

Gyrokinetic predictions of momentum and impurity transport in NSTX

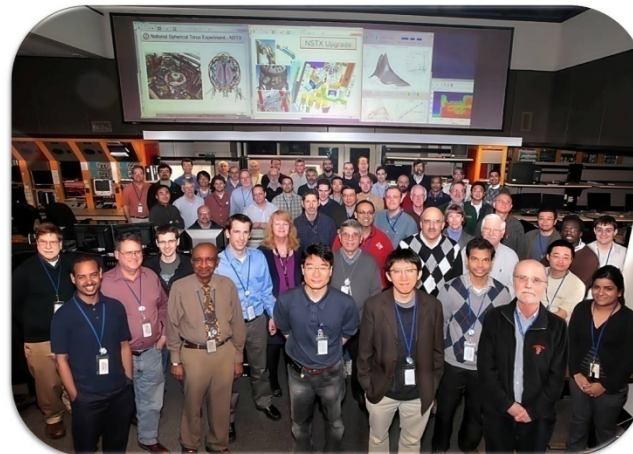
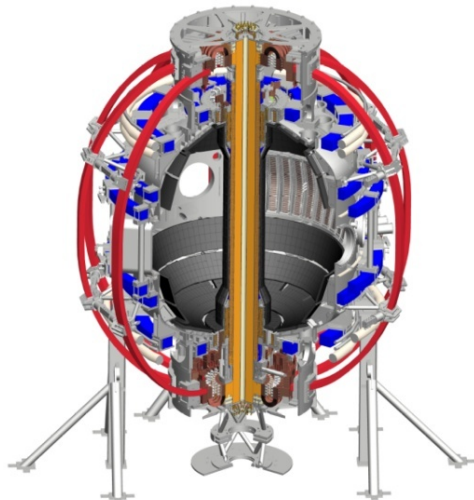
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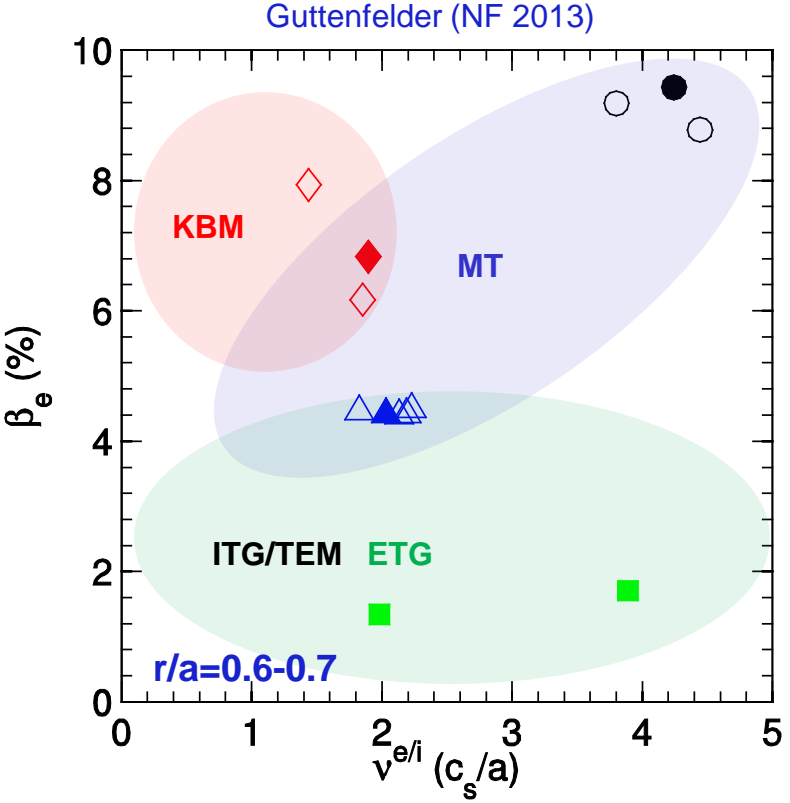
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Broad range of parameters in NSTX requires consideration of many micro-instabilities



Transport Mechanism	Transport channel affected			
	ion energy	electron energy	particle/impurity	momentum
ITG	×	×	×	×
TEM	×	×	×	×
KBM	×	×	×	×
MT		×		
ETG		×		

- All of them of interest for electron thermal transport
 - Only ion scale ballooning instabilities (ITG, TEM, KBM) expected to transport momentum and impurity
- ⇒ Investigate multiple transport channels to help constrain theory

Summary

Momentum transport

- NSTX L-modes governed by ITG/TEM, linear simulations predict:
 - Prandtl number ($Pr \sim 0.2-0.8$) generally consistent with experimental analysis (~ 0.5)
 - Relatively weak inward pinch ($RV_\phi/\chi_\phi \sim -1$) insensitive to many parameters except R/a
- NSTX $n=3$ magnetic perturbation H-mode experiments dominated linearly by microtearing ($r/a=0.6-0.8$)
 - Sub-dominant KBM exists, $Pr \sim 0.3-0.6$ similar to experiment
 - Predicted $RV_\phi/\chi_\phi \sim -1 \rightarrow +2$ small/outward compared to stronger inward experimental values ($-7 \rightarrow -1$), relatively insensitive to parameter variations

Impurity transport

- In lithiated H-mode cases where impurity carbon transport appears to be anomalous:
 - KBM modes (sub-dominant to microtearing) predict inward carbon pinch opposite to experiment

⇒ **Local linear theory does not appear to explain momentum & impurity transport trends in NSTX – future work to investigate nonlinear and global (finite- ρ_*) effects**

Momentum Transport

Interpretation of toroidal angular momentum transport often assumes diffusive and convective components

- Transport equation:
$$\frac{\partial}{\partial t} (n_i m_i \langle R^2 \rangle \Omega) + \nabla \cdot \Pi_\varphi = S_\Omega \rightarrow \sum_s (\dots)$$

- Assumed transport form:
$$\Pi_\varphi = -nmR\chi_\varphi (R\nabla\Omega) + nmV_\varphi (R\Omega)$$

Prandtl number

$$\text{Pr} = \frac{\chi_\varphi}{\chi_i}$$

$$\hat{\Pi}_\varphi = \hat{\chi}_\varphi \left(\hat{u}' + \frac{RV_\varphi}{\chi_\varphi} \hat{u} \right)$$

Pinch parameter

$$\frac{RV_\varphi}{\chi_\varphi}$$

$$\hat{u}' = \frac{-R^2 \nabla \Omega}{c_s} \quad \hat{u} = \frac{R\Omega}{c_s}$$

- Can also have residual stress Π_{RS} contributions (from up-down asymmetric flux surfaces, finite ρ_* profile effects) leading to intrinsic torque \rightarrow intrinsic rotation when $u'=u=0$
 - Perhaps less important in core with large beam torque (co-NBI in NSTX)

Steady state Prandtl numbers $\chi_\phi/\chi_i < 1$ for NSTX L- mode and H-mode discharges

- $Pr = \chi_\phi/\chi_i \approx 0.3-1.0$ over many radii and discharges (assumes $V_\phi=0$)

