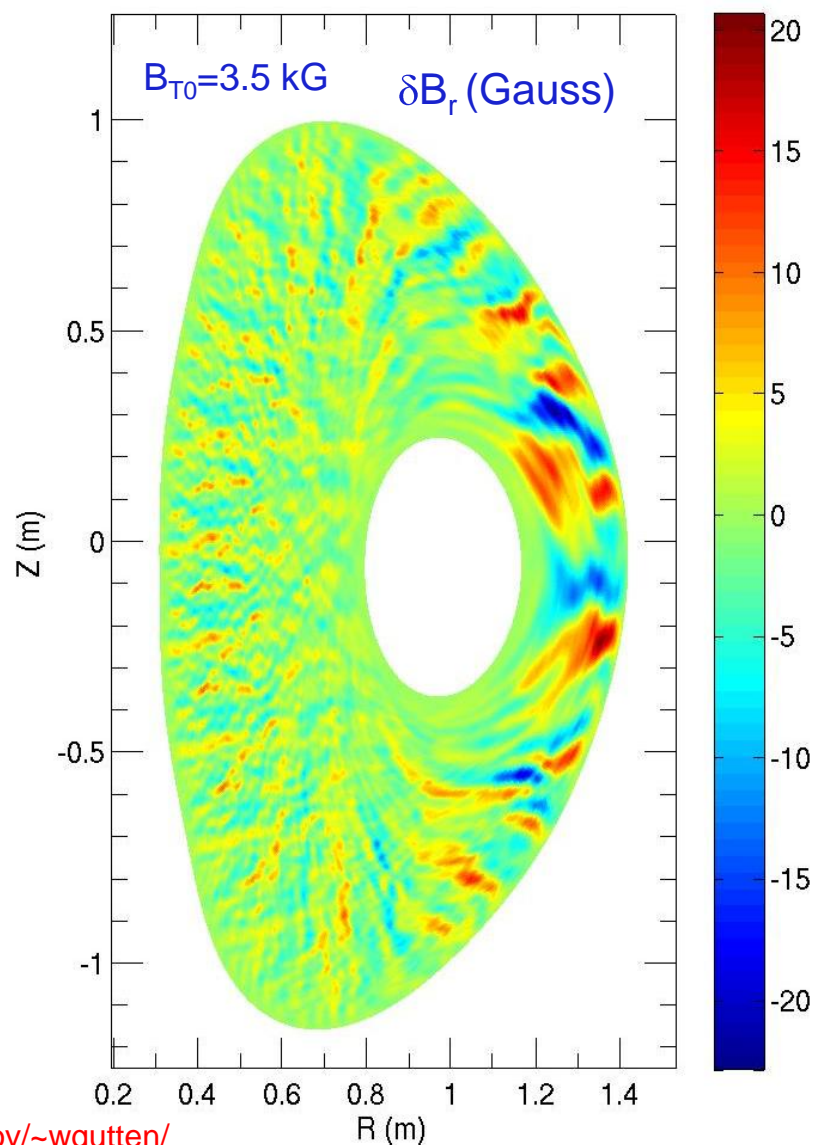
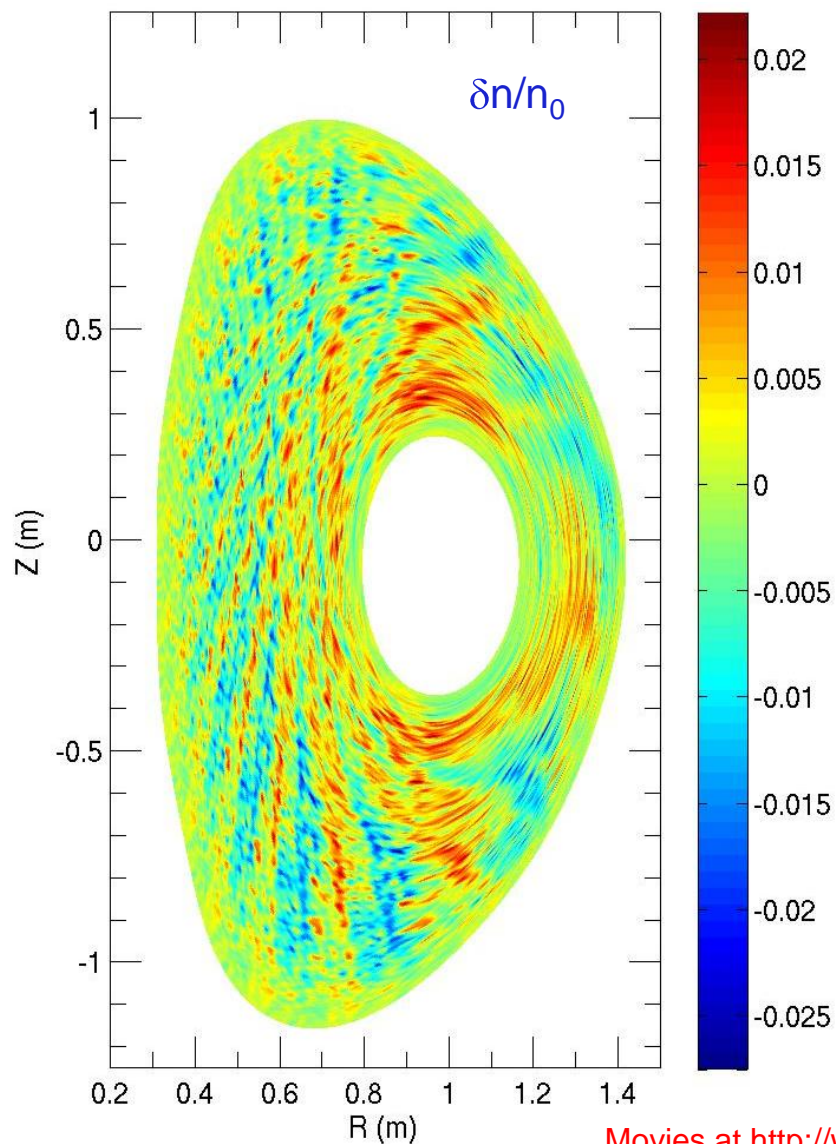
# Nonlinear microtearing transport sensitive to $\gamma_E/\gamma_{lin}$