- $\chi_\phi > \chi_{\phi,NC}$ for both L and H
In L-mode $\chi_i > \chi_{i,NC}$

$$Pr = \frac{\chi_\phi}{\chi_i} \approx \frac{\chi_{\phi,turb}}{\chi_{i,turb}}$$

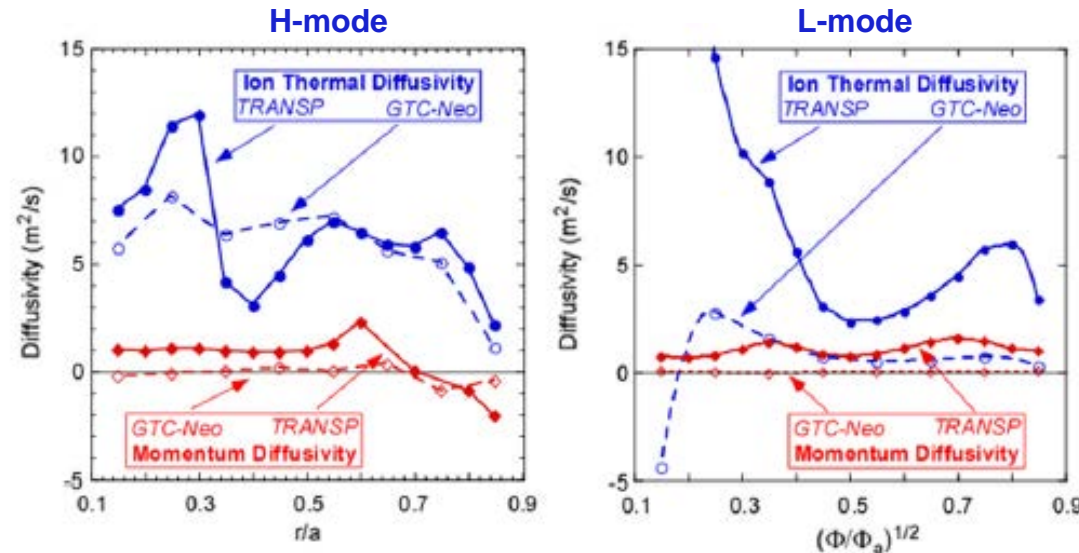
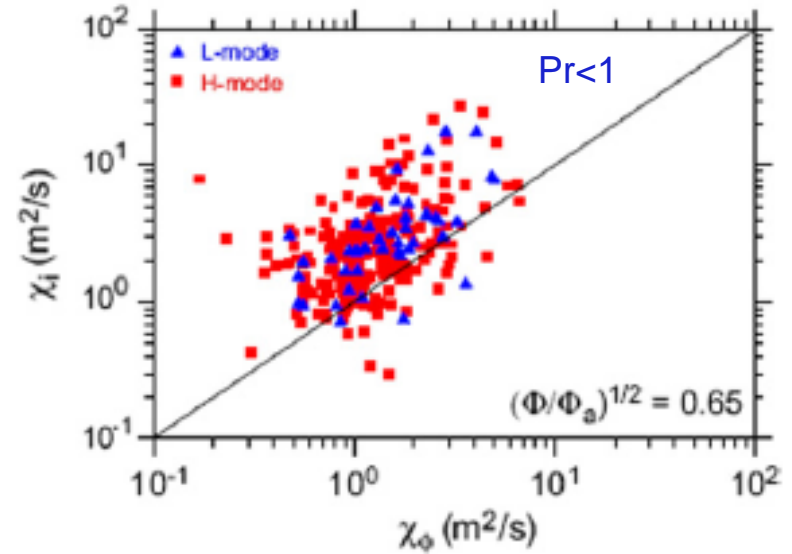
In H-mode $\chi_i \approx \chi_{i,NC}$

$$Pr = \frac{\chi_\phi}{\chi_i} = \frac{\chi_{\phi,turb}}{(\chi_{i,NC} + \chi_{i,turb})} \rightarrow \sim 0$$

⇒ Pr less useful in H-mode?

- RV_ϕ/χ_ϕ less ambiguous

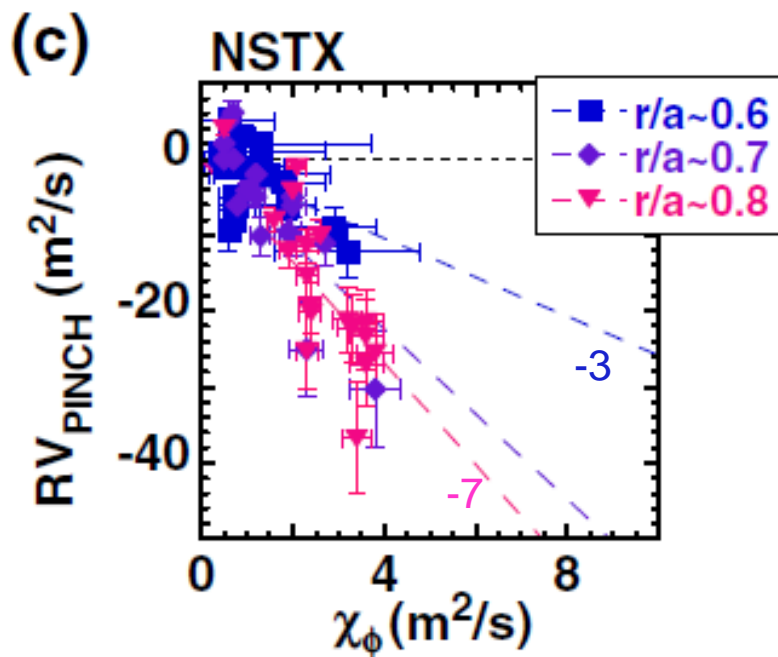
Steady state transport analysis (Kaye et al., 2009)



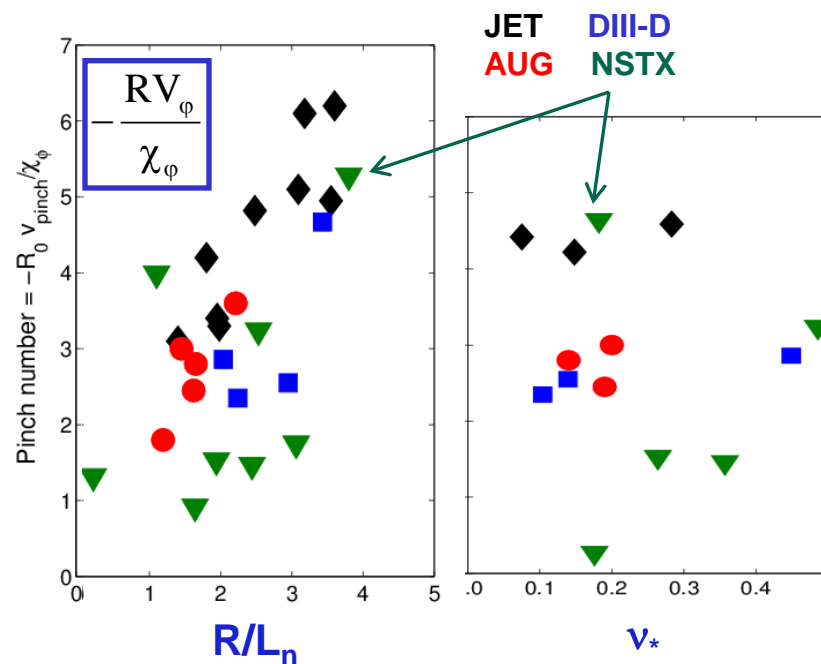
Perturbative H-mode experiments indicate existence of an inward momentum pinch

- $RV_\phi/\chi_\phi \approx -(1-7)$ for many NSTX discharges & radii
 - $Pr \sim 0.3-0.5$, smaller than other machines ($Pr \sim 0.6-2.0$) (Yoshida, NF 2012)
- Possible dependence on density gradient (R/L_n), less clear with collisionality (ν^*), but a lot of scatter

(Solomon et al., PRL 2008, PoP 2010; Yoshida et al., NF 2012)



(Tala et al., IAEA 2012)



Q: What are the relevant momentum transport mechanism(s) in NSTX?

Method for predicting quasi-linear Prandtl (χ_ϕ/χ_i) and Pinch numbers (RV_ϕ/χ_ϕ)

- Local linear GYRO simulations run between $r/a=0.6-0.8$ ($\rho_{\text{tor}}\approx 0.5-0.7$), with
 - deuterium, carbon, electrons
 - $\phi, A_{\parallel}, B_{\parallel}$
 - numerical equilibrium (EFIT/LRDFIT)
 - n_e profiles from averaged inboard/outboard measurements (no centrifugal effects in GYRO)

- Pr and RV_ϕ/χ_ϕ determined using momentum

flux from different combinations of u, u' $\hat{\Pi}_\phi = \hat{\chi}_\phi \hat{u}' + (\hat{R}\hat{V}_\phi + \hat{R}\hat{\Gamma}_p)\hat{u} + \hat{\Pi}_{\phi,RS}$

$$\text{Pr} = \frac{\hat{\chi}_\phi}{\hat{\chi}_i} = \frac{\hat{\Pi}_\phi(0, u') - \hat{\Pi}_\phi(0, 0)}{\hat{u}'} \cdot \frac{a/L_{Ti}}{\hat{Q}_i}$$

Following Peeters et al.
PRL (2007)
Nucl. Fusion (2011)

$$\left(\frac{RV_\phi}{\chi_\phi} \right) = \left[\frac{\hat{\Pi}_\phi(u, 0) - \hat{\Pi}_\phi(0, 0)}{\hat{u}} - \frac{\hat{m}\hat{R}\hat{\Gamma}_p(u, 0)}{\hat{u}} \right] \cdot \frac{\hat{u}'}{\hat{\Pi}_\phi(0, u') - \hat{\Pi}_\phi(0, 0)}$$

- Subtracting particle convection contribution

Transport of toroidal angular momentum calculated from delta-f gyrokinetics (GYRO*)