- Transport reduced when increasing  $\gamma_E$  to local experimental value ( $\gamma_{E,exp} \sim \gamma_{lin,max} \sim 0.17 c_s/a$ )
- Transport partially recovered with **increase in  $\nabla T_e$**



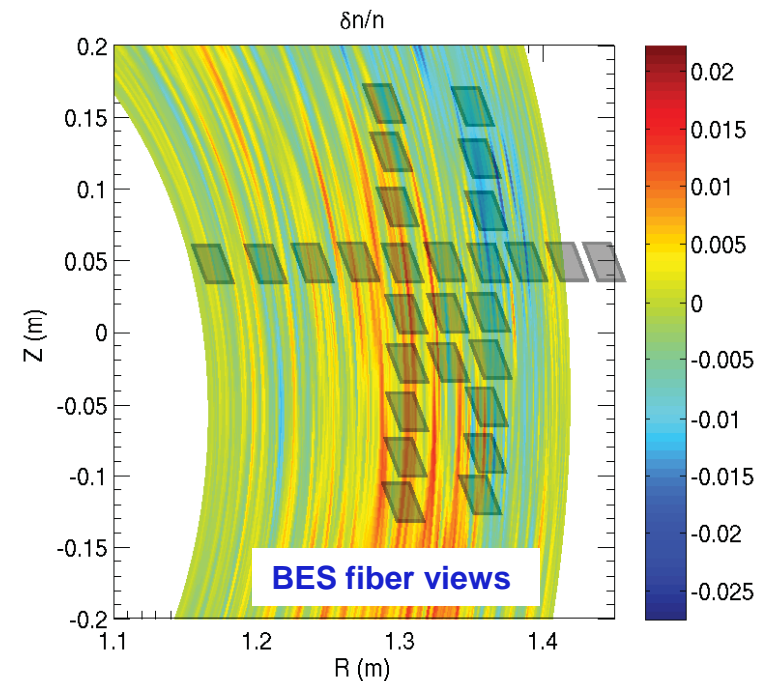
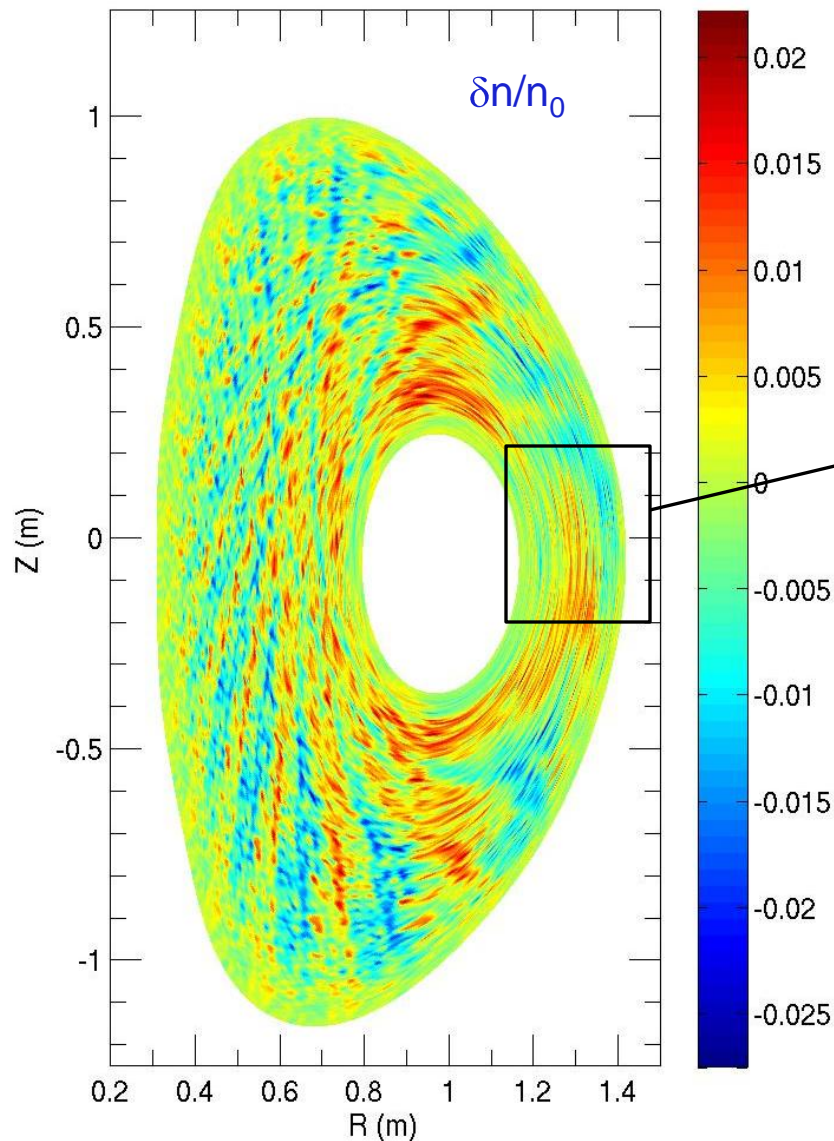
- Higher ionic charge ( $Z_{eff} > 1$ , through adiabatic response) and improved resolution (binormal and radial) could increase transport
- Profile (non-local) effects could also matter -  $\rho_s/a \approx 1/100$  & edge more strongly driven