- Transport calculated for **toroidal momentum** from correlation of **perturbed distribution function** and **effective radial drifts** from all EM fields

$$\delta f_s(\vec{x}) = -\frac{e\delta\phi(\vec{x})}{T_s} F_{s0} + H_s(\vec{R}) \quad (3.22)$$

$$\Pi_s = \oint_{\text{flux surface average}} \int d^3v H_s^*(\vec{R}) \left\langle \underbrace{[m_s \mathbf{R}(\vec{V}_0 + \vec{v}) \cdot \vec{e}_\phi]}_{\text{red line}} \underbrace{\frac{c}{B} \vec{b} \times \nabla \left[\delta\phi(\vec{x}) - \frac{1}{c} (\vec{V}_0 + \vec{v}) \cdot \delta\vec{A}(\vec{x}) \right]}_{\text{blue line}} \cdot \nabla \mathbf{r} \right\rangle_{\text{gyro average}} \quad (3.55)$$

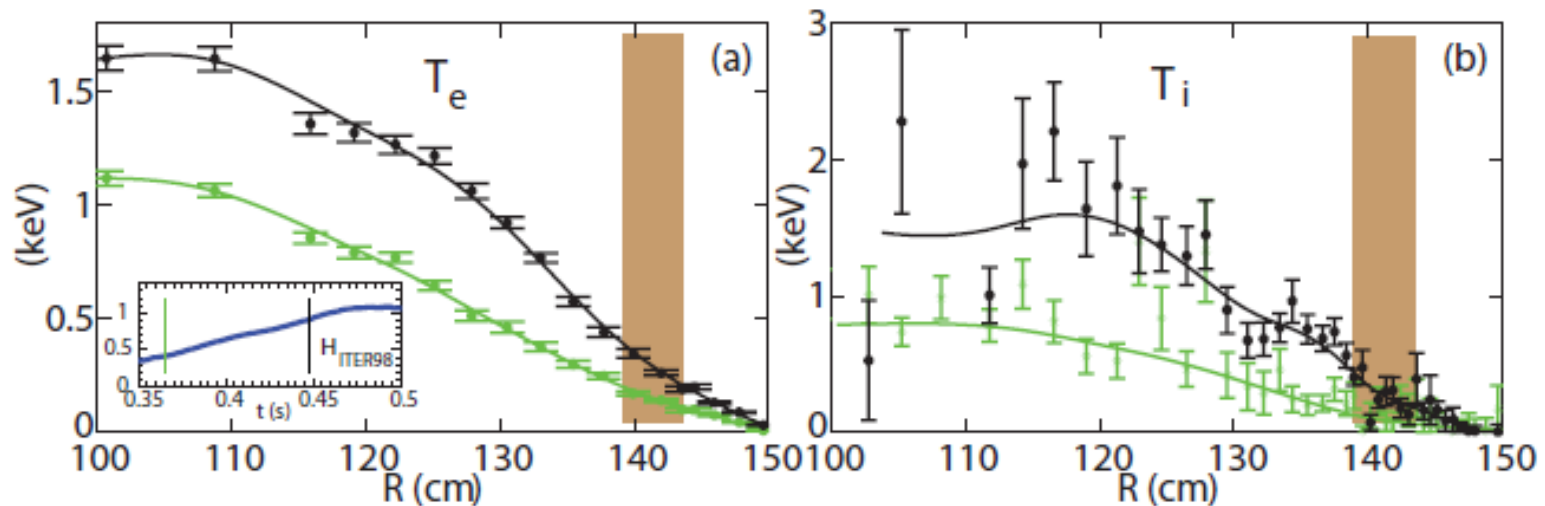
↓ Electrostatic E×B drift
↓ Drifts from shear ($v_{\parallel} \nabla A_{\parallel} \sim v_{\parallel} B_r$) and compressional ($v_{\perp} \nabla A_{\perp} \sim v_{\perp} B_{\parallel}$) magnetic perturbations

- EM contributions are important in NSTX H-modes*

*Candy & Belli, GYRO Technical Guide, <https://fusion.gat.com/theory/Gyro>

Example from NSTX L-mode (Ren, NF 2013)

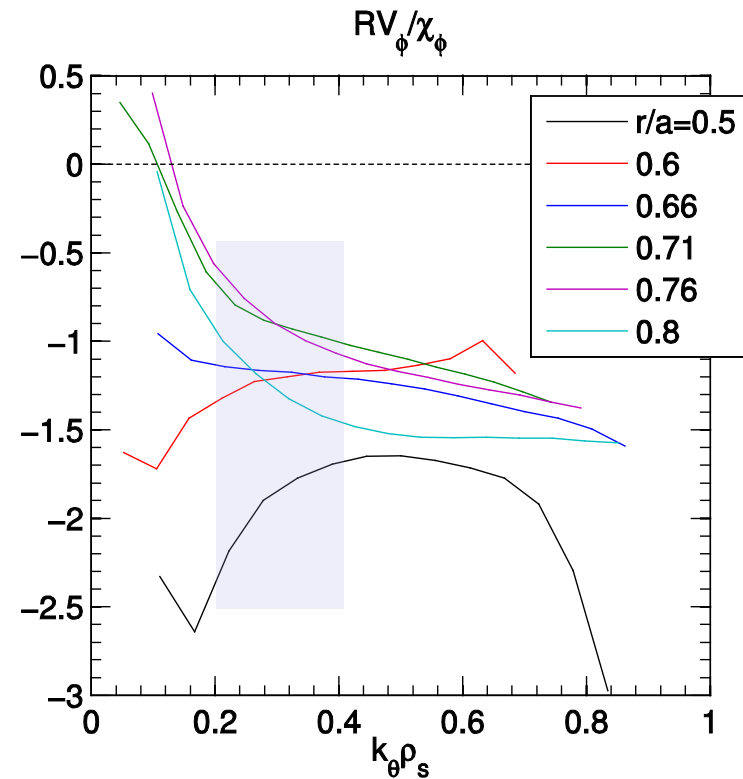
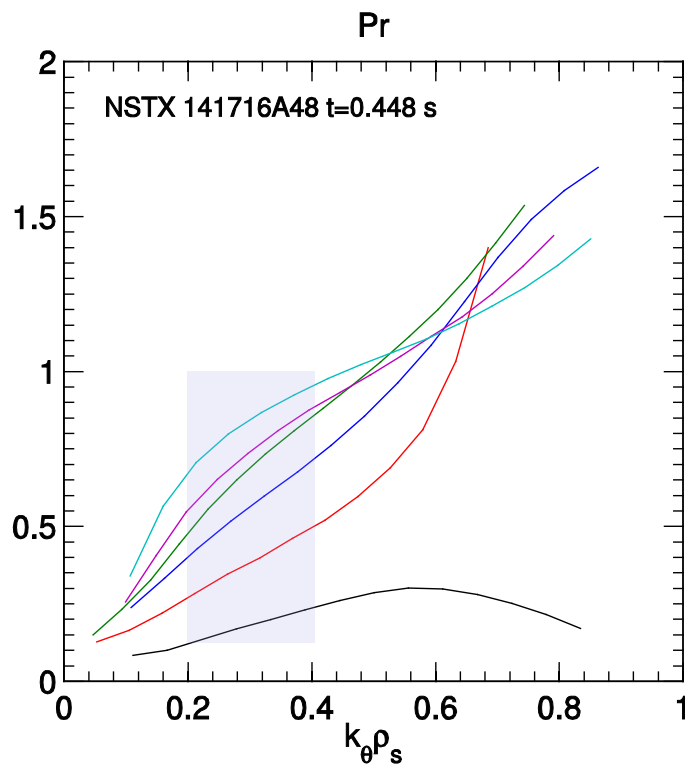
- Low k_θ stability dominated by ITG/TEM
- No perturbative momentum experiments in this case, but it provides a basis for comparing to conventional tokamaks
- MAST perturbative L-mode experiments were run this year (analysis ongoing)



$$B_T = 0.55 \text{ T}, I_p = 0.9 \text{ MA}, P_{\text{NBI}} = 2 \text{ MW}, \langle n \rangle \approx 3 \times 10^{19} \text{ m}^{-3}$$

Quasilinear Prandtl number increases with radius, relatively weak momentum pinch predicted

- Range of $Pr \sim 0.2-0.8$ generally consistent with experiment (~ 0.5)
 - NL spectrum peak around $k_\theta \rho_s \sim 0.3$
 - Small inward pinch $RV_\phi / \chi_\phi \sim -(1-2)$
- \Rightarrow Investigate sensitivity to various parameters