# What hope is there to experimentally identify microtearing modes?



Movies at <http://www.pppl.gov/~wgutten/>

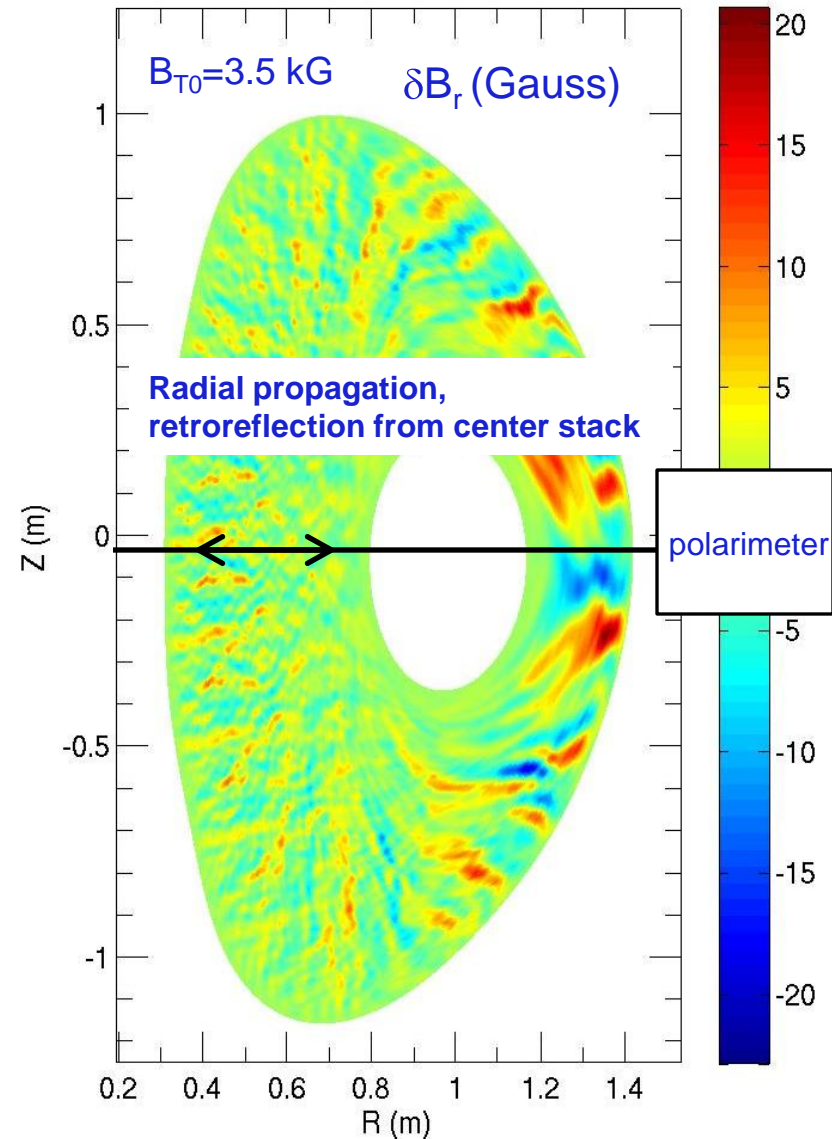
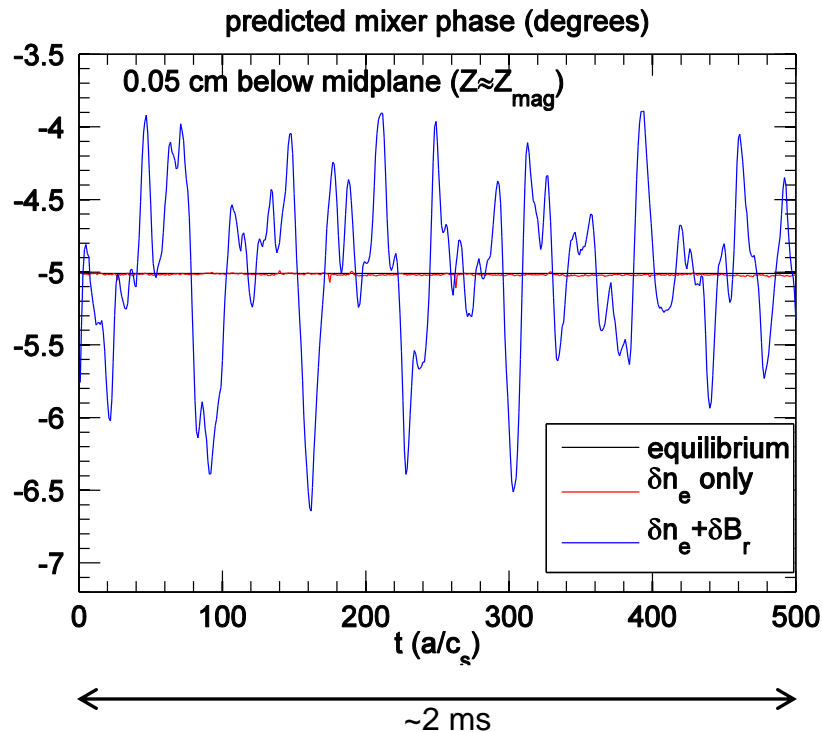
# BES for density fluctuations



- BES suitable for long poloidal scale (U-Wisconsin, Smith et al., RSI 2010)
- May average over narrow radial scale – requires synthetic diagnostic and instrument function (D. Smith, BO4.2)

# Polarimetry for magnetic field fluctuations

- New UCLA polarimetry system (J. Zhang, PP9.71)
- Simulations suggest  $(\delta B/B)_{\text{internal}} \leq 0.1\%$  may be detectable ( $1\text{-}2^\circ$  or  $\sim 0.3^\circ$  rms mixer phase)



# Summary

- Microtearing modes found to be unstable in experimental  $v_*$  scans
  - Scaling of linear growth rates  $\gamma_{\text{lin}} \sim v_e$  – **potential candidate to explain experimental confinement trend**
  - Linear thresholds exist in  $a/L_{Te}$  &  $\beta_e$
  - Ionic charge ( $Z_{\text{eff}}$ ) can enhance instability (opposite to ETG expectations)
- First non-linear microtearing simulations in NSTX
  - Require relatively fine radial resolution ( $\Delta x \approx 0.2\rho_s$ ,  $n_x=400$ ) to capture physics
  - Transport dominated by electromagnetic contribution ( $\delta A_{\parallel}$ )  $\rightarrow$  stochastic field lines
  - **Predicted  $\chi_{e,\text{sim}} \sim v_e^{1.1}$  close to experimental scaling**
  - “Stiff” with  $\nabla T_e$  but suppressible by experimental levels of  $E \times B$  shear

## Acknowledgements

NERSC, OLCF, and support from DOE Contract No's: DE-AC05-00OR22725, DE-AC02-09CH11466, DE-FG03-95ER54309, DE-AC52-07NA27344, DE-FG02-99ER54527