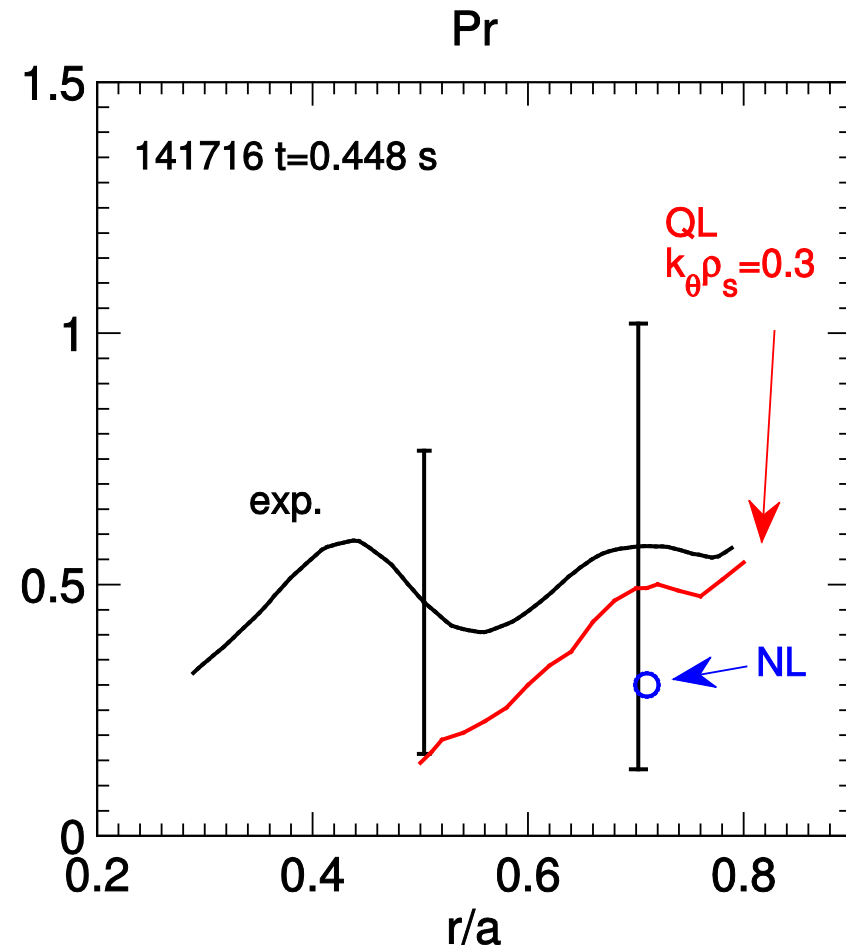


Quasi-linear and non-linear “effective” Pr number in range of experimental values

- Comparing “effective” Pr number (ignoring small pinch which lowers quasilinear Pr)

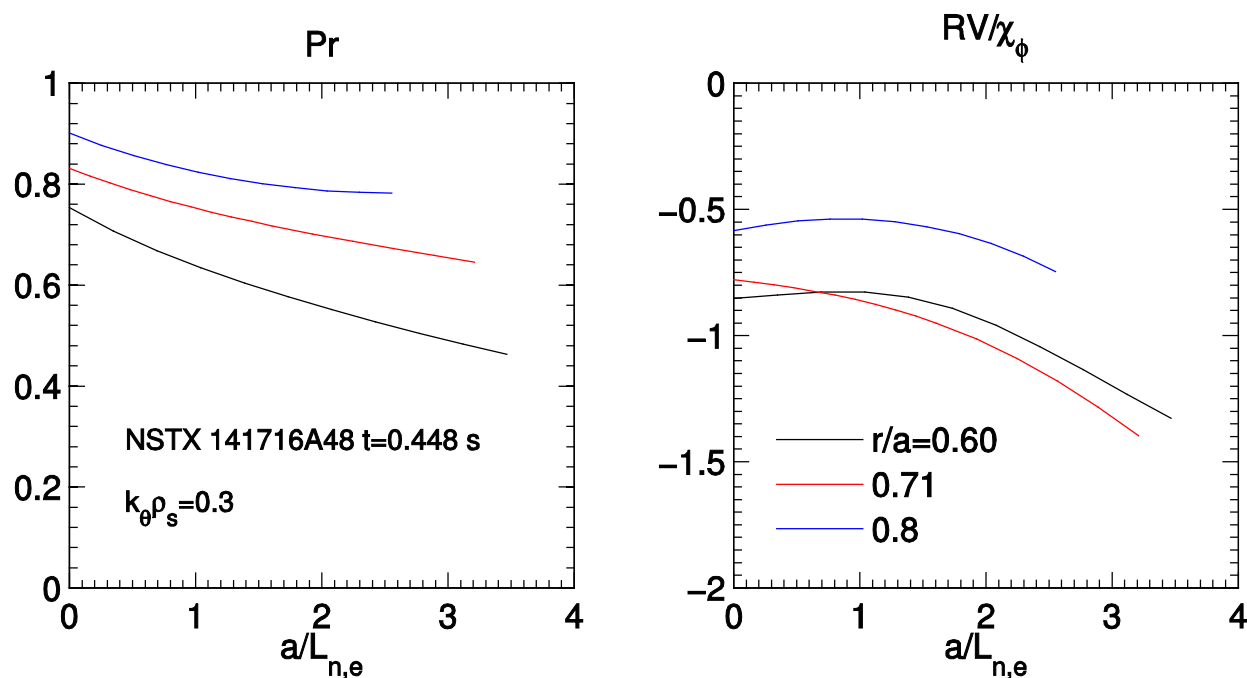
$$\text{Pr}_{\text{eff}} = \frac{\hat{\chi}_{\phi}^{\text{eff}}}{\hat{\chi}_i}, \quad \hat{\chi}_{\phi}^{\text{eff}} = \chi_{\phi} \left(1 + \frac{R V_{\phi}}{\chi_{\phi}} \frac{u}{u'} \right)$$

- Large uncertainty in χ_i due to power balance (e-i coupling, T_i uncertainties)
 - Quasilinear estimates reasonably close
 - Nonlinear simulation predicts lower effective Pr
- ⇒ Investigate sensitivity to various parameters



Pinch remains relatively small even for increased density gradient ($a/L_n = -a\nabla n/n$)

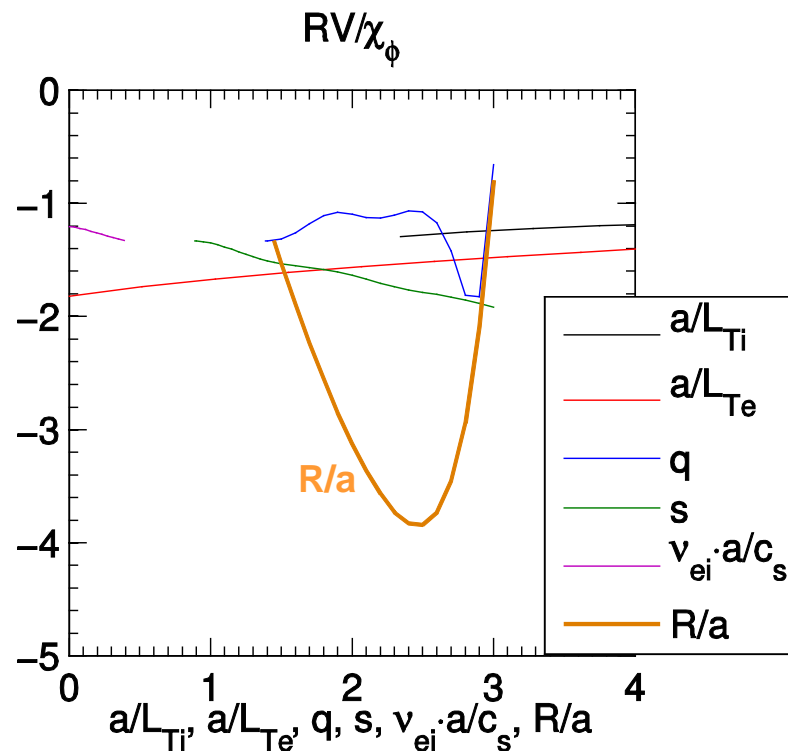
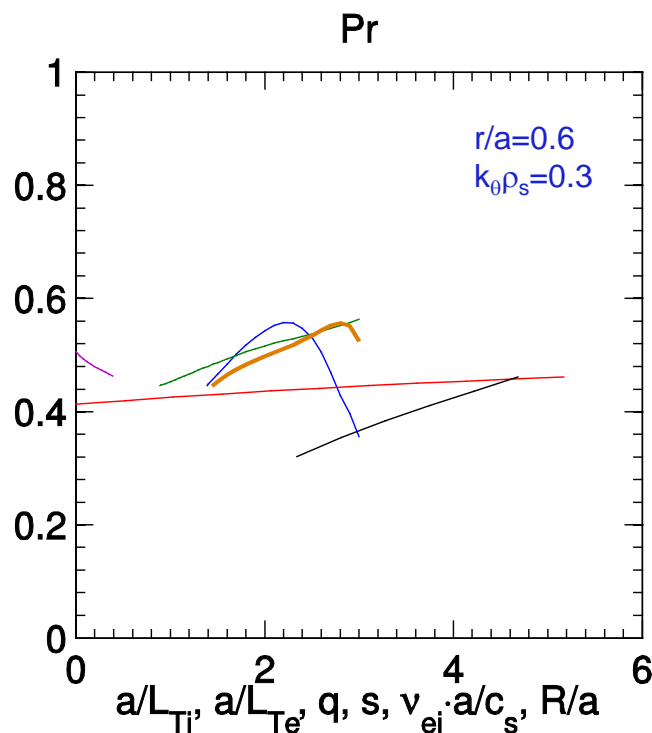
- Weaker dependence than predicted for ITG in conventional tokamaks



- Growth rates at $r/a=0.6$ increase with a/L_n
 - TEM-like at $r/a=0.6$
 - ITG-like at $r/a=0.8$
- Weaker pinch consistent with smaller RV_ϕ/χ_ϕ reported for TEM conditions at higher aspect ratio [Kluy et al., 2009]

Pinch predicted to be weakly dependent on many parameters except aspect ratio (R/a)

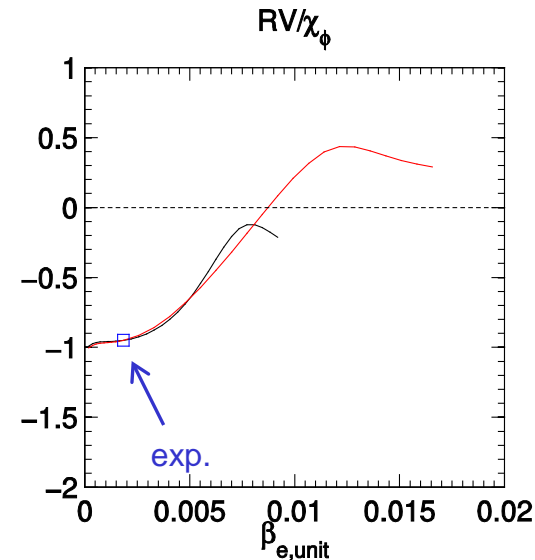
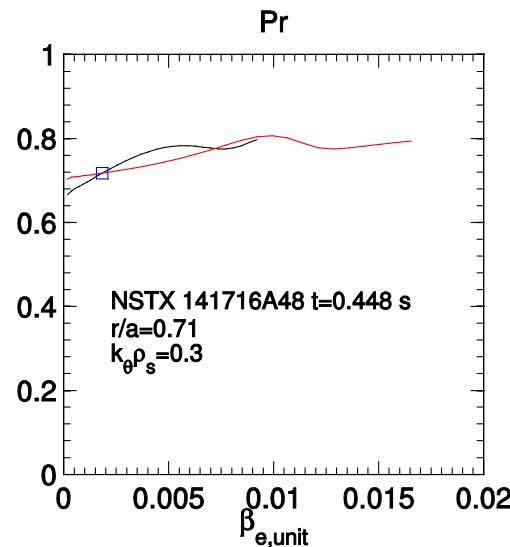
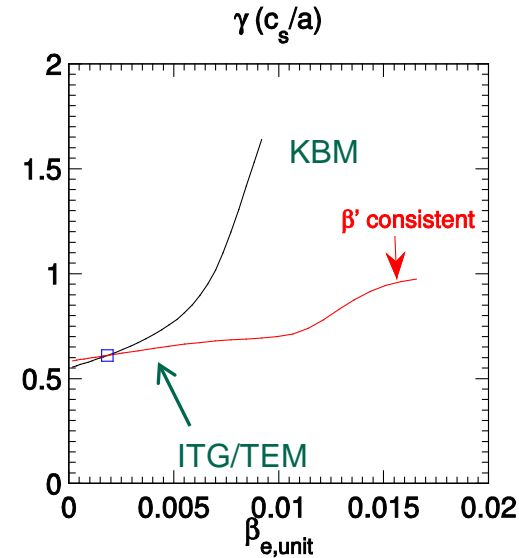
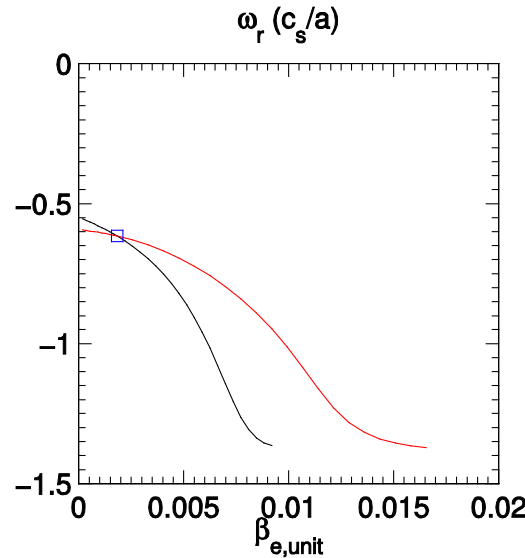
- Prandtl number remains constant $\sim 0.4-0.6$
- RV_ϕ/χ_ϕ relatively insensitive to $a/L_{Ti,e}$, q , s , v_{ei}
- Becomes much larger (inward) for **increased aspect ratio (R/a)**



- q , s , R/a scans using local Miller equilibrium model \Rightarrow not consistent with any particular global equilibrium

Growth rates increase with beta, eventually transition to KBM (preview for H-modes)

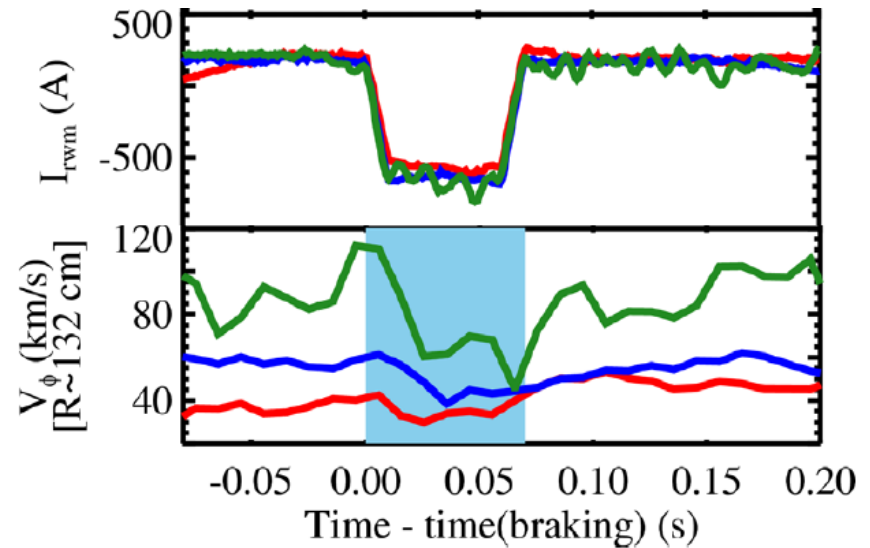
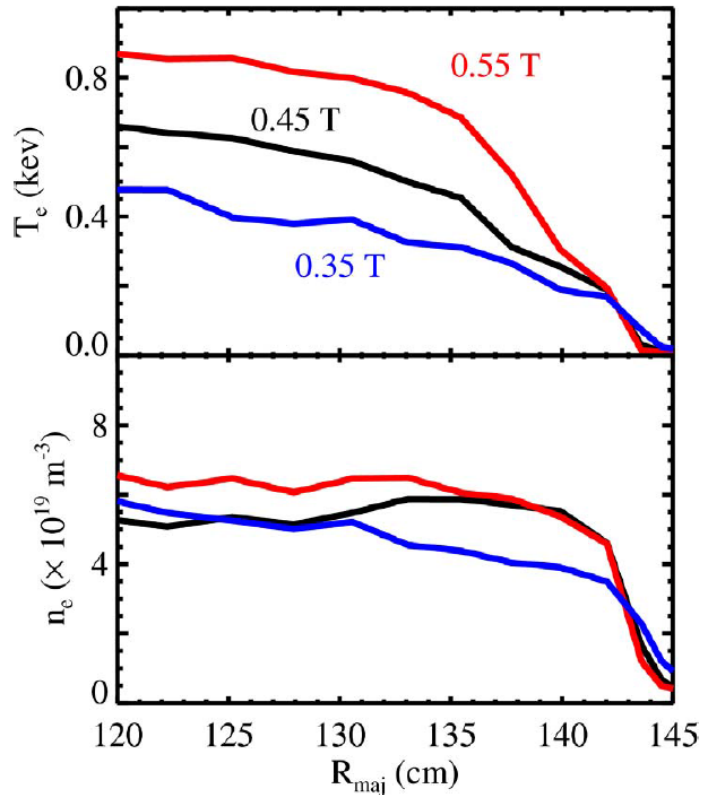
- ITG/TEM growth rates *increase* with β_e , opposite to traditional results (e.g. “cyclone base case”)
- Eventually transitions to KBM (similar to hybrid ITG/KBM [Belli, Candy 2010])
 - Increasing β'_{eq} consistently is stabilizing [Bourdelle, 2003]
- Pr remains ~constant
- Pinch goes toward zero, even positive/outward (depending on β'_{eq})
 - similar to EM behavior predicted in conventional aspect ratio [Hein, 2010]



NSTX H-modes

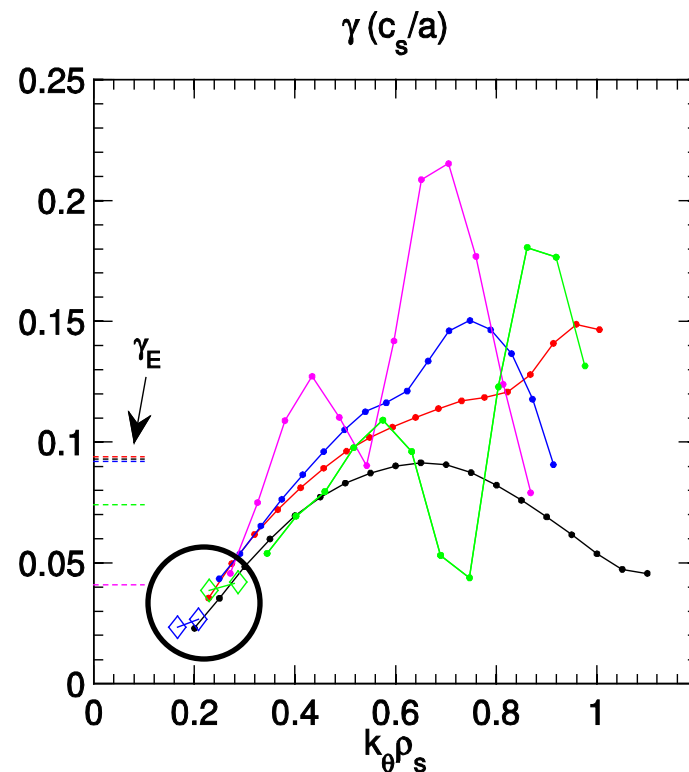
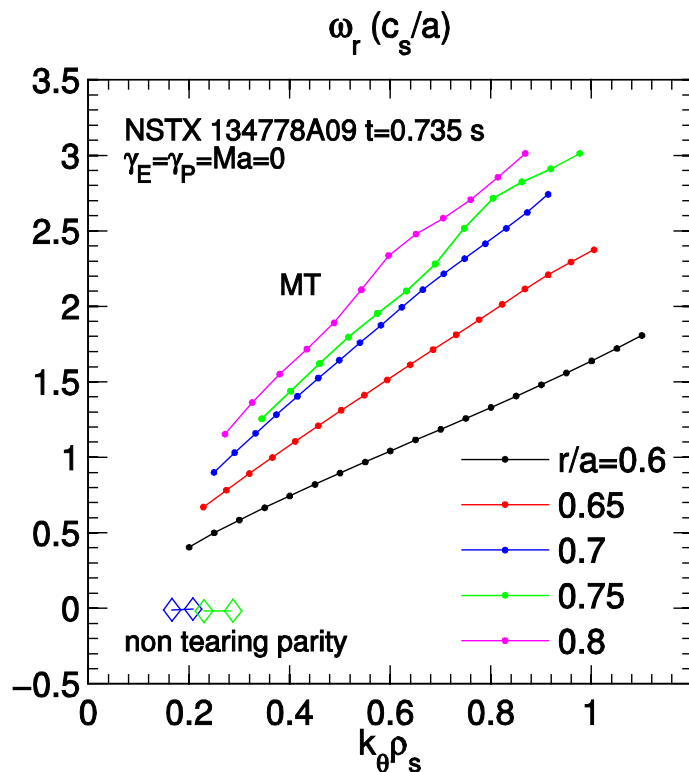
- Simulations run for 7 NBI H-modes with $n=3$ perturbations [Solomon, 2010]

$$B_T = 0.35 - 0.55 \text{ T} \quad I_p = 0.7 - 1.1 \text{ MA}$$
$$P_{\text{NBI}} = 4 - 6 \text{ MW} \quad \langle n \rangle \approx 4 - 6 \times 10^{19} \text{ m}^{-3}$$



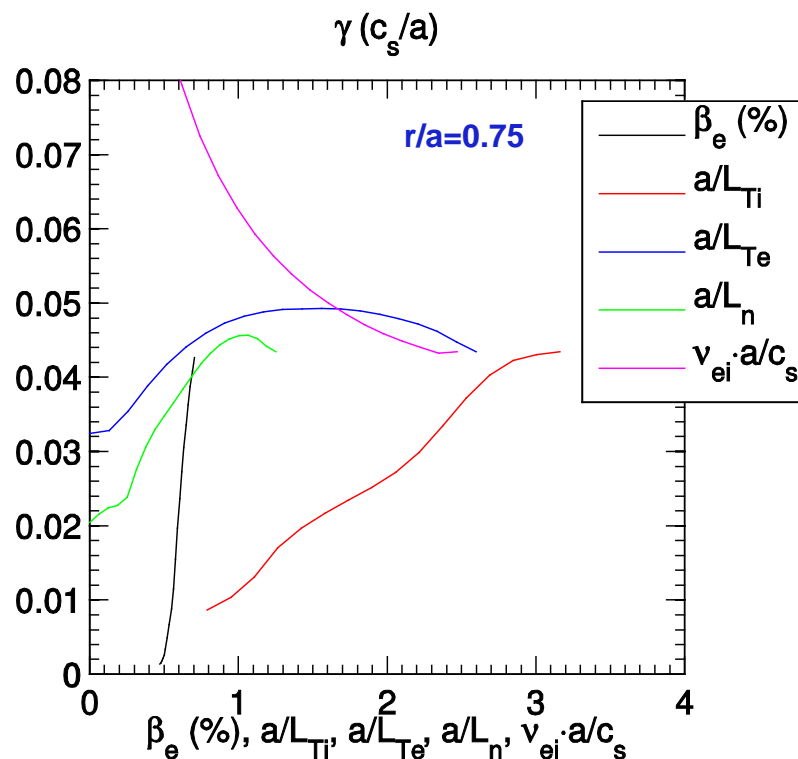
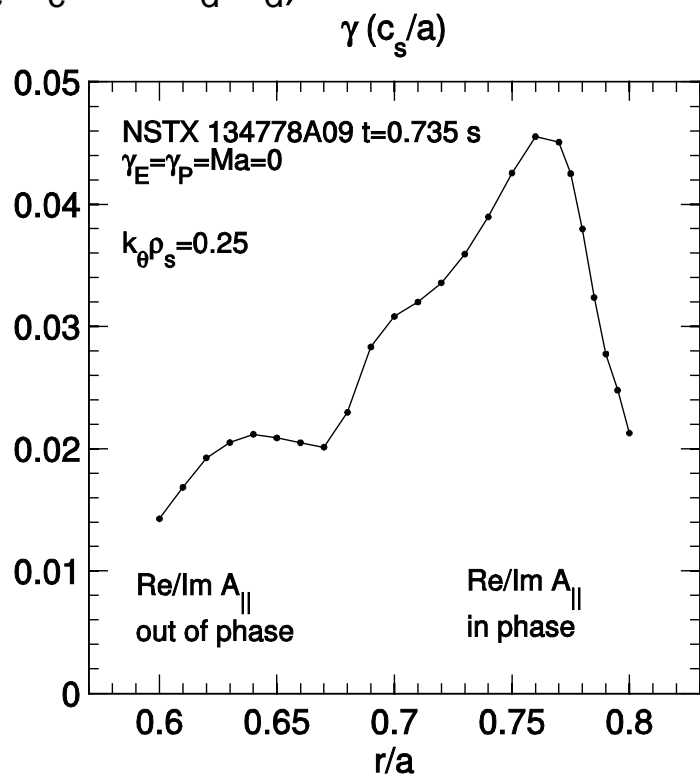
Most cases show broad spectra of microtearing modes

- Apparent in eigenfunctions (not shown) and near linear dispersion $\omega \approx \omega_{*e}$
 - Microtearing only transports electron energy
- Often see hints of subdominant ballooning modes (\diamond)
 - Unknown whether they survive nonlinearly
- $E \times B$ shearing rates comparable to γ_{lin} ($\gamma_{\text{lin}}/\gamma_E \uparrow$ as $r/a \uparrow$)



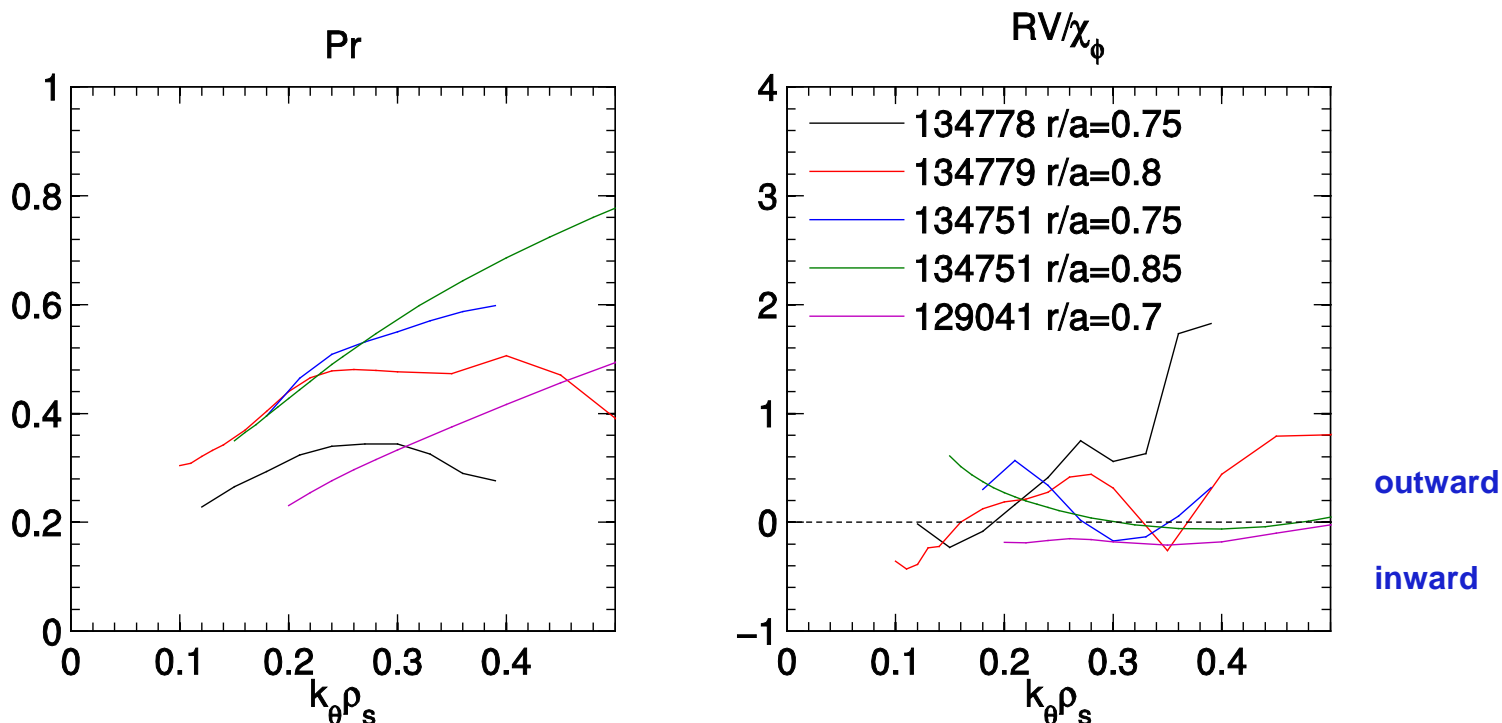
Ballooning modes exist over $0.6 \leq r/a \leq 0.8$, exhibit KBM behavior

- Very sensitive to $\beta_e \rightarrow$ KBM ($\alpha_{\text{MHD,unit}} > 0.6$)
 - Unstable from a/L_{Ti} - similar to hybrid ITG/KBM behavior found by Belli, Candy [2010]
 - Similar hybrid-KBM modes often predicted in NSTX H-modes [Guttenfelder, NF 2013; Canik, NF 2013]
- Transport contributions come from both φ and $B_{||}$; also D and C ($Z_{\text{eff}} \approx 3$, $n_c m_c \sim 0.7 n_d m_d$)



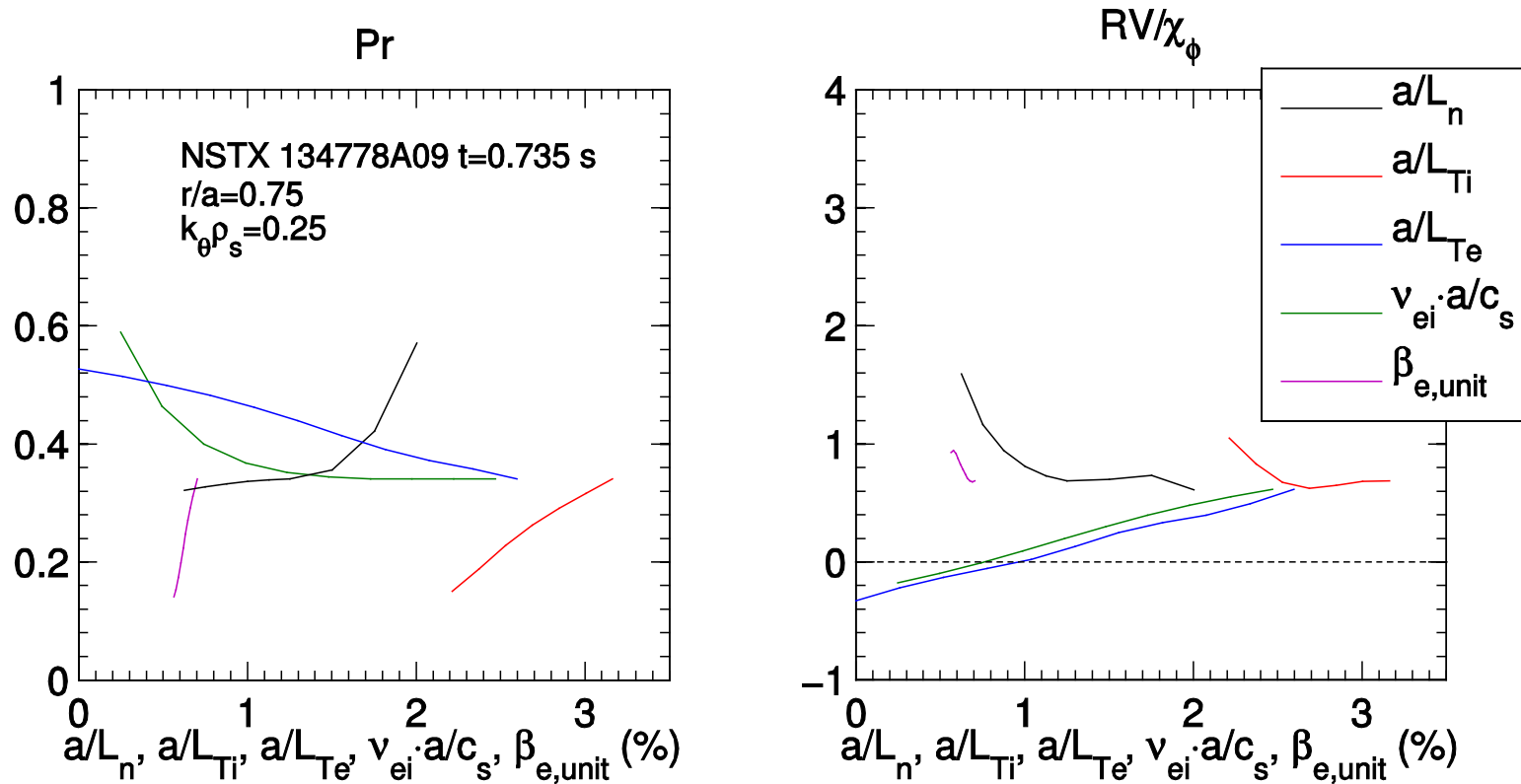
Small Prandtl numbers over KBM range of $k_{\theta}\rho_s$, small/outward Pinch parameter

- Interpreted Pr would be smaller for $\chi_{i,nc} > \chi_{i,turb}$
- Small/outward RV_{ϕ}/χ_{ϕ}
 - consistent with KBM predictions using conventional tokamak parameters [Hein, 2010]
- Small/positive RV_{ϕ}/χ_{ϕ} predicted in multiple cases, never approaches larger inward experimental values (-7)



Pinch parameter shows minor changes with parameters, always remains near zero or outwards

- Never approaches larger inward experimental values (-7)



- What else is missing?

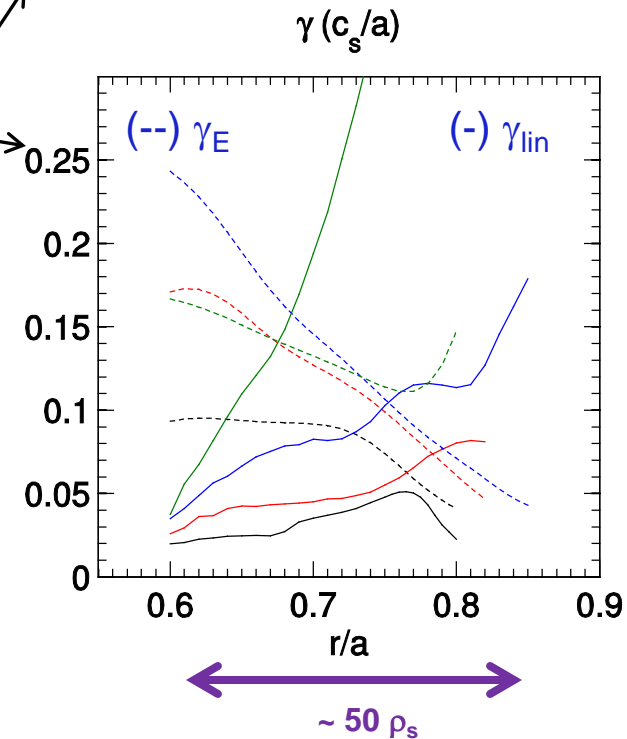
Additional considerations

- Nonlinear transport possibly different from quasilinear (simulations beginning)
- Perpendicular ($E \times B$) flow shear (Dominguez, Casson, Waltz)

$$\Pi_{\varphi} = (\chi_{\varphi} u' + \chi_{\varphi\perp} \gamma_E) + (RV_{\varphi} + R\Gamma_p)u + \Pi_{\varphi,RS}$$

- Influence of particle flux
- Finite ρ_* effects: profile shear, non-local effects, influence from pedestal
- Centrifugal effects on transport and stability
 - GKW work in progress (Buchholtz, Hornsby, Peeters)

⇒ Mechanism(s) for strong observed inward pinch remains unresolved



Impurity Transport

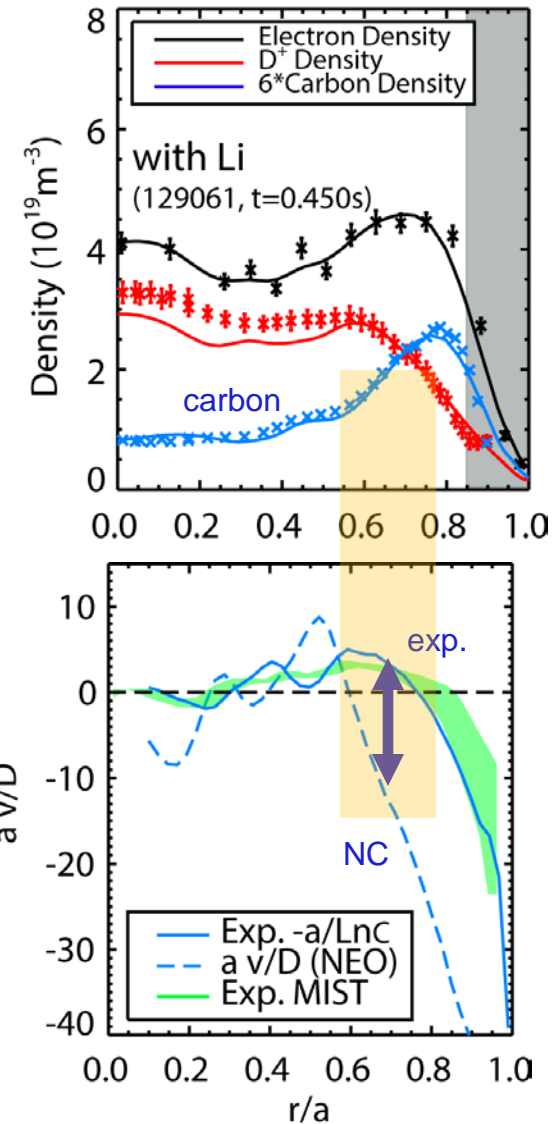
Occasional evidence for non-neoclassical impurity transport in Lithium conditioned H-modes (Scotti, NF 2013)

- Impurity transport often close to neoclassical levels in H-modes [Delgado-Aparicio, NF 2009, 2011; Clayton, PPCF 2012]
- With lithium wall conditioning, ELMs are suppressed and carbon accumulates
 - Lithium does NOT accumulate (better scrape-off layer screening + neoclassical $D_{Li} \gg D_c$)

$$\Gamma_c = -D_c \nabla n_c + V_c n_c \approx 0 \Rightarrow \boxed{\frac{aV_c}{D_c} = -\frac{a}{L_{n,c}}}$$

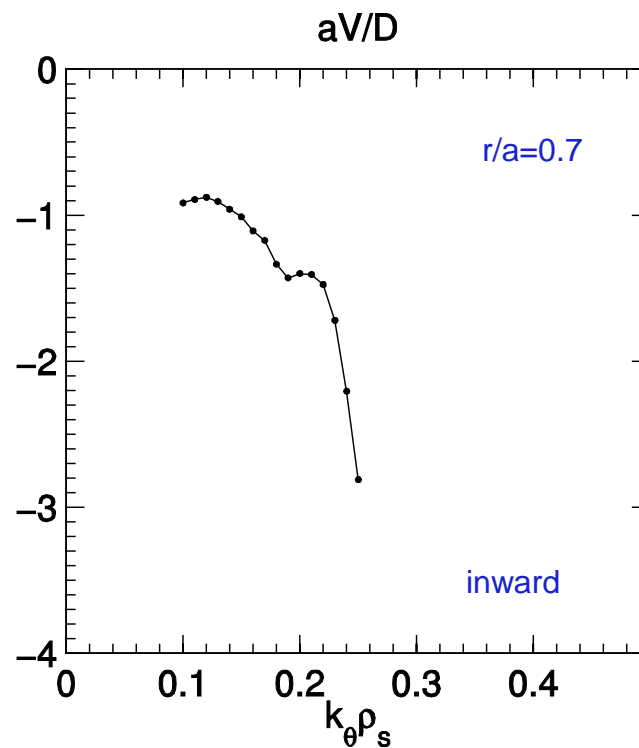
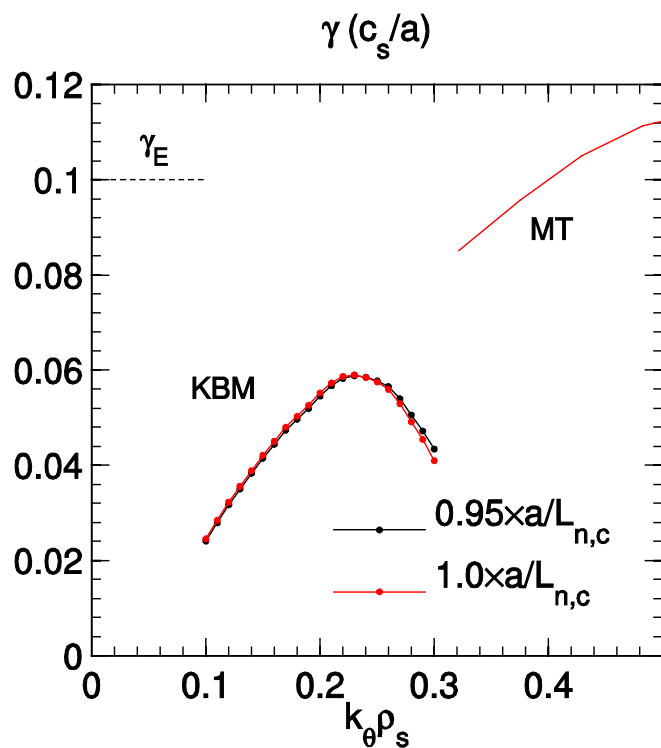
- Profile shape can diverge significantly from neoclassical theory (don't have quantitative source in these cases)

Q: Can ballooning modes influence impurity transport in NSTX H-modes?



Microtearing modes dominate, sub-dominant KBM modes predict inward carbon pinch

- Microtearing dominates (no particle flux)
- Weaker hybrid-KBM ($\gamma_{\text{KBM}} < \gamma_{\text{E}}$) – unknown if this survives nonlinearly
- KBM predicts inward carbon pinch ($r/a=0.6-0.7$)
 - Opposite to experiment, similar to neoclassical



Inward carbon pinch predicted for KBM over a range of parameters

- Direction of carbon convection insensitive to v_{ei} , a/L_{Ti} , a/L_{Te} , and u'
 - Outward carbon convection predicted as beta is reduced and mode transitions to ITG/TEM
- ⇒ Does not appear to reconcile non-neoclassical impurity profile

